Lecture 15 - The pn Junction Diode (I)

I-V CHARACTERISTICS

April 5, 2001

Contents:

- 1. pn junction under bias
- 2. I-V characteristics

Reading assignment:

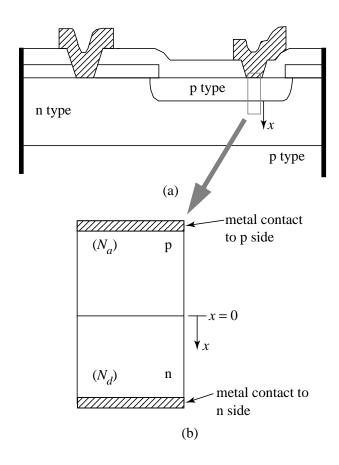
Howe and Sodini, Ch. 6, §§6.1-6.3

Key questions

- Why does the pn junction diode exhibit current rectification?
- Why does the junction current in forward bias increase as $\sim \exp \frac{qV}{kT}$?
- What are the leading dependences of the saturation current (the factor in front of the exponential)?

1. PN junction under bias

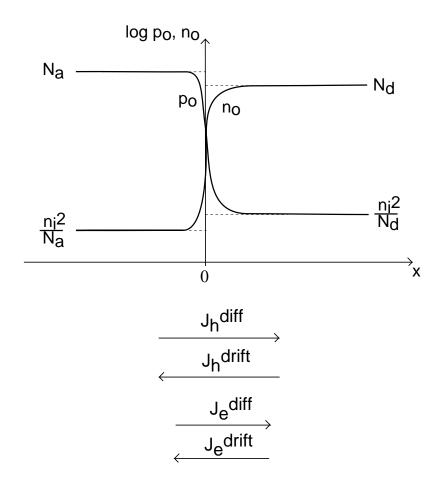
Focus on intrinsic region:



Upon application of voltage:

- electrostatics upset: depletion region widens or shrinks
- current flows (with rectifying behavior)
- carrier charge storage

Carrier profiles in thermal equilibrium:

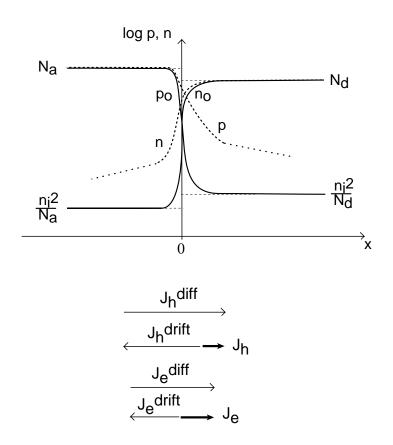


Inside SCR in thermal equilibrium: dynamic balance between drift and diffusion for electrons and holes.

$$|J_{drift}| = |J_{diff}|$$

Carrier concentrations in pn junction under bias:

• for
$$V > 0$$
, $\phi_B - V \downarrow \Rightarrow |E_{SCR}| \downarrow \Rightarrow |J_{drift}| \downarrow$



Current balance in SCR broken:

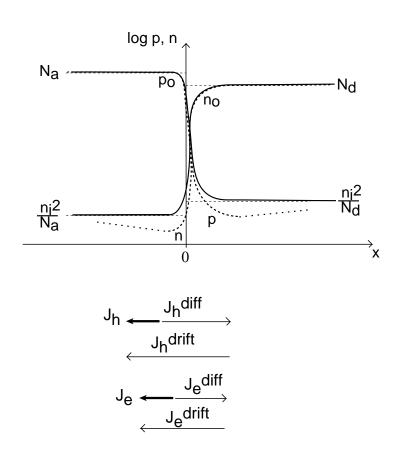
$$|J_{drift}| < |J_{diff}|$$

Net diffusion current in SCR

- \Rightarrow minority carrier *injection* into QNR's
- \Rightarrow excess minority carrier concentrations in QNR's

Lots of majority carriers in QNR's \Rightarrow current can be high.

• for
$$V < 0$$
, $\phi_B - V \uparrow \Rightarrow |E_{SCR}| \uparrow \Rightarrow |J_{drift}| \uparrow$



Current balance in SCR broken:

$$|J_{drift}| > |J_{diff}|$$

Net drift current in SCR

- \Rightarrow minority carrier *extraction* from QNR's
- \Rightarrow defect of minority carrier concentrations in QNR's

Few minority carriers in QNR's \Rightarrow current small.

What happens if minority carrier concentrations in QNR change from equilibrium?

- \Rightarrow Balance between generation and recombination broken
 - In thermal equilibrium: rate of break up of Si-Si bonds balanced by rate of formation of bonds

Si-Si bond
$$\underset{\text{recombination}}{\underbrace{\text{generation}}} n_0 + p_0$$

- If minority carrier injection:
 - ⇒ carrier concentration above equilibrium
 - \Rightarrow recombination prevails

- If minority carrier extraction:
 - ⇒ carrier concentrations below equilibrium
 - \Rightarrow generation prevails

Where does generation and recombination take place?

In modern devices, recombination mainly takes place at *surfaces*:

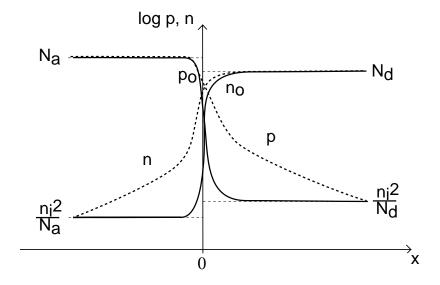
- perfect crystalline periodicity broken at a surface
 ⇒ lots of broken bonds: generation and recombination centers
- modern devices are very small
 ⇒ high area to volume ratio.

High generation and recombination activity at surfaces ⇒ carrier concentrations cannot deviate much from equilibrium values:

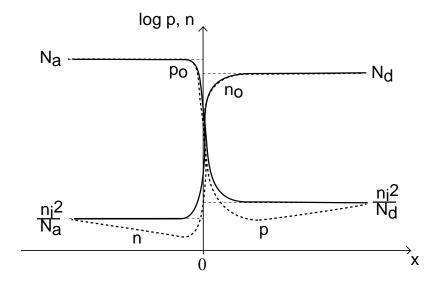
$$n(s) \simeq n_o, \ p(s) \simeq p_o$$

Complete physical picture for pn diode under bias:

 \bullet Forward bias: injected minority carriers diffuse through QNR \Rightarrow recombine at semiconductor surface

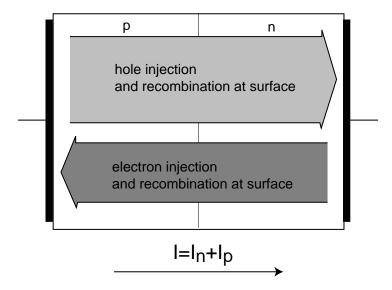


 \bullet Reverse bias: minority carriers extracted by SCR \Rightarrow generated at surface and diffuse through QNR

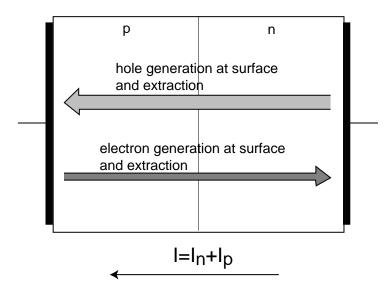


The current view:

• Forward bias:

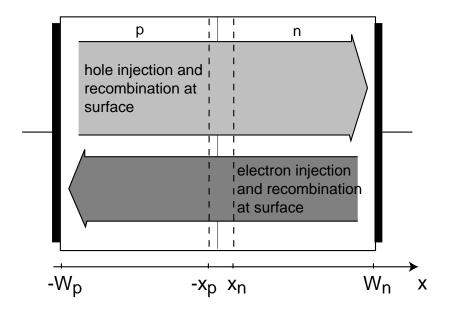


• Reverse bias:



What limits the magnitude of the diode current?

- <u>not</u> generation or recombination rate at surfaces
- <u>not</u> injection or extraction rates through SCR
- diffusion rate through QNR's



Development of analytical current model:

- 1. Calculate concentration of minority carriers at edges of SCR, $p(x_n)$ and $n(-x_p)$
- 2. calculate minority carrier diffusion current in each QNR, I_n and I_p
- 3. sum electron and hole diffusion currents, $I = I_n + I_p$

2. I-V characteristics

□ Step 1: computation of minority carrier boundary conditions at edges of SCR

In thermal equilibrium in SCR, $|J_{drift}| = |J_{diff}|$, and

$$\frac{n_o(x_1)}{n_o(x_2)} = \exp \frac{q[\phi(x_1) - \phi(x_2)]}{kT}$$

and

$$\frac{p_o(x_1)}{p_o(x_2)} = \exp \frac{-q[\phi(x_1) - \phi(x_2)]}{kT}$$

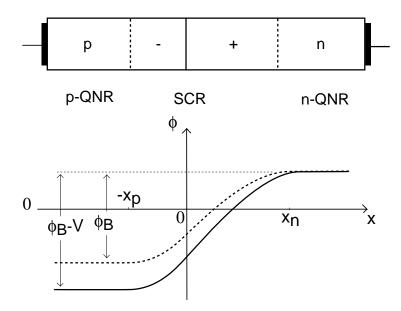
Under bias in SCR, $|J_{drift}| \neq |J_{diff}|$, but if difference small with respect to absolute values of current:

$$\frac{n(x_1)}{n(x_2)} \simeq \exp \frac{q[\phi(x_1) - \phi(x_2)]}{kT}$$

and

$$\frac{p(x_1)}{p(x_2)} \simeq \exp \frac{-q[\phi(x_1) - \phi(x_2)]}{kT}$$

This is called quasi-equilibrium.



At edges of SCR, then:

$$\frac{n(x_n)}{n(-x_p)} \simeq \exp \frac{q[\phi(x_n) - \phi(-x_p)]}{kT} = \exp \frac{q(\phi_B - V)}{kT}$$

and

$$\frac{p(x_n)}{p(-x_p)} \simeq \exp \frac{-q[\phi(x_n) - \phi(-x_p)]}{kT} = \exp \frac{-q(\phi_B - V)}{kT}$$

But:

$$p(-x_p) \simeq N_a$$
 and $n(x_n) \simeq N_d$

This is the *low-level injection* approximation [will discuss in more detail next time].

Then:

$$n(-x_p) \simeq N_d \exp \frac{q(V - \phi_B)}{kT}$$

and

$$p(x_n) \simeq N_a \exp \frac{q(V - \phi_B)}{kT}$$

Built-in potential:

$$\phi_B = \frac{kT}{q} \ln \frac{N_d N_a}{n_i^2}$$

Plug in above and get:

$$n(-x_p) \simeq \frac{n_i^2}{N_a} \exp \frac{qV}{kT}$$

and

$$p(x_n) \simeq \frac{n_i^2}{N_d} \exp \frac{qV}{kT}$$

Voltage dependence:

• Equilibrium (V = 0):

$$n(-x_p) = \frac{n_i^2}{N_a} \qquad p(x_n) = \frac{n_i^2}{N_d}$$

• Forward (V > 0):

$$n(-x_p) \gg \frac{n_i^2}{N_a}$$
 $p(x_n) \gg \frac{n_i^2}{N_d}$

Lots of carriers available for injection:

- $\Rightarrow V \uparrow \rightarrow$ concentration of injected carriers \uparrow
- \Rightarrow forward current can be high.
- Reverse (V < 0):

$$n(-x_p) \ll \frac{n_i^2}{N_a} \qquad p(x_n) \ll \frac{n_i^2}{N_d}$$

Few carriers available for extraction:

 \Rightarrow reverse current is small.

There is limit to how low minority carrier concentrations drop in reverse bias: zero!

 \Rightarrow reverse current saturates.

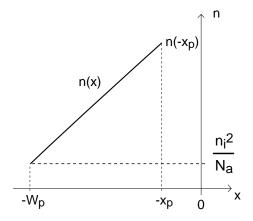
Rectification property of pn diode arises from minority-carrier boundary conditions at edges of SCR.

□ Step 2: Diffusion current in QNR:

Diffusion equation (for electrons in p-QNR):

$$J_n = qD_n \frac{dn}{dx}$$

Inside p-QNR, electrons diffuse to reach and recombine at contact $\Rightarrow J_n$ constant in p-QNR $\Rightarrow n(x)$ linear.



Boundary conditions:

$$n(x = -W_p) = n_o = \frac{n_i^2}{N_a} \qquad n(-x_p) = \frac{n_i^2}{N_a} \exp \frac{qV}{kT}$$

Electron profile:

$$n_p(x) = n_p(-x_p) + \frac{n_p(-x_p) - n_p(-W_p)}{-x_p + W_p}(x + x_p)$$

$$n_p(x) = n_p(-x_p) + \frac{n_p(-x_p) - n_p(-W_p)}{-x_p + W_p}(x + x_p)$$

Electron current density:

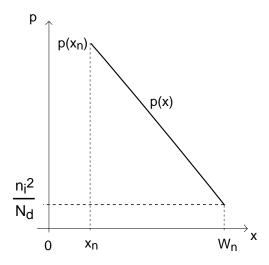
$$J_n = qD_n \frac{dn}{dx} = qD_n \frac{n_p(-x_p) - n_p(-W_p)}{W_p - x_p}$$

$$=qD_n \frac{\frac{n_i^2}{N_a} \exp \frac{qV}{kT} - \frac{n_i^2}{N_a}}{W_p - x_p}$$

or

$$J_n = q \frac{n_i^2}{N_a} \frac{D_n}{W_p - x_p} (\exp \frac{qV}{kT} - 1)$$

Similarly for hole flow in n-QNR:



Hole current density:

$$J_p = q \frac{n_i^2}{N_d} \frac{D_p}{W_n - x_n} (\exp \frac{qV}{kT} - 1)$$

□ Step 3: sum both current components:

$$J = J_n + J_p = q n_i^2 \left(\frac{1}{N_a} \frac{D_n}{W_p - x_p} + \frac{1}{N_d} \frac{D_p}{W_n - x_n} \right) \left(\exp \frac{qV}{kT} - 1 \right)$$

Current:

$$I = qAn_i^2 \left(\frac{1}{N_a} \frac{D_n}{w_p - x_p} + \frac{1}{N_d} \frac{D_p}{w_n - x_n}\right) \left(\exp \frac{qV}{kT} - 1\right)$$

often written as:

$$I = I_o(\exp\frac{qV}{kT} - 1)$$

with

$$I_o \equiv \text{saturation current [A]}$$

B.C.'s contain both forward and reverse bias⇒ equation valid in forward and reverse bias.

[will discuss this result in detail next time]

Key conclusions

- Application of voltage to pn junction results in disruption of balance between drift and diffusion in SCR:
 - in forward bias, minority carriers are *injected* into quasi-neutral regions
 - in reverse bias, minority carriers are *extracted* from quasi-neutral regions
- In forward bias, injected minority carriers recombine at surface.
- In reverse bias, extracted minority carriers are generated at surface.
- Computation of boundary conditions across SCR exploits *quasi-equilibrium*: balance between diffusion and drift in SCR disturbed very little.
- Rate limiting step to current flow: diffusion through quasi-neutral regions.
- I-V characteristics of p-n diode:

$$I = I_o(\exp\frac{qV}{kT} - 1)$$