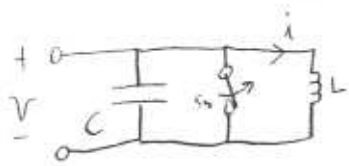


$$3.11 \quad i(t=0^+) = I$$

$$C = 0.1 \mu\text{F}$$

$$L = 40 \text{mH}$$



A) The natural frequency is the roots of the characteristic equation.

Find differential equation for the circuit:

$$\text{We know that } v_L = L \frac{di(t)}{dt} \quad \text{and} \quad v_C = \frac{1}{C} \int i(t) dt$$

Using these equations and KVL in the circuit for  $t > 0$

$$L \frac{di(t)}{dt} + \frac{1}{C} \int i(t) dt = 0$$

differentiate to remove the integral

$$L \frac{d^2 i(t)}{dt^2} + \frac{1}{C} i(t) = 0$$

Now plug in  $i(t) = Ae^{st}$ :

$$L s^2 Ae^{st} + \frac{1}{C} Ae^{st} = 0$$

$$s^2 + \frac{1}{LC} = 0$$

Find the roots of this characteristic equation to determine the natural frequency

$$s = \pm j \omega \frac{\text{rad}}{\text{sec}}$$

where

$$\omega = \sqrt{\frac{1}{LC}} = \sqrt{\frac{1}{(0.1 \mu\text{F})(40 \text{mH})}}$$

$$\omega = 15811.4 \text{ rad/sec}; \quad f = \frac{\omega}{2\pi} = 2.5 \text{ kHz}$$

determine period, first convert  $\omega_0$  into frequency

$$\omega = 2\pi f$$

$$f = \frac{15811.4 \text{ rad/sec}}{2\pi \text{ rad}} = 2516.46 \text{ Hz}$$

$$\text{The period is } T = \frac{1}{f} \Rightarrow T = \frac{1}{2516.46 \text{ Hz}} = 3.97 \times 10^{-4} \text{ sec}$$

B.) As seen in part A.) using KVL, the differential equation for current is.

$$\frac{d^2 i(t)}{dt^2} + \frac{1}{LC} i(t) = 0 \quad \text{for } t > 0$$

C.) To solve the diff. eq.:

$$i(t) = i_h(t) + i_p(t) \quad \text{for this case } i_p(t) = 0$$

Plugging  $i(t) = Ae^{st}$  to solve for  $i_h(t)$  we get:

$$i_h(t) = s^2 Ae^{st} + \frac{1}{LC} Ae^{st} = 0$$

$$s^2 = -\frac{1}{LC}$$

$$s = \pm j\omega \quad \text{where } \omega = \sqrt{\frac{1}{LC}}$$

Therefore

$$i(t) = Ae^{-j\omega t} + Be^{j\omega t}$$

We can switch this to a more intuitive form

$$i(t) = A_1 \sin \omega t + B_1 \cos \omega t$$

Now we use initial conditions to solve for the constants

$$i(t=0) = I \Rightarrow I = A_1 \overset{0}{\sin} + B_1 \overset{1}{\cos}$$

$$\underline{B_1 = I}$$

We can use the voltage to solve for  $A_1$

$$v(t) = \frac{1}{C} \int i(t) dt$$
$$v(t) = \frac{I}{\omega C} \sin \omega t + \frac{A_1}{\omega C} \cos \omega t + K_1 \quad \leftarrow \text{constant of integration}$$

Since there is no impulse into the circuit (3)

$$v(t=0^+) = 0 \Rightarrow 0 = \frac{-I}{\omega C} \sin 0 + \frac{A_1}{\omega C} \cos 0 + K_1$$

$$0 = \frac{A_1}{\omega C} + K \Rightarrow A_1 = -K = 0$$

Therefore the solution is

$$i(t) = I \cos \omega t \quad \text{where } \omega = \sqrt{\frac{1}{LC}}$$

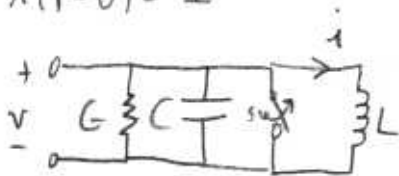
$$v(t) = \frac{-I}{\omega C} \sin \omega t = -I \sqrt{\frac{L}{C}} \sin \omega t$$

or

$$i(t) = I \cos (15811.4 t)$$

$$v(t) = -632.5 I \sin (15811.4 t)$$

8.2]  $i(t=0) = I$



$$C = .1 \mu\text{F}$$

$$L = 40 \text{ mH}$$

$$G = 5 \times 10^4 \text{ Mhos}$$

A) By using KCL at the top node, we obtain a differential equation for  $v(t)$

$$v(t)G + C \frac{dv(t)}{dt} + \frac{1}{L} \int v(t) dt = 0$$

Differentiate to remove integral.

$$G \frac{dv(t)}{dt} + C \frac{d^2v(t)}{dt^2} + \frac{1}{L} v(t) = 0$$

Arrange into standard form

$$\frac{d^2v(t)}{dt^2} + \frac{G}{C} \frac{dv(t)}{dt} + \frac{1}{CL} v(t) = 0$$

B) Plug in  $v(t) = Ae^{st}$  to find the characteristic equation:

$$Ae^{st}Gs + Ae^{st}Cs^2 + \frac{1}{L}Ae^{st} = 0$$

Rearrange:

$$\boxed{s^2 + \frac{G}{C}s + \frac{1}{LC} = 0} \quad \text{- characteristic equation}$$

Now find the roots

$$s = \frac{-\frac{G}{C} \pm \sqrt{\left(\frac{G}{C}\right)^2 - 4\frac{1}{LC}}}{2}$$

Plug in values

$$s = \frac{-5 \times 10^{-4}}{.1 \times 10^{-6}} \pm \sqrt{\frac{5 \times 10^{-4}}{.1 \times 10^{-6}} - 4\left(\frac{1}{(4 \times 10^{-3})(.1 \times 10^{-6})}\right)}$$

$$s = \frac{-5000 \pm \sqrt{(5000)^2 - 1 \times 10^9}}{2}$$

$$\boxed{s = -2500 \pm j 15612.5} \quad \text{- natural frequency}$$

C.) From part A.)

$$\frac{d^2 v(t)}{dt^2} + \frac{G}{C} \frac{dv(t)}{dt} + \frac{1}{LC} v(t) = 0$$

In general

$$\frac{d^2 y}{dt^2} + 2\alpha \frac{dy}{dt} + \omega_0^2 y = 0$$

Therefore

$$2\alpha = \frac{G}{C}$$

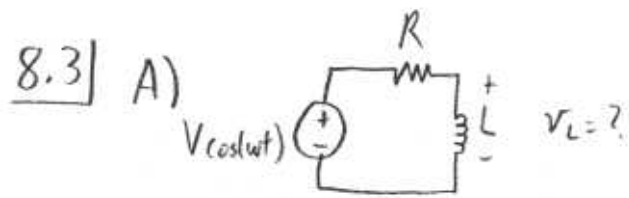
$$\omega_0^2 = \frac{1}{LC}$$

$$\boxed{\alpha = \frac{G}{2C} = 2500}$$

$$\boxed{\omega_0 = \sqrt{\frac{1}{LC}} = 15811.4}$$

$$Q = \frac{\omega_0}{2\alpha} = \frac{\sqrt{\frac{1}{LC}}}{\frac{G}{C}} = \frac{\sqrt{\frac{C}{L}}}{G} = 3.16$$

$$Q = \frac{\sqrt{\frac{L}{C}}}{G} = 3.16$$



Use impedances to solve and to find  $v_L$  use a voltage divider

$$V \cos(\omega t) = V e^{j\omega t}$$

$$\bar{V}_L = \frac{Z_L}{Z_L + Z_R} \bar{V}_i$$

$$\bar{V}_L = \frac{j\omega L}{j\omega L + R}$$

Convert to polar form

$$\bar{V}_L = \frac{\omega L}{\sqrt{R^2 + (\omega L)^2}} e^{j\phi}$$

where  $\phi = \frac{\pi}{2} - \tan^{-1}\left(\frac{\omega L}{R}\right)$

Take real magnitude to get  $v_L$  value

$$v_L(t) = \frac{\omega L}{\sqrt{R^2 + (\omega L)^2}} V \cos(\omega t + \phi)$$

B.)  $i(t) = I \cos(\omega t) = I e^{j\omega t}$

Combine the two parallel impedances

$$Z_{tot} = Z_R \parallel Z_C = \frac{\frac{1}{j\omega C} \cdot R}{\frac{1}{j\omega C} + R}$$

$$Z_{tot} = \frac{\frac{R}{j\omega C}}{1 + Rj\omega C} \Rightarrow Z_{tot} = \frac{R}{1 + j\omega RC}$$

Now we can say

$$\bar{V}_C = \bar{I} Z_{tot}$$

$$\bar{V}_C = \frac{R}{1 + j\omega RC} \bar{I} e^{j\omega t}$$

Put this into polar form

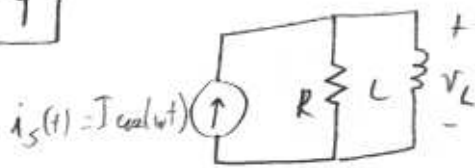
$$\bar{V}_C = \frac{R}{\sqrt{1 + (\omega RC)^2}} e^{-j\phi} \bar{I} e^{j\omega t}$$

where  $\phi = \tan^{-1} \left( \frac{\omega RC}{1} \right)$

Take real magnitude

$$V_C = \frac{R}{\sqrt{1 + (\omega RC)^2}} I \cos(\omega t + \phi)$$

8.4



$$v_L(t) = I \frac{\omega L}{\sqrt{1 + \left(\frac{\omega L}{R}\right)^2}} \cos\left(\omega t + \frac{\pi}{2} - \tan^{-1}\left(\frac{\omega L}{R}\right)\right)$$

A) For small  $\omega$ , the inductor will look close to a short circuit so  $v_L$  will be very small or close to 0

From the equation, the 1 term in the denominator will dominate and we can ignore the  $\left(\frac{\omega L}{R}\right)^2$  term.

Therefore the magnitude is

$I \omega L$  which is approximately zero if  $\omega \ll \frac{R}{L}$

B) At high frequencies the inductor has a high impedance and most all of the current will go to the resistor. Therefore the  $v_L$  will equal  $v_L = IR$

Looking at the equation, the  $(\frac{\omega L}{R})^2$  term dominates the denominator so the magnitude reduces to:

$$V_L = I \frac{\omega L}{\frac{\omega L}{R}} \Rightarrow \underline{V_L = IR}$$

C) The impedances of the resistor and the inductor must be equal or the imaginary and real portions of the denominator must be equal.

$$1 = \frac{\omega L}{R}$$
$$\boxed{\omega = \frac{R}{L}}$$

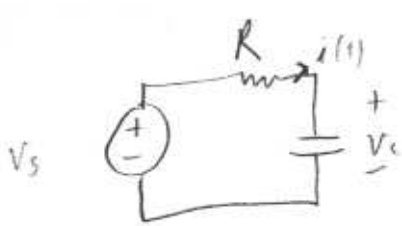
D.) When the imaginary and real parts are equal, the signal is  $\frac{1}{\sqrt{2}}$  times the magnitude with just the resistor. By plugging into the equation we can get this

$$\boxed{|v_L| = \frac{I \left(\frac{R}{L}\right) (L)}{\sqrt{1 + \left(\frac{R}{L}\right) \left(\frac{L}{R}\right)}} = \frac{IR}{\sqrt{2}} = .707 IR}$$

The phase will be  $\boxed{\phi = 45^\circ}$

Using the equation

$$\phi = \tan^{-1}\left(\frac{R}{L} \cdot \frac{L}{R}\right) \Rightarrow \boxed{\phi = \tan^{-1}(1) = 45^\circ = \frac{\pi}{4} \text{ radians}}$$



$$v_c(t) = V \frac{1}{\sqrt{1 + (\omega RC)^2}} \cos\left(\omega t + \frac{\pi}{2} - \tan^{-1}(\omega RC)\right)$$

$$i(t) = \frac{V \omega C}{\sqrt{1 + (\omega RC)^2}} \cos\left(\omega t + \frac{\pi}{2} - \tan^{-1}(\omega RC)\right)$$

E) The condition for low frequencies is  $\omega \ll \frac{1}{RC}$ .

For low frequencies, the capacitor has a large impedance so there will be a current that is almost zero and almost all of the voltage across the capacitor.

$$|v_c| = V \quad |i(t)| = 0$$

Looking at the equations: For  $i(t)$ , the  $(\omega RC)^2$  can be ignored and  $|i(t)| = \frac{\omega C}{1}$ . With small  $\omega$ ,  $|i(t)| = 0$

For the voltage, the  $(\omega RC)^2$  term can be ignored and we have  $|v_c(t)| = \frac{V}{1}$  or all of the voltage across the capacitor.

F) The condition for high frequencies is  $\omega \gg \frac{1}{RC}$ . At high frequencies, the opposite is true. The capacitor has a very small impedance, and looks almost like a short. Therefore, the current is limited by the resistor and almost all the voltage falls across the resistor.

$$|v_c(t)| \approx 0 \quad \text{and} \quad |i(t)| = \frac{V}{R}$$

Looking at the equations, we see the  $(\omega RC)^2$  term dominates the denominator in both.

$$|v_c(t)| = \frac{V}{\omega RC} \Rightarrow \text{with large } \omega \quad |v_c(t)| \approx 0$$

$$|i(t)| = \frac{V \omega C}{\omega RC} \Rightarrow \quad \underline{i(t) = \frac{V}{R}}$$



G) The voltage divides evenly when the impedances of the C and R are equal or when the imaginary and real parts are equal. ①

$$I = \omega RC$$
$$\boxed{\omega = \frac{1}{RC}}$$

H.) As in D.) , the magnitude is  $\frac{1}{\sqrt{2}}$  times the response total response and the phase is  $45^\circ$ .

$$\boxed{|V_C(\omega)| = \frac{V}{\sqrt{1 + (\frac{1}{RC}RC)^2}} = \frac{V}{\sqrt{2}} = .707 V}$$

$$\boxed{|I_C(\omega)| = \frac{V(\frac{1}{RC})C}{\sqrt{1+1}} = \frac{V}{\sqrt{2}R} = .707 \frac{V}{R}}$$

$$\psi = \tan^{-1}\left(\frac{(\frac{1}{RC})RC}{1}\right) = \tan^{-1}\left(\frac{1}{1}\right) \Rightarrow \boxed{\phi = 45^\circ}$$