

Lecture 3

Semiconductor Physics (II)

Carrier Transport

Outline

- Thermal Motion
- Carrier Drift
- Carrier Diffusion

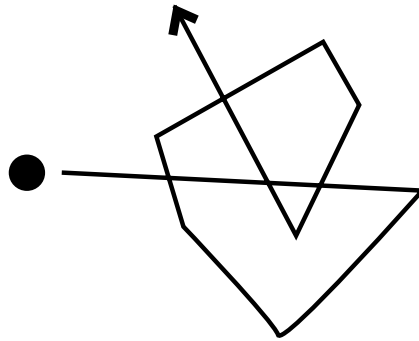
Reading Assignment:

Howe and Sodini; Chapter 2, Sect. 2.4-2.6

1. Thermal Motion

In thermal equilibrium, carriers are not sitting still:

- Undergo collisions with vibrating Si atoms (*Brownian motion*)
- Electrostatically interact with each other and with ionized (charged) dopants



Characteristic time constant of thermal motion:

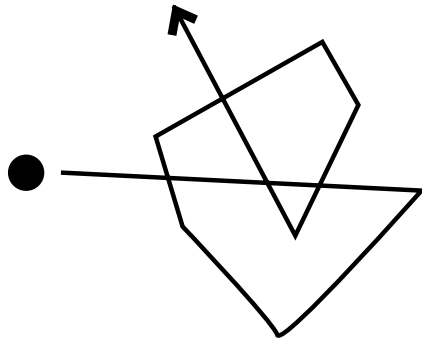
⇒ **mean free time between collisions**

$$\tau_c \equiv \textit{collision time} [s]$$

In between collisions, carriers acquire high velocity:

$$\mathbf{v}_{th} \equiv \textit{thermal velocity} [\text{cms}^{-1}]$$

.... but get nowhere!



Characteristic length of thermal motion:

$$\lambda \equiv \textit{mean free path} \text{ [cm]}$$

$$\lambda = v_{th} \tau_c$$

Put numbers for Si at room temperature:

$$\tau_c \approx 10^{-13} \text{ s}$$

$$v_{th} \approx 10^7 \text{ cm s}^{-1}$$

$$\Rightarrow \lambda \approx 0.01 \text{ } \mu\text{m}$$

For reference, state-of-the-art production MOSFET:

$$L_g \approx 0.1 \text{ } \mu\text{m}$$

\Rightarrow **Carriers undergo many collisions as they travel through devices**

2. Carrier Drift

Apply electric field to semiconductor:

E \equiv electric field [V cm⁻¹]

\Rightarrow net force on carrier

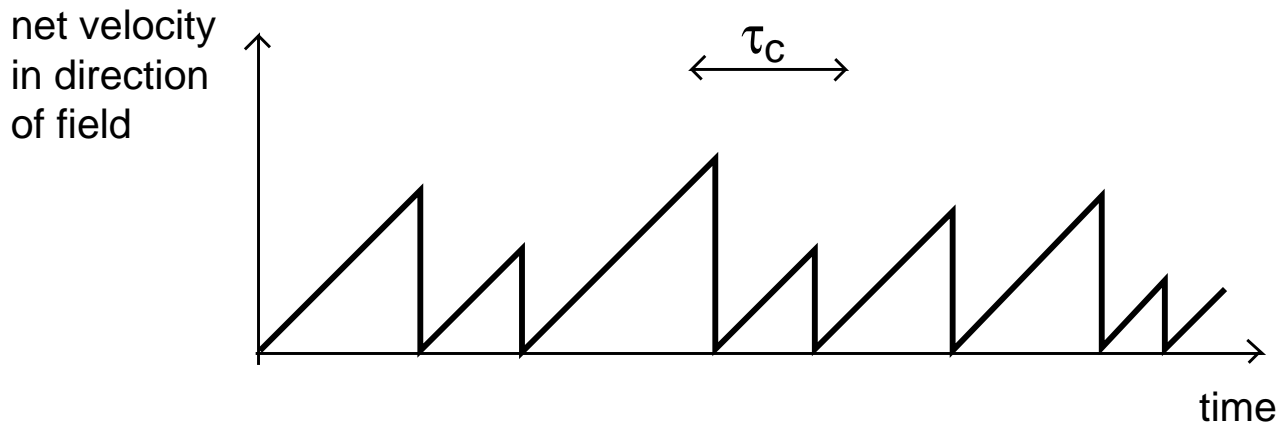
$$F = \pm qE$$



Between collisions, carriers accelerate in the direction of the electrostatic field:

$$\mathbf{v}(\mathbf{t}) = \mathbf{a} \cdot \mathbf{t} = \pm \frac{q\mathbf{E}}{m_{n,p}} \mathbf{t}$$

But there is (on the average) a collision every τ_c and the velocity is randomized:



The average net velocity in direction of the field:

$$\bar{\mathbf{v}} = \mathbf{v}_d = \pm \frac{q\mathbf{E}}{2m_{n,p}} \tau_c = \pm \frac{q\tau_c}{2m_{n,p}} \mathbf{E}$$

This is called **drift velocity** [cm s^{-1}]

Define:

$$\mu_{n,p} = \frac{q\tau_c}{2m_{n,p}} \equiv \text{mobility} [\text{cm}^2 \text{V}^{-1} \text{s}^{-1}]$$

Then, for electrons:

$$\mathbf{v}_{dn} = -\mu_n \mathbf{E}$$

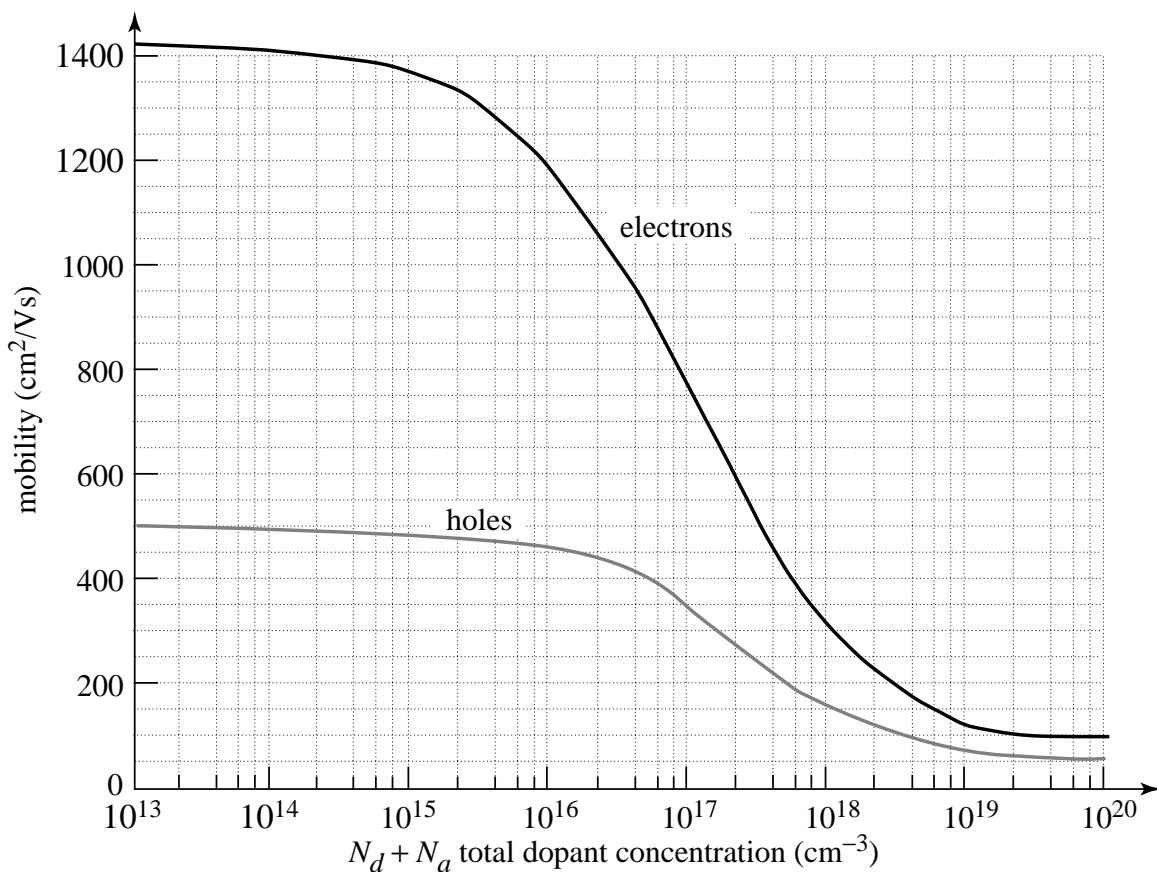
and for holes:

$$\mathbf{v}_{dp} = \mu_p \mathbf{E}$$

Mobility - is a measure of ease of carrier drift

- If $\tau_c \uparrow$, longer time between collisions $\Rightarrow \mu \uparrow$
- If $m \downarrow$, “lighter” particle $\Rightarrow \mu \uparrow$

At room temperature, mobility in Si depends on doping:



- For low doping level, μ is limited by collisions with lattice. As Temp \rightarrow **INCREASES**; $\mu \rightarrow$ **DECREASES**
- For medium doping and high doping level, μ limited by collisions with ionized impurities
- Holes “heavier” than electrons
 - For same doping level, $\mu_n > \mu_p$

Drift Current

Net velocity of charged particles \Rightarrow electric current:

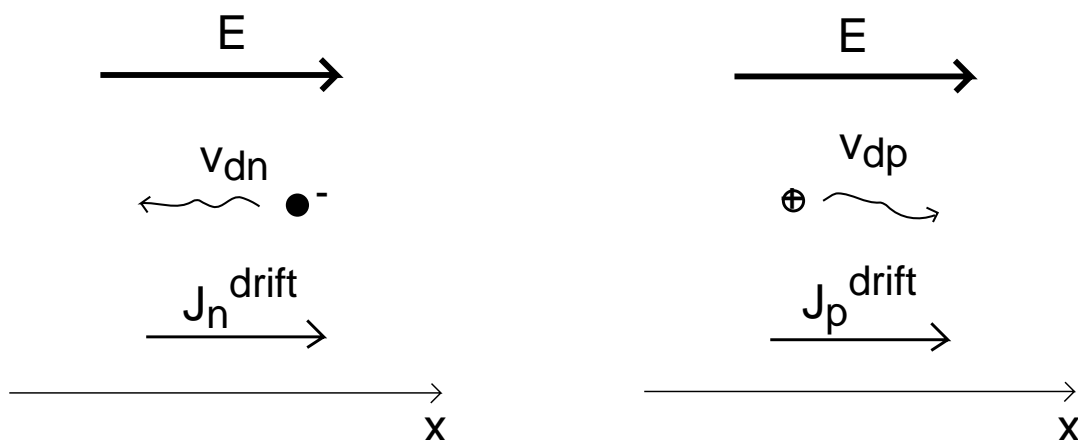
Drift current density \propto *carrier drift velocity*
 \propto *carrier concentration*
 \propto *carrier charge*

Drift current densities:

$$J_n^{drift} = -qnv_{dn} = qn\mu_n E$$

$$J_p^{drift} = qp v_{dp} = qp\mu_p E$$

Check signs:



Total Drift Current Density :

$$\mathbf{J}^{\text{drift}} = \mathbf{J}_n^{\text{drift}} + \mathbf{J}_p^{\text{drift}} = q(n\mu_n + p\mu_p)\mathbf{E}$$

Has the form of *Ohm's Law*

$$\mathbf{J} = \sigma\mathbf{E} = \frac{\mathbf{E}}{\rho}$$

Where:

$\sigma \equiv$ conductivity [$\Omega^{-1} \bullet \text{cm}^{-1}$]

$\rho \equiv$ resistivity [$\Omega \bullet \text{cm}$]

Then:

$$\sigma = \frac{1}{\rho} = q(n\mu_n + p\mu_p)$$

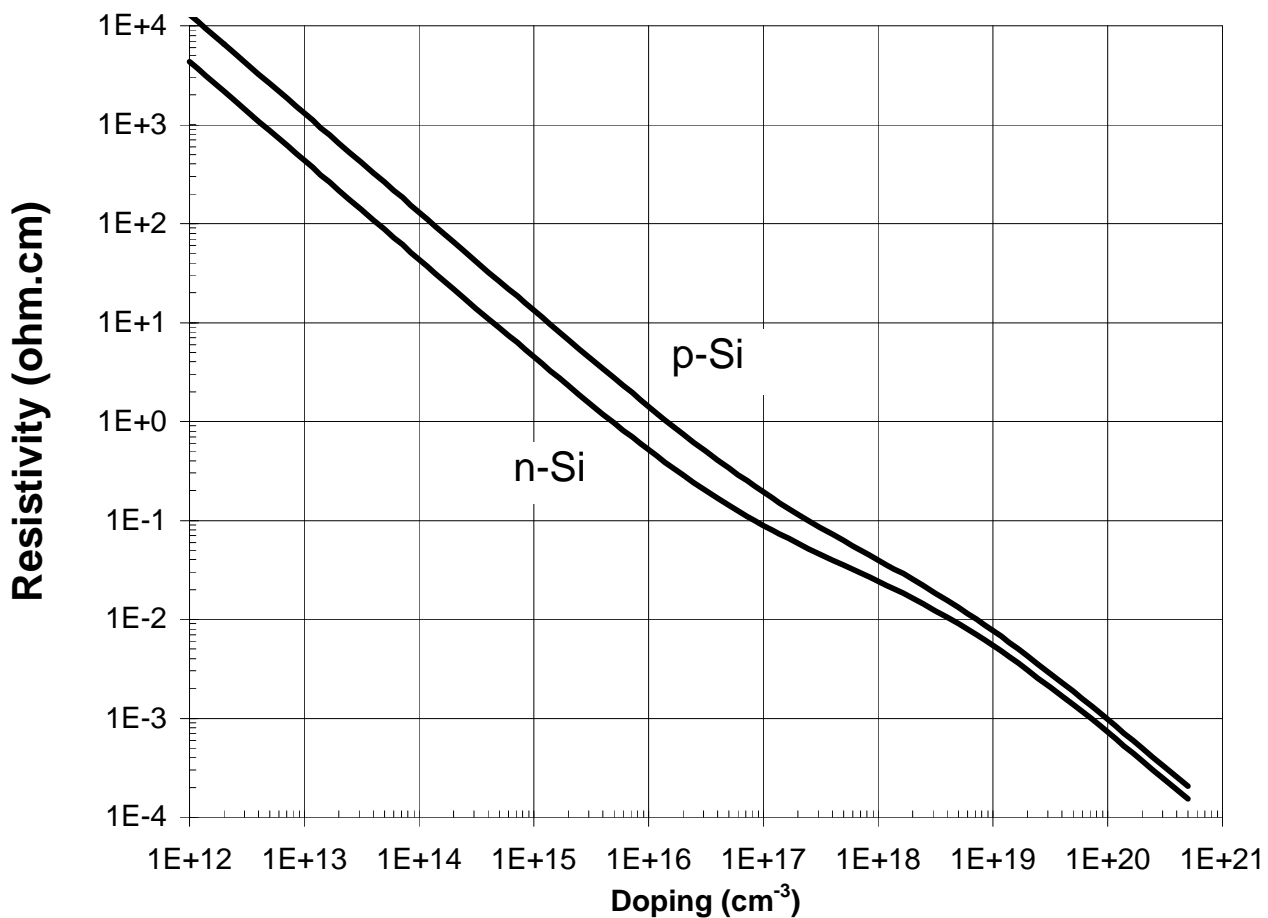
Resistivity is commonly used to specify the doping level

- In n-type semiconductor:

$$\rho_n \approx \frac{1}{qN_d\mu_n}$$

- In p-type semiconductor:

$$\rho_p \approx \frac{1}{qN_a\mu_p}$$



Numerical Example:

Si with $N_d = 3 \times 10^{16} \text{ cm}^{-3}$ at room temperature

$$\mu_n \approx 1000 \text{ cm}^2 / \text{V} \cdot \text{s}$$

$$\rho_n \approx 0.21 \Omega \cdot \text{cm}$$

$$n \approx 3 \times 10^{16} \text{ cm}^{-3}$$

Apply $E = 1 \text{ kV/cm}$

$$|v_{dn}| \approx 10^6 \text{ cm/s} \ll v_{th}$$

$$J_n^{drift} \approx qn v_{dn} = qn \mu_n E = \sigma E = \frac{E}{\rho}$$

$$J_n^{drift} \approx 4.8 \times 10^3 \text{ A/cm}^2$$

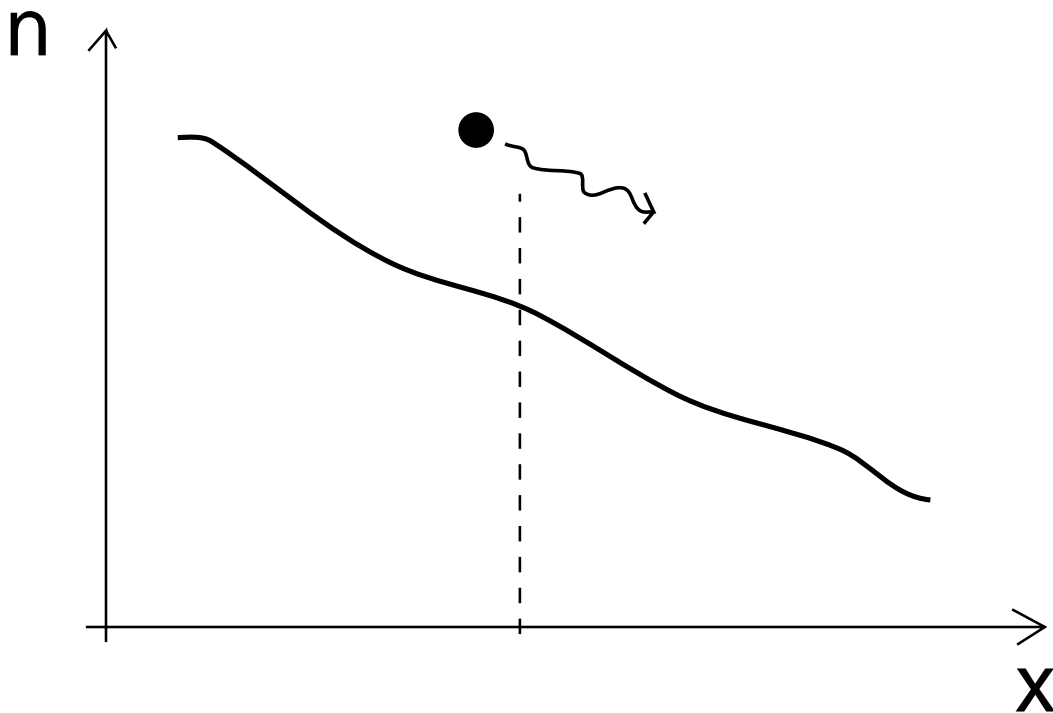
Time to drift through $L = 0.1 \text{ } \mu\text{m}$

$$t_d = \frac{L}{v_{dn}} = 10 \text{ ps}$$

fast!

3. Carrier Diffusion

Diffusion = particle movement (flux) in response to concentration gradient



Elements of diffusion:

- A medium (*Si Crystal*)
- A gradient of particles (*electrons and holes*) inside the medium
- Collisions between particles and medium send particles off in random directions
 - Overall result is to erase gradient

Fick's first law-

Key diffusion relationship

Diffusion flux \propto - concentration gradient

Flux \equiv number of particles crossing a unit area per unit time [$\text{cm}^{-2} \cdot \text{s}^{-1}$]

For Electrons:

$$F_n = -D_n \frac{dn}{dx}$$

For Holes:

$$F_p = -D_p \frac{dp}{dx}$$

$D_n \equiv$ electron diffusion coefficient [$\text{cm}^2 \text{s}^{-1}$]

$D_p \equiv$ hole diffusion coefficient [$\text{cm}^2 \text{s}^{-1}$]

D measures the *ease* of carrier diffusion in response to a concentration gradient: $D \uparrow \Rightarrow F^{\text{diff}} \uparrow$

D limited by vibration of lattice atoms and ionized dopants.

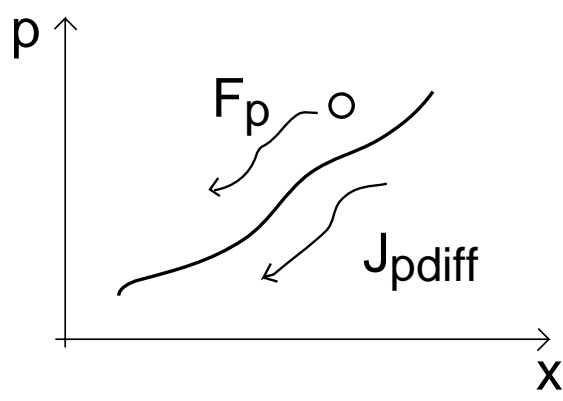
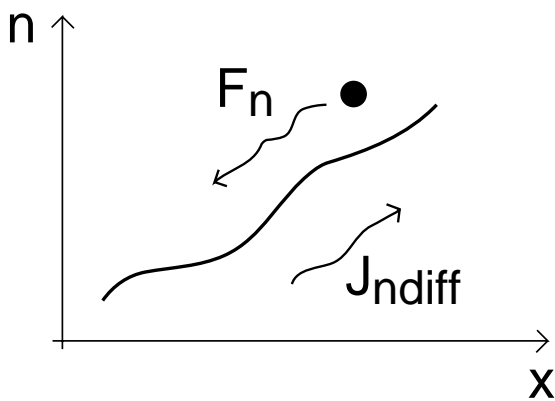
Diffusion Current

Diffusion current density = charge \times carrier flux

$$J_n^{diff} = qD_n \frac{dn}{dx}$$

$$J_p^{diff} = -qD_p \frac{dp}{dx}$$

Check signs:



Einstein relation

At the core of drift and diffusion is same physics:
collisions among particles and medium atoms
⇒ there should be a relationship between D and μ

Einstein relation [will not derive in 6.012]

$$\frac{D}{\mu} = \frac{kT}{q}$$

In semiconductors:

$$\frac{D_n}{\mu_n} = \frac{kT}{q} = \frac{D_p}{\mu_p}$$

$kT/q \equiv$ thermal voltage

At room temperature:

$$\frac{kT}{q} \approx 25 \text{ mV}$$

For example: for $N_d = 3 \times 10^{16} \text{ cm}^{-3}$

$$\mu_n \approx 1000 \text{ cm}^2 / \text{V} \cdot \text{s} \Rightarrow D_n \approx 25 \text{ cm}^2 / \text{s}$$

$$\mu_p \approx 400 \text{ cm}^2 / \text{V} \cdot \text{s} \Rightarrow D_p \approx 10 \text{ cm}^2 / \text{s}$$

Total Current Density

In general, total current can flow by drift and diffusion separately. **Total current density:**

$$\mathbf{J}_n = \mathbf{J}_n^{\text{drift}} + \mathbf{J}_n^{\text{diff}} = qn\mu_n\mathbf{E} + qD_n\frac{dn}{dx}$$

$$\mathbf{J}_p = \mathbf{J}_p^{\text{drift}} + \mathbf{J}_p^{\text{diff}} = qp\mu_p\mathbf{E} - qD_p\frac{dp}{dx}$$

$$\mathbf{J}_{\text{total}} = \mathbf{J}_n + \mathbf{J}_p$$

What did we learn today?

Summary of Key Concepts

- Electrons and holes in semiconductors are mobile and charged
 - \Rightarrow **Carriers of electrical current!**

- **Drift current**: produced by electric field

$$\mathbf{J}^{\text{drift}} \propto \mathbf{E} \quad \mathbf{J}^{\text{drift}} \propto \frac{d\phi}{dx}$$

- **Diffusion current**: produced by concentration gradient

$$\mathbf{J}^{\text{diffusion}} \propto \frac{dn}{dx}, \frac{dp}{dx}$$

- Diffusion and drift currents are sizeable in modern devices
- Carriers move fast in response to fields and gradients