

6.730 Physics for Solid State Applications

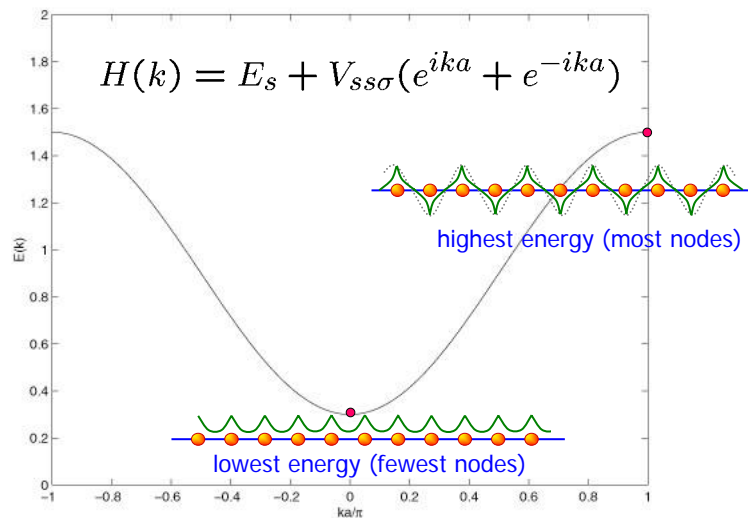
Lecture 15: Electrons in a Periodic Solid

Outline

- Review LCAO for 1-D Crystals
- Preview Problem for 2-D Crystal
- 2-D and 3-D Tight-binding
- Example: 2-D Crystal, single atom basis, 4 orbitals

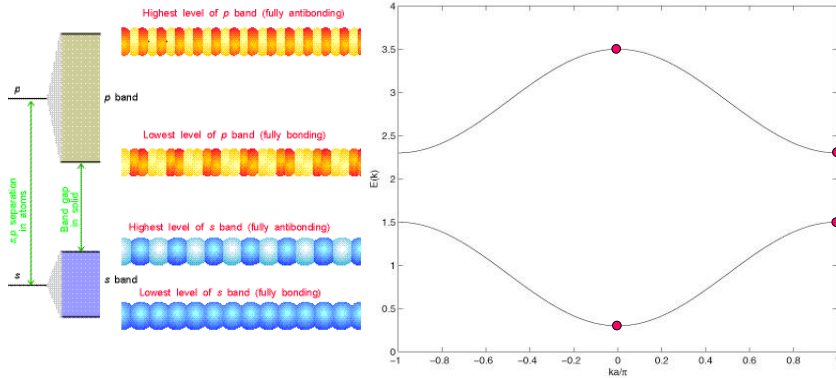
March 8, 2004

Energy Band for 1-D Lattice Single orbital, single atom basis



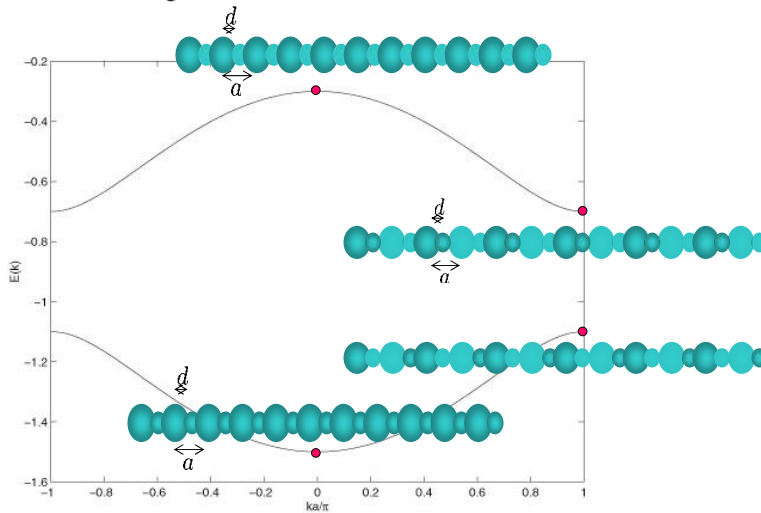
Energy Band for 1-D Lattice Two orbital, single atom basis

$$\mathbf{H}(\mathbf{k}) = \begin{matrix} \langle \phi_s | \\ \langle \phi_{p_x} | \end{matrix} \begin{pmatrix} E_s + V_{ss\sigma} (e^{ika} + e^{-ika}) & V_{sp\sigma} (e^{-ika} - e^{ika}) \\ V_{sp\sigma} (e^{ika} - e^{-ika}) & E_p + V_{pp\sigma} (e^{ika} + e^{-ika}) \end{pmatrix} \begin{matrix} | \phi_s \rangle \\ | \phi_{p_x} \rangle \end{matrix}$$



$$E_s = -12 \text{ eV}, E_p = -6 \text{ eV}, V_{ss\sigma} = -1 \text{ eV}, V_{pp\sigma} = V_{pp\pi} = +1 \text{ eV}$$

Energy Band for 1-D Lattice Single orbital, two atom basis



$$E_s = -0.9 \text{ eV}, V_{s,a} = -0.4 \text{ eV}, V_{s,a-d} = -0.2 \text{ eV}$$

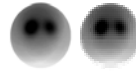
Energy Band for 1-D Lattice

Two orbital, single atom basis
Hamiltonian Matrix

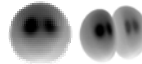
$$E_s = \langle \phi_s(\mathbf{r}) | \hat{\mathcal{H}} | \phi_s(\mathbf{r}) \rangle$$

$$E_p = \langle \phi_p(\mathbf{r}) | \hat{\mathcal{H}} | \phi_p(\mathbf{r}) \rangle$$

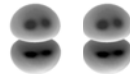
$$V_{ss\sigma} = \langle \phi_s(\mathbf{r}) | \hat{\mathcal{H}} | \phi_s(\mathbf{r} - \mathbf{a}_x) \rangle$$



$$V_{sp\sigma} = \langle \phi_s(\mathbf{r}) | \hat{\mathcal{H}} | \phi_{p_x}(\mathbf{r} - \mathbf{a}_x) \rangle$$



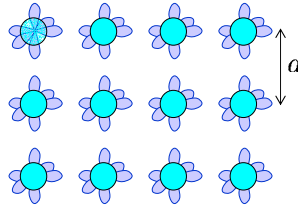
$$V_{pp\sigma} = \langle \phi_{p_x}(\mathbf{r}) | \hat{\mathcal{H}} | \phi_{p_x}(\mathbf{r} - \mathbf{a}_x) \rangle$$



$$V_{pp\pi} = \langle \phi_{p_y}(\mathbf{r}) | \hat{\mathcal{H}} | \phi_{p_y}(\mathbf{r} - \mathbf{a}_x) \rangle$$

Preview Problem: 2D Monatomic Square Crystals

$$\mathbf{c}[\mathbf{R}_n] = \begin{pmatrix} c_s \\ c_{p_x} \\ c_{p_y} \\ c_{p_z} \end{pmatrix}$$



$$\mathbf{H}(\mathbf{k}) = \sum_{\mathbf{R}_p} \tilde{\mathbf{H}}(\mathbf{R}_p) e^{-i\mathbf{k} \cdot \mathbf{R}_p}$$

$$\mathbf{H}(\mathbf{k}) = \begin{matrix} \langle \phi_s | \\ \langle \phi_{p_x} | \\ \langle \phi_{p_y} | \\ \langle \phi_{p_z} | \end{matrix} \begin{pmatrix} |\phi_s\rangle & |\phi_{p_x}\rangle & |\phi_{p_y}\rangle & |\phi_{p_z}\rangle \\ E_s + V_{ss\sigma}g_0 & ? & ? & ? \\ V_{sp\sigma}g_1^* & ? & ? & ? \\ ? & ? & ? & ? \\ ? & ? & ? & ? \end{pmatrix}$$

$$g_0 = e^{-ik_x a} + e^{ik_x a} + e^{-ik_y a} + e^{ik_y a}$$

$$g_1 = e^{-ik_x a} - e^{ik_x a}$$

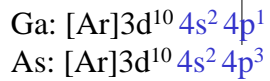
$$g_3 = e^{-ik_y a} + e^{ik_y a}$$

$$g_2 = -e^{-ik_y a} + e^{ik_y a}$$

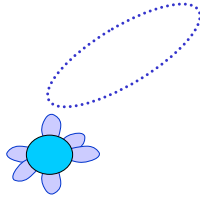
$$g_4 = e^{-ik_x a} + e^{ik_x a}$$

LCAO Basis for FCC Crystals

$$\psi(\mathbf{r}) = \sum_{\alpha} \sum_{\mathbf{n}=\mathbf{R}_{\ell}} c_{\alpha}[\mathbf{R}_{\ell}] \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\ell})$$



$$\mathbf{c}[\mathbf{R}_{\mathbf{n}}] = \begin{pmatrix} c_{s1} \\ c_{p_{x1}} \\ c_{p_{y1}} \\ c_{p_{z1}} \\ c_{s2} \\ c_{p_{x2}} \\ c_{p_{y2}} \\ c_{p_{z2}} \end{pmatrix}$$



Tight-binding for 3-D Crystals

$$\psi(\mathbf{r}) = \sum_{\alpha} \sum_{\mathbf{n}=\mathbf{R}_{\ell}} c_{\alpha}[\mathbf{R}_{\ell}] \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\ell})$$

Best estimate for energy with LCAO basis....

$$\sum_{\alpha} \sum_{\mathbf{R}_{\mathbf{m}}} \tilde{\mathbf{H}}_{\beta,\alpha}(\mathbf{R}_{\mathbf{n}}, \mathbf{R}_{\mathbf{m}}) c_{\alpha}[\mathbf{R}_{\mathbf{m}}] = E \sum_{\alpha} \sum_{\mathbf{R}_{\mathbf{p}}} \tilde{\mathbf{S}}_{\beta,\alpha}(\mathbf{R}_{\mathbf{n}}, \mathbf{R}_{\mathbf{p}}) c_{\alpha}[\mathbf{R}_{\mathbf{p}}]$$

Hamiltonian matrix....

$$\tilde{\mathbf{H}}_{\beta,\alpha}(\mathbf{R}_{\mathbf{n}}, \mathbf{R}_{\mathbf{m}}) = \langle \phi_{\beta}(\mathbf{r} - \mathbf{R}_{\mathbf{n}}) | \hat{\mathcal{H}} | \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\mathbf{m}}) \rangle$$

Overlap matrix....

$$\tilde{\mathbf{S}}_{\beta,\alpha}(\mathbf{R}_{\mathbf{n}}, \mathbf{R}_{\mathbf{p}}) = \langle \phi_{\beta}(\mathbf{r} - \mathbf{R}_{\mathbf{n}}) | \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\mathbf{p}}) \rangle$$

Tight-binding for 3-D Crystals

$$\psi(\mathbf{r}) = \sum_{\alpha} \sum_{\mathbf{n}=\mathbf{R}_{\ell}} c_{\alpha}[\mathbf{R}_{\ell}] \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\ell})$$

Since the probability of finding electrons at each lattice site is equal...

$$c[\mathbf{R}_{\mathbf{n}} + \mathbf{R}_{\ell}] = e^{i\mathbf{k} \cdot \mathbf{R}_{\ell}} c[\mathbf{R}_{\mathbf{n}}] \quad c[\mathbf{R}_{\mathbf{n}}] = e^{i\mathbf{k} \cdot \mathbf{R}_{\mathbf{n}}} \tilde{c}$$

Consequently...

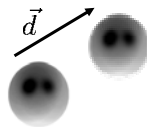
$$\mathbf{H}(\mathbf{k}) \tilde{c} = \mathbf{E} \mathbf{S}(\mathbf{k}) \tilde{c}$$

$$\mathbf{H}(\mathbf{k}) = \sum_{\mathbf{R}_{\mathbf{p}}} \tilde{\mathbf{H}}(\mathbf{R}_{\mathbf{p}}) e^{-i\mathbf{k} \cdot \mathbf{R}_{\mathbf{p}}}$$

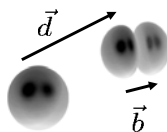
$$\mathbf{S}(\mathbf{k}) = \sum_{\mathbf{R}_{\mathbf{p}}} \tilde{\mathbf{S}}(\mathbf{R}_{\mathbf{p}}) e^{-i\mathbf{k} \cdot \mathbf{R}_{\mathbf{p}}}$$

Orbital Overlaps for 3-D Crystals

$$\langle \phi_{s1} | \hat{\mathcal{H}} | \phi_{s2} \rangle = V_{ss\sigma}$$

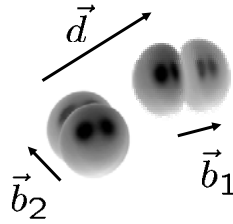


$$\langle \phi_{s1} | \hat{\mathcal{H}} | \phi_{p2} \rangle = V_{sp\sigma} \hat{\mathbf{d}} \cdot \hat{\mathbf{b}}$$



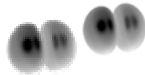
\vec{b} distance from positive to negative lobe of p-orbital

Orbital Overlaps for 3-D Crystals

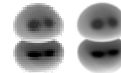


$$\langle \phi_{p1} | \hat{H} | \phi_{p2} \rangle = (\hat{\mathbf{d}} \cdot \hat{\mathbf{b}}_1)(\hat{\mathbf{d}} \cdot \hat{\mathbf{b}}_2) V_{pp\sigma} + [\hat{\mathbf{b}}_1 - \hat{\mathbf{d}}(\hat{\mathbf{b}}_1 \cdot \hat{\mathbf{d}})] \cdot [\hat{\mathbf{b}}_2 - \hat{\mathbf{d}}(\hat{\mathbf{b}}_2 \cdot \hat{\mathbf{d}})] V_{pp\pi}$$

$V_{pp\sigma}$



$V_{pp\pi}$



Orbital Overlaps for 3-D Crystals Diamond and Zincblende

$$V_{\ell\ell'm} = \eta_{\ell\ell'm} \frac{\hbar^2}{2md^2}$$

$$\eta_{ss\sigma} = -1.40 \approx -9\pi^2/64$$

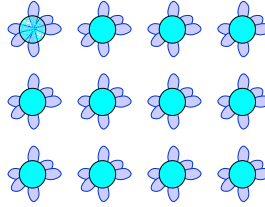
$$\eta_{pp\sigma} = 3.24$$

$$\eta_{sp\sigma} = 1.84$$

$$\eta_{pp\pi} = -0.81$$

2D Monatomic Square Crystals

$$\mathbf{c}[\mathbf{R}_n] = \begin{pmatrix} c_s \\ c_{p_x} \\ c_{p_y} \\ c_{p_z} \end{pmatrix}$$



$$\mathbf{H}(\mathbf{k}) = \begin{pmatrix} \langle \phi_s | & \langle \phi_{p_x} | & \langle \phi_{p_y} | & \langle \phi_{p_z} | \\ \begin{pmatrix} E_s + V_{ss\sigma}g_0 & V_{sp\sigma}g_1 & V_{sp\sigma}g_2^* & 0 \\ V_{sp\sigma}g_1^* & E_p + V_{pp\pi}g_3 + V_{pp\sigma}g_4 & 0 & 0 \\ V_{sp\sigma}g_2 & 0 & E_p + V_{pp\pi}g_4 + V_{pp\sigma}g_3 & 0 \\ 0 & 0 & 0 & E_p + V_{pp\pi}g_0 \end{pmatrix} \end{pmatrix}$$

$$g_0 = e^{-ik_x a} + e^{ik_x a} + e^{-ik_y a} + e^{ik_y a}$$

$$g_1 = e^{-ik_x a} - e^{ik_x a}$$

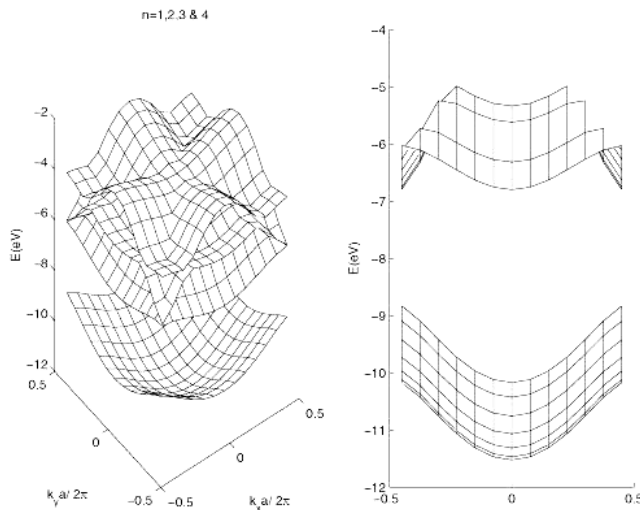
$$g_3 = e^{-ik_y a} + e^{ik_y a}$$

$$g_2 = -e^{-ik_y a} + e^{ik_y a}$$

$$g_4 = e^{-ik_x a} + e^{ik_x a}$$

2D Monatomic Square Crystals

Dispersion Relations

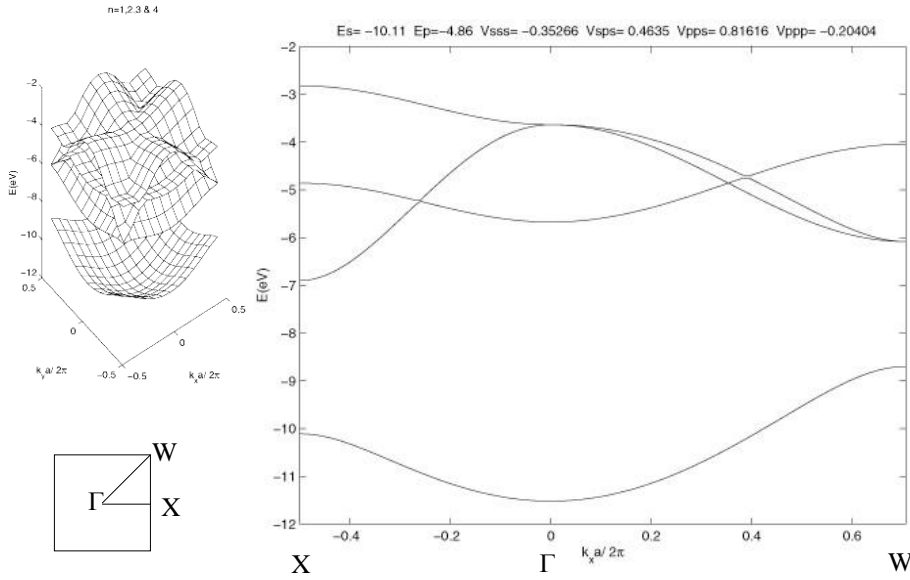


$$E_s = -10.11 \text{ eV}$$

$$E_p = -4.86 \text{ eV}$$

$$a = 5.5 \text{ \AA}$$

2D Monatomic Square Crystals Dispersion Relations



2D Monatomic Square Crystals Dispersion Relations at $\Gamma=0$

$$\mathbf{H}(\mathbf{k}) = \begin{pmatrix} \langle \phi_s | & \langle \phi_{px} | & \langle \phi_{py} | & \langle \phi_{pz} | \\ \left(\begin{array}{cccc} E_s + V_{ss\sigma}g_0 & V_{sp\sigma}g_1 & V_{sp\sigma}g_2^* & 0 \\ V_{sp\sigma}g_1^* & E_p + V_{pp\pi}g_3 + V_{pp\sigma}g_4 & 0 & 0 \\ V_{sp\sigma}g_2 & 0 & E_p + V_{pp\pi}g_4 + V_{pp\sigma}g_3 & 0 \\ 0 & 0 & 0 & E_p + V_{pp\pi}g_0 \end{array} \right) & & & \end{pmatrix}$$

$$g_0 = e^{-ik_x a} + e^{ik_x a} + e^{-ik_y a} + e^{ik_y a} \rightarrow 4$$

$$g_1 = e^{-ik_x a} - e^{ik_x a} \rightarrow 0$$

$$g_2 = -e^{-ik_y a} + e^{ik_y a} \rightarrow 0$$

$$g_3 = e^{-ik_y a} + e^{ik_y a} \rightarrow 2$$

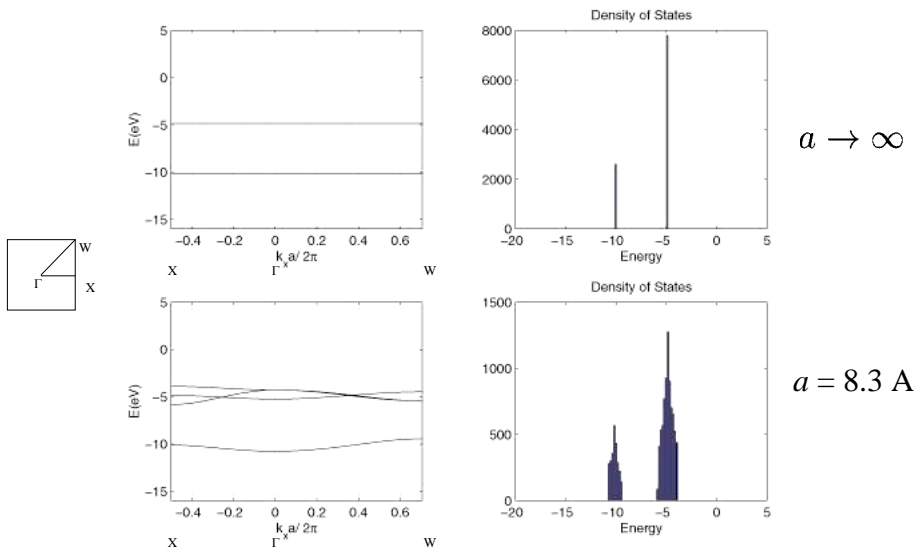
$$g_4 = e^{-ik_x a} + e^{ik_x a} \rightarrow 2$$

2D Monatomic Square Crystals Dispersion Relations at $\Gamma=0$

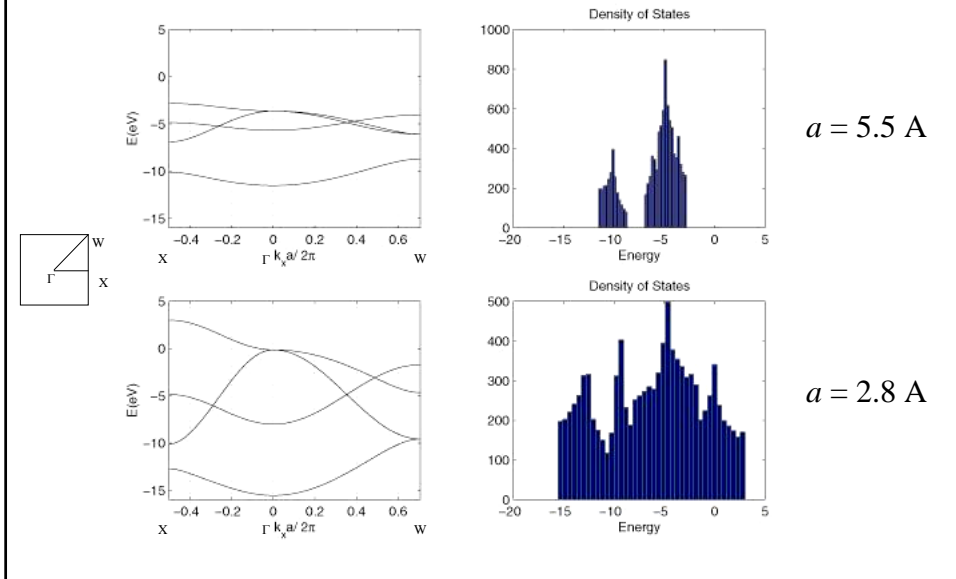
$$\mathbf{H}(\Gamma) = \begin{pmatrix} E_s + 4V_{ss\sigma} & 0 & 0 & 0 \\ 0 & E_p + 2V_{pp\pi} + 2V_{pp\sigma} & 0 & 0 \\ 0 & 0 & E_p + 2V_{pp\pi} + 2V_{pp\sigma} & 0 \\ 0 & 0 & 0 & E_p + 4V_{pp\pi} \end{pmatrix}$$

$$\begin{aligned} E_1(\Gamma) &= E_s + 4V_{ss\sigma} & \vec{\epsilon}_1(\Gamma) &= \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \vec{\epsilon}_3(\Gamma) &= \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \\ E_2(\Gamma) &= E_p + 4V_{pp\pi} \\ E_3(\Gamma) = E_4(\Gamma) &= E_p + 2V_{pp\pi} + 2V_{pp\sigma} & \vec{\epsilon}_4(\Gamma) &= \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix} & \vec{\epsilon}_2(\Gamma) &= \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} \end{aligned}$$

2D Monatomic Square Crystals Variations with Lattice Constant



2D Monatomic Square Crystals Dispersion Relations



2D Monatomic Square Crystals Fermi Energy

How many states per band ?

$$2 \cdot \frac{\text{Area of BZ}}{\text{Area per state}} = \frac{2 \cdot \left(\frac{2\pi}{a}\right)^2}{\left(\frac{2\pi}{L}\right)^2} = 2n$$

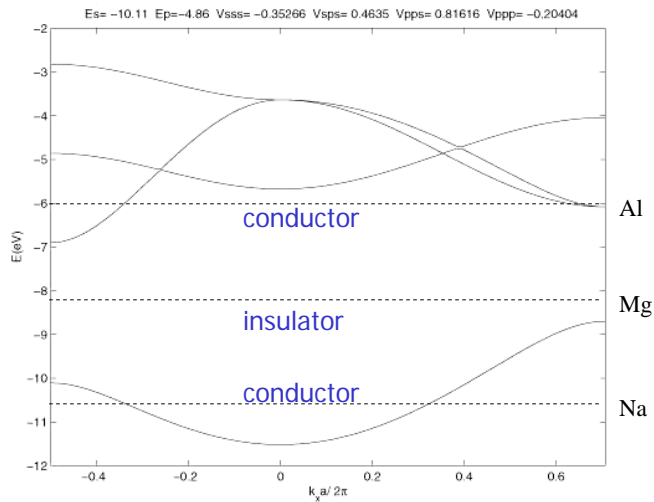
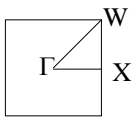
where n is the areal density of atoms

To estimate Fermi energy we need to know the number of outermost valence electrons each atom has...

I	II	III	IV	V	VI	VII	VIII
11	12	13	14	15	16	17	18
Na	Mg	Al	Si	P	S	Cl	Ar

2D Monatomic Square Crystals Dispersion Relations

$a = 5.5 \text{ \AA}$



I	II	III	IV	V	VI	VII	VIII
11	12	13	14	15	16	17	18
Na	Mg	Al	Si	P	S	Cl	Ar

Reducing a , makes Mg a conductor !

Name: _____

$$\mathbf{H}(\mathbf{k}) = \begin{pmatrix} \langle \phi_s | & \langle \phi_{p_x} | & \langle \phi_{p_y} | & \langle \phi_{p_z} | \\ \left(\begin{array}{cccc} E_s + V_{ss\sigma} g_0 & ? & ? & ? \\ V_{sp\sigma} g_1^* & ? & ? & ? \\ ? & ? & ? & ? \\ ? & ? & ? & ? \end{array} \right) & & & \end{pmatrix}$$

Matrix element (s- p_x) _____

Matrix element (s- p_y) _____

Matrix element (p_x - p_x) _____

Matrix element (p_x - p_y) _____