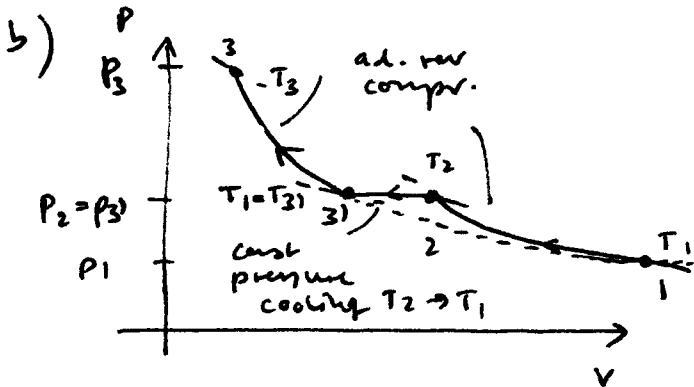


concepts:

- 1st law
- shaft and flow work
- adiabatic rev. process
- p-v diagrams

Assume: adiab rev. compressors, ideal gas, steady flow, no KE, PE



c) compressor I:



$$u_2 - u_1 = -(-p_1 v_1 - w_{sI} + p_2 v_2)$$

$$h_2 - h_1 = w_{sI} = c_p (T_2 - T_1)$$

adiab. rev: $\frac{T_2}{T_1} = \left(\frac{p_2}{p_1}\right)^{\frac{\gamma-1}{\gamma}} \rightarrow w_{sI} = c_p T_1 \left[\left(\frac{p_2}{p_1}\right)^{\frac{\gamma-1}{\gamma}} - 1\right]$

compressor II: similar by $\rightarrow w_{sII} = c_p T_3 \left[\left(\frac{p_3}{p_2}\right)^{\frac{\gamma-1}{\gamma}} - 1\right] = c_p T_1 \left[\left(\frac{p_3}{p_2}\right)^{\frac{\gamma-1}{\gamma}} - 1\right]$

so $w_s = w_{sI} + w_{sII} \rightarrow w_s = c_p T_1 \left[\left(\frac{p_2}{p_1}\right)^{\frac{\gamma-1}{\gamma}} + \left(\frac{p_3}{p_2}\right)^{\frac{\gamma-1}{\gamma}} - 2 \right]$

1) find minimum: $\frac{dw_s}{dp_2} = 0$

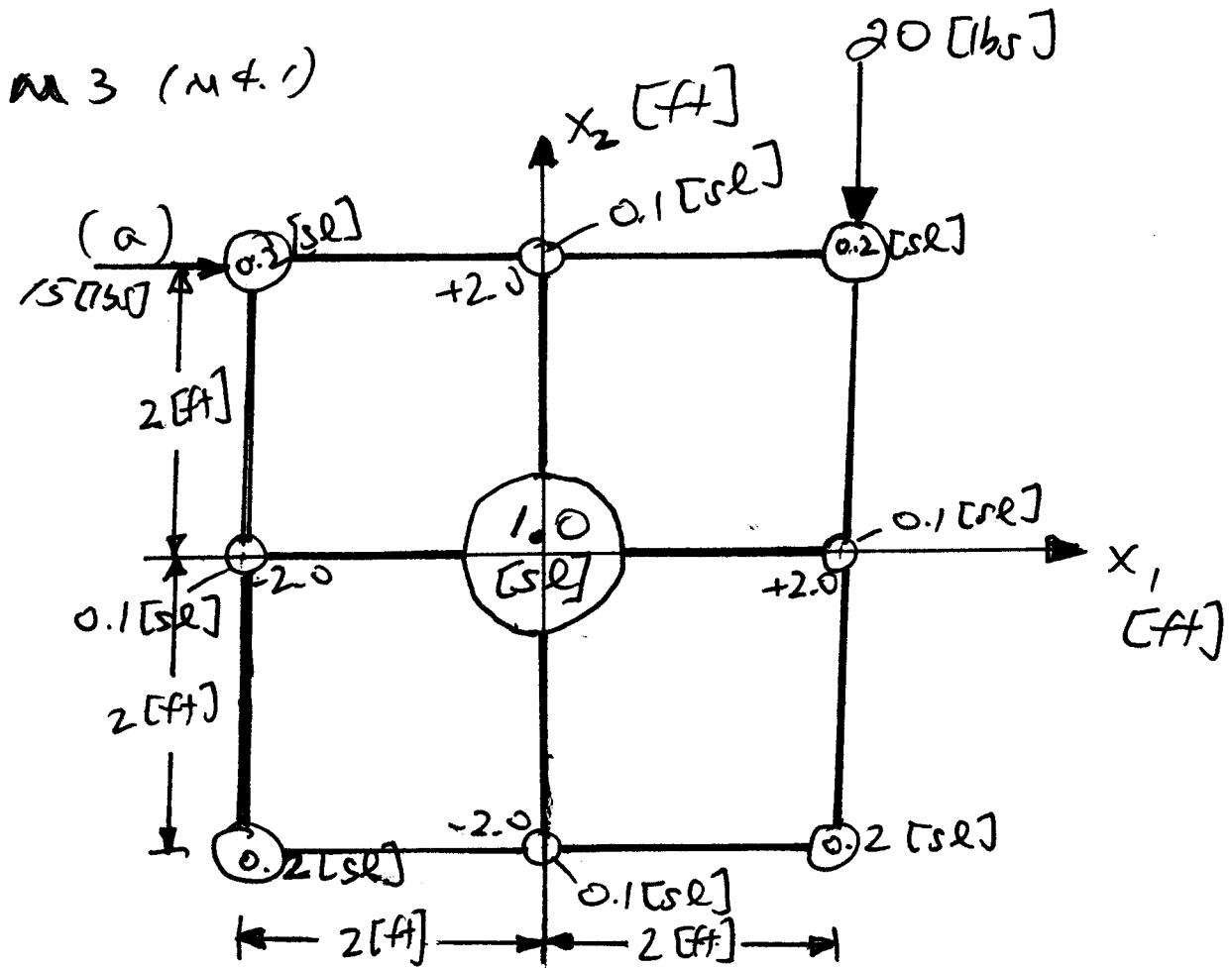
$$0 = c_p T_1 \left[\frac{\gamma-1}{\gamma} \left(\frac{p_2}{p_1}\right)^{\frac{\gamma-1}{\gamma}-1} \cdot \frac{1}{p_1} + \frac{\gamma-1}{\gamma} \left(\frac{p_3}{p_2}\right)^{\frac{\gamma-1}{\gamma}-1} \cdot \left(-\frac{p_3}{p_2^2}\right) \right]$$

$$0 = \left(\frac{p_2}{p_1}\right)^{\frac{\gamma-1}{\gamma}} \cdot \frac{1}{p_2} - \left(\frac{p_3}{p_2}\right)^{\frac{\gamma-1}{\gamma}} \cdot \frac{1}{p_2} \rightarrow \frac{p_2}{p_1} = \frac{p_3}{p_2}$$

find for minimum shaft work

$$P_{2, \min, w_s} = \sqrt{P_1 P_3}$$

Unified Engineering
 Homework SOLUTIONS
 Week #4; Problem Set #3



[sl] = [sl/ft]

— = rigid, massless rods

(b) This system is not in equilibrium. One can see that both its net force and net moment are non-zero.

Call the 20 lb force \underline{F}_1 and the 15 lb force \underline{F}_2 . Express each as vectors:

$$\underline{F}_1 = -20[\text{lb}] \hat{i}_2 \quad \text{at } (2\text{ft}, 2\text{ft})$$

$$\underline{F}_2 = +15[\text{lb}] \hat{i}_1 \quad \text{at } (-2\text{ft}, 2\text{ft})$$

Thus, the net force, $\underline{F}_{\text{net}}$, is:

$$\underline{F}_{\text{net}} = \underline{F}_1 + \underline{F}_2 = +15[\text{lb}] \hat{i}_1 + (-20)[\text{lb}] \hat{i}_2$$

The acceleration of the rigidly connected system is:

$$\underline{a} = \frac{\underline{F}_{\text{net}}}{m}$$

where m is the total mass = $(4 \times 0.1[\text{sl}])$
 $+ (4 \times 0.2[\text{sl}]) + 1.0[\text{sl}]$
 $= 2.2[\text{slugs}]$

$$\text{So: } \underline{a} = \frac{15\text{ lb}}{2.2\text{ slugs}} \hat{i}_1 - \frac{20\text{ lb}}{2.2\text{ slugs}} \hat{i}_2$$

Note the relationship of force in [lbs] to mass in [slugs]

Thus: $\underline{a} = 6.82 \text{ [ft/s}^2\text{]} \hat{i}_1 - 9.09 \text{ [ft/s}^2\text{]} \hat{i}_2$

So the system has a linear acceleration of $6.82 \text{ [ft/s}^2\text{]}$ in the $+x_1$ -direction and $9.09 \text{ [ft/s}^2\text{]}$ in the $-x_2$ -direction

There will also be a clockwise rotation about the origin

This is due to the net moment:

$$\underline{M}_{\text{net}} = \sum (\underline{r}_i \times \underline{F}_i)$$

where \underline{r}_i = position vector to force vector from origin

One can find:

$$\underline{M}_{\text{net}} = (2 \text{ ft}) \hat{i}_1 \times (-20 \text{ lbs}) \hat{i}_2 + (2 \text{ ft}) \hat{i}_2 \times (+15 \text{ lbs}) \hat{i}_1$$

$$\text{Also noting: } \hat{i}_1 \times \hat{i}_2 = +\hat{i}_3$$

$$\hat{i}_2 \times \hat{i}_1 = -\hat{i}_3$$

$$\Rightarrow \underline{M}_{\text{net}} = (-40 \text{ [ft}\cdot\text{lbs]})\hat{i}_3 + (30 \text{ [ft}\cdot\text{lbs]})(-\hat{i}_3)$$

$$= -70 \text{ [ft}\cdot\text{lbs] } \hat{i}_3$$

- \Rightarrow clockwise

(c) By applying other forces and moments, one can achieve equilibrium if the net force and moment are zero.

Thus, by applying a force, one can get the net force to be zero. The applied force must be equal in magnitude and opposite in direction to the current net force (or it exists)

$$\Rightarrow \text{apply } \underline{F}_{\text{applied}} = (-15 \text{ [lb]})\hat{i}_1 + 20 \text{ [lb]}\hat{i}_2$$

and the linear acceleration will be zero.

BUT, by applying a force at the origin one cannot cause a moment about the origin

as the position vector is: $\underline{r} = \underline{0}$. Thus the net moment cannot become zero.

\Rightarrow NO

equilibrium
cannot be
achieved

(d) The opposite occurs in this case. By applying a moment about the origin, one can cause the net moment to be zero (by applying a moment that is equal in magnitude and opposite in direction)

$$\Rightarrow \underline{M}_{\text{applied}} = +70 \text{ [ft}\cdot\text{lb]} \hat{i}_3$$

however, one cannot get the net force to be zero and thus there will be linear acceleration and motion

\Rightarrow NO

equilibrium
cannot be
achieved

(e) Go back to the reasoning of parts (c) and (d). Add these two (superposition applies), and equilibrium can be

achieved.

$$\Rightarrow \boxed{YES}$$

Apply:

$$\begin{aligned} \underline{F}_{\text{applied}} &= (-15 [15]) \hat{i}_1 + (20 [15]) \hat{i}_2 \\ \underline{M}_{\text{applied}} &= +70 [15] \hat{i}_3 \end{aligned}$$

(f) The same reasoning applies as in the answer to part (d). A couple results in a net moment but no net force. So a couple can cause the rotation not to occur, but the total net force will still be non-zero and thus there will be linear acceleration and motion.

$$\Rightarrow \boxed{NO}$$

equilibrium cannot be achieved

(g) A force can achieve a desired moment by proper placement. Thus, using a similar reasoning as in part (e), a force of the proper magnitude and direction can be properly placed to give the needed opposite

moment about the origin.

→ The force must be as in part (c) to achieve a net force of zero:

$$\underline{F}_{\text{applied}} = (-15 \text{ [lb]})\hat{i}_1 + (20 \text{ [lb]})\hat{i}_2$$

→ One must have a net moment, or per part (d), of:

$$\underline{M}_{\text{applied}} = +70 \text{ [ft. lbs]} \hat{i}_3$$

This is achieved via:

$$\underline{M}_{\text{applied}} = \underline{r}_{\text{applied}} \times \underline{F}_{\text{applied}}$$

where $\underline{r}_{\text{applied}}$ is the position vector to the point (x_1, x_2) where $\underline{F}_{\text{applied}}$ is applied.

Write this as:

$$\underline{r}_{\text{applied}} = x_1 \hat{i}_1 + x_2 \hat{i}_2$$

dot this gives via the moment expression:

$$+70 \text{ [ft. lbs]} \hat{i}_3 = (x_1 \hat{i}_1 + x_2 \hat{i}_2) \times \left\{ (-15 \text{ [lb]})\hat{i}_1 + (20 \text{ [lb]})\hat{i}_2 \right\}$$

$$\Rightarrow (-15 \times 2 [15s]) (-\hat{i}_3) + (20 \times 1 [15s]) (\hat{i}_3)$$

$$= +70 [15s] \hat{i}_3$$

$$\Rightarrow \frac{3}{14} x_2 + \frac{2}{7} x_1 = 1 [ft]$$

So this defines a line along which

$\underline{F} = (-15 [15s]) \hat{i}_1 + (20 [15s]) \hat{i}_2$
 can be placed to get equilibrium

So:

YES

equilibrium
 can be achieved

(NOTE: There are 4
 points of interaction
 with the world of the
 system)

Solutions F4

1. From the continuity equation: $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0$

we have $\boxed{\nabla \cdot \vec{v} = 0}$ for a steady, incompressible flow.

The velocity vector is $\vec{v} = u \hat{i} + v \hat{j}$

So the condition $\nabla \cdot \vec{v} = 0$ reduces to $\frac{\partial u}{\partial x} = -\frac{\partial v}{\partial y}$

If I know $u = u(x) \Rightarrow \boxed{v = - \int \frac{\partial u}{\partial x} dy}$

u comes from our measurements: $u = v_0 \left(1 + \left(\frac{x}{x_0}\right)^2\right) e^{-y/y_0}$

then: $\frac{\partial u}{\partial x} = \frac{2v_0 x}{x_0^2} e^{-y/y_0}$

Integrate in $y \Rightarrow v = - \int \frac{2v_0 x}{x_0^2} e^{-y/y_0} dy = + \frac{2v_0 x y_0}{x_0^2} e^{-y/y_0} + f(x) + A$

\swarrow function of x
 \nwarrow constant

Evaluate at $y=0 \Rightarrow v_{y=0}(x) = \frac{2v_0 x y_0}{x_0^2} + f(x) + A$

therefore $f(x) + A = v_{y=0}(x) - \frac{2v_0 x y_0}{x_0^2}$

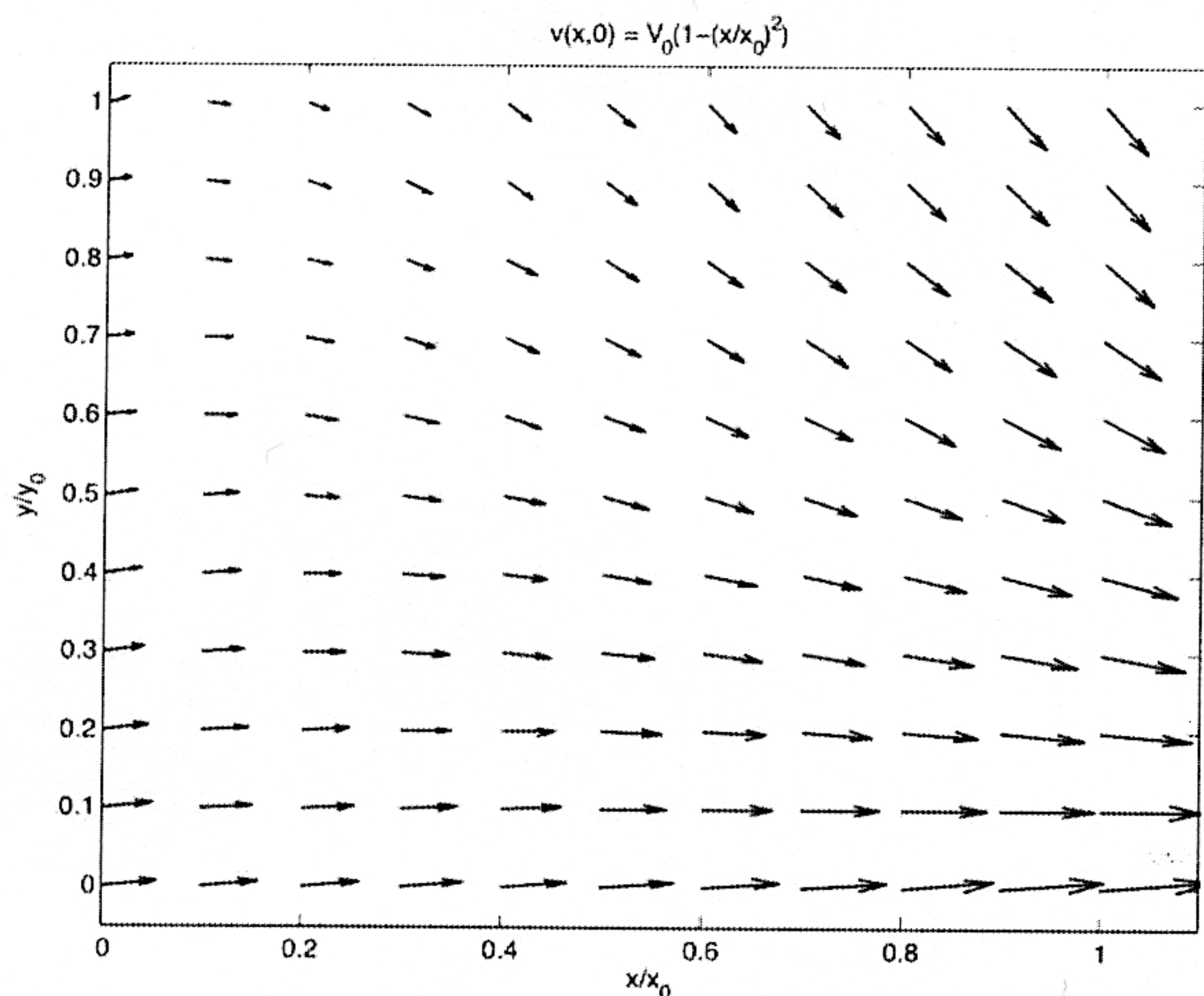
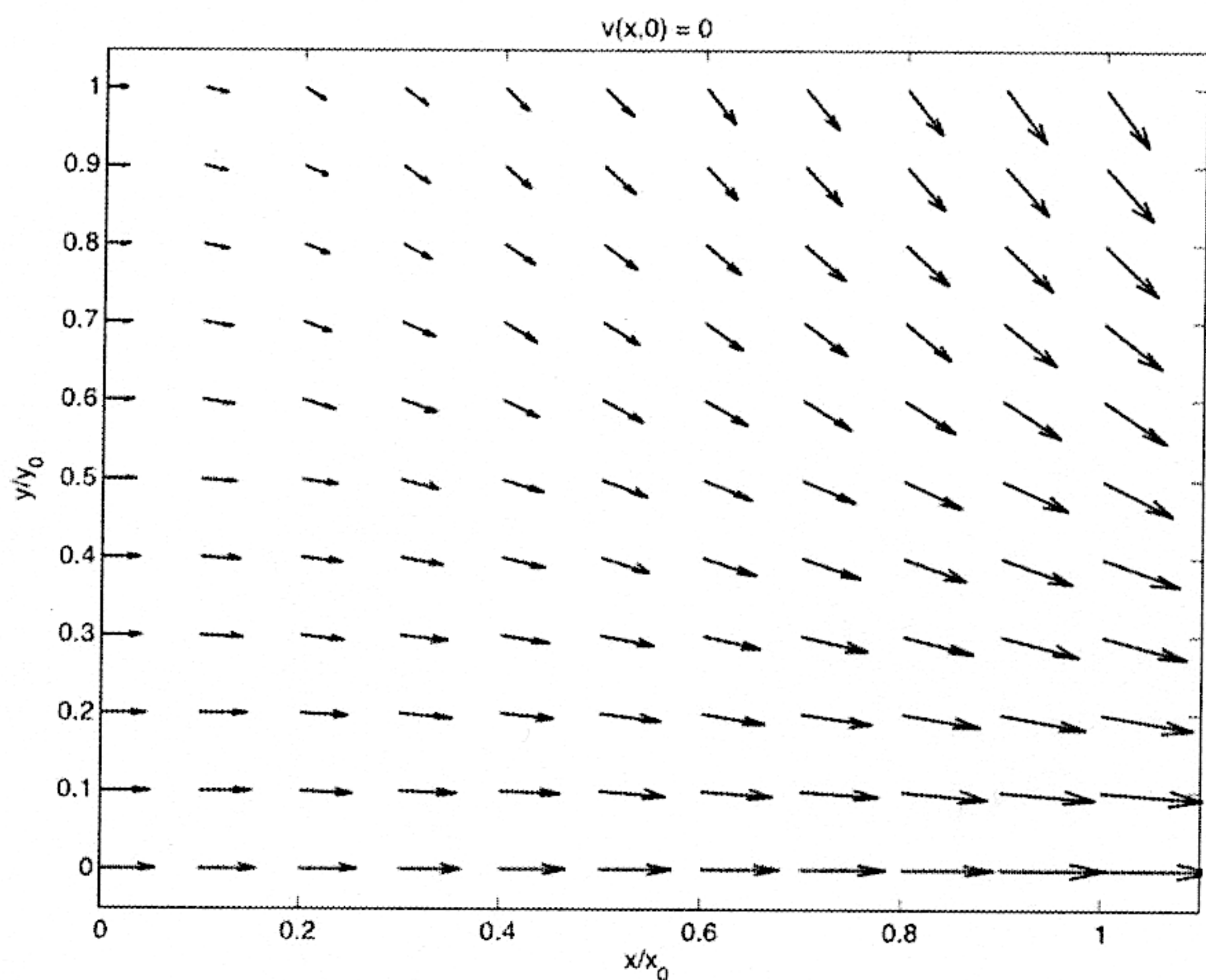
the velocity vector is then:

$$\frac{\vec{v}}{v_0} = \underbrace{\left[1 + \left(\frac{x}{x_0}\right)^2\right] e^{-y/y_0}}_{f_x} \hat{i} + \underbrace{\left\{ \frac{2xy_0}{x_0^2} (e^{-y/y_0} - 1) + \frac{v_{y=0}(x)}{v_0} \right\}}_{f_y} \hat{j}$$

We have a free parameter, the function $v_{y=0}(x)$, and we don't have information about it. To perform the plots, we need to make an

assumption. We will do 2 limit cases: $v_{y=0}(x) = v_0 \left(1 + \left(\frac{x}{x_0}\right)^2\right) \leftarrow$ same as u

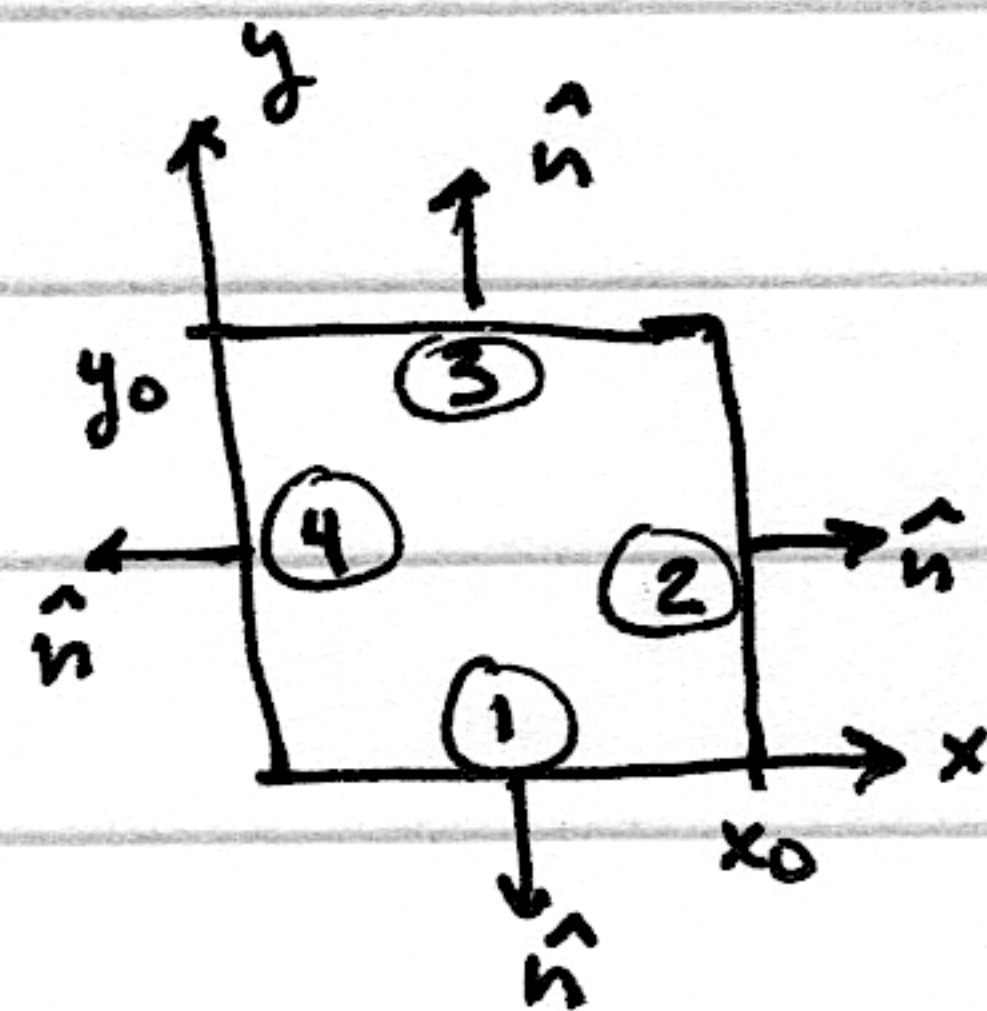
and $v_{y=0}(x) = 0$



The general shape of the field is very similar in both cases

3. Evaluate the flow integral: $\int \rho \vec{v} \cdot \hat{n} dA$ over the 4 faces
p. unit depth

const $\rightarrow \rho V_0 \int \vec{v} \cdot \hat{n} dA$



Face 1: $\hat{n} = -\hat{j}$ then $-\int_0^{x_0} \frac{v_{y=0}(x)}{V_0} dx$
($y=0$)

Face 2: $\hat{n} = \hat{i}$ then $\int_0^{y_0} 2e^{-y/y_0} dy = -2y_0(e^{-1} - 1)$
($x=x_0$)

Face 3: $\hat{n} = \hat{j}$ then $\int_0^{x_0} \left[\frac{2xy_0}{x_0^2}(e^{-1} - 1) + \frac{v_{y=0}(x)}{V_0} \right] dx = y_0(e^{-1} - 1) + \int_0^{x_0} \frac{v_{y=0}(x)}{V_0} dx$
($y=y_0$)

Face 4: $\hat{n} = -\hat{i}$ then $-\int_0^{y_0} e^{-y/y_0} dy = y_0(e^{-1} - 1)$
($x=0$)

Adding all together: $-\int_0^{x_0} \frac{v_{y=0}(0)}{V_0} dx + \int_0^{x_0} \frac{v_{y=0}(x)}{V_0} dx - 2y_0(e^{-1} - 1) + 2y_0(e^{-1} - 1) = 0$

as it should be.

$$\nabla \cdot \vec{V} = 0 \iff \int \vec{V} \cdot \hat{n} dA$$

it is consistent.