
 Problem Set No. 4

 Problem 1

$$[I]_G = \begin{bmatrix} 9 & 4 & 0 \\ & 3 & 0 \\ & & 5 \end{bmatrix} \quad G: \text{origin of the coordinate system}$$

Principal moments of inertia are the eigenvalues of $[I]_G$:

To find the eigenvalues I of $[I]_G$,

$$\det \begin{bmatrix} 9-I & 4 & 0 \\ & 3-I & 0 \\ & & 5-I \end{bmatrix} = 0 \quad \Rightarrow \quad (9-I)(3-I)(5-I) - 16(5-I) = 0$$

$$\Rightarrow \quad \underline{I_1 = 11}, \quad \underline{I_2 = 5}, \quad \underline{I_3 = 1} \quad \text{principal moments of inertia}$$

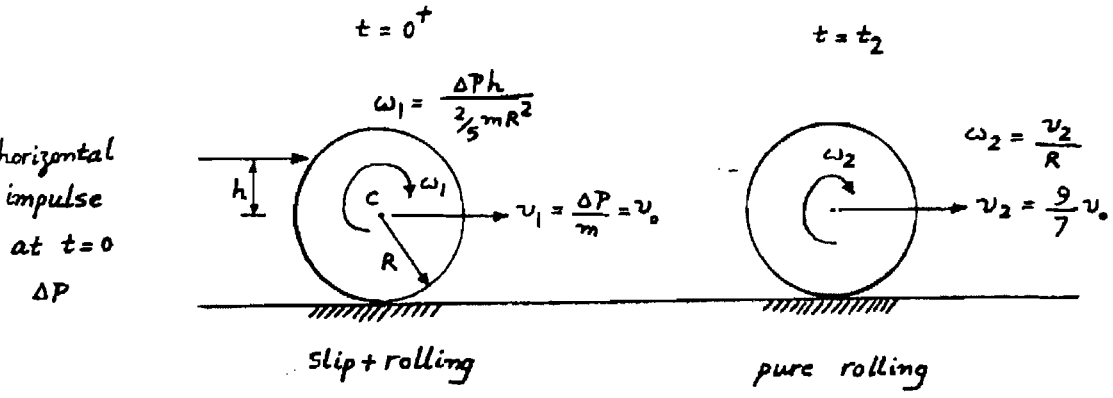
To find direction of principal axis,

$$\text{for } I_1 = 11, \quad \begin{bmatrix} 9-11 & 4 & 0 \\ & 3-11 & 0 \\ & & 5-11 \end{bmatrix} \begin{Bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{Bmatrix}_1 = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \Rightarrow \begin{matrix} \omega_x = 2\omega_y \\ \omega_z = 0 \end{matrix} \Rightarrow \underline{\underline{\omega_1 = \begin{Bmatrix} \frac{2}{\sqrt{5}} \\ \frac{1}{\sqrt{5}} \\ 0 \end{Bmatrix}}}}$$

$$\text{for } I_2 = 5, \quad \begin{bmatrix} 9-5 & 4 & 0 \\ & 3-5 & 0 \\ & & 5-5 \end{bmatrix} \begin{Bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{Bmatrix}_2 = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \Rightarrow \begin{matrix} \omega_x = \omega_y = 0 \\ \omega_z \neq 0 \end{matrix} \Rightarrow \underline{\underline{\omega_2 = \begin{Bmatrix} 0 \\ 0 \\ 1 \end{Bmatrix}}}}$$

$$\text{for } I_3 = 1, \quad \begin{bmatrix} 9-1 & 4 & 0 \\ & 3-1 & 0 \\ & & 5-1 \end{bmatrix} \begin{Bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{Bmatrix}_3 = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \Rightarrow \begin{matrix} \omega_y = -2\omega_x \\ \omega_z = 0 \end{matrix} \Rightarrow \underline{\underline{\omega_3 = \begin{Bmatrix} \frac{1}{\sqrt{5}} \\ -\frac{2}{\sqrt{5}} \\ 0 \end{Bmatrix}}}}$$

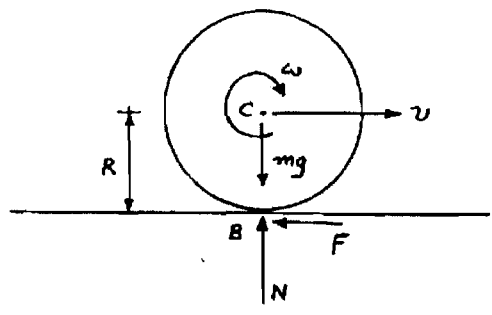
Problem 2



horizontal impulse at $t=0$ ΔP , $\Delta P = \text{change in linear momentum} = mv_1 = mv_0$

$$\Delta P \cdot h = \text{change in angular momentum about } C = I\omega_1 = \frac{2}{5}mR^2\omega_1$$

Apply angular momentum principle about (moving) contact point B:



$$\underline{\tau}_B = \frac{d}{dt} \underline{H}_B + \underline{v}_B \times \underline{P}$$

$$\left. \begin{array}{l} \underline{P} = m\underline{v} \\ \underline{v}_B \parallel \underline{v} \end{array} \right\} \Rightarrow \underline{v}_B \times \underline{P} = \underline{0}$$

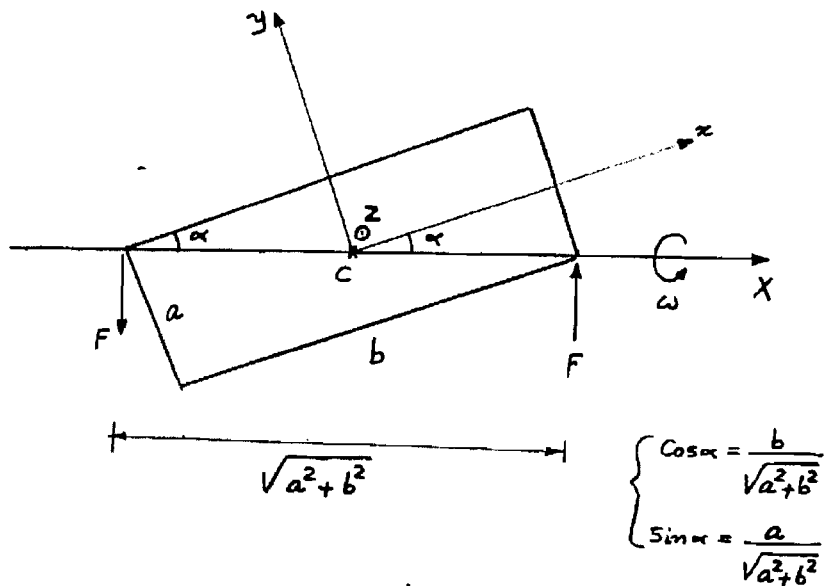
$$\underline{\tau}_B = 0$$

$\therefore \frac{d}{dt} \underline{H}_B = 0 \rightarrow \underline{H}_B$ is conserved. $\rightarrow \underline{H}_B|_{t=0^+} = \underline{H}_B|_{t=t_2}$

$$\underline{H}_B = \underline{H}_C + \underline{b}_C \times \underline{P} = \frac{2}{5}mR^2\omega + mRv$$

$$\left. \begin{array}{l} t=0^+, \quad \underline{H}_B|_{t=0^+} = \frac{2}{5}mR^2\omega_1 + mRv_1 = \Delta P \cdot h + mRv_0 = mv_0(h+R) \\ t=t_2, \quad \underline{H}_B|_{t=t_2} = \frac{2}{5}mR^2\omega_2 + mRv_2 = \left(\frac{2}{5}mR + mR\right)\frac{9}{7}v_0 = \frac{9}{5}v_0mR \end{array} \right\} \Rightarrow h+R = \frac{9}{5}R \Rightarrow h = \frac{4}{5}R$$

Introducing xyz coordinate system which is fixed to the plate and rotates with ω about X axis:

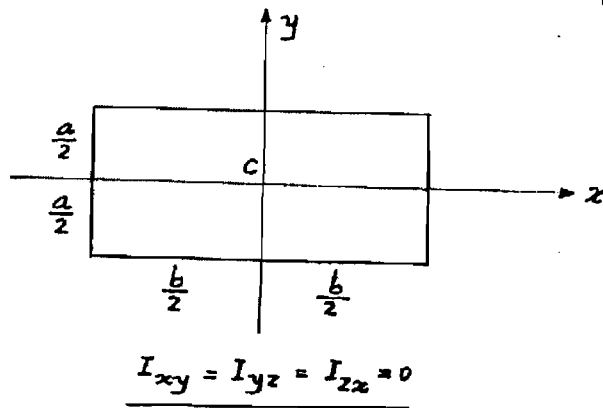


(a) Since $\underline{v}_C = 0$, forces on the bearings are equal and in opposite directions.

$$\begin{cases} \cos \alpha = \frac{b}{\sqrt{a^2 + b^2}} \\ \sin \alpha = \frac{a}{\sqrt{a^2 + b^2}} \end{cases}$$

$$I_x = \int (\rho dV) (y^2 + z^2) \quad z \approx 0 \text{ (thin plate)}$$

$$= \int_{-\frac{a}{2}}^{\frac{a}{2}} \int_{-\frac{b}{2}}^{\frac{b}{2}} \rho dx dy (y^2) = \rho b \frac{a^3}{12} = \frac{1}{12} M a^2$$



$$I_{xy} = I_{yz} = I_{zx} = 0$$

$$I_y = \frac{1}{12} M b^2 \quad I_z = \frac{1}{12} M (a^2 + b^2)$$

$$\underline{\omega}_{plate} = \omega \hat{e}_x = \omega \cos \alpha \hat{e}_x - \omega \sin \alpha \hat{e}_y \quad \rightarrow \quad \begin{cases} \omega_x = \omega \cos \alpha \\ \omega_y = -\omega \sin \alpha \\ \omega_z = 0 \end{cases}$$

$$\underline{H}_C = [I]_C \underline{\omega} \quad \Rightarrow \quad \underline{H}_C = I_x \omega_x \hat{e}_x + I_y \omega_y \hat{e}_y$$

$$\underline{\tau}_C = \frac{d\underline{H}_C}{dt}$$

$$\frac{d\underline{H}_C}{dt} = I_x \omega_x \frac{d\hat{e}_x}{dt} + I_y \omega_y \frac{d\hat{e}_y}{dt} \quad \begin{matrix} \omega_x \hat{e}_x = \omega \sin \alpha \hat{e}_z \\ \omega_x \hat{e}_y = \omega \cos \alpha \hat{e}_z \end{matrix}$$

$$= \left(\frac{1}{12} M a^2 \omega^2 \cos \alpha \sin \alpha - \frac{1}{12} M b^2 \omega^2 \sin \alpha \cos \alpha \right) \hat{e}_z = \frac{1}{12} M \omega^2 \sin \alpha \cos \alpha (a^2 - b^2) \hat{e}_z$$

$$\underline{\tau}_C = F \sqrt{a^2 + b^2} \hat{e}_z$$

Problem 3

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$$\therefore F\sqrt{a^2+b^2} = \frac{1}{12} M\omega^2 \frac{ab}{a^2+b^2} (a^2-b^2)$$

$$\Rightarrow F = \frac{1}{12} M\omega^2 ab \frac{a^2-b^2}{(a^2+b^2)^{3/2}}$$

Note that force F rotates about X axis and is always in xy plane.

(b)

$$KE = \frac{1}{2} \{\omega\}^t [I]_c \{\omega\} + \frac{1}{2} M \cancel{v_c \cdot v_c} \quad \begin{array}{l} \nearrow \\ 0 \end{array}$$

$$= \frac{1}{2} (I_x \omega_x^2 + I_y \omega_y^2)$$

$$= \frac{1}{2} \left(\frac{1}{12} M a^2 \omega^2 \cos^2 \alpha + \frac{1}{12} M b^2 \omega^2 \sin^2 \alpha \right)$$

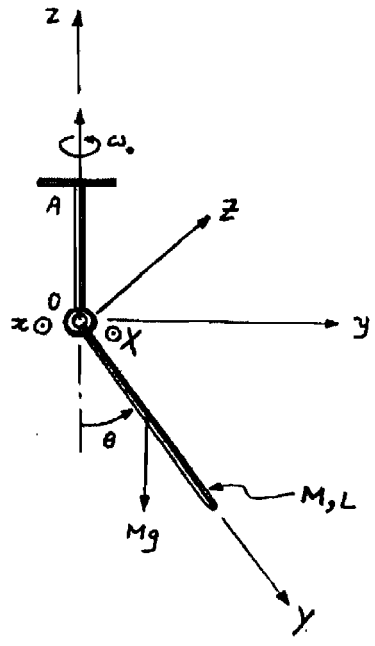
$$= \frac{1}{12} M \omega^2 \frac{a^2 b^2}{a^2 + b^2}$$

Kinetic energy of the rotating plate

Problem 4

xyz coordinate system rotates about z with ω_z so that the rod is always in yz plane.

XYZ coordinate system is fixed to the rod.



(a)

Point O is stationary, so we have

$$\underline{H}_O = [I]_O \underline{\omega}_{rod} \quad , \quad \underline{\tau}_O = \frac{d}{dt} \underline{H}_O$$

$$\underline{\omega}_{rod} = \omega_z \hat{e}_z + \dot{\theta} \hat{e}_x = \omega_z (-\cos\theta \hat{e}_y + \sin\theta \hat{e}_z) + \dot{\theta} \hat{e}_x$$

To find $[I]_O$ for the rod,

$$\begin{cases} I_X = \int \rho dV (y^2 + z^2) = \int_0^L \frac{M}{L} dy y^2 = M \frac{L^2}{3} \\ I_Y = \int \rho dV (x^2 + z^2) = 0 \\ I_Z = \int \rho dV (y^2 + x^2) = \int_0^L \frac{M}{L} dy y^2 = M \frac{L^2}{3} \\ I_{XY} = I_{YZ} = I_{ZX} = 0 \end{cases} \quad (x=0, z=0)$$

$$\therefore [I]_O = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix} \frac{ML^2}{3}$$

$$\underline{H}_O = [I]_O \begin{Bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{Bmatrix} = \frac{ML^2}{3} \dot{\theta} \hat{e}_x + \frac{ML^2}{3} \omega_z \sin\theta \hat{e}_z = \frac{ML^2}{3} \left[\dot{\theta} \hat{e}_x + \omega_z \sin\theta (\cos\theta \hat{e}_y + \sin\theta \hat{e}_z) \right]$$

$$\frac{d}{dt} \underline{H}_O = \frac{ML^2}{3} \left[\ddot{\theta} \hat{e}_x + \dot{\theta} \frac{d\hat{e}_x}{dt} + \omega_z \dot{\theta} \cos 2\theta \hat{e}_y + \omega_z \frac{\sin 2\theta}{2} \frac{d\hat{e}_y}{dt} + \omega_z \dot{\theta} \sin 2\theta \hat{e}_z \right]$$

$\omega_z \hat{e}_z \times \hat{e}_x = \omega_z \hat{e}_y$
 $\omega_z \hat{e}_z \times \hat{e}_y = -\omega_z \hat{e}_x$

$$\underline{\tau}_0 = -Mg \frac{L}{2} \sin \theta \hat{e}_x + M_y \hat{e}_y + M_z \hat{e}_z$$

$$\underline{\tau}_0 = \frac{d}{dt} \underline{H}_0$$

$$\therefore -Mg \frac{L}{2} \sin \theta = \frac{ML^2}{3} (\ddot{\theta} - \omega_0^2 \sin \theta \cos \theta)$$

$$\Rightarrow \quad \underline{L\ddot{\theta} + \left(\frac{3g}{2} - L\omega_0^2 \cos \theta\right) \sin \theta = 0} \quad \text{equation of motion for } \theta(t)$$

(b)

For stationary angle θ , $\dot{\theta} = \ddot{\theta} = 0$

$$\ddot{\theta} = 0 \quad \Rightarrow \quad \left(\frac{3g}{2} - L\omega_0^2 \cos \theta\right) \sin \theta = 0 \quad \Rightarrow \quad \begin{cases} \theta_0 = 0 \\ \theta_0 = \cos^{-1}\left(\frac{3g}{2L\omega_0^2}\right), \quad \frac{3g}{2L\omega_0^2} < 1 \end{cases}$$

$$\text{Stationary angles } \theta_0 : \quad \begin{cases} \theta_0 = 0, & \omega_0^2 < \frac{3g}{2L} \\ \theta_0 = \cos^{-1}\left(\frac{3g}{2L\omega_0^2}\right) \quad \& \quad \theta_0 = 0, & \omega_0^2 > \frac{3g}{2L} \end{cases}$$

It can be shown that $\theta_0 = 0$ is unstable when $\omega_0^2 > \frac{3g}{2L}$.

Problem 5

xyz rotates with Ω about Z axis.

$$\begin{aligned} \underline{\omega}_{\text{Cone}} &= -\Omega \cot \alpha \hat{e}_{OA} \\ &= -\Omega \cot \alpha (-\cos \alpha \hat{e}_z + \sin \alpha \hat{e}_y) \\ &= -\Omega \cos \alpha \hat{e}_y + \Omega \frac{\cos^2 \alpha}{\sin \alpha} \hat{e}_z \end{aligned}$$

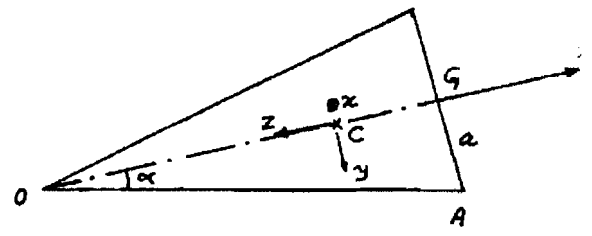
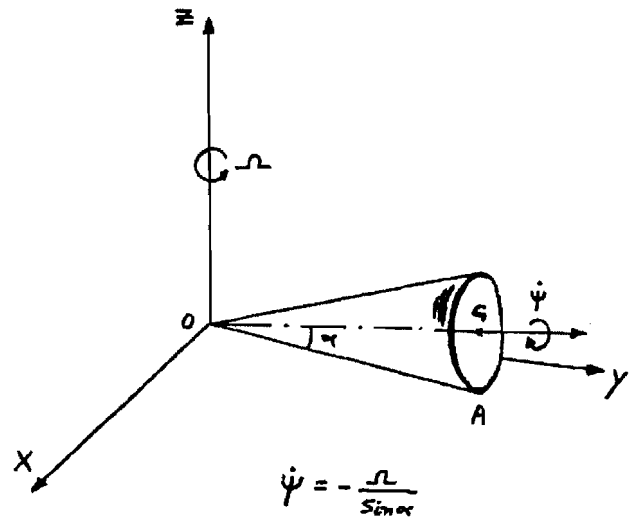
(a)

Angular momentum of the cone about the tip O : $\underline{H}_O = [I]_O \underline{\omega}$

$$[I]_O = \begin{bmatrix} I_1 & 0 & 0 \\ 0 & I_1 & 0 \\ 0 & 0 & I_3 \end{bmatrix}$$

$$\underline{\omega} = \begin{cases} \omega_x = 0 \\ \omega_y = -\Omega \cos \alpha \\ \omega_z = \Omega \frac{\cos^2 \alpha}{\sin \alpha} \end{cases}$$

$$\therefore \underline{H}_O = [I]_O \underline{\omega} = -I_1 \Omega \cos \alpha \hat{e}_y + I_3 \Omega \frac{\cos^2 \alpha}{\sin \alpha} \hat{e}_z$$



$$a = \sqrt{\frac{10}{3} \frac{I_3}{M}}$$

(b)

$$\underline{P} = M \underline{v}_C = M \Omega \frac{3}{4} a \frac{\cos^2 \alpha}{\sin \alpha} \hat{e}_x$$

$$\underline{F} = \frac{d\underline{P}}{dt} = \frac{3}{4} M \Omega a \frac{\cos^2 \alpha}{\sin \alpha} \frac{d\hat{e}_x}{dt}$$

$$\begin{aligned} \Omega \hat{e}_z \times \hat{e}_x &= \Omega (-\cos \alpha \hat{e}_y - \sin \alpha \hat{e}_z) \times \hat{e}_x \\ &= \Omega (\cos \alpha \hat{e}_z - \sin \alpha \hat{e}_y) = \Omega \hat{e}_{AO} \end{aligned}$$

$$\underline{F} = \frac{3}{4} M \Omega^2 a \frac{\cos^2 \alpha}{\sin \alpha} \hat{e}_{AO}$$

total required force

Problem 5

(b)

$$\dot{z}_0 = \frac{dH_0}{dt}$$

$$(z_0 = 0)$$

$$\frac{dH_0}{dt} = -I_1 \Omega \cos \alpha \frac{d\hat{e}_y}{dt} + I_3 \Omega \frac{\cos \alpha}{\sin \alpha} \frac{d\hat{e}_z}{dt}$$

$$\frac{d\hat{e}_y}{dt} = \Omega \hat{e}_z \times \hat{e}_y = \Omega (-\cos \alpha \hat{e}_y - \sin \alpha \hat{e}_z) \times \hat{e}_y = \Omega \sin \alpha \hat{e}_x$$

$$\frac{d\hat{e}_z}{dt} = \Omega \hat{e}_z \times \hat{e}_z = \Omega (-\cos \alpha \hat{e}_y - \sin \alpha \hat{e}_z) \times \hat{e}_z = -\Omega \cos \alpha \hat{e}_x$$

$$\therefore \dot{z}_0 = -\Omega^2 \left(I_1 \sin \alpha \cos \alpha + I_3 \frac{\cos^3 \alpha}{\sin \alpha} \right) \hat{e}_x$$

total required torque
about point O