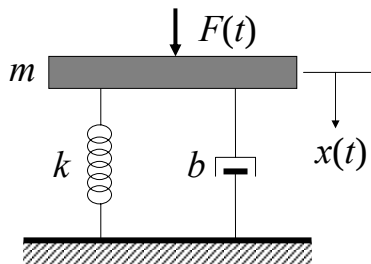


Linear shift invariant systems in the time domain



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Linear shift invariant systems

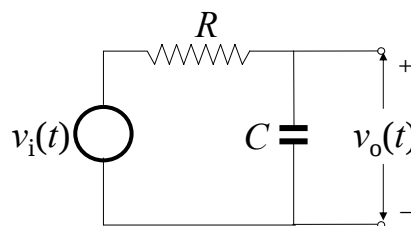


$$m\ddot{x} + b\dot{x} + kx = F(t)$$

$$H(\omega) = \frac{1}{-\omega^2 m - i\omega b + k}$$

$$|H(\omega)|^2 = \frac{1}{m^2} \frac{1}{(\omega^2 - \omega_r^2)^2 + 4\zeta^2 \omega_r^2 \omega^2}$$

$$\omega_r = \sqrt{\frac{k}{m}} \quad \zeta = \frac{b}{2\sqrt{km}}$$



$$RC\dot{v}_o + v_o = v_i(t)$$

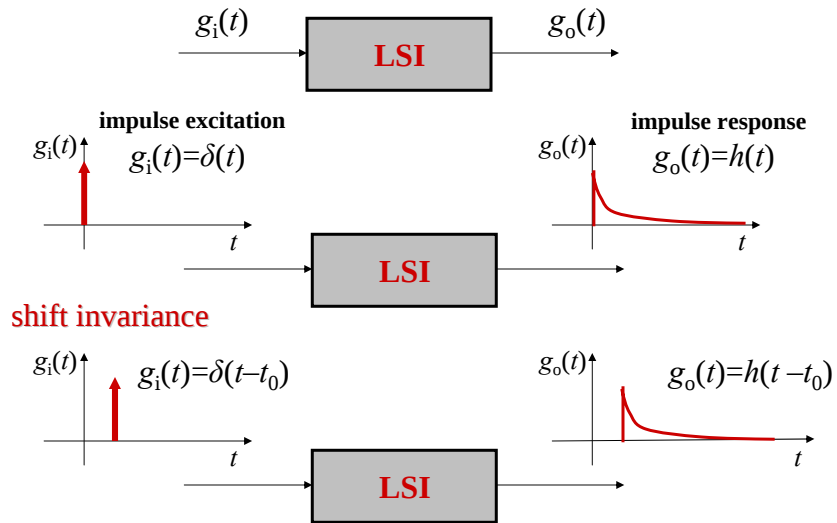
$$H(\omega) = \frac{1}{-i\omega\tau + 1}$$

$$|H(\omega)|^2 = \frac{1}{(\omega\tau)^2 + 1}$$

$$\tau = RC$$

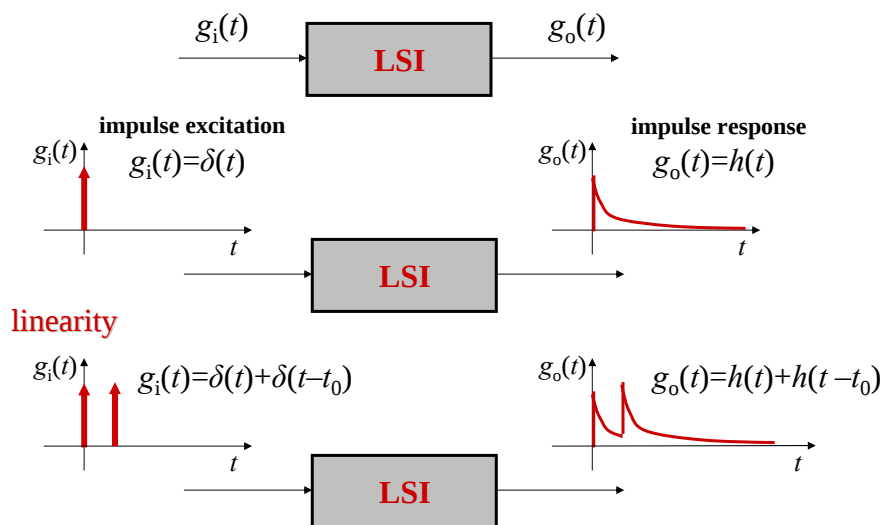
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Linear shift invariant systems



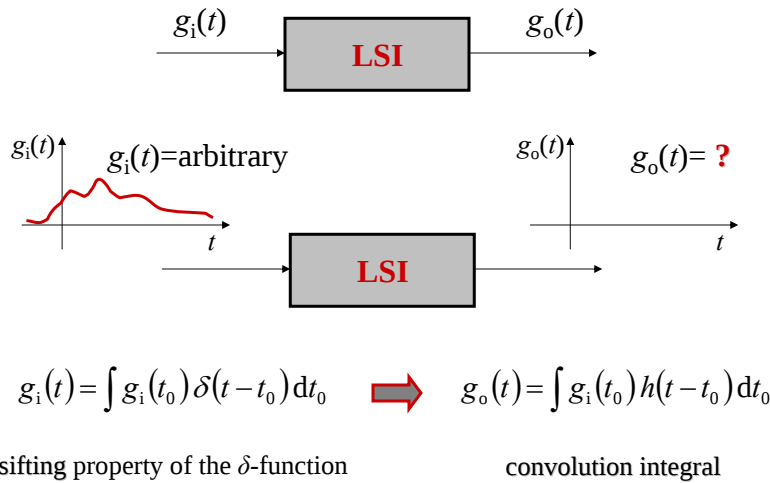
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Linear shift invariant systems



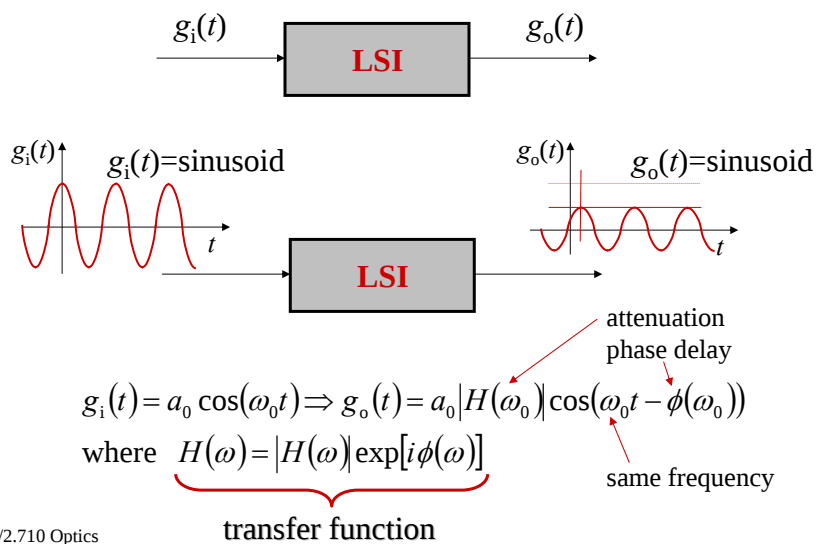
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Linear shift invariant systems



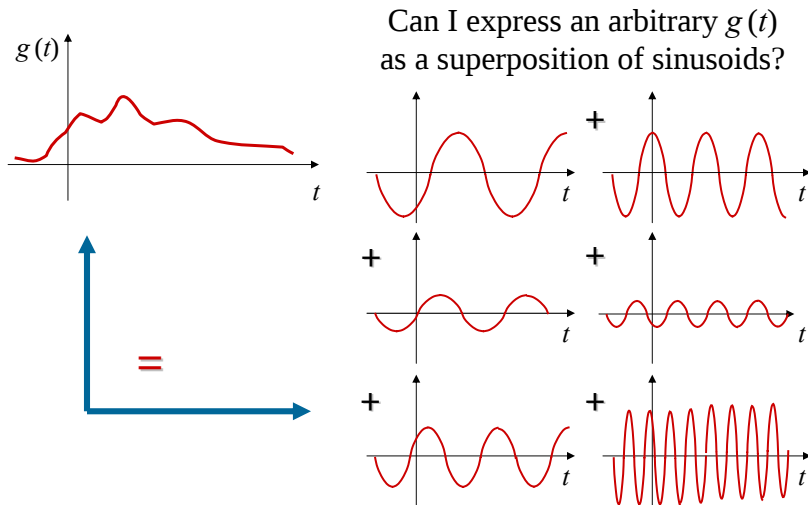
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Linear shift invariant systems



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From time to frequency: Fourier analysis



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The Fourier integral

(aka **inverse Fourier transform**)

$$g(t) = \int G(\nu) e^{+i2\pi\nu t} d\nu$$

superposition

sinusoids

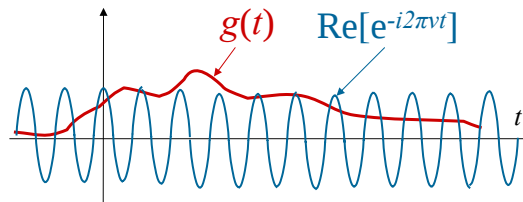
complex weight,
expresses relative amplitude
(magnitude & phase)
of superposed sinusoids

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The Fourier transform

The complex weight coefficients $G(\nu)$,
aka **Fourier transform** of $g(t)$
are calculated from the integral

$$G(\nu) = \int g(t) e^{-i2\pi\nu t} dt$$

$$\text{Re}[G(\nu)] = \int \text{Re}[g(t) e^{-i2\pi\nu t}] dt$$


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Fourier transform *pairs*

$$g(t) = \delta(t) \quad \leftrightarrow \quad G(\nu) = 1$$

$$g(t) = \exp(i2\pi\nu_0 t) \quad \leftrightarrow \quad G(\nu) = \delta(\nu - \nu_0)$$

$$g(t) = \cos(2\pi\nu_0 t) \quad \leftrightarrow \quad G(\nu) = \frac{1}{2} [\delta(\nu + \nu_0) + \delta(\nu - \nu_0)]$$

$$g(t) = \sin(2\pi\nu_0 t) \quad \leftrightarrow \quad G(\nu) = \frac{1}{2i} [-\delta(\nu + \nu_0) + \delta(\nu - \nu_0)]$$

$$g(t) = \text{rect}\left(\frac{t}{T}\right) \quad \leftrightarrow \quad G(\nu) = T \text{sinc}(T\nu)$$

$$g(t) = \exp\left(-\frac{t^2}{2T^2}\right) \quad \leftrightarrow \quad G(\nu) = T \exp(-2T^2\nu^2)$$

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