

2.111J/18.435J Quantum Computation Problem Set 7

(Due: Tuesday, November 15, 2005)

1) Consider a 2 spin system possessing an Ising interaction $\hbar\Gamma(\sigma_z^A \otimes \sigma_z^B)/2$ that is placed in a magnetic field $B(t) = B \cos(\omega t)\hat{x} + B \sin(\omega t)\hat{y} + B_0\hat{z}$. Let $\omega_A = \gamma_A B_0$ and $\omega_B = \gamma_B B_0$ denote the respective Larmor frequencies of spins A and B around the z -axis. Viewed in a frame co-rotating with the circularly polarized field, the system's equation of motion is

$$\begin{aligned} \frac{d|\chi(t)\rangle}{dt} &= -\frac{i}{\hbar} H_{rot} |\chi(t)\rangle \\ &= -\frac{i}{2} [(\omega_A - \omega)\sigma_z^A + (\omega_B - \omega)\sigma_z^B + \Gamma(\sigma_z^A \otimes \sigma_z^B) + \gamma_A B \sigma_x^A + \gamma_B B \sigma_x^B] |\chi(t)\rangle. \end{aligned}$$

Evaluate $e^{-itH_{rot}/\hbar}$ for the case where the time-dependent part of the B field is applied at frequency $\omega = \omega_A - \Gamma$ for a time $t = \pi/(\gamma_A B)$. To what type of quantum logic operation does such time evolution correspond?

2) Compare the effect of the sequence of rotations

$$e^{i\pi\sigma_y^A/4} e^{-i\pi(\sigma_z^A \otimes \sigma_z^B)/4} e^{-i\pi\sigma_x^A/4}$$

to that of a standard CNOT gate.

[*Hint:* The formula $e^{ikM} = (\cos k)\mathbb{I} + i(\sin k)M$ holds for any matrix M such that $M^2 = \mathbb{I}$, not just 2×2 matrices.]

3) Consider the 3 spin Hamiltonian

$$H = -\frac{\hbar}{2} \left[\omega_A \sigma_z^A + \omega_B \sigma_z^B + \omega_C \sigma_z^C + \Gamma_{AB}(\sigma_z^A \otimes \sigma_z^B) + \Gamma_{AC}(\sigma_z^A \otimes \sigma_z^C) + \Gamma_{BC}(\sigma_z^B \otimes \sigma_z^C) \right].$$

The presence of the interactions splits each of the individual NMR peaks (“singlets”) at ω_A , ω_B , and ω_C that would occur in the noninteracting case into 4 peaks (a “doublet of doublets”). Derive formulas for the size of splittings of the ω_A peak in terms of Γ_{AB} , Γ_{AC} , and/or Γ_{BC} .

4) Consider the quantized simple harmonic oscillator,

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2.$$

Define the annihilation operator a and creation operator a^\dagger ,

$$a \equiv \frac{1}{\sqrt{2m\hbar\omega}}(m\omega x + ip); \quad a^\dagger \equiv \frac{1}{\sqrt{2m\hbar\omega}}(m\omega x - ip).$$

Verify the following four facts.

- $[a, a^\dagger] = \mathbb{I}$.
 - $[a, a^\dagger a] = a$.
 - $[a^\dagger, a^\dagger a] = -a^\dagger$.
 - $H = \frac{\hbar\omega}{2}(aa^\dagger + a^\dagger a) = \hbar\omega \left(a^\dagger a + \frac{1}{2} \right)$.
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5) Let $|0\rangle$ denote the vacuum state for the quantized simple harmonic oscillator.

$$|0\rangle \equiv \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} \int dx e^{-m\omega x^2/(2\hbar)} |x\rangle = \left(\frac{1}{\pi\hbar m\omega} \right)^{1/4} \int dp e^{-p^2/(2m\hbar\omega)} |p\rangle.$$

Verify that $a|0\rangle = 0$.

6) Define the non-normalized n th excited state of the quantized simple harmonic oscillator to be $|n\rangle \equiv (a^\dagger)^n |0\rangle$. Verify the following four facts.

- $a|n\rangle = \sqrt{n}|n-1\rangle$.
 - $a^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle$.
 - $H|n\rangle = \hbar\omega \left(n + \frac{1}{2} \right) |n\rangle$.
 - $|n\rangle = \frac{1}{\sqrt{n!}}(a^\dagger)^n |0\rangle$ is the proper normalization.
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7) Consider the following form of the Jaynes-Cummings Hamiltonian,

$$H = \hbar(C + D) + \frac{\hbar\omega}{2},$$

where $C \equiv \omega \left(a^\dagger a + \frac{1}{2} \sigma_z \right)$ and $D \equiv \kappa(a^\dagger \sigma_- + a \sigma_+) - \frac{\omega - \omega_0}{2} \sigma_z$. Verify that $[C, D] = 0$.