

Lenses and Imaging (Part II)

- Review from Part I
- Thin lens
 - matrix ray-tracing
 - imaging condition
 - real and virtual images
- Thick lens
- Generalized optical systems:
 - principal planes
 - effective, front and back focal lengths
- Reflective optics

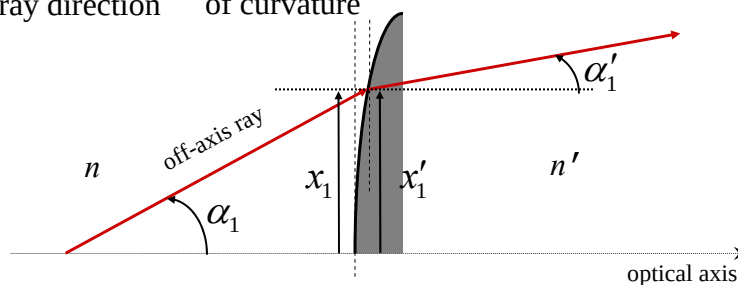
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Refraction at spherical surface

Refraction at positive spherical surface:
$$\begin{cases} x'_1 = x_1 \\ \alpha'_1 = \frac{n}{n'} \alpha_1 - \left[\frac{(n' - n)}{nR} \right] x_1 \end{cases}$$

x : ray height R : radius
 α : ray direction of curvature

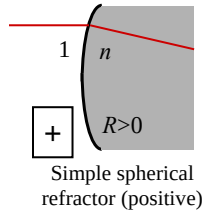
Power



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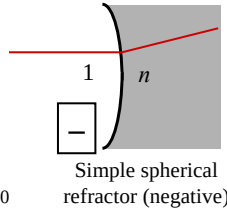
The power of surfaces

- Positive power bends rays “inwards” (*converging bundle*)



$$P \equiv \frac{(n_{\text{right}} - n_{\text{left}})}{R} = \frac{n - 1}{R} \frac{(+)}{(+)} > 0$$

- Negative power bends rays “outwards” (*diverging bundle*)

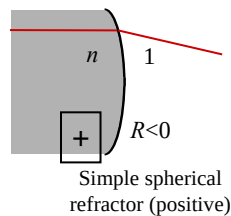


$$P \equiv \frac{(n_{\text{right}} - n_{\text{left}})}{R} = \frac{n - 1}{R} \frac{(+)}{(-)} < 0$$

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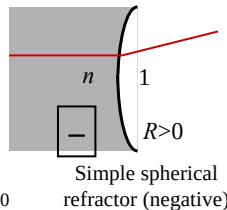
The power of surfaces

- Positive power bends rays “inwards” (*converging bundle*)



$$P \equiv \frac{(n_{\text{right}} - n_{\text{left}})}{R} = \frac{1 - n}{R} \frac{(-)}{(-)} > 0$$

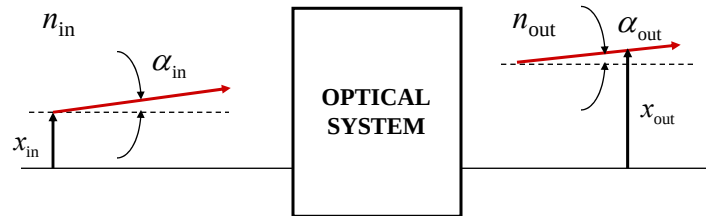
- Negative power bends rays “outwards” (*diverging bundle*)



$$P \equiv \frac{(n_{\text{right}} - n_{\text{left}})}{R} = \frac{1 - n}{R} \frac{(-)}{(+)} < 0$$

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(Paraxial) Ray-tracing in general



For an arbitrary ray entering with

- lateral departure displacement x_{in} wrt the optical axis
- angle of departure α_{in} wrt the optical axis
- in a departure space of refractive index n_{in}

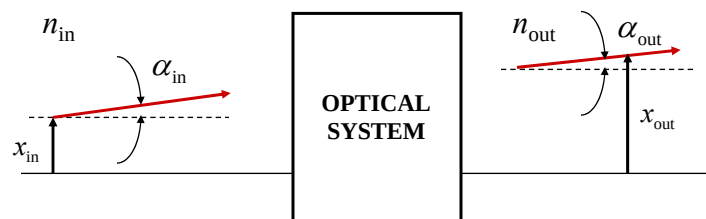
determine the ray's

- lateral arrival displacement x_{out} wrt the optical axis
- angle of arrival α_{out} wrt the optical axis
- in an arrival space of refractive index n_{out}

upon exiting the optical system

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Matrix ray-tracing in general

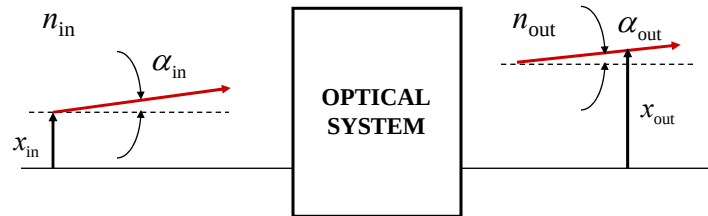


$$\begin{pmatrix} n_{\text{out}} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} n_{\text{in}} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

$\mathbf{M} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix}$ $\Rightarrow$ is the “system matrix,” found as **the product, in reverse order,** of the matrices of the elements (free space, spherical surfaces) comprising the system

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Matrix ray-tracing

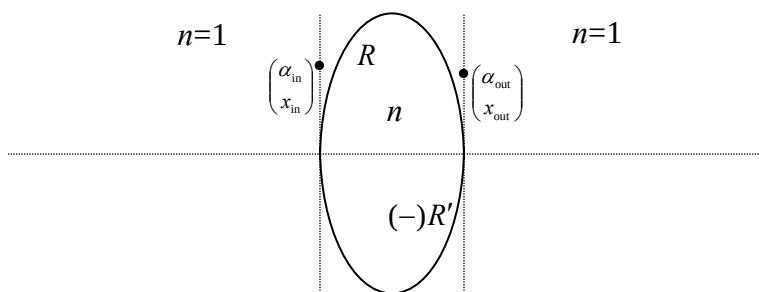


$$\begin{pmatrix} n_{\text{out}} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} n_{\text{in}} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

- Power M_{12}
- Imaging condition $M_{21} = 0$
- Lateral magnification $M_{21}|_{M_{21}=0}$
- Angular magnification $M_{11}|_{M_{21}=0}$

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Thin lens in air

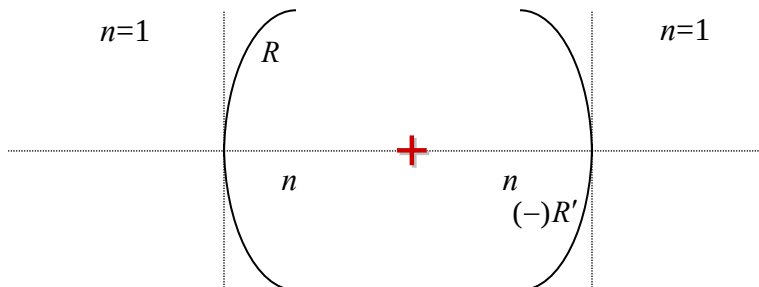


$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} ? & ? \\ ? & ? \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

Objective: specify input-output relationship

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Thin lens in air

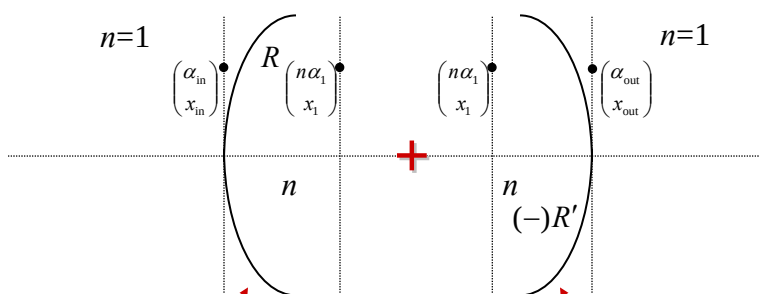


Model: refraction from first (positive) surface
+ refraction from second (negative) surface

Ignore space in-between (*thin* lens approx.)

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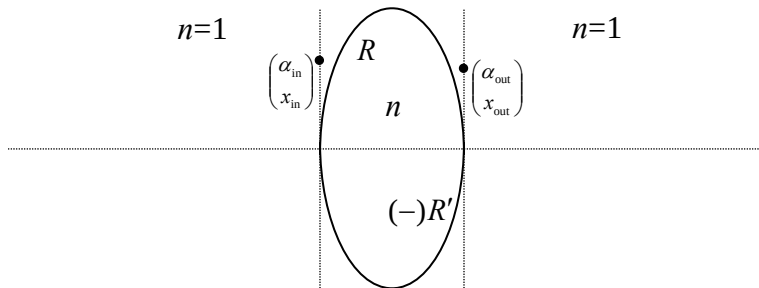
Thin lens in air



$$\begin{pmatrix} n\alpha_1 \\ x_1 \end{pmatrix} = \begin{pmatrix} 1 & -\frac{n-1}{R} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix} \quad \begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & -\frac{1-n}{R'} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} n\alpha_1 \\ x_1 \end{pmatrix}$$

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Thin lens in air

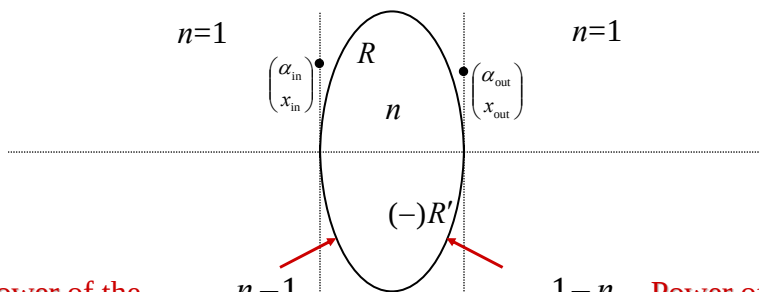


$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & -\frac{1-n}{R'} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & -\frac{n-1}{R} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & -(n-1)\left(\frac{1}{R} - \frac{1}{R'}\right) \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

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Thin lens in air



Power of the
first surface

$$P = \frac{n-1}{R}$$

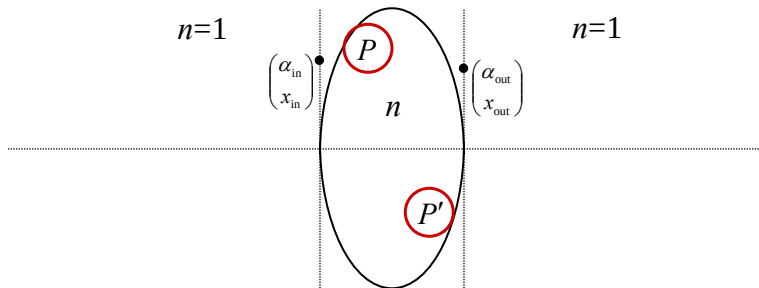
$$P' = \frac{1-n}{R'}$$

Power of the
second surface

$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & -(n-1)\left(\frac{1}{R} - \frac{1}{R'}\right) \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

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Thin lens in air

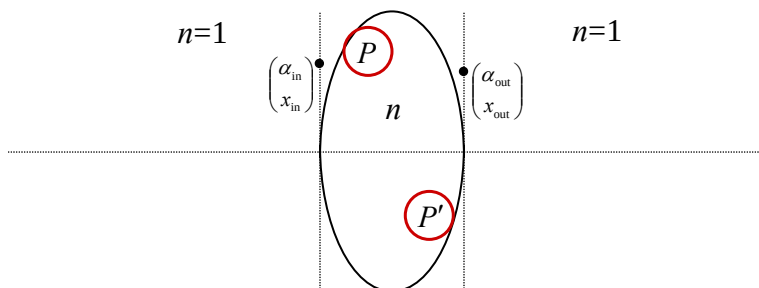


$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & -P_{\text{thin lens}} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

$$P_{\text{thin lens}} = P + P' = \frac{n-1}{R} + \frac{1-n}{R'} = \boxed{(n-1) \left(\frac{1}{R} - \frac{1}{R'} \right)} \quad \text{Lens-maker's formula}$$

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Thin lens in air



$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & -P_{\text{thin lens}} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix} \Rightarrow$$

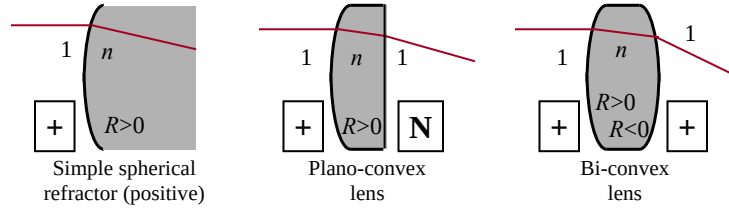
$$\alpha_{\text{out}} = \alpha_{\text{in}} - \boxed{P_{\text{thin lens}} x_{\text{in}}} \quad \text{Ray bending is proportional to the distance from the axis}$$

$$x_{\text{out}} = x_{\text{in}}$$

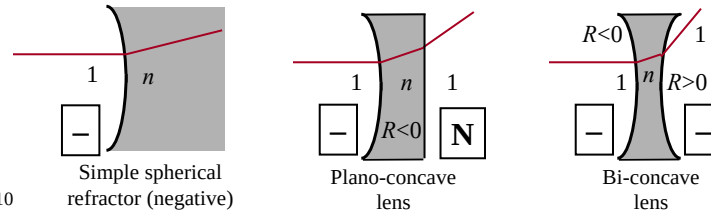
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The power of lenses

- Positive power bends rays “inwards” (*converging bundle*)

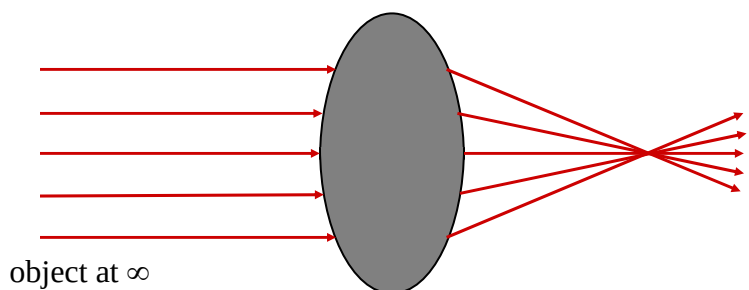


- Negative power bends rays “outwards” (*diverging bundle*)



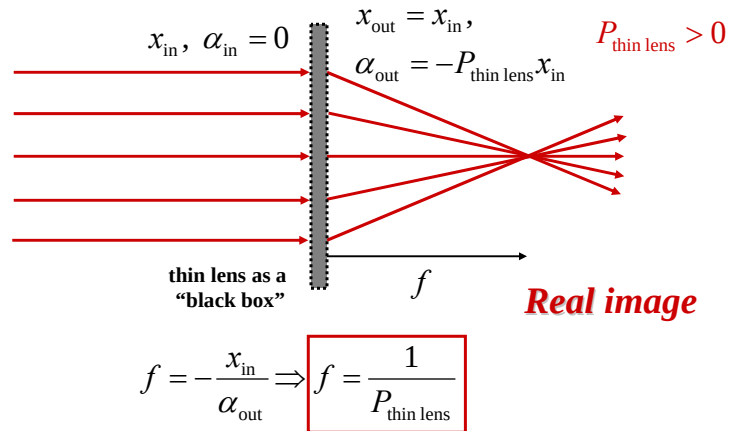
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Positive thin lens in air



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Positive thin lens in air

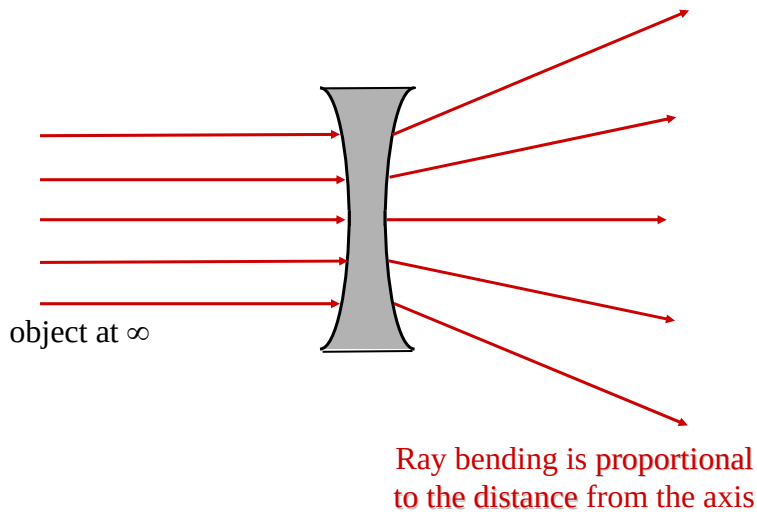


Focal point = image of an object at ∞

Focal length = distance between lens & focal point

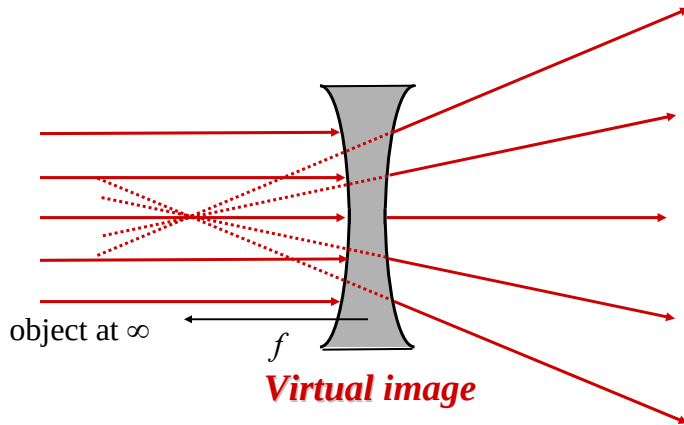
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Negative thin lens in air



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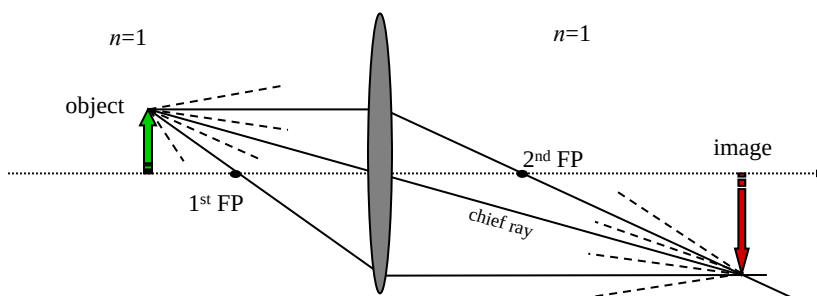
Negative thin lens in air



$$f = \frac{1}{P_{\text{thin lens}}} \quad \text{still applies, now with} \quad P_{\text{thin lens}} < 0 \Rightarrow f < 0 \quad (\text{to the "left"})$$

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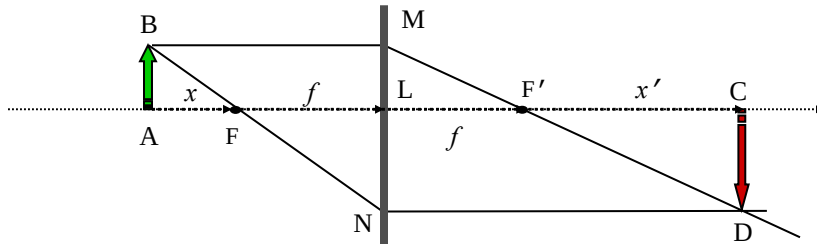
Imaging condition: ray-tracing



- Image point is located at the common intersection of **all** rays which emanate from the corresponding object point
- The two rays passing through the two focal points and the chief ray can be ray-traced directly

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Imaging condition: ray-tracing



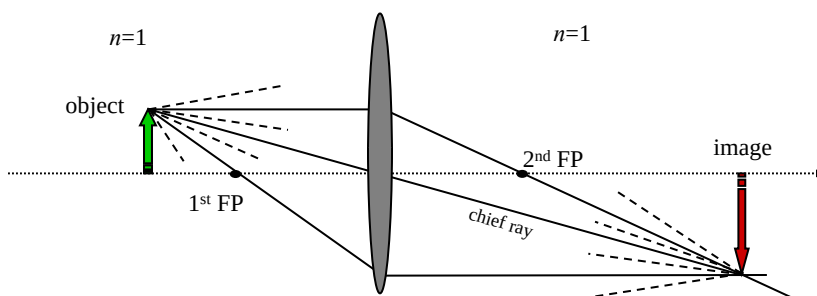
- $(ABF) \sim (FLN)$ and $(F'CD) \sim (MLF')$ are pairs of similar triangles

$$\frac{(AB)}{(AF)} = \frac{(LN)}{(FL)} \quad \frac{(LM)}{(LF')} = \frac{(CD)}{(F'C)} \quad (AB) = (ML) \quad (LN) = (CD)$$

$$(AF)(F'C) = (FL)(LF') \quad xx' = f^2$$

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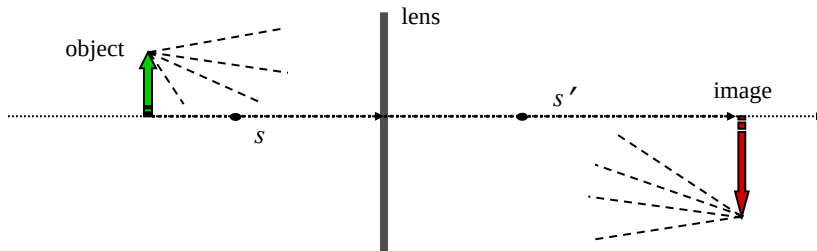
Imaging condition: matrix method



- Location of image point must be independent of ray departure angle at the object

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Imaging condition: matrix method

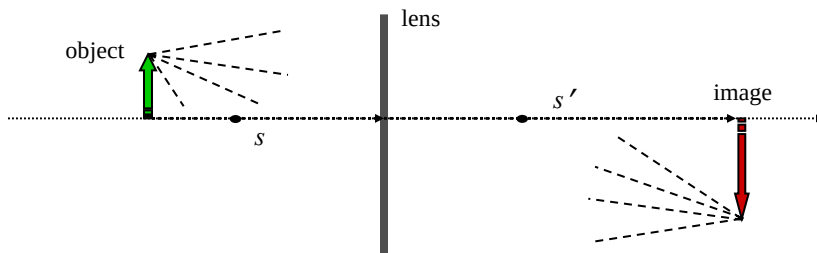


$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ s' & 1 \end{pmatrix} \begin{pmatrix} 1 & -\frac{1}{f} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ s & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix} = \begin{pmatrix} 1 - \frac{s}{f} & -\frac{1}{f} \\ s + s' - \frac{ss'}{f} & 1 - \frac{s'}{f} \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

$= 0$

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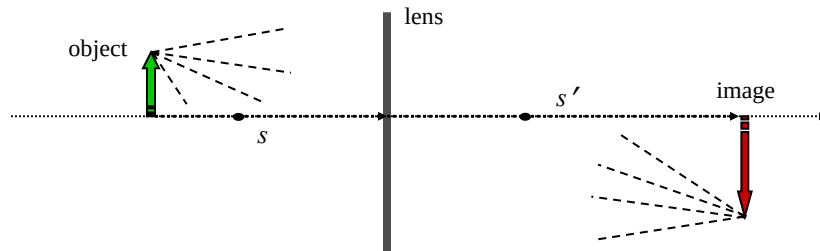
Imaging condition: matrix method



$$s + s' - \frac{ss'}{f} = 0 \quad \Leftrightarrow \quad \boxed{\frac{1}{s} + \frac{1}{s'} = \frac{1}{f}} \quad \text{Imaging condition (aka Lens Law)}$$

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Imaging condition: matrix method



$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 - \frac{s}{f} & -\frac{1}{f} \\ 0 & 1 - \frac{s'}{f} \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix} \Rightarrow x_{\text{out}} = \left(1 - \frac{s'}{f}\right) x_{\text{in}}$$

Lateral magnification :

$$M_T = \frac{x_{\text{out}}}{x_{\text{in}}} \\ M_T = 1 - \frac{s'}{f} = -\frac{s'}{s}$$

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Real & virtual images

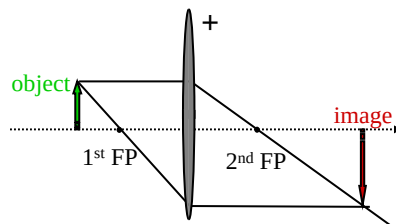


image: real & inverted; $M_T < 0$

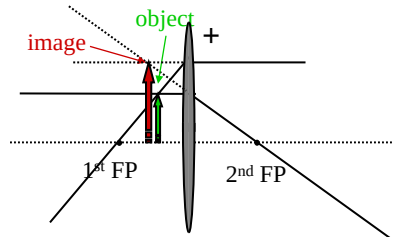


image: virtual & erect; $M_T > 1$

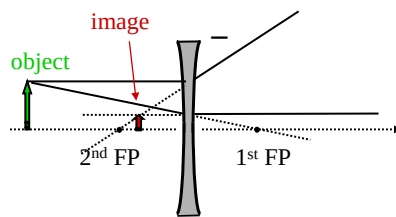


image: virtual & erect; $0 < M_T < 1$

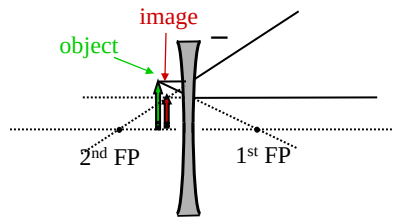
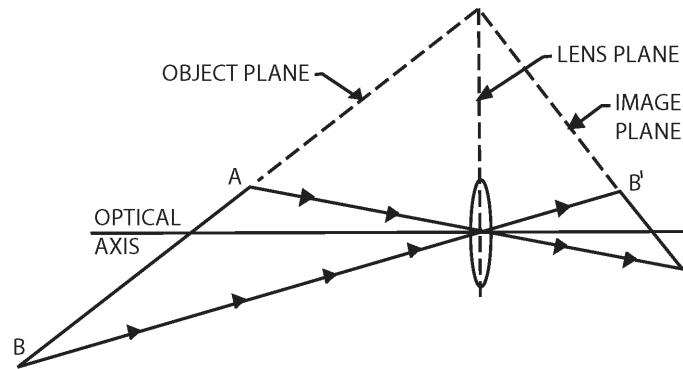


image: virtual & erect; $0 < M_T < 1$

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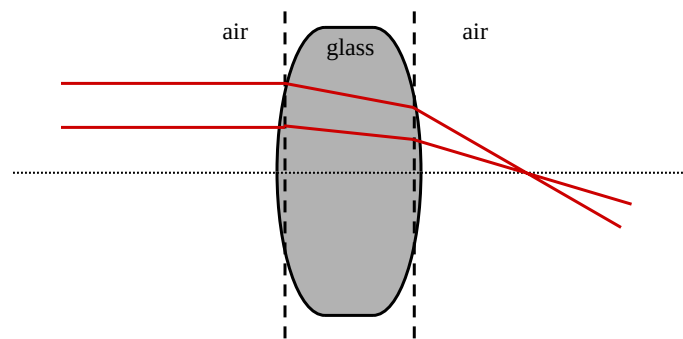
The Scheimpflug condition



The object plane and the image plane intersect at the plane of the lens.

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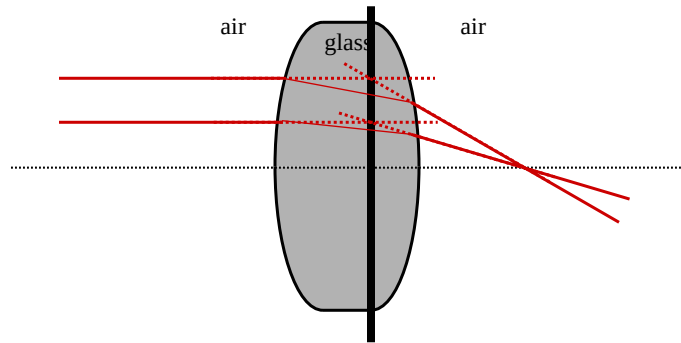
The thick lens



Rays bend in "two steps"

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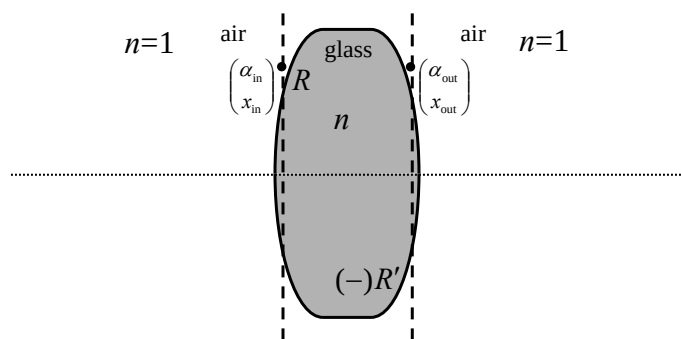
The thick lens



Equivalent to a thin lens placed “somewhere” within the thick element.
The location of this “equivalent thin lens” is the **Principal Plane** of the thick element

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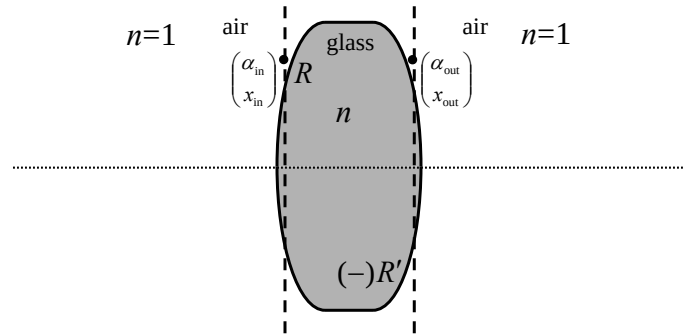
The thick lens



$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 & -\frac{1-n}{R'} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \frac{1}{d} & 0 \\ \frac{n}{1} & 1 \end{pmatrix} \begin{pmatrix} 1 & -\frac{n-1}{R} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

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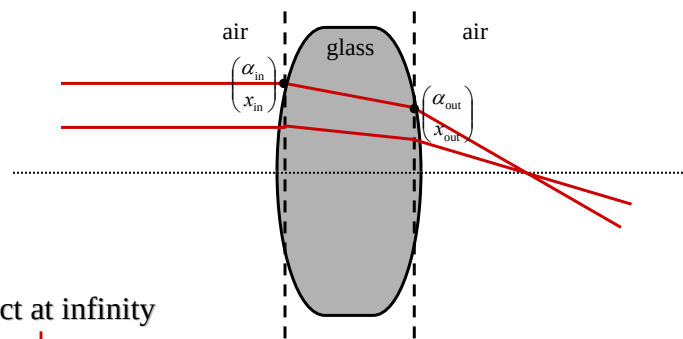
The thick lens



$$\begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} 1 + \frac{n-1}{n} \frac{d}{R'} - \left[(n-1) \left(\frac{1}{R} - \frac{1}{R'} \right) + \frac{(n-1)^2 d}{n R R'} \right] \\ \frac{d}{n} \\ 1 - \frac{n-1}{n} \frac{d}{R} \end{pmatrix} \begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix}$$

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The thick lens: power

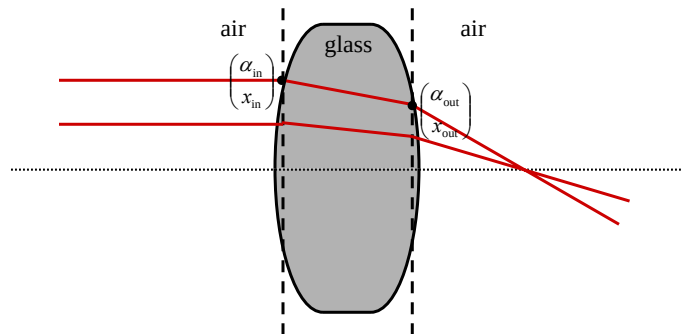


Object at infinity

$$\begin{pmatrix} \alpha_{\text{in}} \\ x_{\text{in}} \end{pmatrix} = \begin{pmatrix} 0 \\ x \end{pmatrix} \Rightarrow \begin{pmatrix} \alpha_{\text{out}} \\ x_{\text{out}} \end{pmatrix} = \begin{pmatrix} - \left[(n-1) \left(\frac{1}{R} - \frac{1}{R'} \right) + \frac{(n-1)^2 d}{n R R'} \right] x \\ \left(1 - \frac{n-1}{n} \frac{d}{R} \right) x \end{pmatrix} \Rightarrow \alpha_{\text{out}} = -P x_{\text{in}}$$

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The thick lens: power



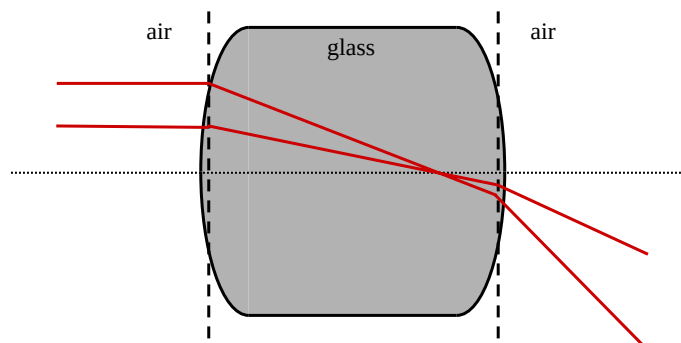
$$P = \left[(n-1) \left(\frac{1}{R} - \frac{1}{R'} \right) + \frac{(n-1)^2 d}{n R R'} \right] \quad \frac{1}{f} = \left[(n-1) \left(\frac{1}{R} - \frac{1}{R'} \right) + \frac{(n-1)^2 d}{n R R'} \right]$$

Power

f : Effective Focal Length

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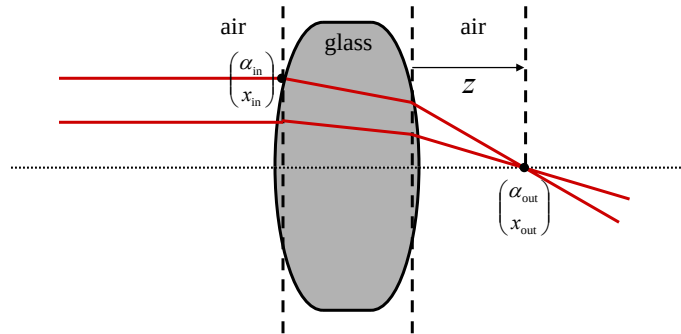
The very thick lens



Funny things happening:
rays diverge upon exiting from the element, *i.e.*
too much positive power leading to a negative element!

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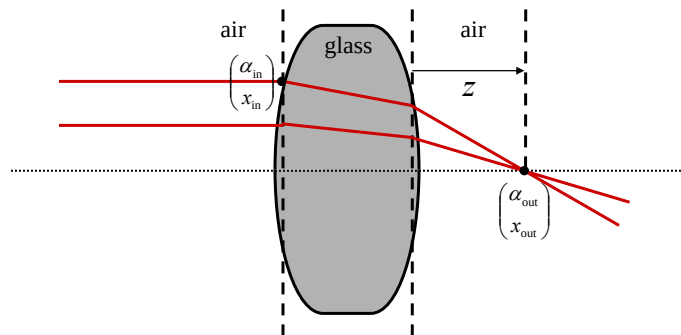
The thick lens: back focal length



$$\begin{pmatrix} \alpha_{out} \\ x_{out} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ z & 1 \end{pmatrix} \begin{pmatrix} 1 + \frac{n-1}{n} \frac{d}{R'} & -\frac{1}{f} \\ \frac{d}{n} & 1 - \frac{n-1}{n} \frac{d}{R} \end{pmatrix} \begin{pmatrix} \alpha_{in} \\ x_{in} \end{pmatrix}$$

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The thick lens: back focal length

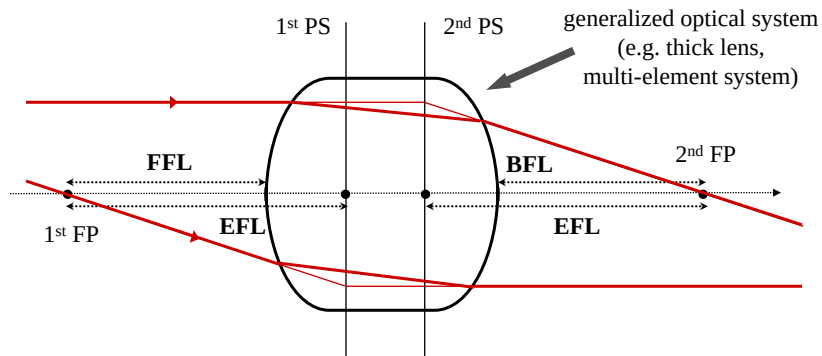


$$x_{out} = 0 \Rightarrow z = f \left(1 - \frac{n-1}{n} \frac{d}{R} \right)$$

z : Back Focal Length

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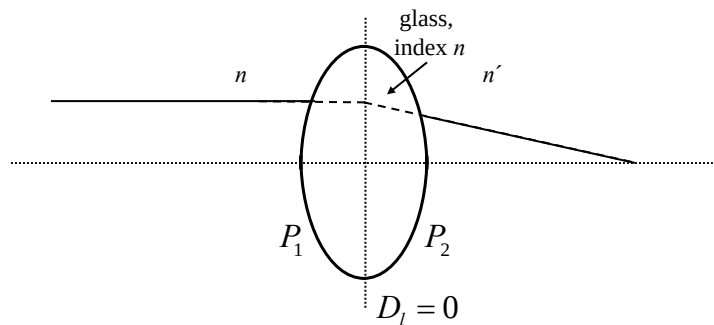
Focal Lengths & Principal Planes



EFL: Effective Focal Length (or simply “focal length”)
FFL: Front Focal Length
BFL: Back Focal Length
FP: Focal Point/Plane **PS:** Principal Surface/Plane

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PSs and FLs for thin lenses

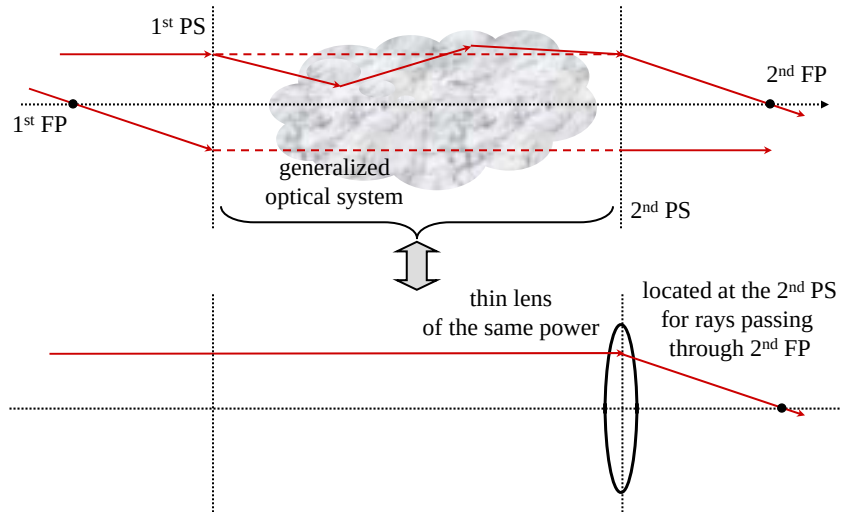


$$\frac{1}{(\text{EFL})} \equiv P = P_1 + P_2 \quad (\text{BFL}) = (\text{EFL}) = (\text{FFL})$$

- The principal planes coincide with the (collocated) glass surfaces
- The rays bend precisely at the thin lens plane (=collocated glass surfaces & PP)

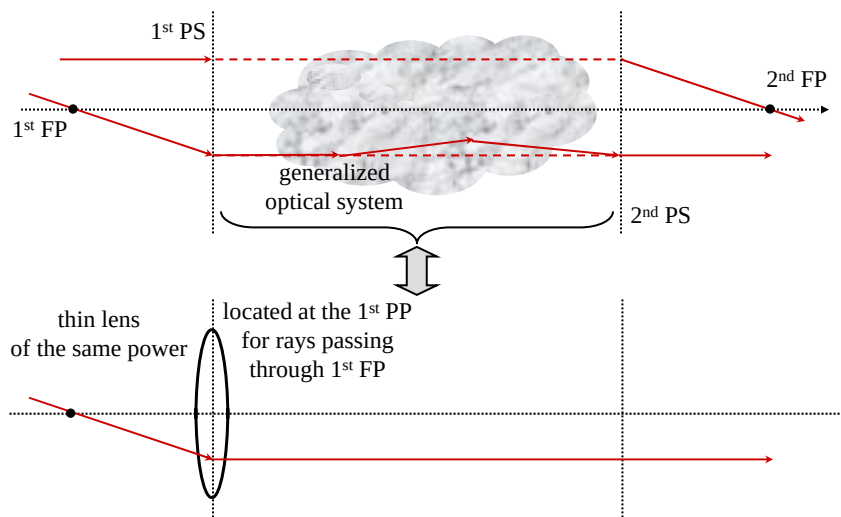
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The significance of principal planes /1



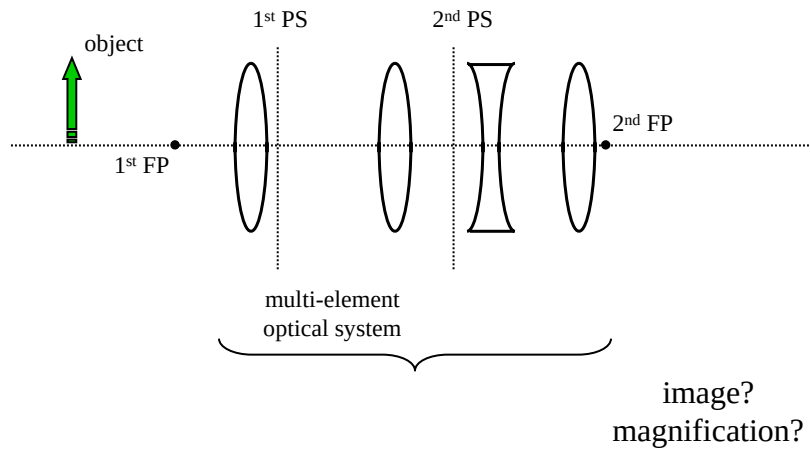
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The significance of principal planes /2



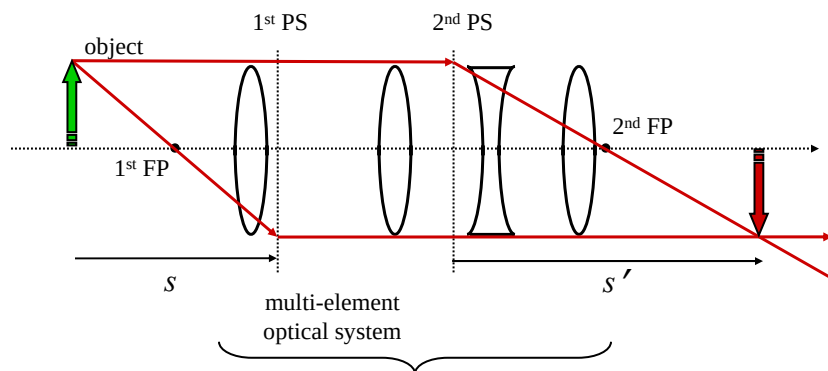
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The significance of principal planes /3



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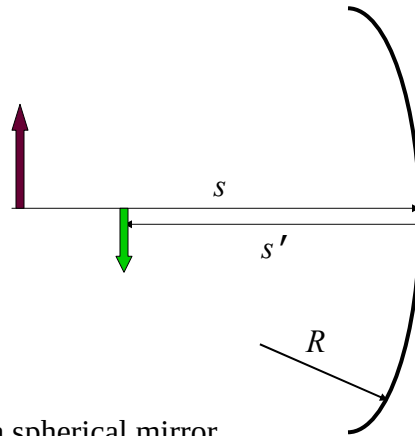
The significance of principal planes /3



$$\frac{1}{s} + \frac{1}{s'} = \frac{1}{f}, \quad m_{\text{lateral}} = -\frac{s'}{s} \quad \text{hold, where } f = (\text{EFL})$$

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Reflective Optics



Example:
imaging by a spherical mirror

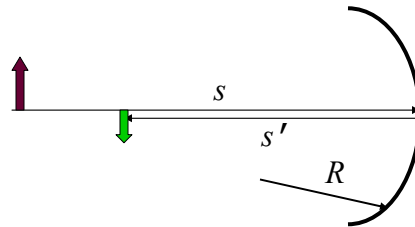
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Sign conventions for reflective optics

- Light travels from left to right *before reflection* and from right to left *after reflection*
- A radius of curvature is positive if the surface is convex towards the left
- Longitudinal distances *before reflection* are positive if pointing to the right; *longitudinal distances after reflection* are positive if pointing to the left
- Longitudinal distances are positive if pointing up
- Ray angles are positive if the ray direction is obtained by rotating the +z axis counterclockwise through an acute angle

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Reflective optics formulae



Imaging condition $\frac{1}{s} + \frac{1}{s'} = -\frac{2}{R}$

Focal length $f = -\frac{R}{2}$

Magnification $m_x = -\frac{s'}{s} \quad m_\alpha = -\frac{s}{s'}$

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