

NOTES FOR 6.002 LECTURE #15, APRIL 3, 2003

READ 14.1-14.1

SINUSOIDAL STEADY STATE ANALYSIS OF LINEAR CIRCUITS

ACTUAL INPUT:  $v(t) = V \cos(\omega t + \phi)$

REPLACE WITH  $\tilde{v}(t) = \tilde{V} e^{j\omega t}$  WHERE  $\tilde{V} = V e^{j\phi}$

USING THE ALGEBRA OF COMPLEX NUMBERS FIND AN OUTPUT OF THE FORM

$\tilde{v}_{out}(t) = \tilde{V}_{out} e^{j\omega t}$  THE REAL PART OF WHICH IS THE ACTUAL

OUTPUT  $v_{out} = V_{out} \cos(\omega t + \phi + \theta)$  WHERE  $\tilde{V}_{out} = V_{out} e^{j(\phi + \theta)}$

CONSIDER THE BASIC PASSIVE ELEMENTS WITH EXPONENTIAL EXCITATION

LET  $\tilde{v}(t) = \tilde{V} e^{st}$ ,  $\tilde{i}(t) = \tilde{I} e^{st}$  WHERE  $s = j\omega$  FOR CONVENIENCE

$\tilde{V}$  AND  $\tilde{I}$  ARE COMPLEX AMPLITUDES

RESISTANCE

$$\tilde{v}(t) = R \tilde{i}(t)$$

$$\tilde{V} = R \tilde{I}$$

$$\tilde{Z}_R = \frac{\tilde{V}}{\tilde{I}} = R$$

$$\tilde{Y}_R = 1/R = G$$

CAPACITANCE

$$\tilde{I} e^{st} = C \frac{d}{dt} \tilde{V} e^{st}$$

$$= C s \tilde{V} e^{st}$$

$$\tilde{Y}_C = \frac{\tilde{I}}{\tilde{V}} = C s$$

$$\tilde{Z}_C = \frac{1}{C s}$$

INDUCTANCE

$$\tilde{V} e^{st} = L \frac{d}{dt} \tilde{I} e^{st}$$

$$\tilde{V} = L s \tilde{I}$$

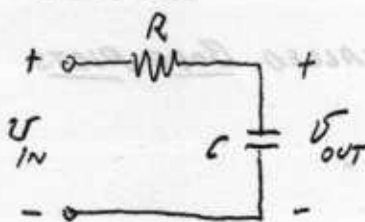
$$\tilde{Z}_L = \frac{\tilde{V}}{\tilde{I}} = L s$$

$$\tilde{Y}_L = 1/L s$$

$\tilde{Z}$  IS THE IMPEDANCE OF THE ELEMENT;  $\tilde{Y}$  IS THE ADMITTANCE.

ONCE THE ELEMENTS ARE REPRESENTED BY THEIR IMPEDANCES OR ADMITTANCE THE CIRCUIT CAN BE ANALYZED JUST LIKE A PURELY RESISTING NETWORK! THE ALGEBRA OF COMPLEX NUMBERS IS EMPLOYED.

EXAMPLE I



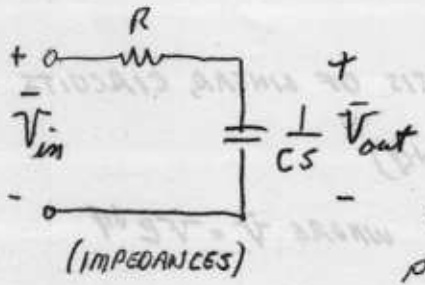
$$v_{in} = V_m \cos \omega t$$

FIND  $v_{out}$

EXPLORE FREQUENCY DEPENDENCE OF  $v_{out}$

FIRST CONSIDER THE ASYMPTOTES

ASSUME AN EXCITATION OF  $\bar{V}_{in} e^{st}$  WHERE  $\bar{V}_{in} = V e^{j0} = V$



JUST LIKE A RESISTIVE VOLTAGE DIVIDER:

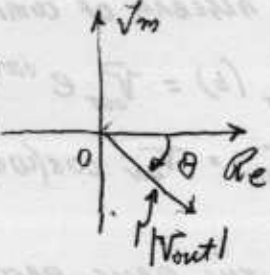
$$\bar{V}_{out} = \bar{V}_{in} \frac{\frac{1}{Cs}}{R + \frac{1}{Cs}} = \bar{V}_{in} \frac{1}{1 + RCs}$$

LET  $s = j\omega$  AND

PUT  $\bar{V}_{out}$  IN POLAR FORM:  $\bar{V}_{out} = V_{in} \frac{1}{\sqrt{1 + (RC\omega)^2}} e^{-j\theta}$

THEN

$$v_{out} = \text{Re} \left[ V e^{j0} \frac{1}{\sqrt{1 + (RC\omega)^2}} e^{-j\theta} e^{j\omega t} \right]$$



$$\theta = \tan^{-1}(RC\omega)$$

$$v_{out} = \frac{V}{\sqrt{1 + (RC\omega)^2}} \cos(\omega t - \theta)$$

THE OUTPUT IS, OF COURSE, A COSINE OF FREQUENCY  $\omega$ , DELAYED WITH RESPECT TO THE INPUT BY  $\theta$  AND ATTENUATED BY  $1/\sqrt{1 + (RC\omega)^2}$ . BOTH  $\theta$  AND THE AMPLITUDE  $V/\sqrt{1 + (RC\omega)^2}$  ARE FUNCTIONS OF  $\omega$ .

THE DEPENDENCE OF THE ATTENUATION AND PHASE LAG ON  $\omega$  IS BEST EXPLORED BY FOCUSING ON THE VOLTAGE TRANSFER RATIO  $\bar{A}$

$$\bar{A}(\omega) = \frac{\bar{V}_{out}}{\bar{V}_{in}} = \frac{1}{1 + jRC\omega} = \frac{1}{\sqrt{1 + (RC\omega)^2}} e^{-j\theta} \quad \theta = \tan^{-1}(RC\omega)$$

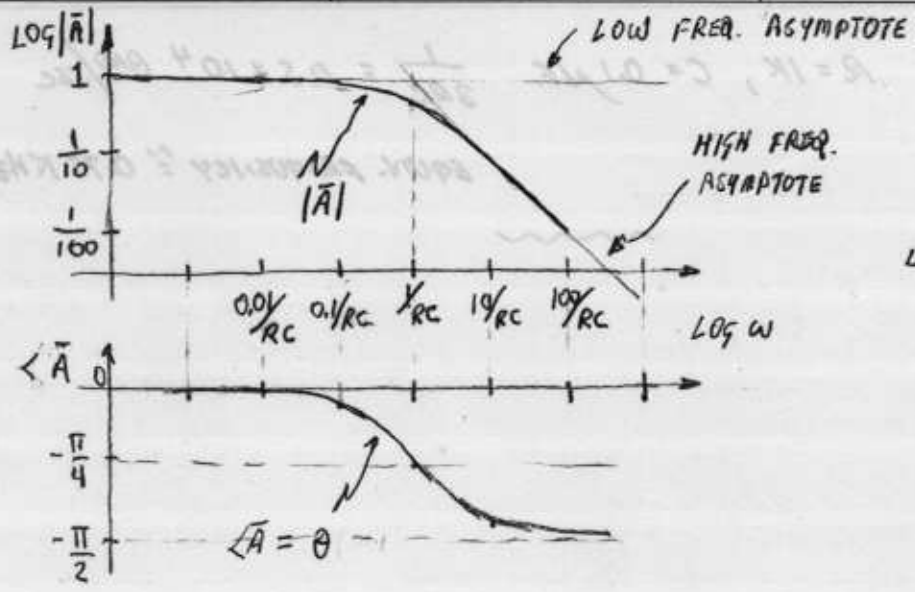
FOR LOW FREQUENCIES ( $\omega \ll \frac{1}{RC}$ )  $\bar{A}(\omega) \approx 1$  WHY?

FOR HIGH FREQUENCIES ( $\omega \gg \frac{1}{RC}$ )  $\bar{A}(\omega) \approx \frac{1}{RC\omega} e^{-j\pi/2}$  WHY?

BECAUSE  $\omega$  IN MANY CASES MAY VARY OVER MANY DECADES IT IS CONVENIENT TO PLOT  $|\bar{A}(\omega)|$  ON A LOGARITHMIC SCALE FOR  $\omega$ .

FOR REASONS THAT WILL SOON BE CLEAR,  $|\bar{A}(\omega)|$  IS ALSO PLOTTED ON A LOGARITHMIC SCALE WHILE  $\theta = \angle \bar{A}(\omega)$  IS PLOTTED ON A LINEAR SCALE.

THE RESULTING DIAGRAMS ARE CALLED BODE PLOTS



FOR HIGH FREQUENCIES  
 $|A| \approx \frac{1}{RC} \times \frac{1}{w}$   
 $\log |A| = \log \frac{1}{RC} - \log w$   
 A STRAIGHT LINE WITH  
 A SLOPE OF -1 ON  
 $\log |A|$  VS  $\log w$   
 COORDINATES  
 WHEN  $w = \frac{1}{RC}$   $|A| = 1$   
 WHICH LOCATES THE  
 ASYMPTOTE

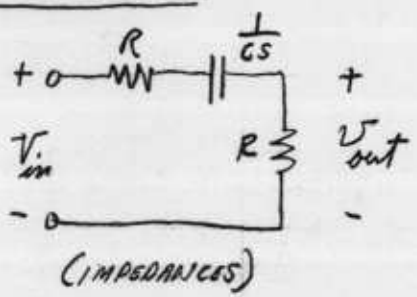
DEMO  $R = 1K\Omega$   $C = 0.1\mu F$

$1/RC = 10^4$  RAD/SEC WHERE THE ASYMPTOTES INTERSECT

THE EQUIVALENT FREQUENCY IS  $\frac{10^4}{2\pi} \approx 1.5$  KHZ

A LOW-PASS FILTER

EXAMPLE II



AGAIN A VOLTAGE DIVIDER:

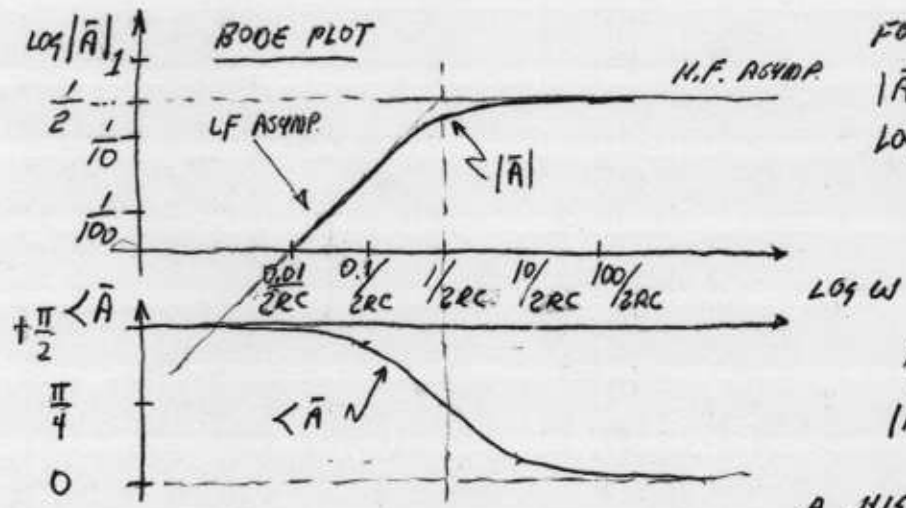
$$\bar{A}(w) = \frac{R}{R + R + \frac{1}{cS}} = \frac{RCs}{1 + 2RCs}$$

$$|\bar{A}(w)| = \frac{RCw}{\sqrt{1 + (2RCw)^2}} \quad \angle \bar{A}(w) = \frac{\pi}{2} - \theta$$

$$\theta = \tan^{-1}(2RCw)$$

FOR LOW FREQUENCIES  $|\bar{A}| \rightarrow 0$  WHY?

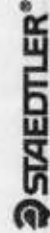
FOR HIGH FREQUENCIES  $|\bar{A}| \approx \frac{1}{2}$  WHY?



FOR LOW FREQUENCIES  
 $|\bar{A}| = RC \times w$   
 $\log |\bar{A}| = \log RC + \log w$   
 STRAIGHT LINE OF SLOPE +1

FOR HIGH FREQUENCIES  
 $|\bar{A}| \approx \frac{1}{2}$

A HIGH-PASS FILTER

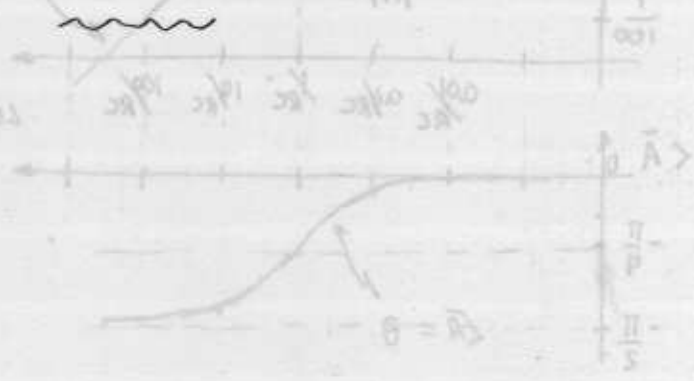


DEMO

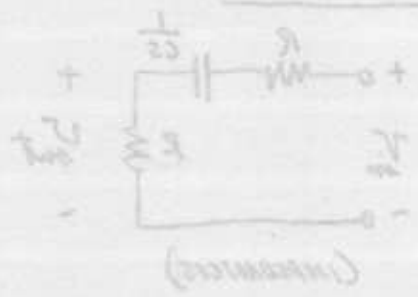
$R = 1K, C = 0.1 \mu F$

$\frac{1}{2RC} = 0.5 \times 10^4 \text{ RAD/SEC}$

EQUIV. FREQUENCY  $\approx 0.75 \text{ KHz}$



Example II

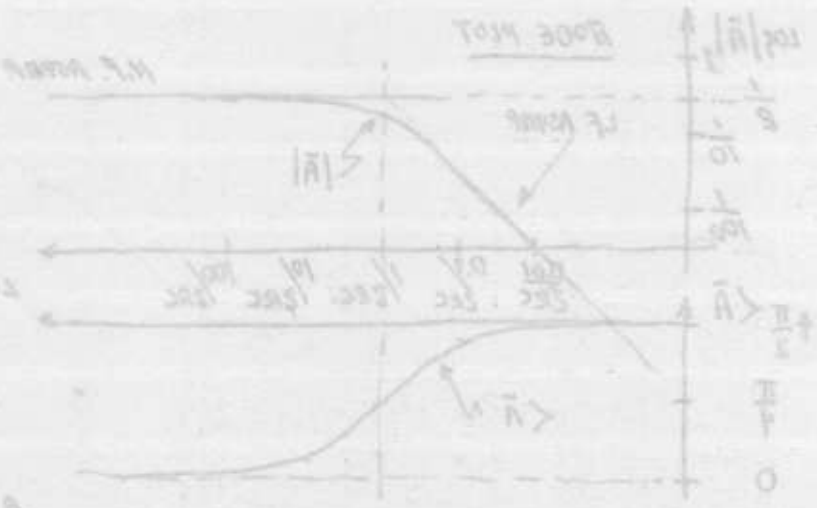


$$A(\omega) = \frac{R}{R + \frac{1}{j\omega C}} = \frac{j\omega RC}{1 + j\omega RC}$$

$$|A(\omega)| = \frac{\omega RC}{\sqrt{1 + (\omega RC)^2}}$$

$$\theta = \tan^{-1}(\omega RC)$$

For low frequencies  $|A| \approx 1$  what?  
 For high frequencies  $|A| \approx \frac{1}{\omega RC}$  what?



For low frequencies  
 $|A| = RC \omega$   
 For high frequencies  
 $|A| = \frac{1}{\omega RC}$

For high frequencies  
 $|A| \approx \frac{1}{\omega RC}$   
 A HIGH-PASS FILTER