

## **Lecture 3 - Semiconductor Physics (II)**

### **CARRIER TRANSPORT**

February 13, 2001

#### **Contents:**

1. Thermal motion
2. Carrier drift
3. Carrier diffusion

#### **Reading assignment:**

Howe and Sodini, Ch. 2, §§2.4-2.6

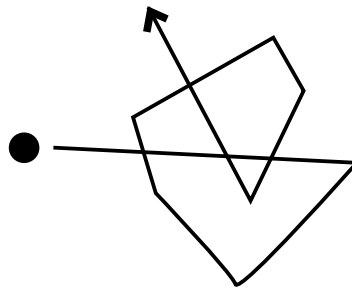
## Key questions

- What are the physical mechanisms responsible for current flow in semiconductors?
- How do electrons and holes in a semiconductor behave in an electric field?
- How do electrons and holes in a semiconductor behave if their concentration is non-uniform?

# 1. Thermal Motion

In thermal equilibrium, carriers are not sitting still:

- undergo collisions with vibrating Si atoms (*Brownian motion*)
- electrostatically interact with charged dopants and with each other



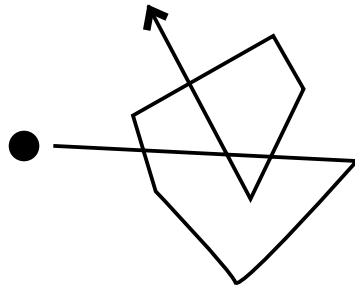
Characteristic time constant of thermal motion - mean free time between collisions:

$$\tau_c \equiv \text{collision time [s]}$$

In between collisions, carriers acquire high velocity:

$$v_{th} \equiv \text{thermal velocity [cm/s]}$$

...but get nowhere!



Characteristic length of thermal motion:

$$\lambda \equiv \text{mean free path [cm]}$$

$$\lambda = v_{th} \tau_c$$

Put numbers for Si at room temperature:

$$\tau_c \simeq 10^{-13} \text{ s}$$

$$v_{th} \simeq 10^7 \text{ cm/s}$$

$$\Rightarrow \lambda \simeq 0.01 \text{ } \mu\text{m} = 10 \text{ nm}$$

For reference, state-of-the-art MOSFETs today:

$$L_g \simeq 0.1 \text{ } \mu\text{m}$$

$\Rightarrow$  carriers undergo many collisions in modern devices

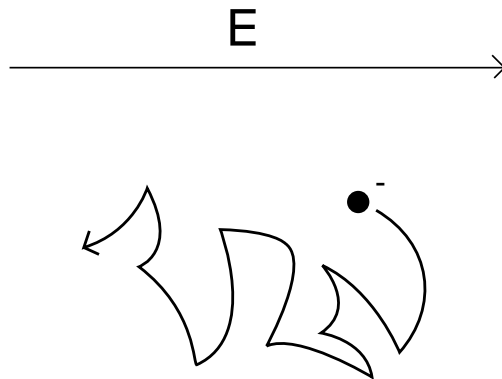
## 2. Carrier Drift

Apply electric field to semiconductor:

$$E \equiv \text{electric field [V/cm]}$$

$\Rightarrow$  net force on carrier

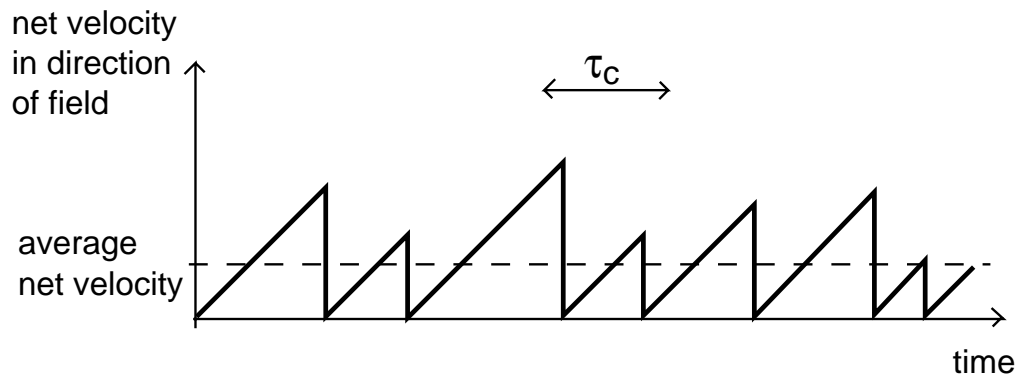
$$F = \pm qE$$



Between collisions, carriers accelerate in direction of field:

$$v(t) = at = \pm \frac{qE}{m_{n,p}} t$$

But velocity randomized every  $\tau_c$  (on average):



Then, average net velocity in direction of field:

$$\bar{v} = v_d = \pm \frac{qE}{2m_{n,p}} \tau_c = \pm \frac{q\tau_c}{2m_{n,p}} E$$

This is called *drift velocity* [cm/s].

Define:

$$\mu = \frac{q\tau_c}{2m_{n,p}} \equiv \text{mobility} [cm^2/V \cdot s]$$

Then, for electrons:

$$v_{dn} = -\mu_n E$$

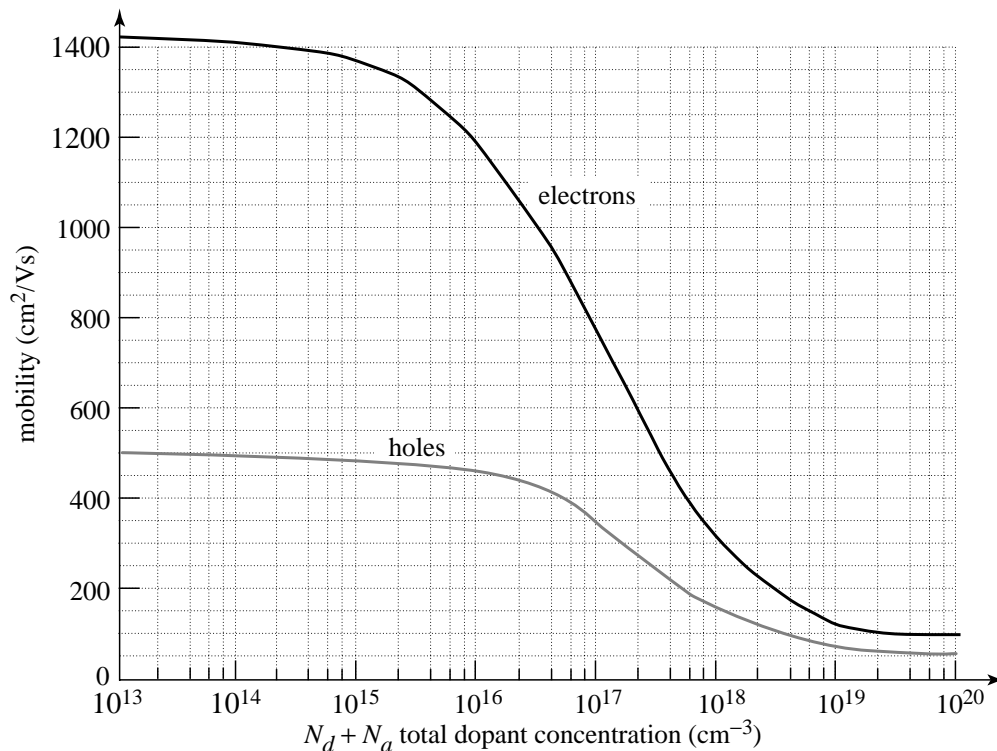
for holes:

$$v_{dp} = \mu_p E$$

Mobility is measure of *ease* of carrier drift:

- if  $\tau_c \uparrow$ , longer time between collisions  $\rightarrow \mu \uparrow$
- if  $m \downarrow$ , "lighter" particle  $\rightarrow \mu \uparrow$

Mobility in Si at room temperature, depends on doping:



- for low doping level,  $\mu$  limited by collisions with lattice
- for medium and high doping level,  $\mu$  limited by collisions with ionized impurities
- holes "heavier" than electrons:
  - $\rightarrow$  for same doping level,  $\mu_n > \mu_p$

## Drift current

Net velocity of charged particles  $\Rightarrow$  electric current:

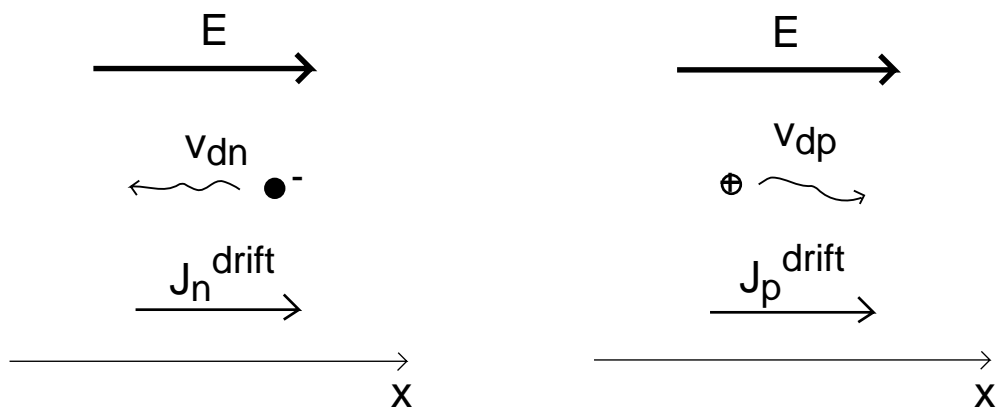
$$\begin{aligned} \text{Drift current density} &\propto \text{carrier drift velocity} \\ &\propto \text{carrier concentration} \\ &\propto \text{carrier charge} \end{aligned}$$

Drift currents:

$$J_n^{drift} = -qn v_{dn} = qn \mu_n E$$

$$J_p^{drift} = qp v_{dp} = qp \mu_p E$$

Check signs:



Total drift current:

$$J^{drift} = J_n^{drift} + J_p^{drift} = q(n\mu_n + p\mu_p)E$$

Has the shape of *Ohm's Law*:

$$J = \sigma E = \frac{E}{\rho}$$

Where:

$$\sigma \equiv \text{conductivity} [\Omega^{-1} \cdot \text{cm}^{-1}]$$

$$\rho \equiv \text{resistivity} [\Omega \cdot \text{cm}]$$

Then:

$$\rho = \frac{1}{\sigma} = \frac{1}{q(n\mu_n + p\mu_p)}$$

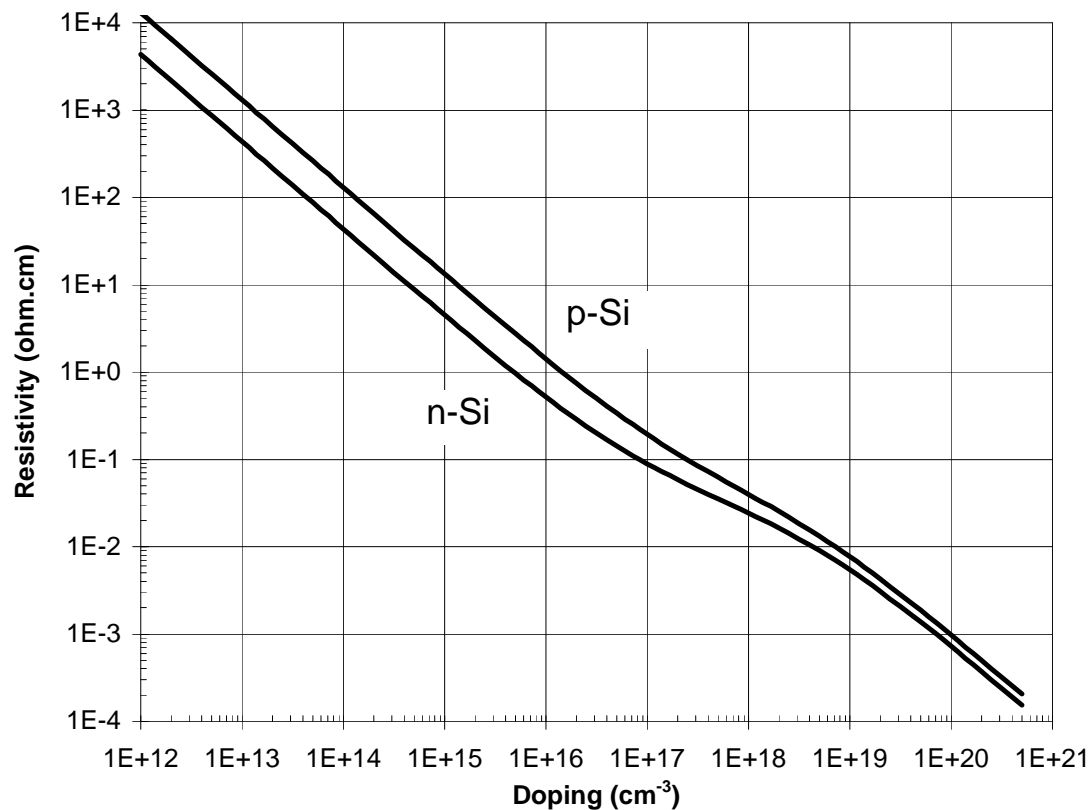
Resistivity commonly used to specify doping level.

- In n-type semiconductor:

$$\rho_n \simeq \frac{1}{qN_d\mu_n}$$

- In p-type semiconductor:

$$\rho_p \simeq \frac{1}{qN_a\mu_p}$$



Numerical example:

- Si with  $N_d = 3 \times 10^{16} \text{ cm}^{-3}$  at room temperature

$$\mu_n \simeq 1000 \text{ cm}^2/\text{V} \cdot \text{s}$$

$$\rho_n \simeq 0.21 \text{ } \Omega \cdot \text{cm}$$

- apply  $E = 1 \text{ kV/cm}$

$$v_{dn} \simeq -10^6 \text{ cm/s} \ll v_{th}$$

$$J_n^{drift} \simeq 4.8 \times 10^3 \text{ A/cm}^2$$

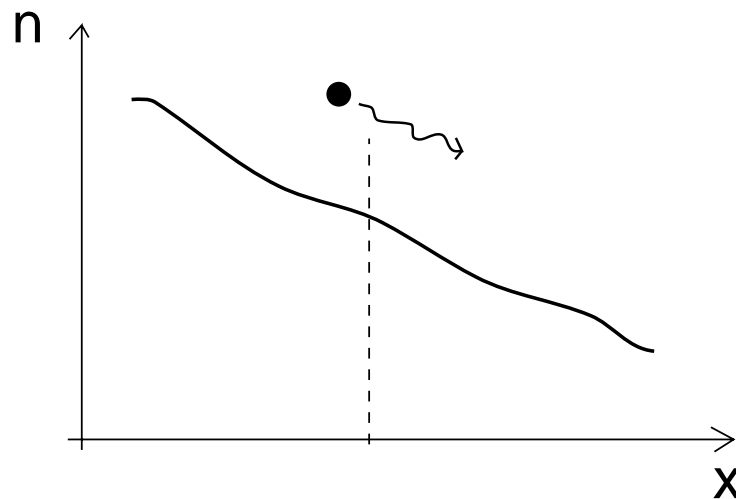
- time to drift through  $L = 0.1 \text{ } \mu\text{m}$ :

$$t_d = \frac{L}{v_{dn}} = 10 \text{ ps}$$

fast!

### 3. Carrier diffusion

Diffusion: particle movement in response to concentration gradient.



Elements of diffusion:

- a medium (*Si crystal*)
- a gradient of particles (*electrons and holes*) inside the medium
- collisions between particles and medium send particles off in random directions:
  - overall result is to erase gradient

Key diffusion relationship (*Fick's first law*):

*Diffusion flux*  $\propto$  - *concentration gradient*

Flux  $\equiv$  number of particles crossing unit area per unit time  $[cm^{-2} \cdot s^{-1}]$

For electrons:

$$F_n = -D_n \frac{dn}{dx}$$

For holes:

$$F_p = -D_p \frac{dp}{dx}$$

$D_n \equiv$  electron diffusion coefficient  $[cm^2/s]$

$D_p \equiv$  hole diffusion coefficient  $[cm^2/s]$

$D$  measures the ease of carrier diffusion in response to a concentration gradient:  $D \uparrow \Rightarrow F^{diff} \uparrow$ .

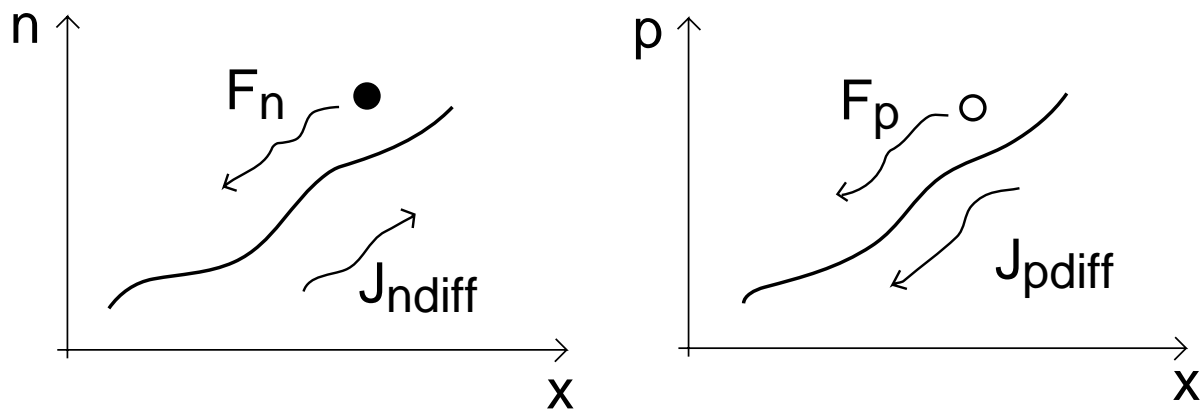
$D$  limited by vibrating lattice atoms and ionized dopants

*Diffusion current density = charge  $\times$  carrier flux*

$$J_n^{diff} = qD_n \frac{dn}{dx}$$

$$J_p^{diff} = -qD_p \frac{dp}{dx}$$

Check signs:



## Einstein relation

At the core of diffusion and drift is same physics: collisions among particles and medium atoms

$\Rightarrow$  there should be a relationship between  $D$  and  $\mu$

Einstein relation [don't derive in 6.012]:

$$\frac{D}{\mu} = \frac{kT}{q}$$

In semiconductors:

$$\frac{D_n}{\mu_n} = \frac{D_p}{\mu_p} = \frac{kT}{q}$$

$$\frac{kT}{q} \equiv \text{thermal voltage [V]}$$

At room temperature:

$$\frac{kT}{q} \simeq 25 \text{ mV}$$

For example: for  $N_d = 3 \times 10^{16} \text{ cm}^{-3}$ :

$$\begin{aligned} \mu_n &\simeq 1000 \text{ cm}^2/\text{V} \cdot \text{s} \rightarrow D_n \simeq 25 \text{ cm}^2/\text{s} \\ \mu_p &\simeq 400 \text{ cm}^2/\text{V} \cdot \text{s} \rightarrow D_p \simeq 10 \text{ cm}^2/\text{s} \end{aligned}$$

## Total current

In general, current can flow by drift and diffusion separately. Total current:

$$J_n = J_n^{drift} + J_n^{diff} = qn\mu_n E + qD_n \frac{dn}{dx}$$

$$J_p = J_p^{drift} + J_p^{diff} = qp\mu_p E - qD_p \frac{dp}{dx}$$

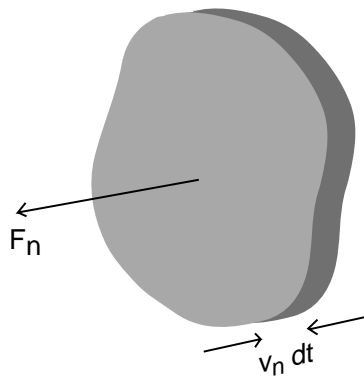
And

$$J_{total} = J_n + J_p$$

## Summary: relationship between $v$ , $F$ , and $J$

In semiconductors: charged particles move  
 $\Rightarrow$  particle flux  $\Rightarrow$  electrical current density

*Particle flux*: number of particles that cross surface of unit area placed normal to particle flow every unit time



Relationship between particle flux and velocity:

$$F_n = nv_n \quad F_p = pv_p$$

*Current density*: amount of charge that crosses surface of unit area placed normal to particle flow every unit time

$$J_n = -qF_n = -qnv_n \quad J_p = qF_p = qp v_p$$

whether carriers move by drift or diffusion.

## Key conclusions

- Electrons and holes in semiconductors are mobile and charged  $\Rightarrow$  *carriers* of electrical current!
- *Drift current*: produced by electric field

$$J^{drift} \propto E$$

- *Diffusion current*: produced by concentration gradient

$$J^{diff} \propto \frac{dn}{dx}, \frac{dp}{dx}$$

- Carriers move fast in response to fields and gradients
- Diffusion and drift currents are sizable in modern devices