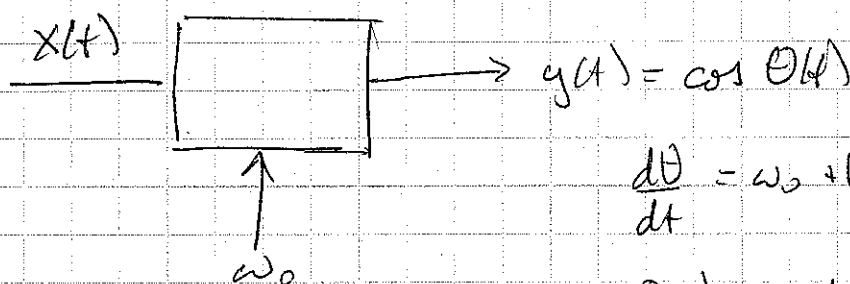


FM modulation

P. Hagel



$$\frac{d\theta}{dt} = \omega_0 + \alpha x(t)$$

$$\theta(t) = \omega_0 t + \alpha \int_{-\infty}^t x(t') dt'$$

Two limits:

(a) large modulation index

$$\theta(t) - \omega_0 t \gg 1$$

⇒ large bandwidth

(b) small modulation index

$$\theta(t) - \omega_0 t \ll 1$$

⇒ Bandwidth like AM case

If large modulation index, then need to work to do demodulation.

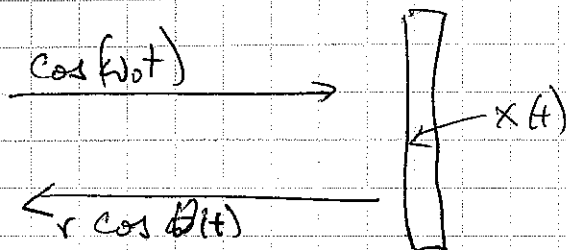
- track phase

- differentiation

- ~~Quadrature~~ Quadrature demodulation (low index)

$$\begin{aligned} \frac{dy}{dt} &= -\sin \theta \frac{d\theta}{dt} \\ &= -\sin \theta [\omega_0 + \alpha x(t)] \end{aligned}$$

## Example



Doppler shift:

When a train coming toward us sounds a whistle, the frequency is shifted up. When it is going away, the frequency is shifted down.

This is termed the Doppler effect. Light reflected from a moving surface is Doppler shifted. The amount of the shift is proportional to the velocity.

$$\frac{\Delta\omega}{\omega_0} = \frac{v}{c}$$

for small velocities  
toward the beam

So if we look at  $\frac{d\theta}{dt}$ , we can get an estimate of the velocity signal

$$\frac{d\theta}{dt} = \omega_0 + \omega_0 \left( \frac{v(t)}{c} \right)$$

OK - how much change in frequency?

For red light  $h\nu_0 \approx 2 \text{ eV}$

$$f_0 \approx 2 \cdot (2.4 \times 10^{14} \text{ Hz})$$

$$c \approx 3 \times 10^{10} \frac{\text{cm}}{\text{sec}}$$

Suppose audio frequency  $\sim 10^3 \text{ Hz}$

Distances  $\sim 1 \text{ m}$

$$v \sim 1 \text{ m} \cdot 10^3 \text{ Hz} \sim 10^{-7} \text{ cm} \cdot 10^3 \text{ sec}^{-1}$$

$$\sim 10^{-4} \frac{\text{cm}}{\text{sec}}$$

$$\Delta \omega \sim \omega_0 \cdot \frac{v}{c}$$

$$\sim 2\pi \cdot (4.8 \times 10^{14} \text{ Hz}) \cdot \frac{10^{-4} \text{ cm/sec}}{3 \times 10^{10} \text{ cm/sec}}$$

$$= 2\pi (4.8 \times 10^{14}) \frac{10^{-4}}{3}$$

$$= 10 \text{ Rad/sec} \rightarrow 1.6 \text{ Hz}$$

Thinking about the reflected signal

$$r \cos \theta(t) = r \cos \left( \int^t \omega_0 + \frac{v(t')}{c} \omega_0 dt' \right)$$

$$= r \cos \left( \omega_0 t + \omega_0 \int^t \frac{v(t')}{c} dt' \right)$$

$$= r \cos(\omega_0 t) \cos \left[ \omega_0 \int^t \frac{v(t')}{c} dt' \right]$$

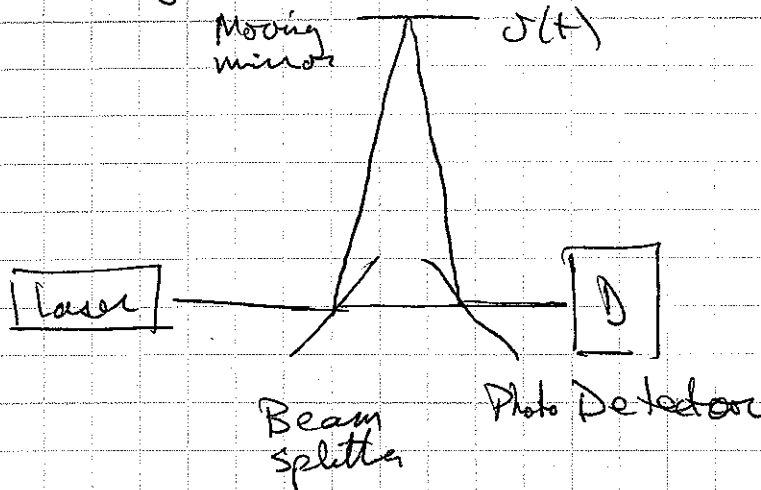
$$- r \sin(\omega_0 t) \sin \left[ \omega_0 \int^t \frac{v(t')}{c} dt' \right]$$

$\sin \omega_0 \int^t \frac{v(t')}{c} dt'$  looks small here

$$\approx r \cos(\omega_0 t) - r \sin(\omega_0 t) \sin \left( \omega_0 \int^t \frac{v(t')}{c} dt' \right)$$

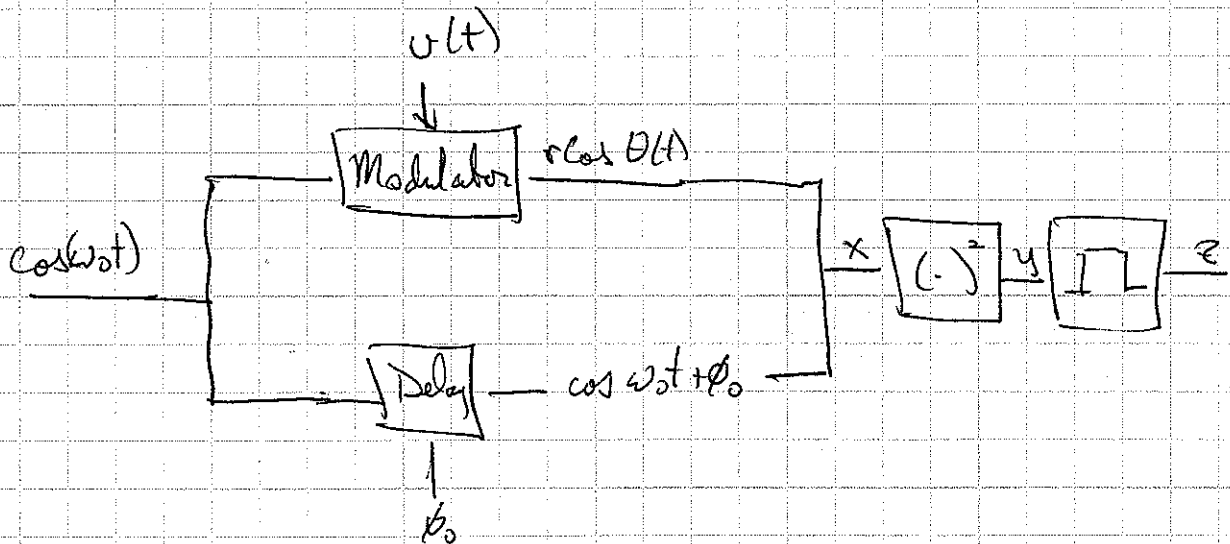
$$\approx r \cos(\omega_0 t) - r \sin(\omega_0 t) \left[ \omega_0 \int^t \frac{v(t')}{c} dt' \right]$$

# Homodyne detection scheme



Detector is a low pass filter on the square of the signal.

We can make a 6.003-type circuit for it.



$$x(t) = r \cos \theta(t) + \cos[\omega_0 t + \phi_0]$$

$$y(t) = x^2(t) = r^2 \cos^2 \theta(t) + \cos^2[\omega_0 t + \phi_0] + 2r \cos \theta(t) \cos[\omega_0 t + \phi_0]$$

$$\begin{aligned}
 y(t) &= r^2 \left( \frac{e^{j\theta} + e^{-j\theta}}{2} \right)^2 + r^2 \left( \frac{e^{j(\omega t + \phi_0)} + e^{-j(\omega t + \phi_0)}}{2} \right)^2 \\
 &\quad + 2r \left( \frac{e^{j\theta} + e^{-j\theta}}{2} \right) \left( \frac{e^{j(\omega t + \phi_0)} + e^{-j(\omega t + \phi_0)}}{2} \right) \\
 &= \frac{r^2 e^{2j\theta} + 2 + e^{-2j\theta}}{2} \\
 &\quad + \frac{r^2 e^{2j(\omega t + \phi_0)} + 2 + e^{-2j(\omega t + \phi_0)}}{2} \\
 &\quad + \frac{2r e^{j(\theta + \omega t + \phi_0)} + e^{j(\theta - \omega t - \phi_0)} + e^{-j(\theta - \omega t - \phi_0)} + e^{-j(\theta + \omega t + \phi_0)}}{4}
 \end{aligned}$$

The low pass filter part of this is

$$z(t) = \frac{r^2}{2} + \frac{1}{2} + \frac{r}{2} \left[ e^{j(\theta - \omega t - \phi_0)} + e^{-j(\theta - \omega t - \phi_0)} \right]$$

$$= \left[ \frac{r^2}{2} + r \cos \right] \theta(t) - \omega t - \phi_0 \quad \leftarrow \text{Detector sees the difference frequency}$$

Now, if  $\phi_0 = \frac{\pi}{2}$ , this is

$$z(t) \rightarrow \left[ \frac{r^2}{2} + r \sin \right] \theta(t) - \omega t$$

Recall that  $\theta(t) = \omega_0 t + \omega_0 \int^t \frac{v(t')}{c} dt'$

$$z(t) \rightarrow \frac{Lr^2}{2} + v \sin \left[ \omega_0 \int^t \frac{v(t')}{c} dt' \right]$$

$$\approx \frac{Lr^2}{2} + v \omega_0 \int^t \frac{v(t')}{c} dt'$$

Since  $\frac{dx}{dt} = v$   $x =$  position of surface

$$x = \int^t v dt'$$

$$z(t) \rightarrow \frac{Lr^2}{2} + v \frac{\omega_0}{c} x(t)$$

$$\frac{Lr^2}{2} + v \frac{2\pi}{\lambda} x(t)$$

### Applications:

Spying - listening to conversations in a room by bouncing a ~~light~~ laser beam off of the window

Gravitational wave detector - can see very tiny motions less than 1 Å easily.

