

Alternative Objectives

Frequently, an engineering design problem will involve competing objectives. For example, an alternative objective to minimize flight power (or maximize payload), is to maximize the maximum-attainable flight speed. The level-flight thrust relations still apply to this flight condition, but now the thrust is a known quantity, and equal to the maximum available thrust from the powerplant.

$$T_{\max} = \frac{1}{2} \rho V_{\max}^2 S \left[\frac{CDA_0}{S} + c_d(C_{L_{\min}}; Re_{\max}) + \frac{C_{L_{\min}}^2}{\pi e AR} \right] \quad (1)$$

where $C_{L_{\min}} = \frac{2W/S}{\rho V_{\max}^2}$

This can be solved for V_{\max} , or equivalently $C_{L_{\min}}$, by numerical means if necessary. If one can assume that at V_{\max} the induced drag is negligible, and c_d is some constant, then we have

$$V_{\max}(AR, S) \simeq \left[\frac{2T_{\max}}{\rho(CDA_0 + S c_d)} \right]^{1/2} \quad (2)$$

which is a suitable alternative objective function.

The figure shows the isolines of V_{\max} versus the design variables, calculated using (1). The slight upturn in the isolines for decreasing AR is the Re effect in c_d . The sharp downturn near $AR = 0$ is due to the induced drag term. Isolines computed using the approximation (2) would be level.

Now the stiffness-constrained optimum results in an unreasonably small wing, which means that other more practical constraints are likely to come into play.

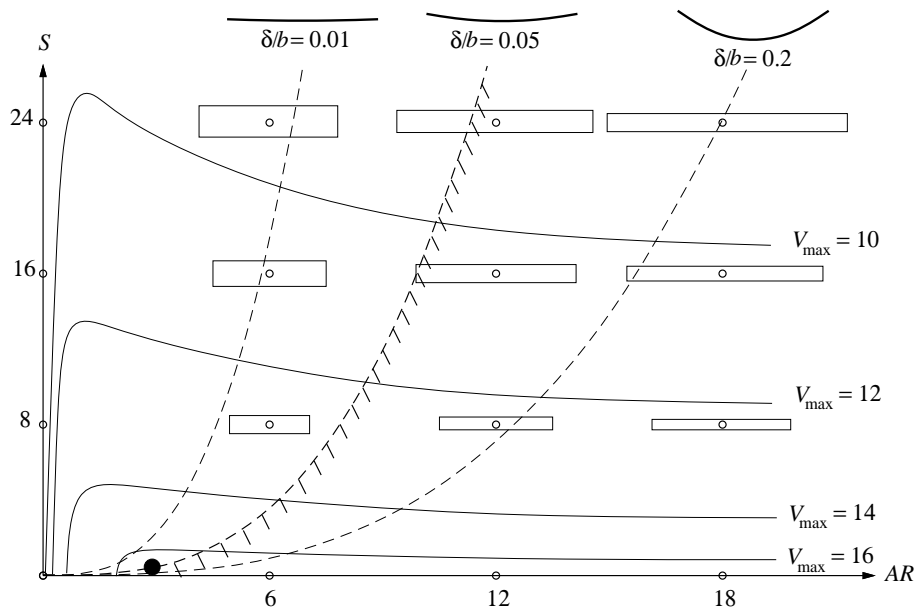


Figure 1: Objective function contours (isolines) in design space of a rectangular wing. Dot shows the stiffness-constrained optimum point.

Iterative Calculation of V_{\max}

Although equation (1) cannot be explicitly solved for V_{\max} , it is possible to solve it by a reasonably simple iterative procedure.

We first use equation (2) to give a first approximation \tilde{V}_{\max} .

$$\tilde{V}_{\max}(AR, S) = \left[\frac{2T_{\max}}{\rho(CDA_0 + S c_d)} \right]^{1/2} \quad (3)$$

This is then used to compute an estimated \tilde{C}_L and \tilde{Re} .

$$\tilde{C}_L(AR, S) = \frac{2W/S}{\rho \tilde{V}_{\max}^2} \quad (4)$$

$$\tilde{Re}(AR, S) = \frac{\tilde{V}_{\max} c}{\nu} \quad (5)$$

These then allow an improved estimation of C_{D_i} , or c_d , or both.

$$\tilde{C}_{D_i}(AR, S) = \frac{\tilde{C}_L^2}{\pi e AR} \quad (6)$$

$$\tilde{c}_d(AR, S) = c_d(\tilde{C}_L, \tilde{Re}) \quad (7)$$

Note that this new C_{D_i} is much better than the first estimate, which merely assumed $C_{D_i} = 0$. The new c_d may or may not be better than the initial estimate of assuming c_d was some constant. In any case, the new \tilde{C}_{D_i} and/or new \tilde{c}_d are then used in equation (1) to get an improved V_{\max} estimate.

$$V_{\max}(AR, S) = \left[\frac{2T_{\max}}{\rho(CDA_0 + S \tilde{c}_d + S \tilde{C}_{D_i})} \right]^{1/2} \quad (8)$$

The iteration can be repeated by starting again at (4), although the additional accuracy gain is likely to be minimal.

The V_{\max} calculated by this procedure will fully include the C_{D_i} contribution to the total drag, and will also include Reynolds number effects if the $c_d(C_L, Re)$ function is used. The resulting V_{\max} isolines should resemble those in Figure 1.