

# Lecture 6

## PN Junction and MOS Electrostatics(III) Metal-Oxide-Semiconductor Structure

### Outline

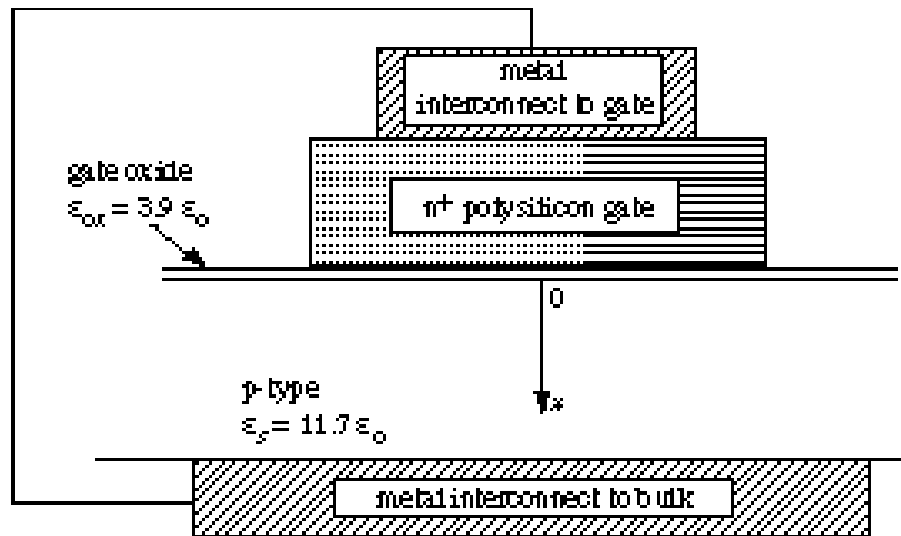
1. Introduction to MOS structure
2. Electrostatics of MOS in thermal equilibrium
3. Electrostatics of MOS with applied bias

#### **Reading Assignment:**

Howe and Sodini, Chapter 3, Sections 3.7-3.8

# 1. Introduction

## Metal-Oxide-Semiconductor structure

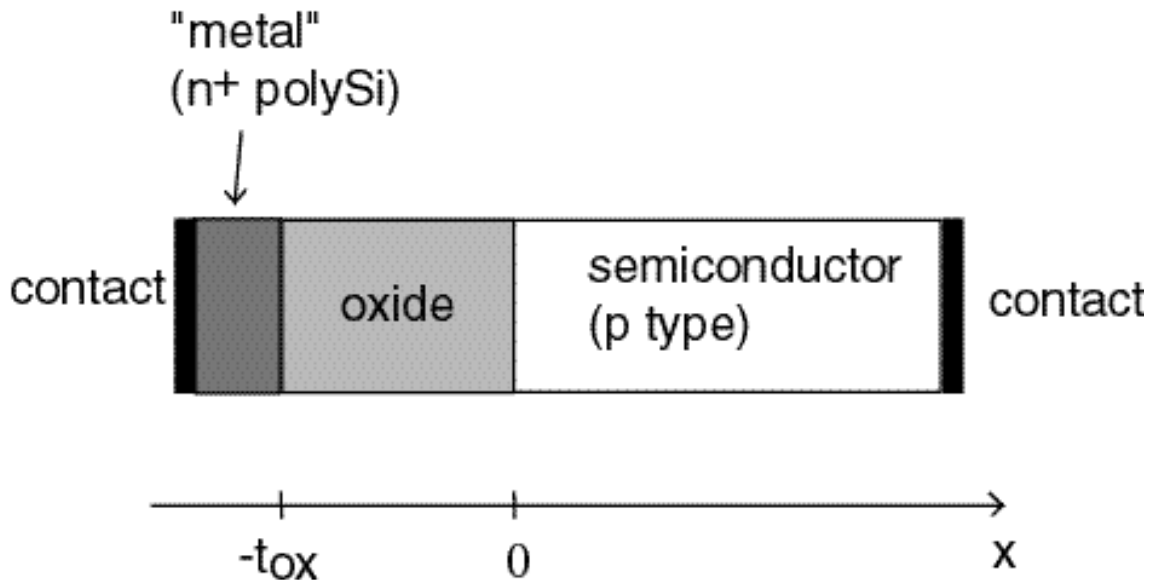


MOS at the heart of the electronics revolution:

- Digital and analog functions
  - Metal-Oxide-Semiconductor Field-Effect Transistor (MOSFET) is key element of Complementary Metal-Oxide-Semiconductor (CMOS) circuit family
- Memory function
  - Dynamic Random Access Memory (DRAM)
  - Static Random Access Memory (SRAM)
  - Non-Volatile Random Access Memory (NVRAM)
- Imaging
  - Charge Coupled Device (CCD) and CMOS cameras
- Displays
  - Active Matrix Liquid Crystal Displays (AMLCD)

## 2. MOS Electrostatics in equilibrium

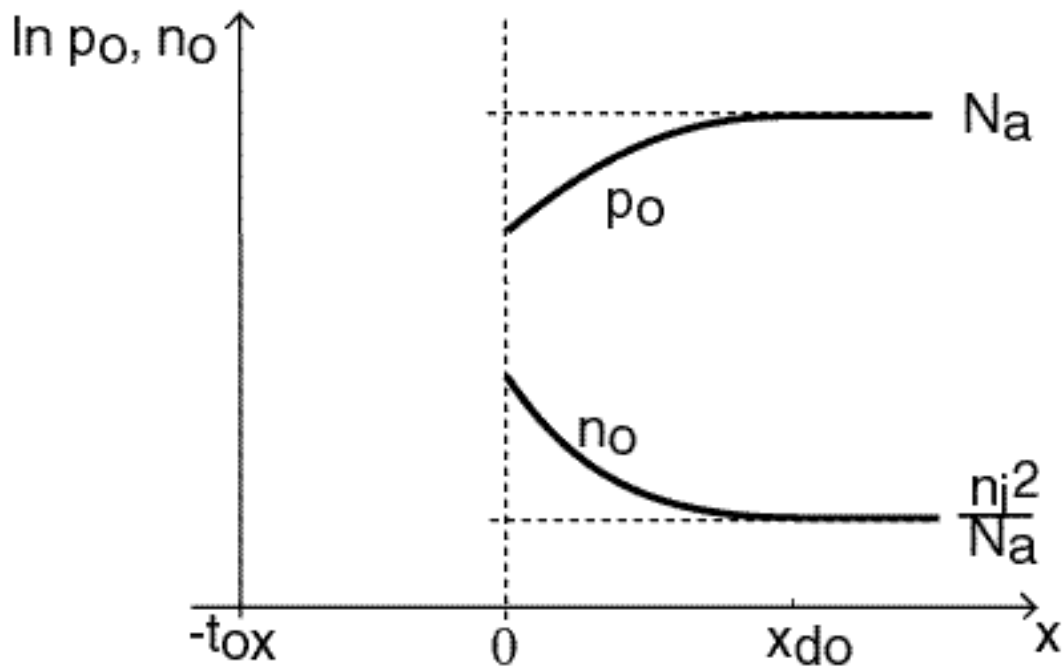
### Idealized 1D structure:



- **Metal:** does not tolerate volume charge
  - $\Rightarrow$  charge can only exist at its surface
- **Oxide:** insulator and does not have volume charge
  - $\Rightarrow$  no free carriers, no dopants
- **Semiconductor:** can have volume charge
  - $\Rightarrow$  Space charge region (SCR)

In thermal equilibrium we assume Gate contact is shorted to Bulk contact. (i. e,  $V_{GB} = 0V$ )

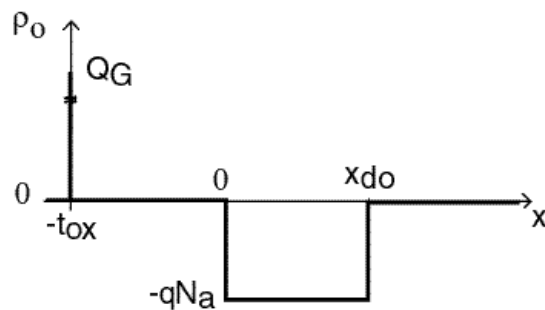
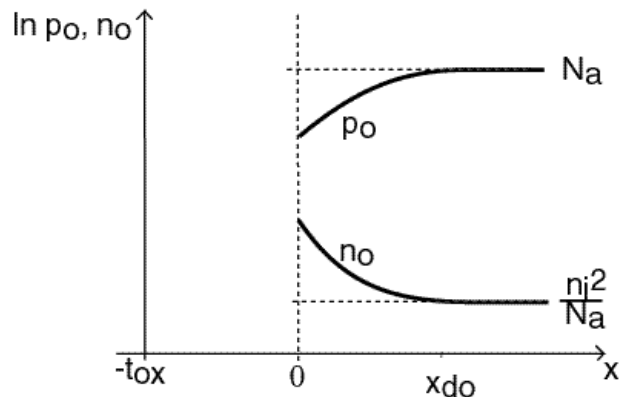
For most metals on p-Si, equilibrium achieved by electrons flowing from metal to semiconductor and holes from semiconductor to metal:



Remember:  $n_o p_o = n_i^2$

**Fewer holes near Si / SiO<sub>2</sub> interface**  
**⇒ ionized acceptors exposed (*volume charge*)**

# Space Charge Density



- In semiconductor: space-charge region close Si /SiO<sub>2</sub> interface
  - can use *depletion approximation*
- In metal: sheet of charge at metal /SiO<sub>2</sub> interface
- Overall charge neutrality

$$x = -t_{ox};$$

$$\sigma = Q_G$$

$$-t_{ox} < x < 0;$$

$$\rho_0(x) = 0$$

$$0 < x < x_{do};$$

$$\rho_0(x) = -qN_a$$

$$x > x_{do};$$

$$\rho_0(x) = 0$$

# Electric Field

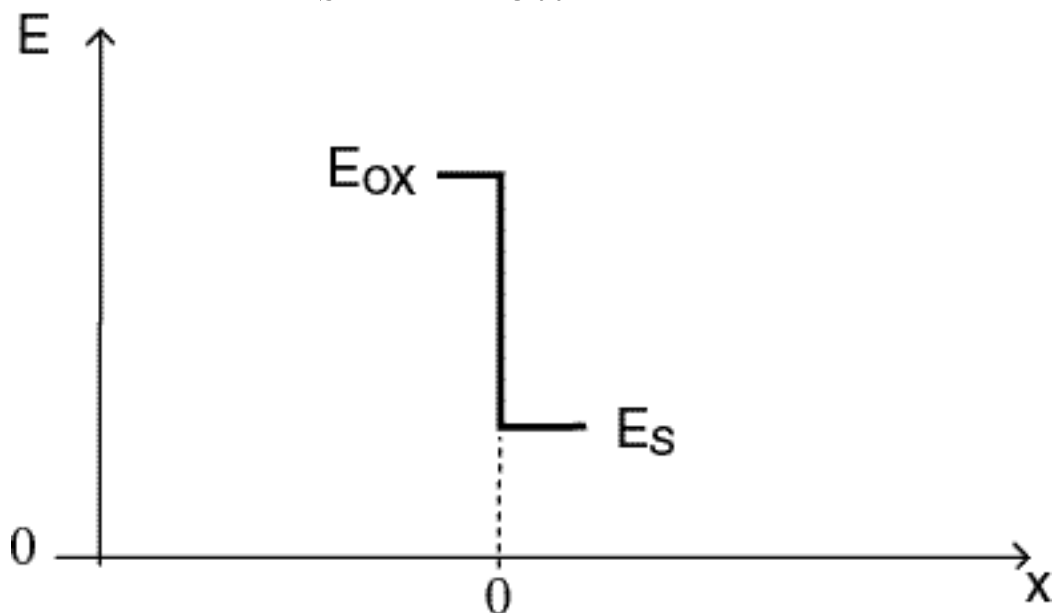
Integrate Poisson's equation

$$E_o(x_2) - E_o(x_1) = \frac{1}{\epsilon} \int_{x_1}^{x_2} \rho(x') dx'$$

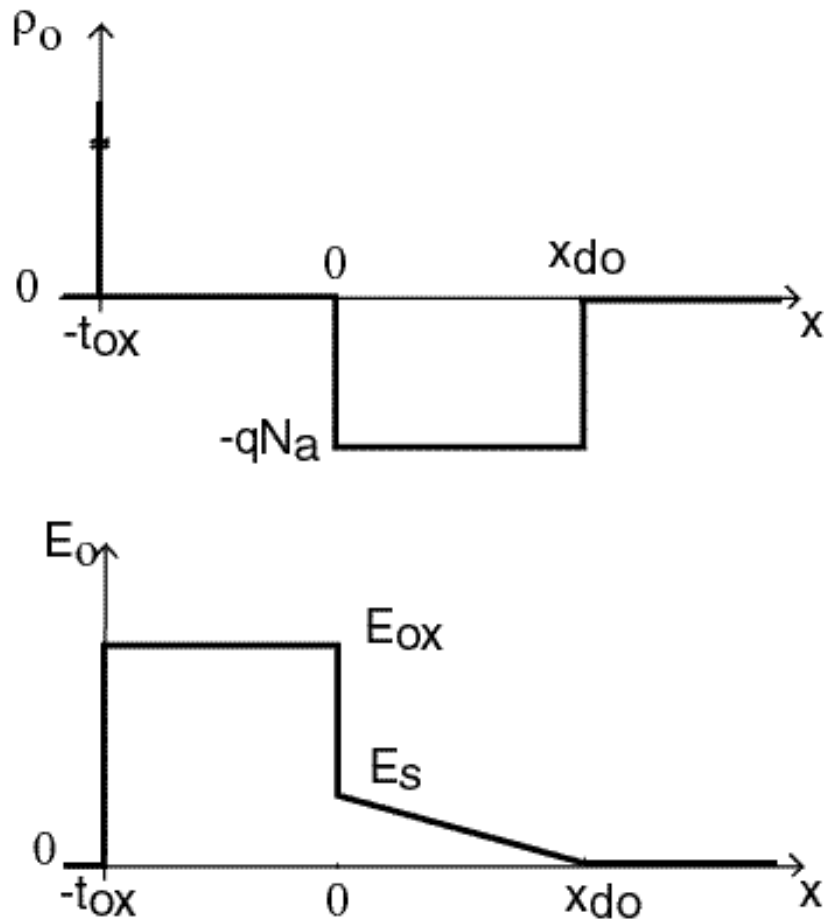
At interface between oxide and semiconductor, there is a change in **permittivity**  $\Rightarrow$  change in electric field

$$\epsilon_{ox} E_{ox} = \epsilon_s E_s$$

$$\frac{E_{ox}}{E_s} = \frac{\epsilon_s}{\epsilon_{ox}} \approx 3$$



**Start integrating from deep inside semiconductor:**



$$x > x_{do}; \quad E_0(x) = 0$$

$$0 < x < x_{do}; \quad E_0(x) - E_0(x_{do}) = \frac{1}{\epsilon_s} \int_{x_{do}}^x -qN_a dx' = -\frac{qN_a}{\epsilon_s} (x - x_{do})$$

$$-t_{ox} < x < 0; \quad E_0(x) = \frac{\epsilon_s}{\epsilon_{ox}} E_0(x = 0^+) = \frac{qN_a x_{do}}{\epsilon_{ox}}$$

$$x < -t_{ox}; \quad E(x) = 0$$

# Electrostatic Potential

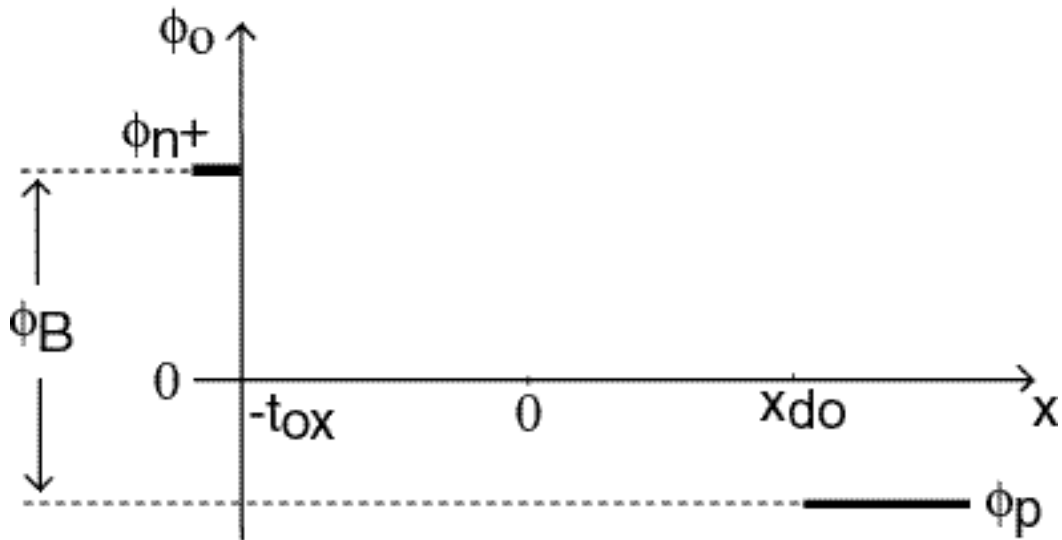
(with  $\phi = 0$  @  $n_o = p_o = n_i$ )

$$\phi = \frac{kT}{q} \cdot \ln \frac{n_o}{n_i} \quad \phi = -\frac{kT}{q} \cdot \ln \frac{p_o}{n_i}$$

In QNRs,  $n_o$  and  $p_o$  are known  $\Rightarrow$  can determine  $\phi$

$$\text{in p-QNR: } p_o = N_a \Rightarrow \phi_p = -\frac{kT}{q} \cdot \ln \frac{N_a}{n_i}$$

$$\text{in } n^+\text{-gate: } n_o = N_d^+ \Rightarrow \phi_g = \phi_{n^+}$$



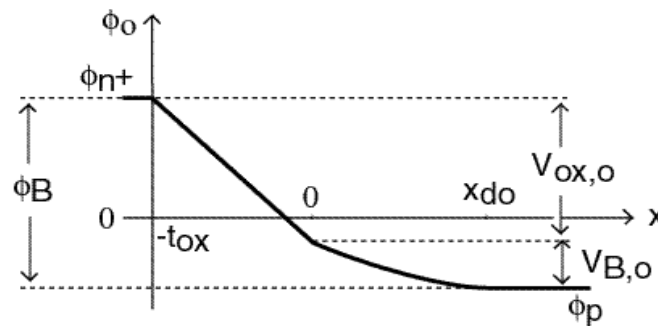
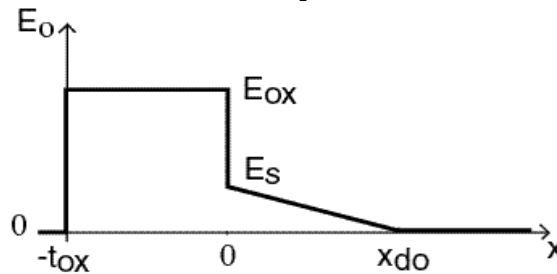
Built-in potential:

$$\phi_B = \phi_g - \phi_p = \phi_{n^+} + \frac{kT}{q} \cdot \ln \frac{N_a}{n_i}$$



To obtain  $\phi_o(x)$ , integrate  $E_o(x)$ ; start from deep inside semiconductor bulk:

$$\phi_o(x_2) - \phi_o(x_1) = - \int_{x_1}^{x_2} E_o(x') dx'$$



$$x > x_{do};$$

$$\phi_o(x) = \phi_p$$

$$0 < x < x_{do};$$

$$\phi_o(x) - \phi_o(x_{do}) = - \int_{x_{do}}^x - \frac{qN_a}{\epsilon_s} (x' - x_{do}) dx'$$

$$\phi_o(x) = \phi_p + \frac{qN_a}{2\epsilon_s} (x - x_{do})^2$$

$$AT \quad x = 0 \quad \phi_o(x) = \phi_p + \frac{qN_a}{2\epsilon_s} (x_{do})^2$$

$$-t_{ox} < x < 0;$$

$$\phi_o(x) = \phi_p + \frac{qN_a x_{do}^2}{2\epsilon_s} + \frac{qN_a x_{do}}{\epsilon_{ox}} (-x)$$

$$x < -t_{ox};$$

$$\phi_o(x) = \phi_{n^+}$$

Almost done ....

**Still do not know  $x_{do} \Rightarrow$  need one more equation**

Potential difference across structure has to add up to  $\phi_B$ :

$$\phi_B = V_{B,o} + V_{ox,o} = \frac{qN_a x_{do}^2}{2\epsilon_s} + \frac{qN_a x_{do} t_{ox}}{\epsilon_{ox}}$$

Solve quadratic equation:

$$x_{do} = \frac{\epsilon_s}{\epsilon_{ox}} t_{ox} \left[ \sqrt{1 + \frac{2\epsilon_{ox}^2 \phi_B}{q\epsilon_s N_a t_{ox}^2}} - 1 \right]$$
$$= \frac{\epsilon_s}{C_{ox}} \left[ \sqrt{1 + \frac{2C_{ox}^2 \phi_B}{q\epsilon_s N_a}} - 1 \right]$$

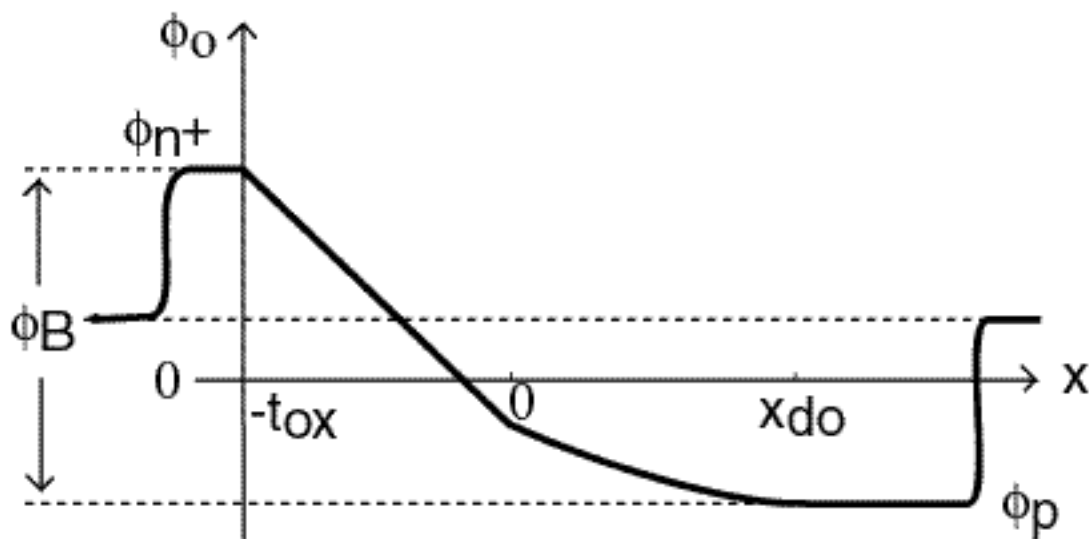
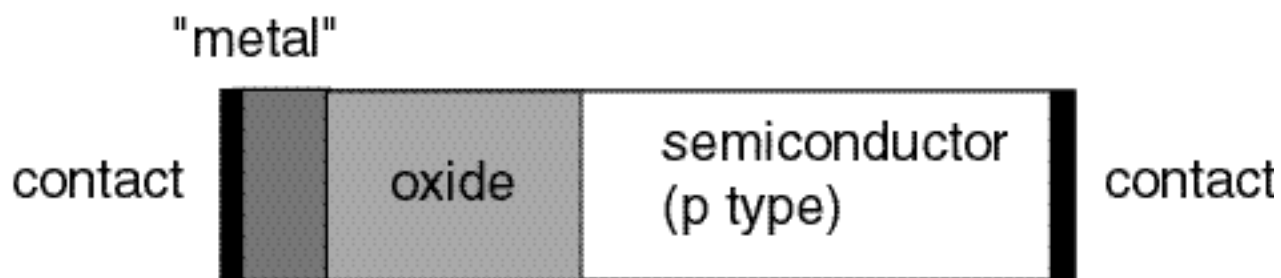
where  $C_{ox}$  is the capacitance per unit area of oxide

$$C_{ox} = \frac{\epsilon_{ox}}{t_{ox}}$$

**Now problem is completely solved!**

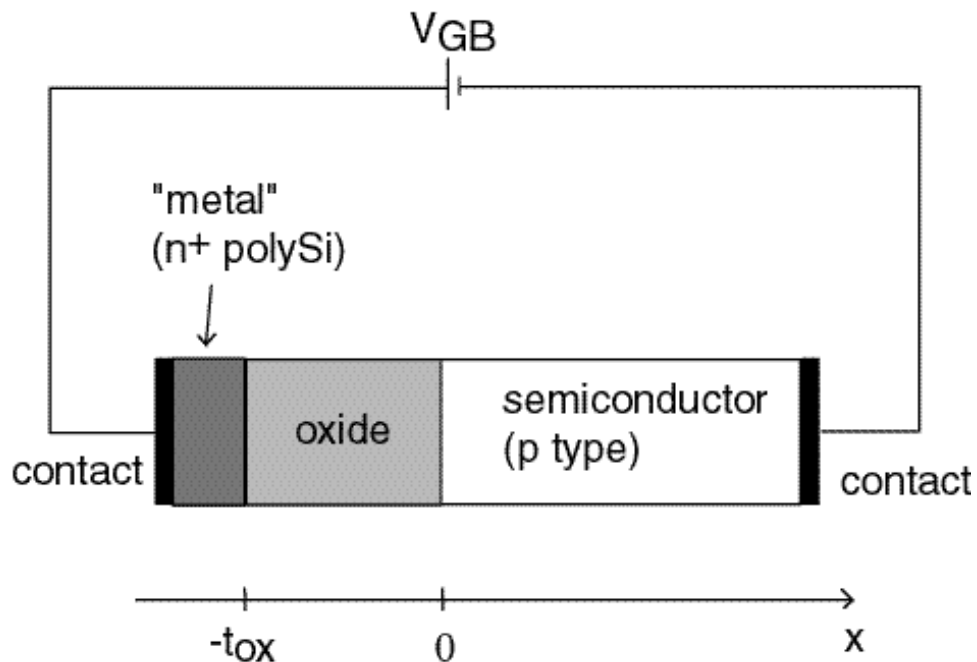
## There are also contact potentials

⇒ total potential difference from contact to contact is zero!



### 3. MOS with applied bias $V_{GB}$

Apply voltage to gate with respect to semiconductor:

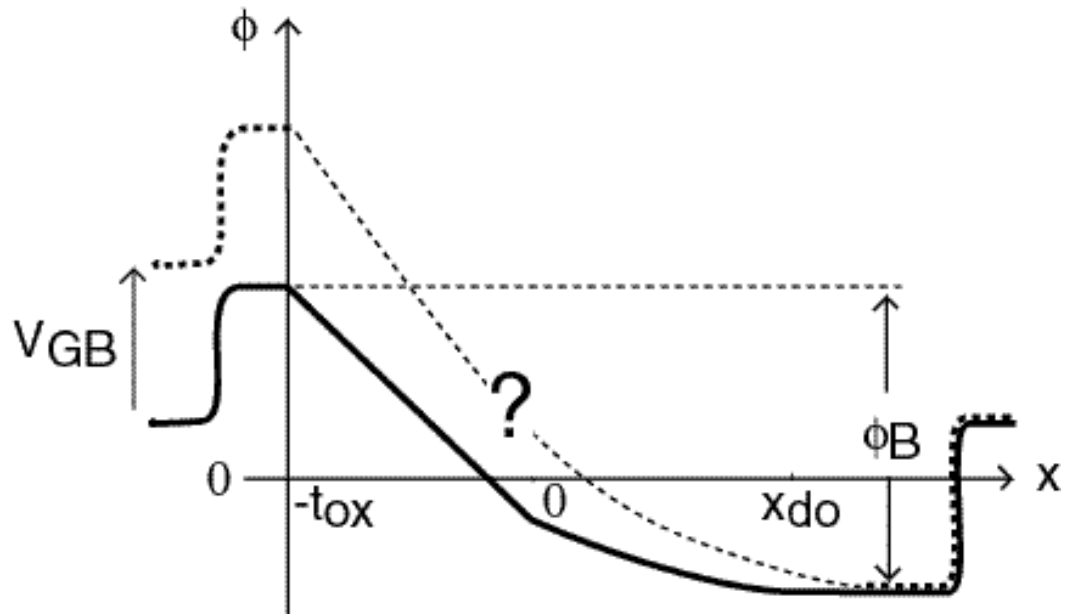


Electrostatics of MOS structure affected

$\Rightarrow$  potential difference across entire structure now  $\neq 0$

**How is potential difference accommodated?**

## Potential difference shows up across oxide and SCR in semiconductor

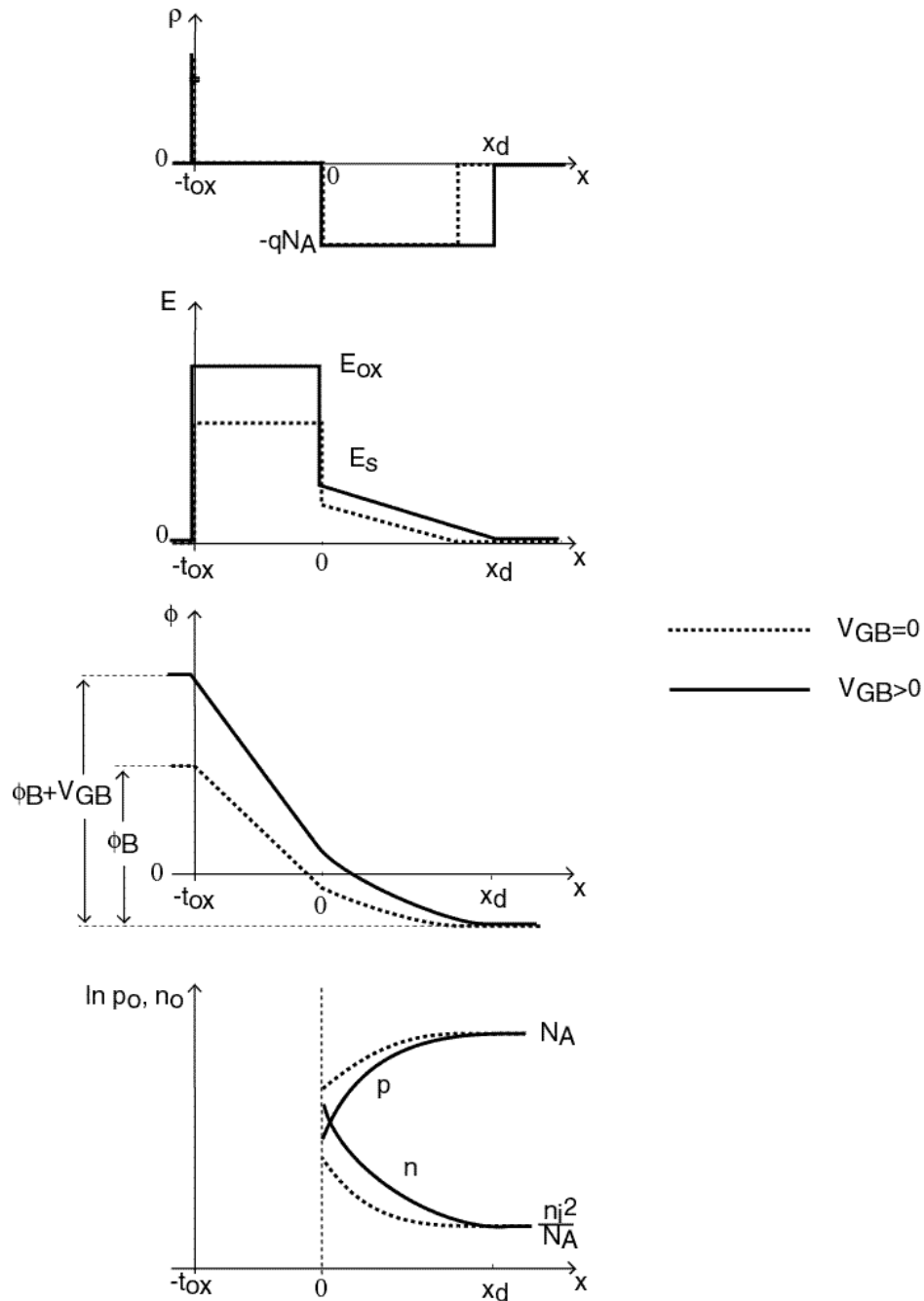


Oxide is an insulator  $\Rightarrow$  no current anywhere in structure

In SCR, quasi-equilibrium situation prevails  
 $\Rightarrow$  New balance between drift and diffusion

- Electrostatics qualitatively identical to thermal equilibrium (*but amount of charge redistribution is different*)
- **$np = n_i^2$**

**Apply  $V_{GB} > 0$ :** potential difference across structure increases  $\Rightarrow$  need larger charge dipole  $\Rightarrow$  SCR expands into semiconductor substrate:



**Simple way to remember:**

with  $V_{GB} > 0$ , gate attracts electrons and repels holes.

Qualitatively, physics unaffected by application of  $V_{GB} > 0$ . Use mathematical formulation in thermal equilibrium, but:

$$\phi_B \rightarrow \phi_B + V_{GB}$$

For example, to determine  $x_d(V_{BG})$ :

$$\begin{aligned} \phi_B + V_{GB} &= V_B(V_{GB}) + V_{ox}(V_{GB}) \\ &= \frac{qN_a x_d^2(V_{GB})}{2\epsilon_s} + \frac{qN_a x_d(V_{GB}) t_{ox}}{\epsilon_{ox}} \end{aligned}$$

$$x_d(V_{GB}) = \frac{\epsilon_s}{C_{ox}} \left[ \sqrt{1 + \frac{2C_{ox}^2 (\phi_B + V_{GB})}{\epsilon_s qN_a}} - 1 \right]$$

$$\phi(0) = \phi_s = \phi_p + \frac{qN_a x_d^2(V_{GB})}{2\epsilon_s}$$

$\phi_s$  gives n & p concentration at the surface

# What did we learn today?

## Summary of Key Concepts

- Charge redistribution in MOS structure in thermal equilibrium
  - SCR in semiconductor
  - $\Rightarrow$  **built-in potential across MOS structure.**
- In most cases, we can use depletion approximation in semiconductor SCR
- Application of voltage modulates depletion region width in semiconductor
  - **No current flows**