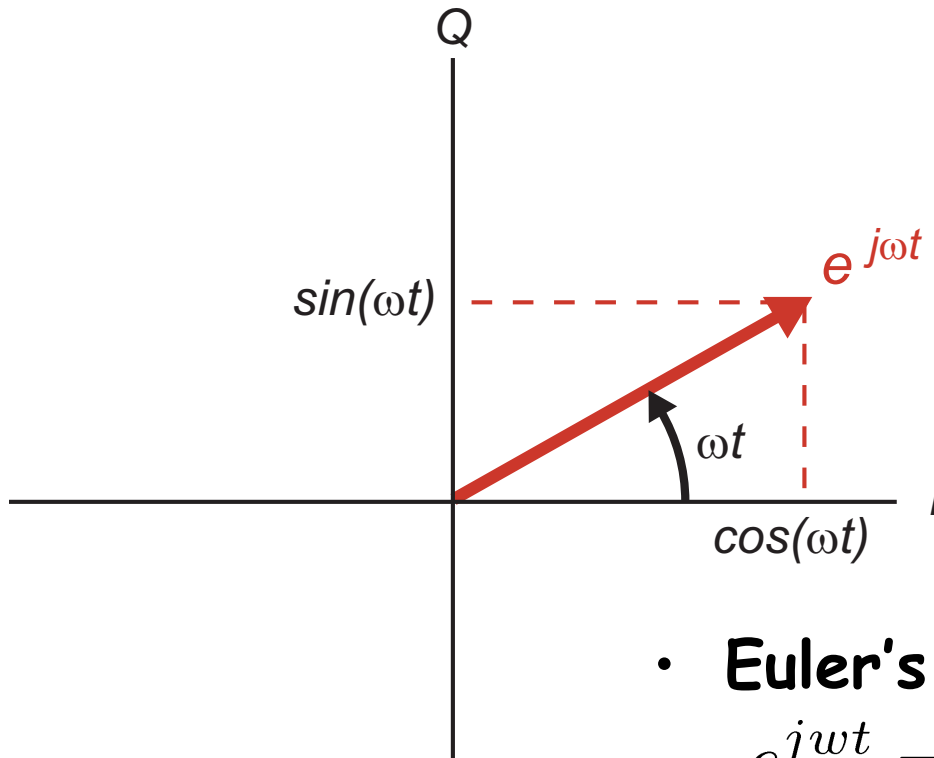


Fourier Series and Fourier Transform

- Complex exponentials
- Complex version of Fourier Series
- Time Shifting, Magnitude, Phase
- Fourier Transform

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The Complex Exponential as a Vector



Note:

$$j = \sqrt{-1}$$

- **Euler's Identity:**

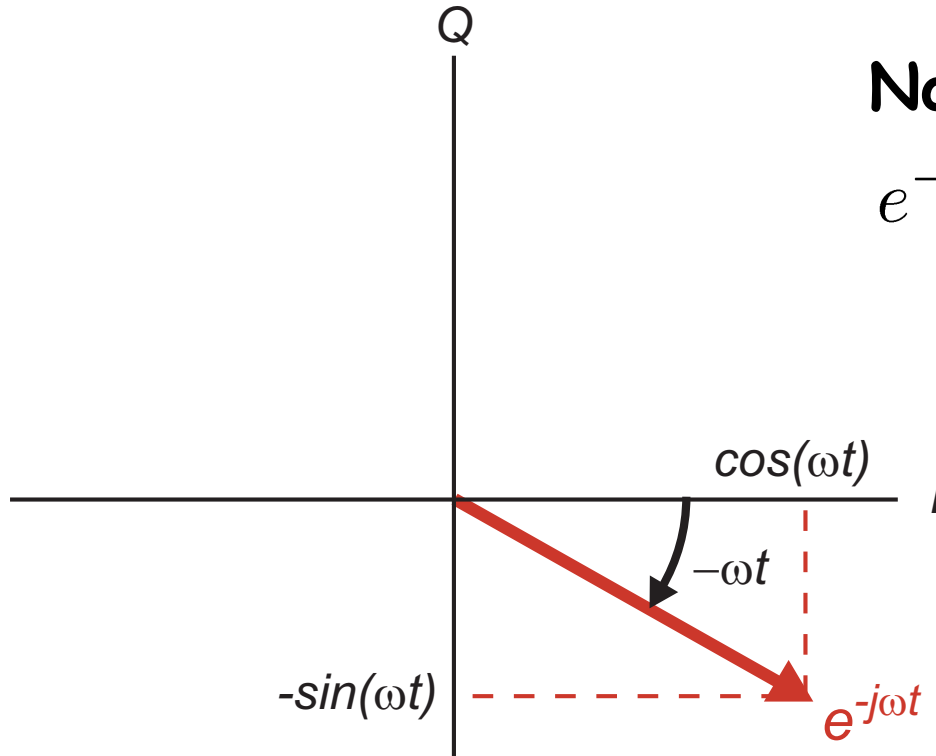
$$e^{j\omega t} = \cos(\omega t) + j \sin(\omega t)$$

- Consider I and Q as the *real* and *imaginary* parts
 - As explained later, in communication systems, I stands for *in-phase* and Q for *quadrature*
- As t increases, vector rotates *counterclockwise*
 - We consider $e^{j\omega t}$ to have *positive* frequency

The Concept of Negative Frequency

Note:

$$e^{-j\omega t} = \cos(\omega t) - j \sin(\omega t)$$

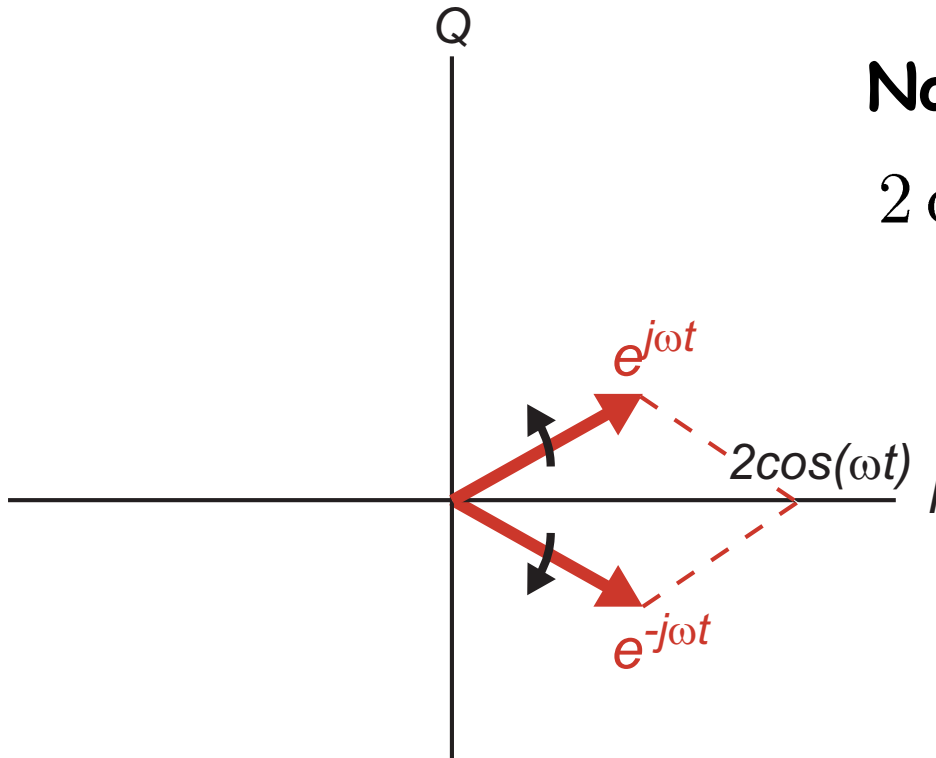


- As t increases, vector rotates *clockwise*
 - We consider $e^{-j\omega t}$ to have *negative* frequency
- Note: $A - jB$ is the *complex conjugate* of $A + jB$
 - So, $e^{-j\omega t}$ is the complex conjugate of $e^{j\omega t}$

Add Positive and Negative Frequencies

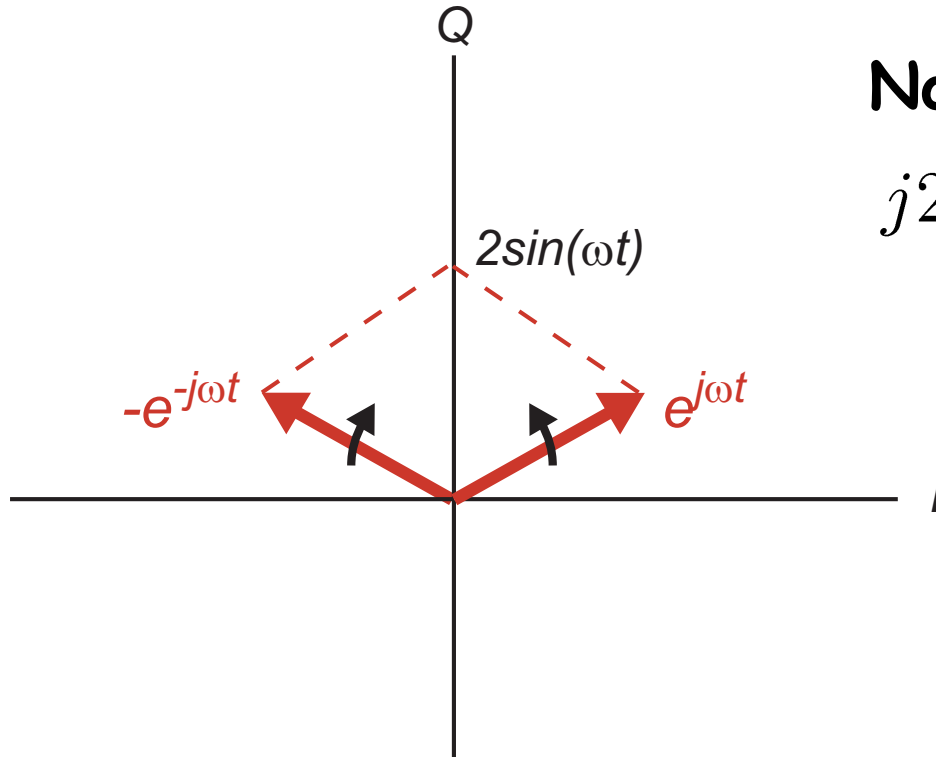
Note:

$$2 \cos(\omega t) = e^{j\omega t} + e^{-j\omega t}$$



- As t increases, the *addition of positive and negative frequency complex exponentials leads to a cosine wave*
 - Note that the resulting cosine wave is purely *real* and considered to have a *positive frequency*

Subtract Positive and Negative Frequencies

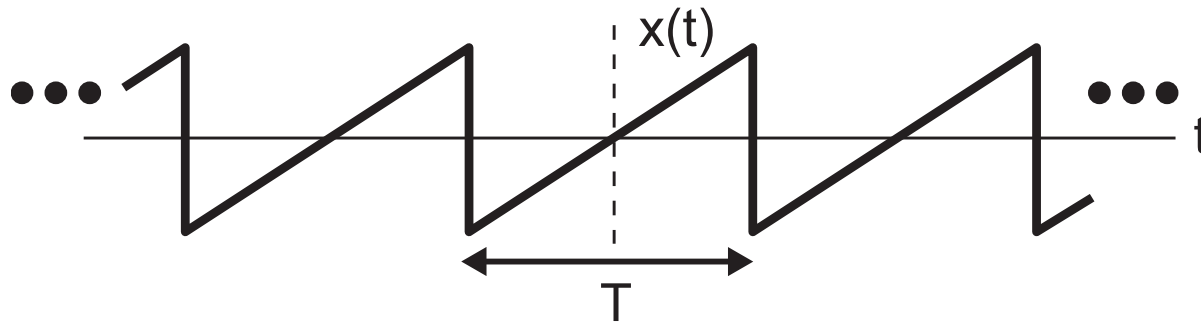


Note:

$$j2 \sin(\omega t) = e^{j\omega t} - e^{-j\omega t}$$

- As t increases, the *subtraction of positive and negative frequency complex exponentials* leads to a *sine wave*
 - Note that the resulting sine wave is purely *imaginary* and considered to have a *positive frequency*

Fourier Series



- The Fourier Series is compactly defined using complex exponentials

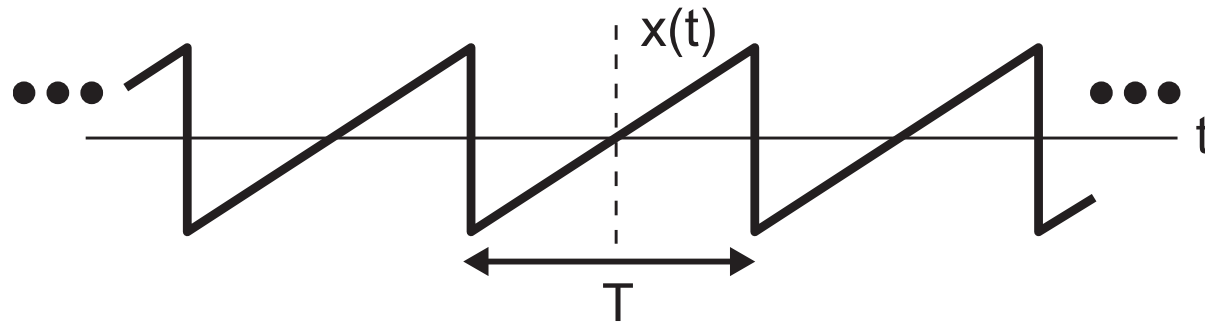
$$x(t) = \sum_{n=-\infty}^{\infty} \hat{X}_n e^{jn\omega_0 t}$$

$$\hat{X}_n = \frac{1}{T} \int_{t_0}^{t_0+T} x(t) e^{-jn\omega_0 t} dt$$

- Where:

$$\omega_0 = \frac{2\pi}{T} \qquad \hat{X}_n = A_n + jB_n$$

From The Previous Lecture



- The Fourier Series can also be written in terms of cosines and sines:

$$x(t) = a_0 + \sum_{n=1}^{\infty} a_n \cos(n\omega_0 t) + b_n \sin(n\omega_0 t)$$

where for $n > 0$:

$$a_n = \frac{2}{T} \int_{t_0}^{t_0+T} x(t) \cos(n\omega_0 t) dt, \quad b_n = \frac{2}{T} \int_{t_0}^{t_0+T} x(t) \sin(n\omega_0 t) dt$$

and where :

$$\omega_0 = \frac{2\pi}{T}, \quad a_0 = \frac{1}{T} \int_{t_0}^{t_0+T} x(t) dt$$

Compare Fourier Definitions

- Let us assume the following: $\hat{X}_n = A_n + jB_n$

$$A_n = A_{-n} \quad B_n = -B_{-n} \quad B_0 = 0$$

- Then:

$$\begin{aligned} x(t) &= \sum_{n=-\infty}^{\infty} \hat{X}_n e^{jnw_0 t} = \sum_{n=-\infty}^{\infty} A_n e^{jnw_0 t} + \sum_{n=-\infty}^{\infty} jB_n e^{jnw_0 t} \\ &= A_0 + \sum_{n=1}^{\infty} A_n (e^{jnw_0 t} + e^{-jnw_0 t}) + \sum_{n=1}^{\infty} jB_n (e^{jnw_0 t} - e^{-jnw_0 t}) \\ &= A_0 + \sum_{n=1}^{\infty} 2A_n \cos(nw_0 t) + \sum_{n=1}^{\infty} -2B_n \sin(nw_0 t) \end{aligned}$$

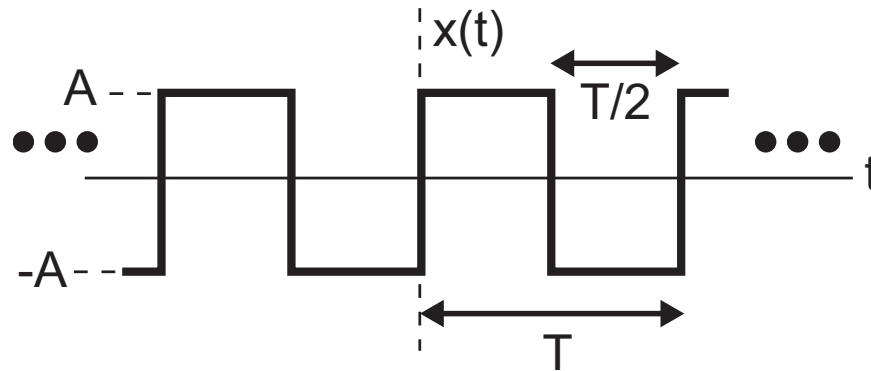
- So:

$$A_0 = a_0$$

$$2A_n = a_n$$

$$-2B_n = b_n$$

Square Wave Example



$$\hat{X}_n = \frac{1}{T} \int_{t_0}^{t_0+T} x(t) e^{-jn\omega_0 t} dt$$

$$= \frac{1}{T} \int_{-T/2}^0 -A e^{-jn\omega_0 t} dt + \frac{1}{T} \int_0^{T/2} A e^{-jn\omega_0 t} dt$$

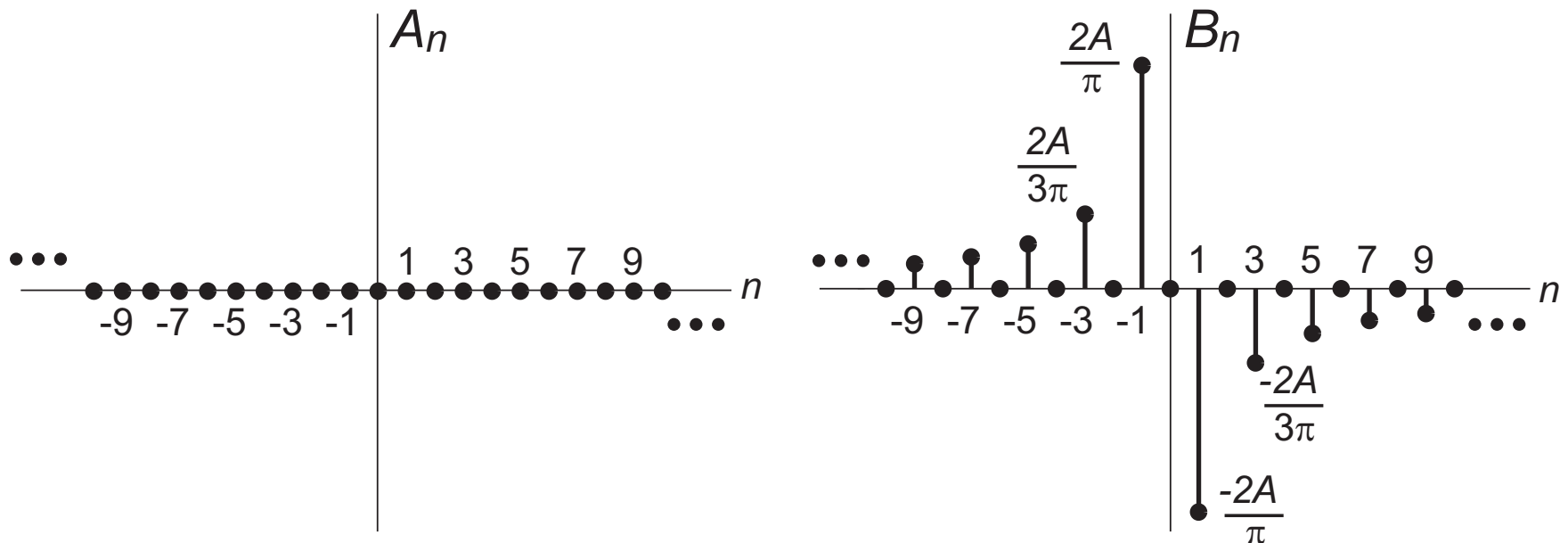
$$= \frac{1}{T} \frac{-A}{-jn\omega_0} (1 - e^{jn\omega_0 T/2}) + \frac{1}{T} \frac{A}{-jn\omega_0} (e^{-jn\omega_0 T/2} - 1)$$

$$= \frac{1}{T} \frac{2A}{jn\omega_0} (1 - \cos(n\omega_0 T/2)) = -j \frac{A}{n\pi} (1 - \cos(n\pi))$$

Graphical View of Fourier Series

- As in previous lecture, we can plot Fourier Series coefficients
 - Note that we now have *positive* and *negative* values of n
- Square wave example:

$$\hat{X}_n = A_n + jB_n = -j \frac{A}{n\pi} (1 - \cos(n\pi)) = \begin{cases} 0 & (\text{even } n) \\ j \frac{-2A}{n\pi} & (\text{odd } n) \end{cases}$$



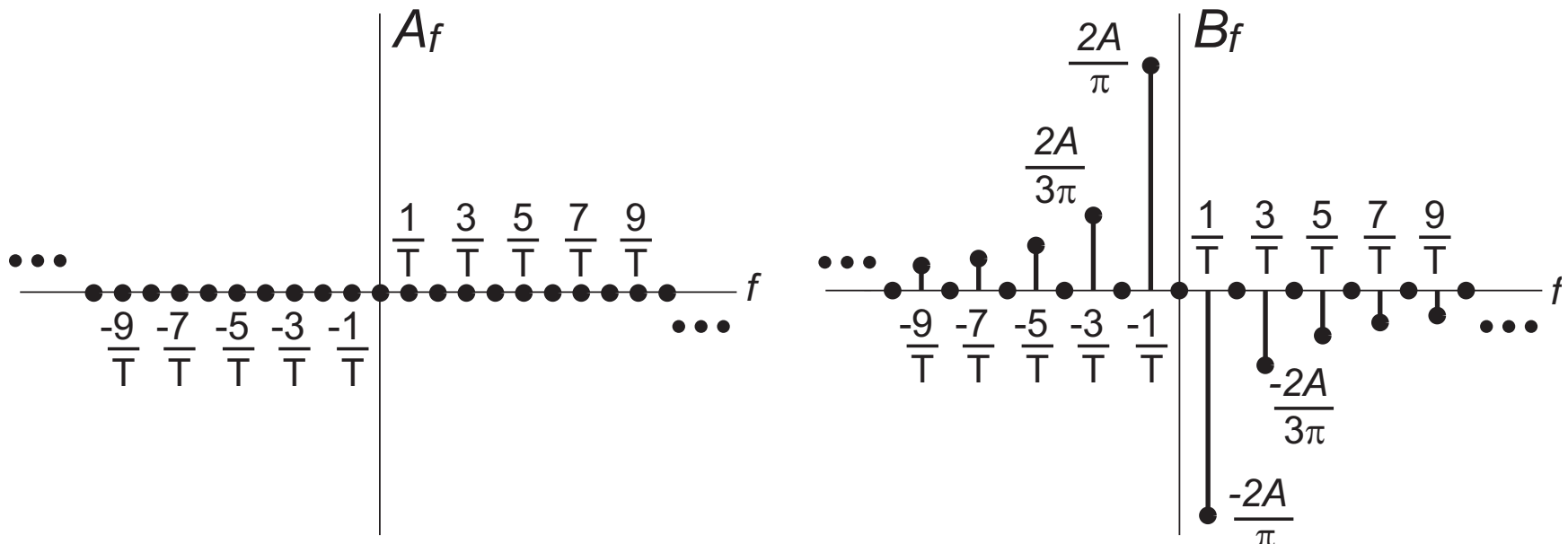
Indexing in Frequency

- A given Fourier coefficient, \hat{X}_n , represents the weight corresponding to frequency $n\omega_0$

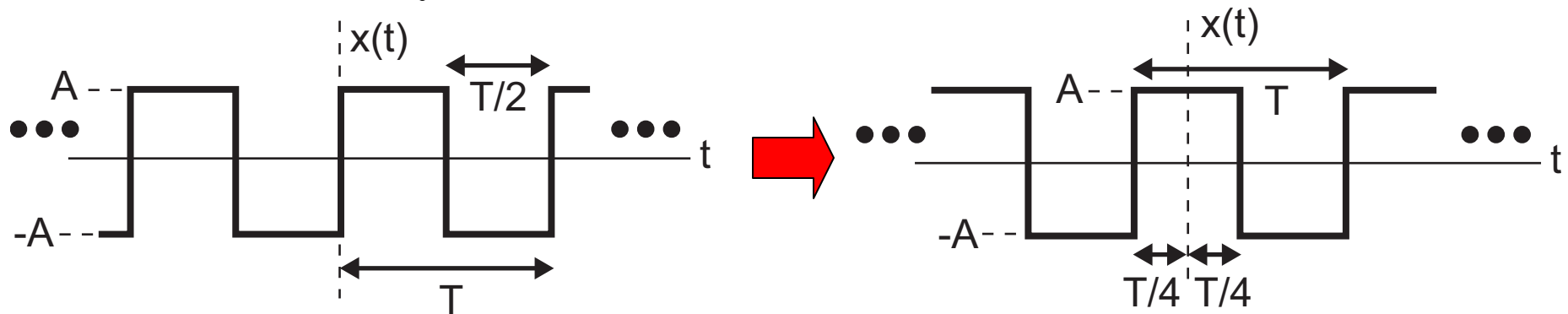
$$x(t) = \sum_{n=-\infty}^{\infty} \hat{X}_n e^{jn\omega_0 t}$$

- It is often convenient to index in *frequency (Hz)*

$$n\omega_0 = 2\pi(nf_0) = 2\pi \left(n \frac{1}{T} \right)$$



The Impact of a Time (Phase) Shift



- Consider shifting a signal $x(t)$ in time by T_d

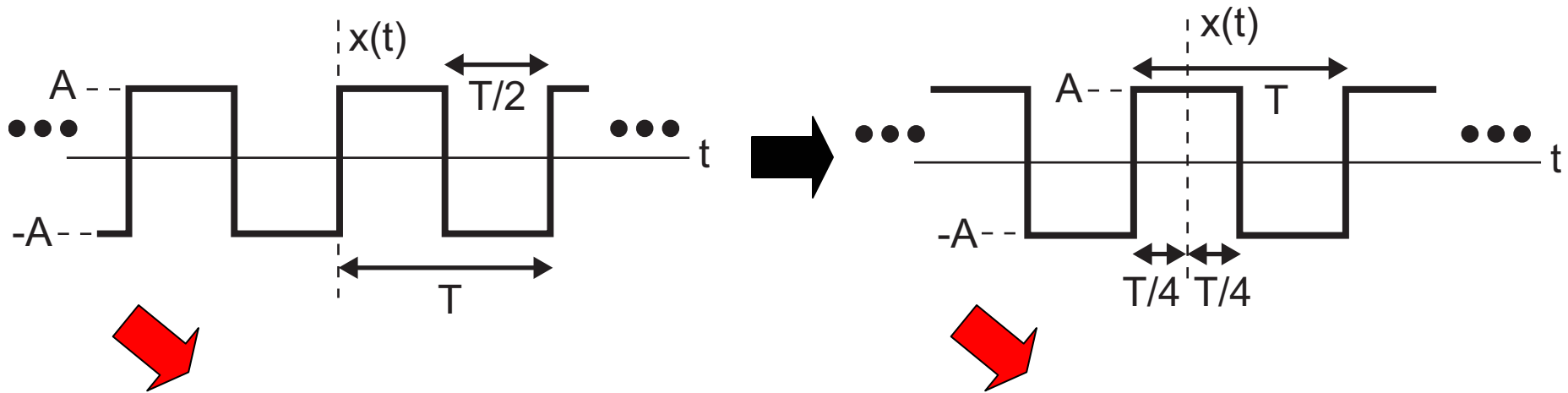
$$\hat{Y}_n = \frac{1}{T} \int_{t_o}^{t_o+T} x(t - T_d) e^{-jn\omega_o t} dt$$

- Define: $\tau = t - T_d \Rightarrow d\tau = dt$

- Which leads to: $\hat{Y}_n = \frac{1}{T} \int_{t_o+T_d}^{t_o+T+T_d} x(\tau) e^{-jn\omega_o(\tau+T_d)} d\tau$

$$= e^{-jn\omega_o T_d} \left(\frac{1}{T} \int_{t_o}^{t_o+T} x(\tau) e^{-jn\omega_o \tau} d\tau \right) = \boxed{e^{-jn\omega_o T_d} \hat{X}_n}$$

Square Wave Example of Time Shift



$$\hat{X}_n = -j \frac{A}{n\pi} (1 - \cos(n\pi))$$

$$= \begin{cases} 0 & (\text{even } n) \\ -j \frac{2A}{n\pi} & (\text{odd } n) \end{cases}$$

$$\hat{Y}_n = e^{-jn\omega_o T_d} \hat{X}_n$$

$$= e^{-jn\omega_o (-T/4)} \hat{X}_n$$

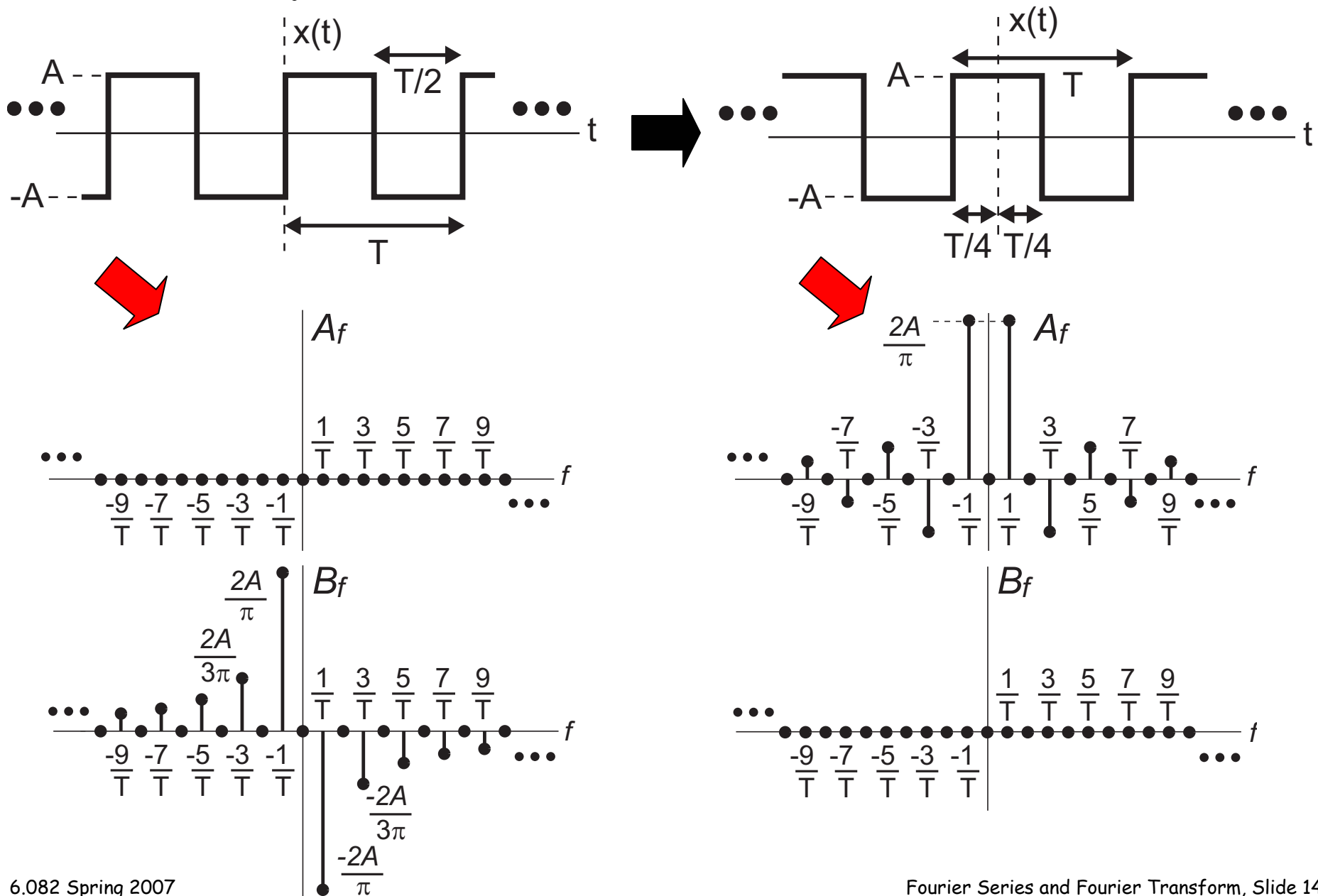
$$= e^{-jn(2\pi/T)(-T/4)} \hat{X}_n$$

$$= e^{jn\pi/2} \hat{X}_n$$

- To simplify, note that $\hat{X}_n = 0$ **except for odd n**

$$\Rightarrow \hat{Y}_n = \begin{cases} 0 & (\text{even } n) \\ j \sin(n\pi/2) \left(-j \frac{2A}{n\pi}\right) = \sin(n\pi/2) \frac{2A}{n\pi} & (\text{odd } n) \end{cases}$$

Graphical View of Fourier Series



Magnitude and Phase

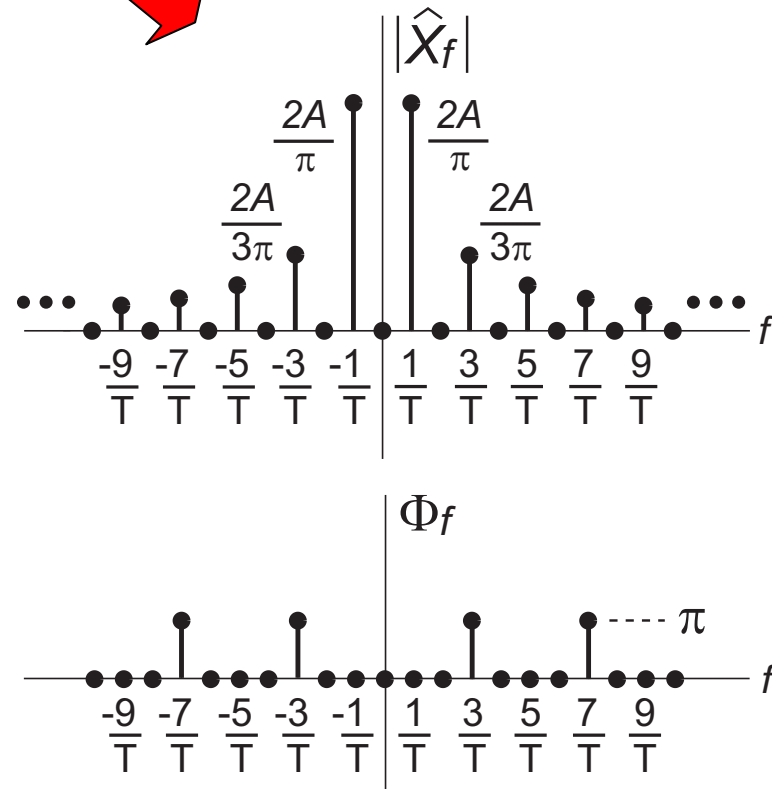
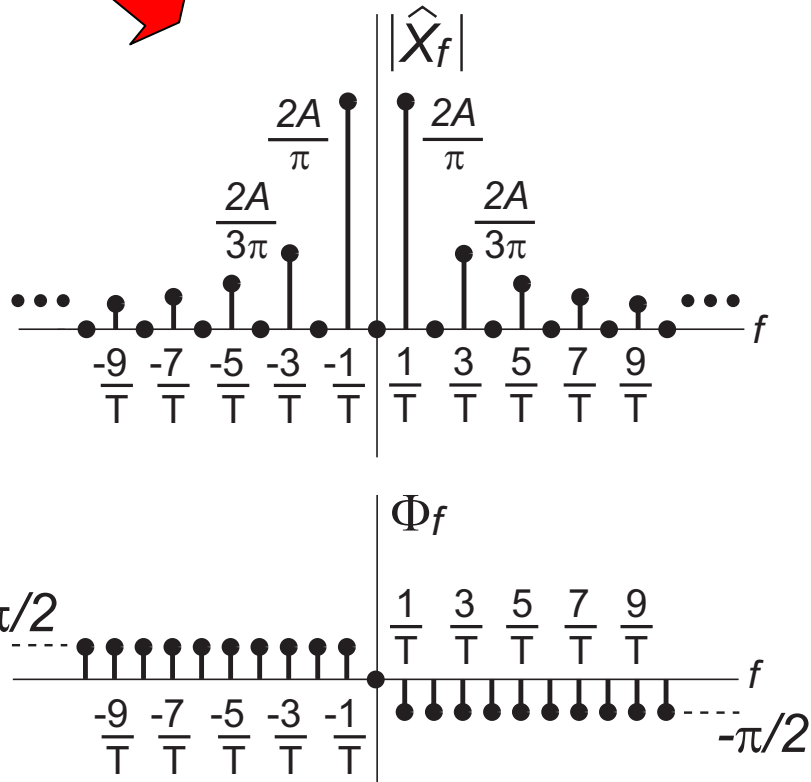
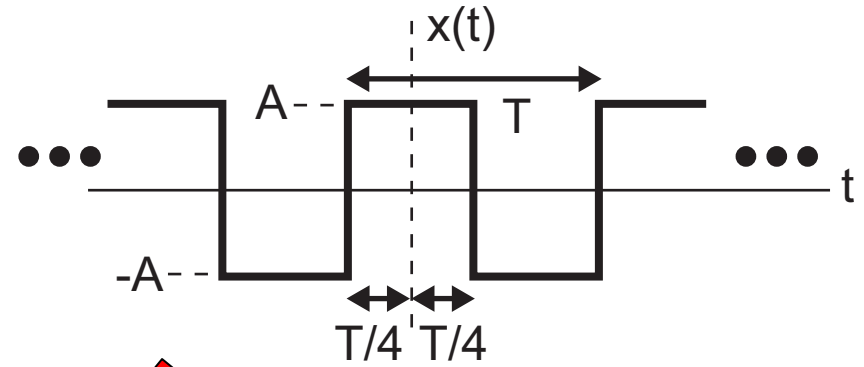
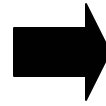
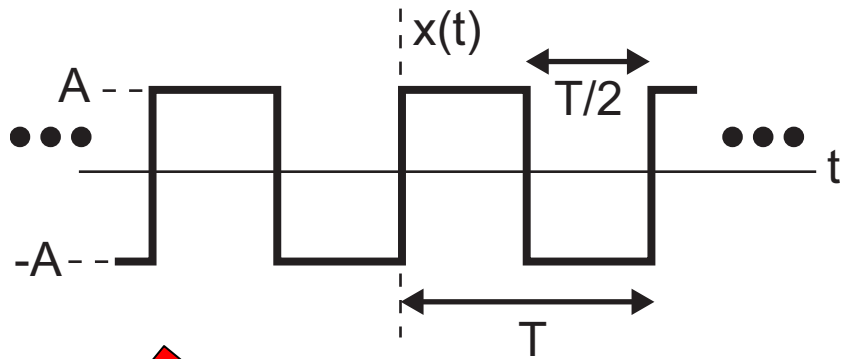
- We often want to ignore the issue of time (phase) shifts when using Fourier analysis
 - Unfortunately, we have seen that the A_n and B_n coefficients are very sensitive to time (phase) shifts
- The Fourier coefficients can also be represented in term of magnitude and phase

$$\hat{X}_n = A_n + jB_n = |\hat{X}_n|e^{j\Phi_n}$$

- where:

$$|\hat{X}_n| = \sqrt{A_n^2 + B_n^2} \quad \Phi_n = \tan^{-1} \left(\frac{B_n}{A_n} \right)$$

Graphical View of Magnitude and Phase



Does Time Shifting Impact Magnitude?

- Consider a waveform $x(t)$ along with its Fourier Series

$$x(t) \Leftrightarrow \hat{X}_n$$

- We showed that the impact of time (phase) shifting $x(t)$ on its Fourier Series is

$$x(t - T_d) \Leftrightarrow e^{-jn\omega_o T_d} \hat{X}_n$$

- We therefore see that time (phase) shifting does *not* impact the Fourier Series magnitude

$$\left| e^{-jn\omega_o T_d} \hat{X}_n \right| = \left| e^{-jn\omega_o T_d} \right| \left| \hat{X}_n \right| = \left| \hat{X}_n \right|$$

Parseval's Theorem

- The squared magnitude of the Fourier Series coefficients indicates *power* at corresponding frequencies

- *Power* is defined as:

$$\frac{1}{T} \int_{t_0}^{t_0+T} x^2(t) dt$$

$$\frac{1}{T} \int_{t_0}^{t_0+T} x^2(t) dt = \frac{1}{T} \int_{t_0}^{t_0+T} x(t) \sum_{n=-\infty}^{\infty} \hat{X}_n e^{jn\omega_0 t} dt$$

Note:

*** means**

complex

conjugate

$$= \sum_{n=-\infty}^{\infty} \hat{X}_n \frac{1}{T} \int_{t_0}^{t_0+T} x(t) e^{jn\omega_0 t} dt$$

$$= \sum_{n=-\infty}^{\infty} \hat{X}_n \hat{X}_n^* =$$

$$\sum_{n=-\infty}^{\infty} |\hat{X}_n|^2$$

The Fourier Transform

- The Fourier Series deals with *periodic* signals

$$x(t) = \sum_{n=-\infty}^{\infty} \hat{X}_n e^{jn\omega_0 t}$$

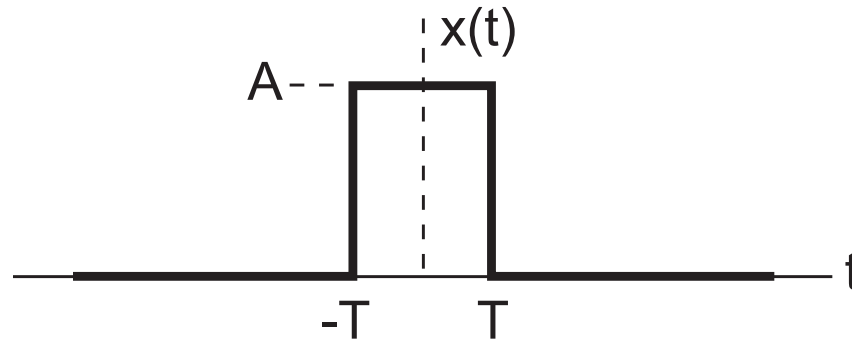
$$\hat{X}_n = \frac{1}{T} \int_{t_0}^{t_0+T} x(t) e^{-jn\omega_0 t} dt$$

- The Fourier Transform deals with *non-periodic* signals

$$x(t) = \int_{-\infty}^{\infty} X(j2\pi f) e^{j2\pi f t} df$$

$$X(j2\pi f) = \int_{-\infty}^{\infty} x(t) e^{-j2\pi f t} dt$$

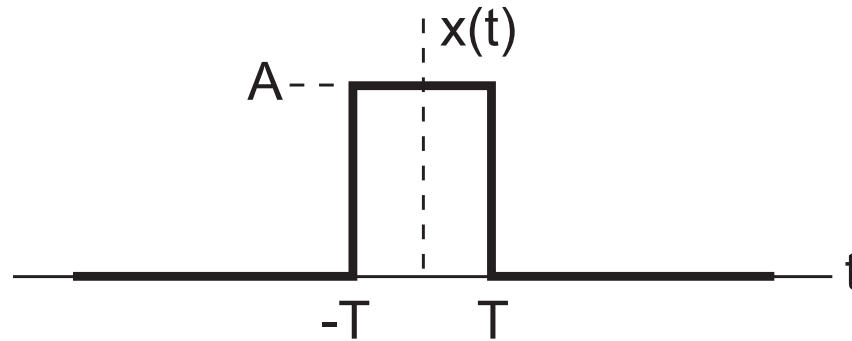
Fourier Transform Example



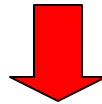
- Note that $x(t)$ is *not* periodic
- Calculation of Fourier Transform:

$$\begin{aligned} X(j2\pi f) &= \int_{-\infty}^{\infty} x(t)e^{-j2\pi ft} dt \\ &= \int_{-T}^T Ae^{-j2\pi ft} dt = \frac{A}{-j2\pi f} e^{-j2\pi ft} \Big|_{-T}^T \\ &= \frac{A \sin(2\pi fT)}{\pi f} \end{aligned}$$

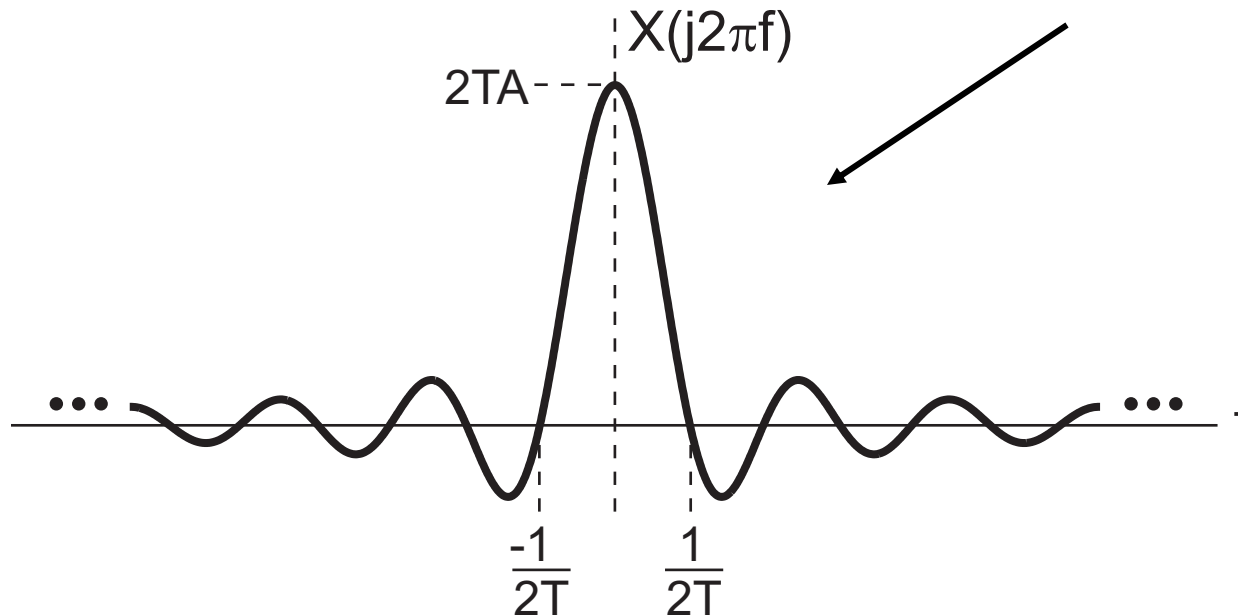
Graphical View of Fourier Transform



$$X(j2\pi f) = \frac{A \sin(2\pi fT)}{\pi f}$$



This is called a *sinc* function



Summary

- The Fourier Series can be formulated in terms of complex exponentials
 - Allows convenient mathematical form
 - Introduces concept of positive and negative frequencies
- The Fourier Series coefficients can be expressed in terms of magnitude and phase
 - Magnitude is independent of time (phase) shifts of $x(t)$
 - The magnitude squared of a given Fourier Series coefficient corresponds to the power present at the corresponding frequency
- The Fourier Transform was briefly introduced
 - Will be used to explain modulation and filtering in the upcoming lectures
 - We will provide an intuitive comparison of Fourier Series and Fourier Transform in a few weeks ...