

Lecture 1: General Introduction to Quantum Computing and Superconducting Devices

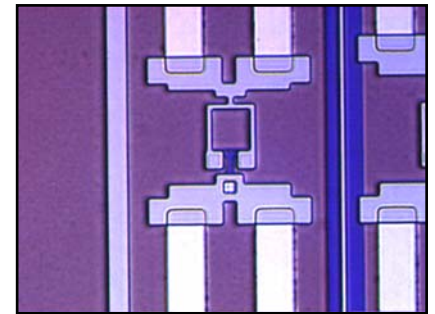
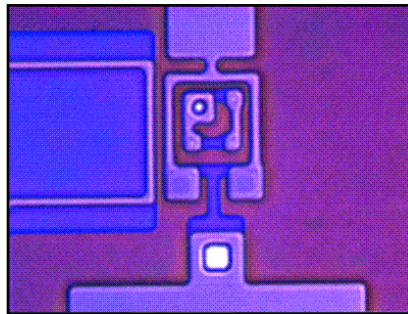
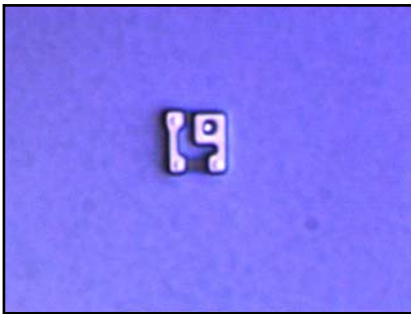
Terry P. Orlando

6.975

EECS MIT

February 4, 2003

Orlando Group Website: <http://rleweb.mit.edu/superconductivity>



Some viewgraphs from Janice Lee, Ken Segall, Donald Crankshaw, Daniel Nakada, and Yang Yu.

Outline

1. Introduction to Quantum Computation

- a. The Unparalleled Power of a Quantum Computer
- b. Two State Systems: Qubits
- c. Types of Qubits

2. Background on Superconductors

- a. What is a Superconductor
- b. Uses for Superconductors

3. Quantum Circuits

4. Building a Quantum Computer with Superconductors

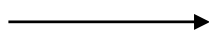
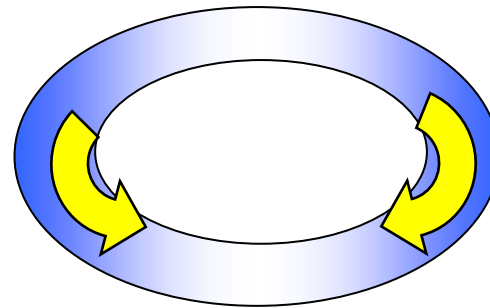
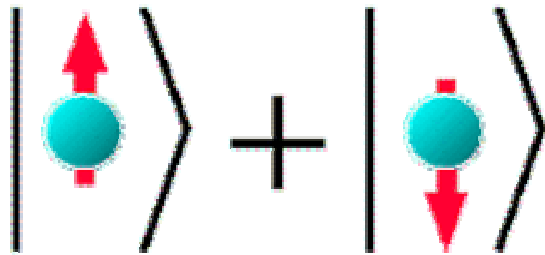
- a. Types of Superconducting Qubits
- b. Experiments on Superconducting Qubits
 - 1. Charge qubits
 - 2. Phase/Flux qubits
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- c. Advantages of superconductors as qubits

5. Outline of Class



The “Magic” of Quantum Mechanics

States 0 and 1 are stored and processed AT THE SAME TIME



Parallel
Computation



Exponential
Speedup to
get Answers

Qubits: Quantum Bits

- Qubits are two level systems
 - a) Spin states can be true two level systems, or
 - b) Any two quantum energy levels can also be used
- We will call the lower energy state $|0\rangle$ and the higher energy state $|1\rangle$
- In general, the wave function can be in a superposition of these two states

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$



Computing with Quantum States

- Consider two qubits, each in superposition states

$$|\psi\rangle_a = |0\rangle_a + |1\rangle_a \quad |\psi\rangle_b = |0\rangle_b + |1\rangle_b$$

- We can rewrite these states as a single state of the 2 spin system

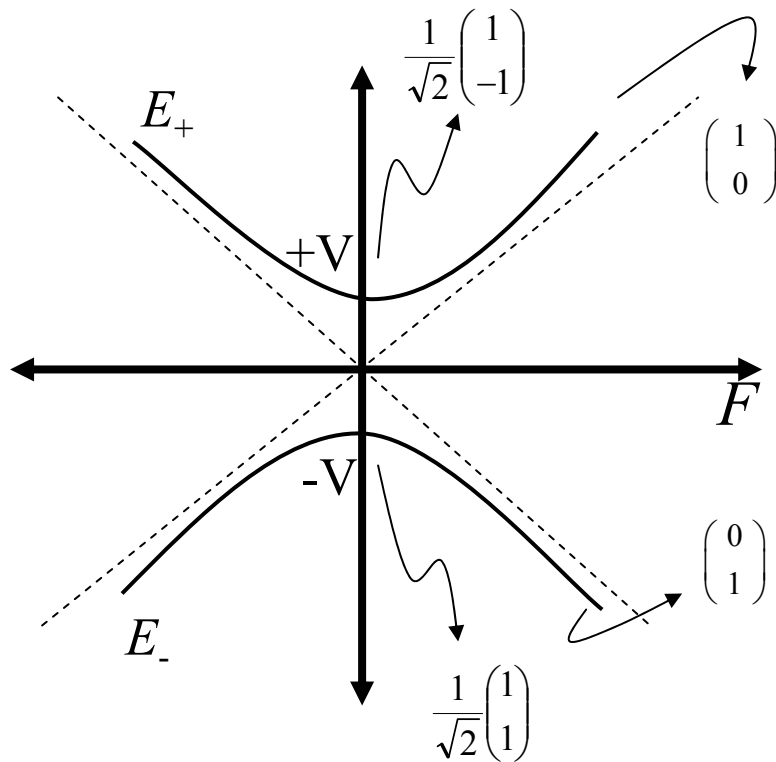
$$\begin{aligned} |\Psi\rangle &= |\psi\rangle_a \otimes |\psi\rangle_b = (|0\rangle_a + |1\rangle_a) \otimes (|0\rangle_b + |1\rangle_b) \\ &= |0\rangle_a |0\rangle_b + |0\rangle_a |1\rangle_b + |1\rangle_a |0\rangle_b + |1\rangle_a |1\rangle_b \\ &= |00\rangle + |01\rangle + |10\rangle + |11\rangle \end{aligned}$$

- All four “numbers” (0, 1, 2, & 3 in binary notation) exist simultaneously
- Algorithm designed so that states interfere to give one “number” with high probability



Two Level Systems

$$H = \begin{pmatrix} -F & -V \\ -V^* & F \end{pmatrix}, \quad |\psi\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$



Eigenenergies $E = \sqrt{F^2 + V^2}$

At $F=0$, let $|\psi(t=0)\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

$$\begin{aligned} |\psi(t)\rangle &= \frac{1}{2} \begin{pmatrix} 1 \\ 1 \end{pmatrix} e^{i\frac{V}{\hbar}t} + \frac{1}{2} \begin{pmatrix} 1 \\ -1 \end{pmatrix} e^{-i\frac{V}{\hbar}t} \\ &= \begin{pmatrix} 1 \\ 0 \end{pmatrix} \cos \frac{V}{\hbar}t + i \begin{pmatrix} 0 \\ 1 \end{pmatrix} \sin \frac{V}{\hbar}t \end{aligned}$$

System oscillates between $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$ and $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$ with period $T = \frac{h}{2V}$



Rabi Oscillations

Drive the system with $V(t)=V_0 e^{i\omega t}$ at the resonant frequency $\omega = E_+ - E_-$

$$\text{If } |\psi(t=0)\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \text{ then } |\psi(t)\rangle = \cos\frac{V_0 t}{\hbar} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + i \sin\frac{V_0 t}{\hbar} e^{i\omega t} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Oscillations between states can be controlled by V_0 and the time of AC drive, with period

$$T = \frac{h}{2V_0}$$



The Promise of a Quantum Computer

A Quantum Computer ...

- Offers exponential improvement in *speed* and *memory* over existing computers
- Capable of *reversible computation*
- e.g. Can factorize a 250-digit number in seconds while an ordinary computer will take 800 000 years!



Evolution of Computing Technology

Vacuum Tubes

- Slow
- Power consuming
- Huge in size



Transistors

- Improved size and efficiency
- Heat dissipation



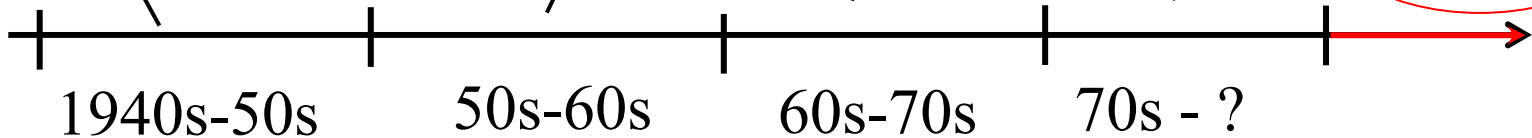
Integrated Circuits

- Fitting everything on a chip

VLSI, ULSI

- Yet smaller sizes

Quantum Computer





Quantum Information Science and Technology Roadmapping Project

1. Quantum Computing Roadmap Overview
2. Nuclear Magnetic Resonance Approaches
3. Ion Trap Approaches
4. Neutral Atom Approaches
5. Optical Approaches
6. Solid State Approaches
7. Superconducting Approaches
8. “Unique” Qubit Approaches
9. The Theory Component of the Quantum Information Processing and Quantum Computing Roadmap

<http://qist.lanl.gov>

Massachusetts Institute of Technology



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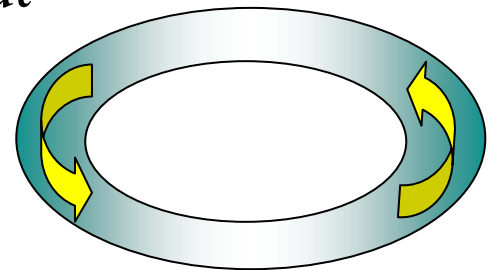
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What is a Superconductor?

“A *Superconductor* has *ZERO* electrical resistance *BELOW* a certain critical temperature. Once set in motion, a persistent electric current will flow in the superconducting loop *FOREVER* without any power loss.”



Magnetic Levitation

A *Superconductor EXCLUDES* any magnetic fields that come near it.

How “Cool” are Superconductors?

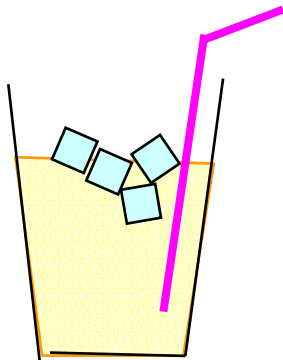
Below **77 Kelvin** (-200 °C):

- Some Copper Oxide Ceramics superconduct

Below **4 Kelvin** (-270 °C):

- Some Pure Metals e.g. Lead, Mercury, **Niobium** superconduct

Keeping at 0 °C



Keeping at 77 K

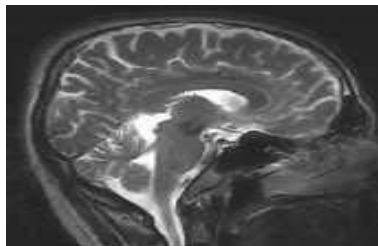


Keeping at 4K



Uses for Superconductors

- Magnetic Levitation allows trains to “float” on strong superconducting magnets (MAGLEV in Japan, 1997)



- To generate Huge Magnetic field e.g. for Magnetic Resonance Imaging (MRI)

- ➔ • A SQUID (Superconducting Quantum Interference Device) is the most sensitive magnetometer. (sensitive to 100 billion times weaker than the Earth’s magnetic field)
- ➔ • Quantum Computing

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Circuits for Qubits

- Need to find dissipationless circuits which have two “good” energy levels
- Need to be able to “manipulate” qubits and couple them together



Quantization of Circuits

- Find the energy of the circuit
- Change the energy into the Hamiltonian of the circuit by identifying the canonical variables
- Quantize the Hamiltonian
- Usually we can make it look like a familiar quantum system



Harmonic Oscillator



$$H = \frac{1}{2} m v^2 + \frac{1}{2} m \omega^2 x^2$$

$$v = \frac{dx}{dt}$$

$$H = \frac{1}{2} m \left(\frac{dx}{dt} \right)^2 + \frac{1}{2} m \omega^2 x^2$$

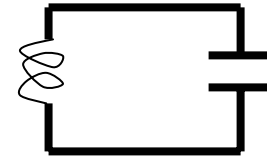
$$p = m \frac{dx}{dt}$$

Quantum Mechanically

$$\Delta x \Delta p \geq \hbar / 2$$

$$E = \hbar \omega \left(n + \frac{1}{2} \right)$$

LC Circuit



$$H = \frac{1}{2} C V^2 + \frac{1}{2} L I^2$$

$$v = \frac{d\Phi}{dt} \text{ and } I = \frac{\Phi}{L}$$

$$H = \frac{1}{2} \underbrace{C}_{\overline{M}} \left(\frac{d\Phi}{dt} \right)^2 + \frac{1}{2} \underbrace{C}_{\overline{M}} \frac{1}{\underbrace{LC}_{\omega^2}} \Phi^2$$

$$p = C \frac{d\Phi}{dt} = C V$$

Quantum Mechanically

$$L C \Delta I \Delta V \geq \hbar / 2$$

$$E = \hbar \omega \left(n + \frac{1}{2} \right)$$

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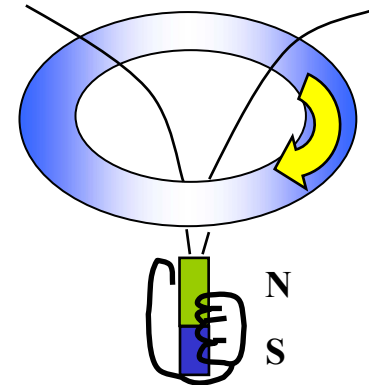
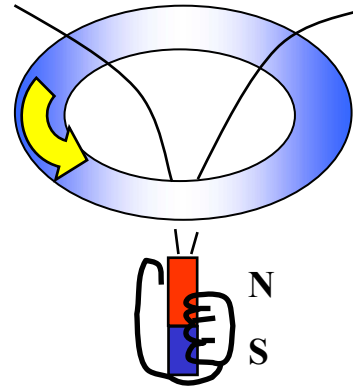
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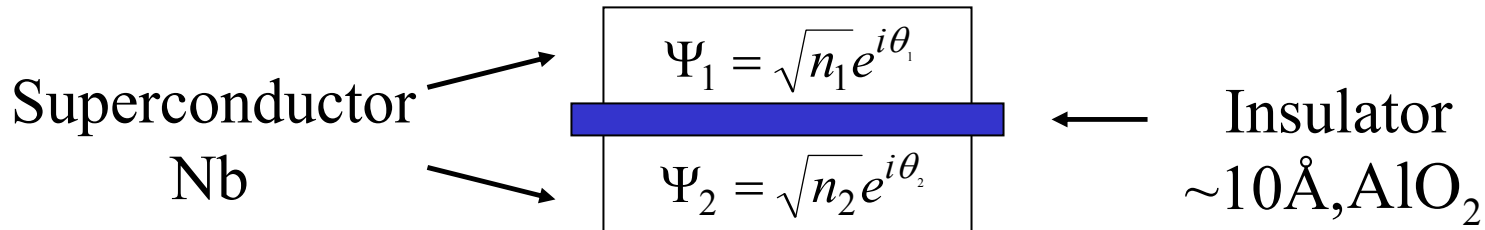


The Superconducting “Quantum Bit”

- An External Magnet can induce a current in a superconducting loop
- The induced current can be in the opposite direction if we carefully choose a *different* magnetic field this time
- To store and process information as a computer bit, we assign:



Josephson Junction



- Josephson relations:

$$I = I_c \sin \varphi$$

$$V = \frac{\Phi_0}{2\pi} \frac{d\varphi}{dt}$$

$$\varphi = \theta_2 - \theta_1$$

$$-\frac{2\pi}{\Phi_0} \int \vec{A}(r, t) \cdot d\vec{l}$$

- Behaves as a nonlinear inductor:

$$V = L_J \frac{dI}{dt},$$

where $L_J = \frac{\Phi_0}{2\pi I_c \cos \varphi}$

$\Phi_0 =$ flux quantum

483.6 GHz / mV

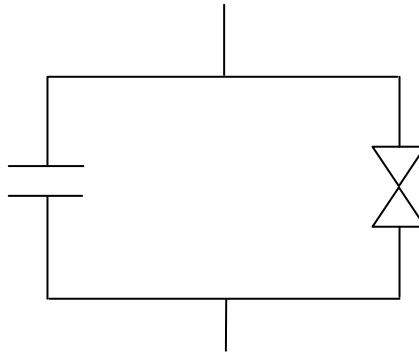


Quantization of a Josephson Junction

Charging Energy

$$U_C = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} CV^2$$
$$= \frac{1}{2} \left(\frac{\Phi_0}{2\pi} \right)^2 C \left(\frac{\partial \varphi}{\partial t} \right)^2$$

$$E_C = \frac{e^2}{2C}$$



Josephson Energy

$$U_J = \frac{\Phi_0 I_c}{2\pi} (1 - \cos \varphi)$$

$$E_J = \frac{\Phi_0 I_c}{2\pi}$$

Hamiltonian:
$$H = \frac{1}{2} \left(\frac{\Phi_0}{2\pi} \right)^2 C \left(\frac{\partial \varphi}{\partial t} \right)^2 + \frac{\Phi_0 I_c}{2\pi} (1 - \cos \varphi)$$

Circuit behaves just like a physical pendulum.

For Al-Al₂O₃-Al junction with an area of 100x100 nm²

C = 1fF and I_c=300 nA, which gives E_C=10μeV and E_J=600μeV

To see quantization, Temperature < 300 mK

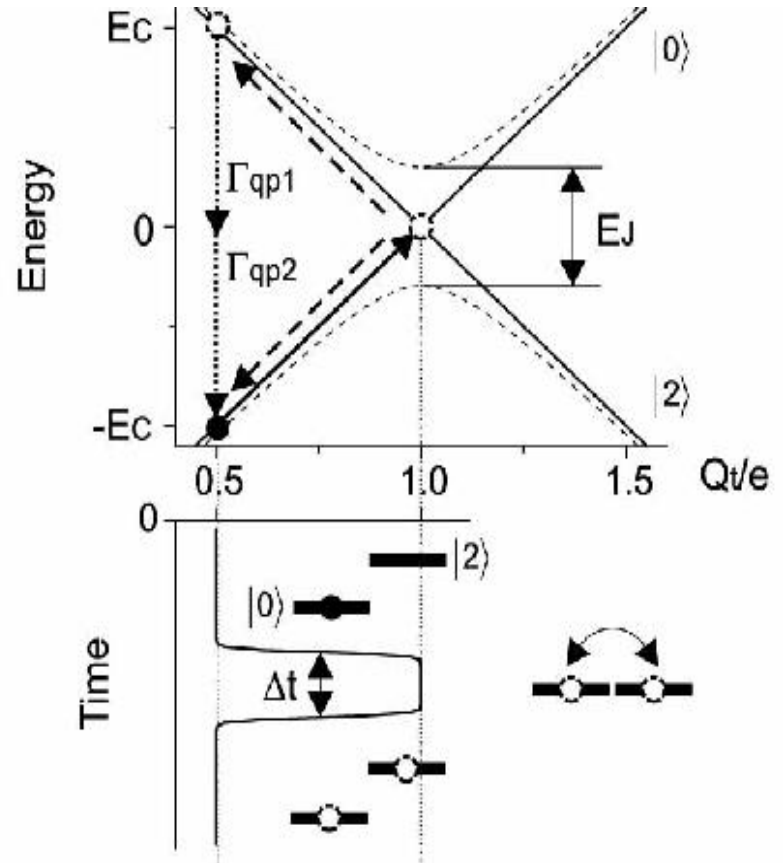
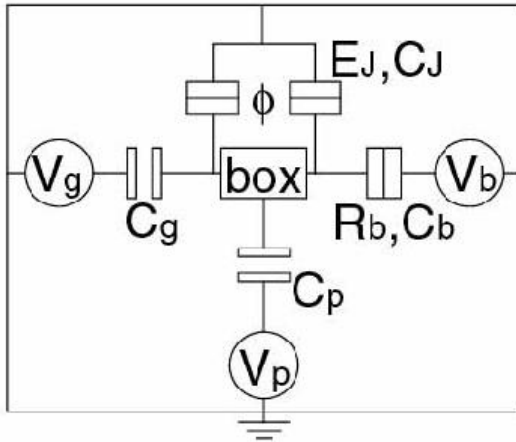


Types of Superconducting Qubits

- Charge-state Qubits (voltage-controlled)
 - Cooper pair boxes
- Flux/Phase-state qubits (flux-current control)
 - Persistent Current Qubits
 - RF SQUID Qubits
 - Phase Qubits (single junction)
- Hybrid Charge-Phase Qubits

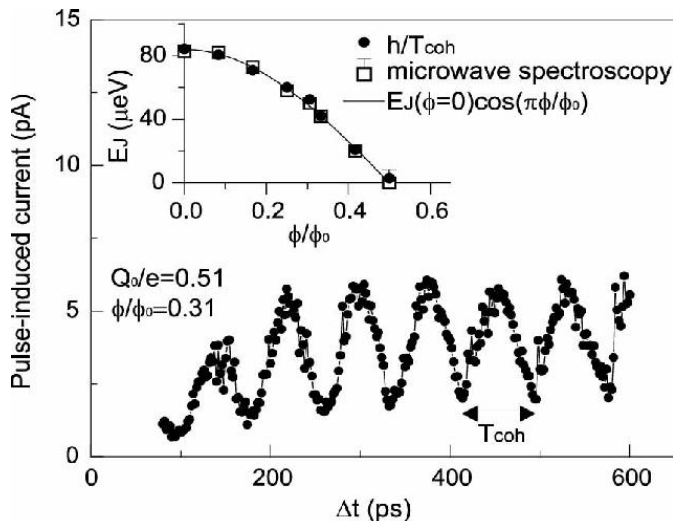


Charge-State Superconducting Qubit



TOP: "Electrical schematic"

BOTTOM: "Evidence of Rabi Oscillations"



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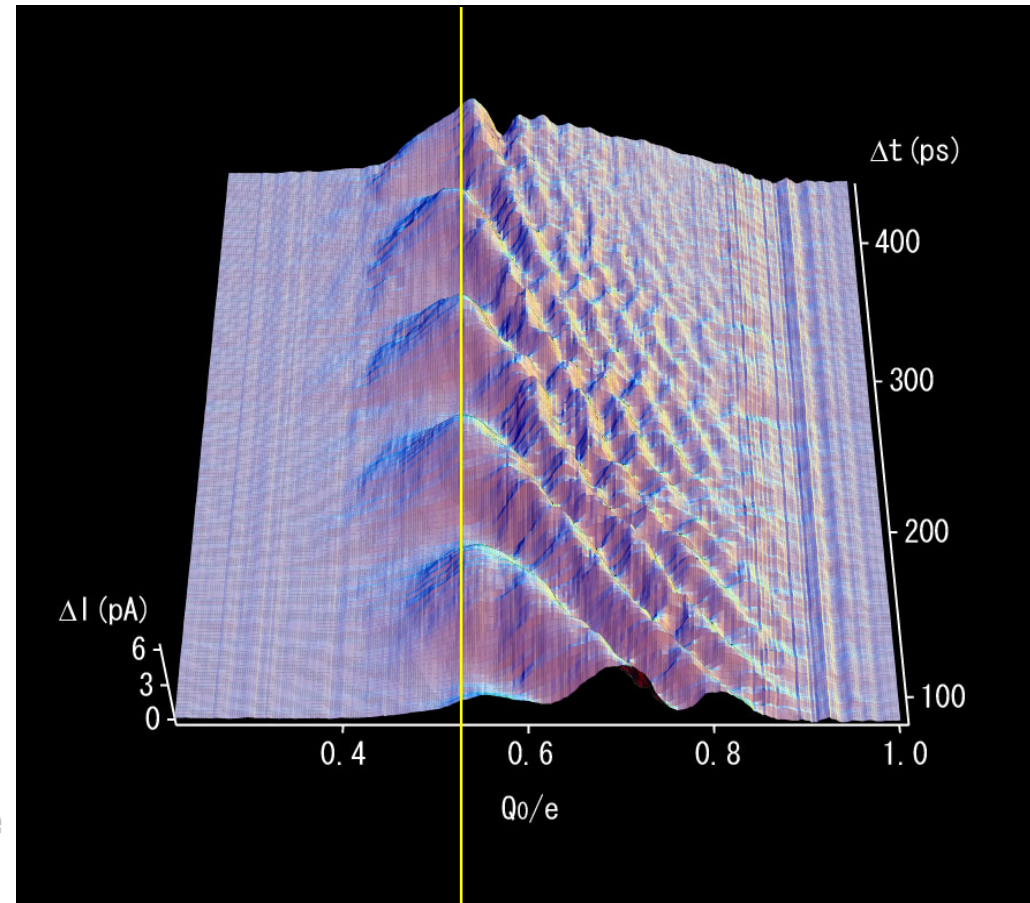
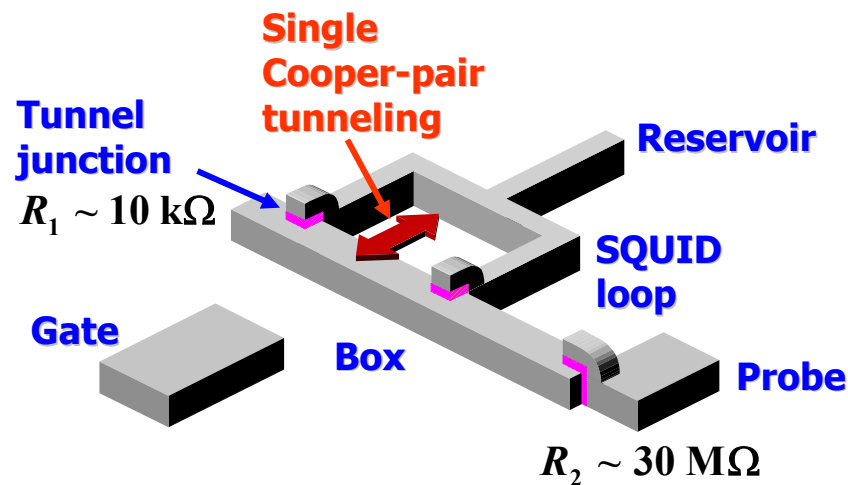
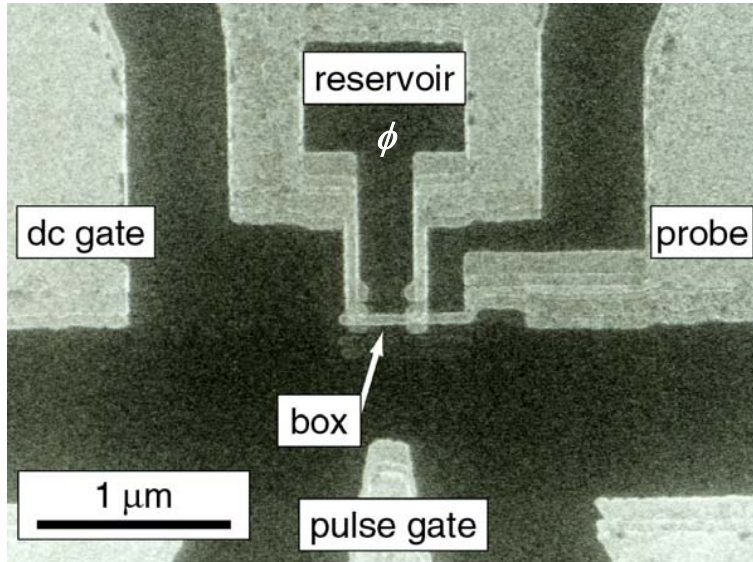
Y. Nakamura, Yu A. Pashkin, and J.S. Tsai, *Nature* **398**, 786 (1999).



Charge qubit

a Cooper-pair box $E_J / E_C \sim 0.3$

Coherence up to ~ 5 ns, presently limited by background charge noise (dephasing) and by readout process (relaxation)



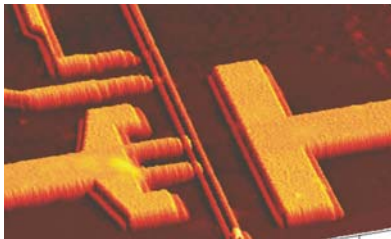
Y.Nakamura *et al.*, Nature **398**, 786 (1999).



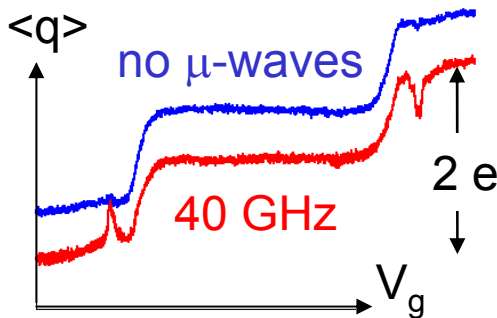
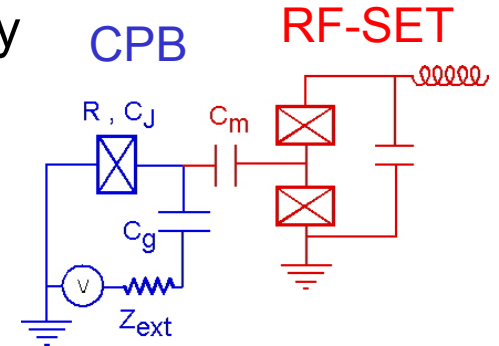
Quantum Coherence in the Cooper-Pair Box Measured with an RF-SET



Schoelkopf Group, Yale University



- Superconducting charge qubit: the Cooper-pair box (CPB)
- Fast, single-charge readout: the RF single-electron transistor



• **Quantum coherence of Cooper-pair box qubit observed by CW microwave spectroscopy**



- Transition frequency ~ 40 GHz: $Q_\phi = \omega_{10} T_\phi \sim 250$ & $Q_1 = \omega_{10} T_1 > 10^5$
- Ensemble decoherence time from linewidth: $T_\phi = 1$ ns
- Lifetime from time-resolved decay of photon peaks: $T_1 = 1.6$ μ s

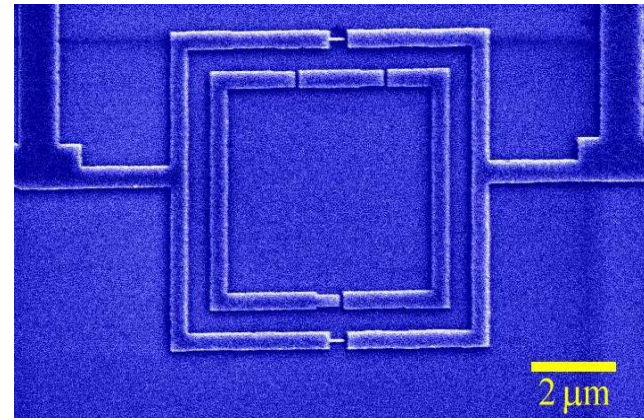
Next steps:

- Operate @ charge-noise insensitive point to reduce decoherence
- Perform time-gated, single-shot measurements

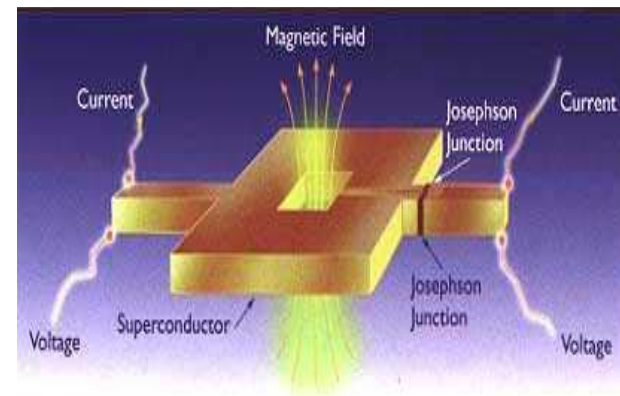
Funding: ARO/ARDA/NSA

Our Persistent Current Qubit

- Depending on the direction of the current, state $|0\rangle$  and state $|1\rangle$  will *add* a different magnetic field to the external magnet



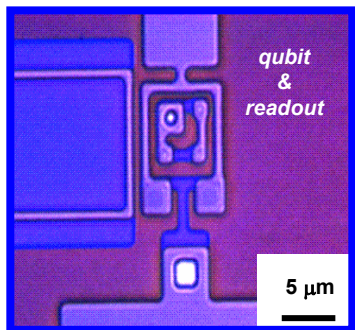
- This difference is very small but can be distinguished by the extremely sensitive SQUID sensor



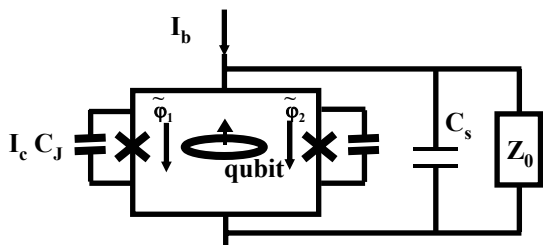
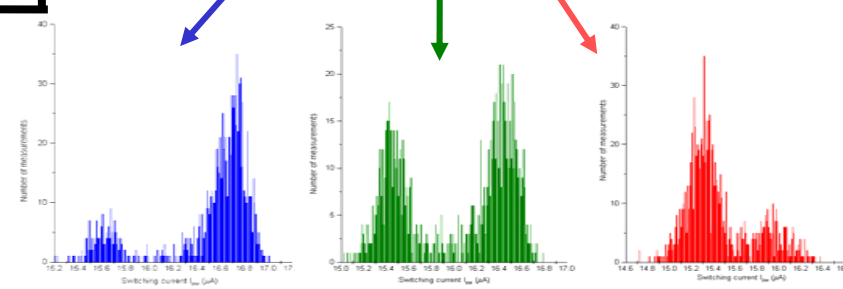
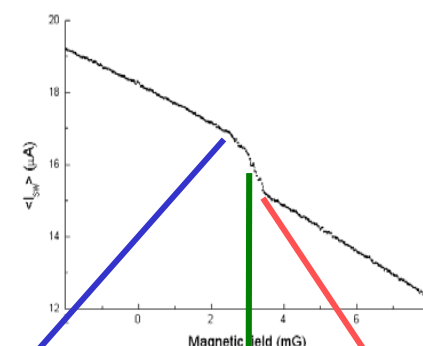
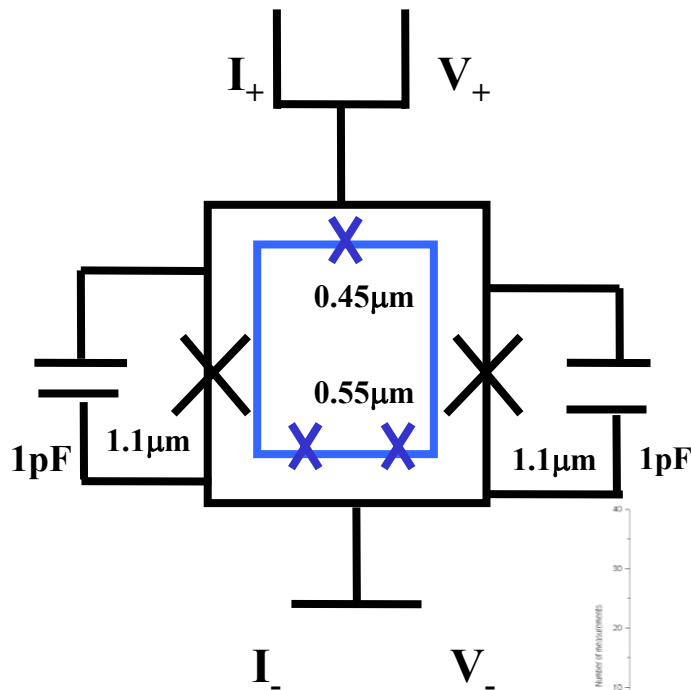


Quantum Computation with Superconducting Quantum Devices

T.P. Orlando, S. Lloyd, L. Levitov, J.E. Mooij - MIT
M. Tinkham – Harvard; M. Bocko, M. Feldman – U. of Rochester
in collaboration with K. Berggren, MIT Lincoln Laboratory



Fabrication modeling, and measurements



- Persistent current qubit fabricated in Nb with submicron junctions
- Two states seen in measurement (thermal activations and energy levels)



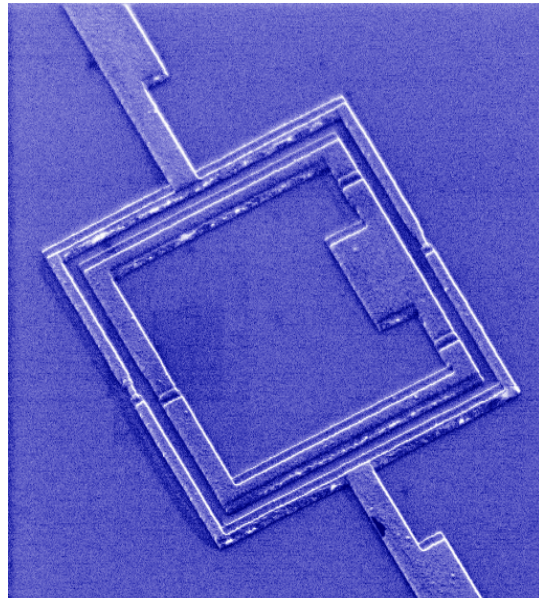
Macroscopic quantum superposition in a Josephson junction loop

Delft University of Technology & DIMES The Netherlands

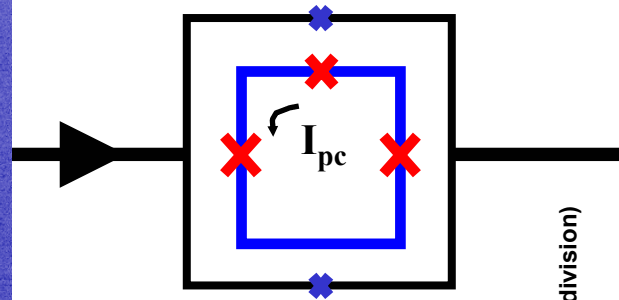
MIT Cambridge

Caspar van der Wal, A. Ter Haar, Kees Harmans, Hans Mooij

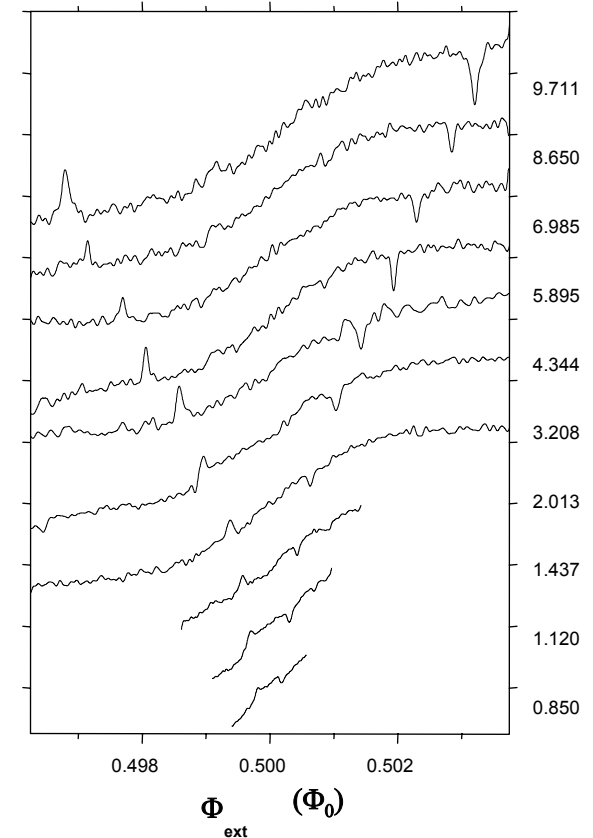
T. Orlando, L. Levitov, S. Lloyd



3 μm



\bar{I}_{sw} (0.4 nA per division)



- *Superposition of states observed*
- *Relaxation time 5 μsec ,*
- *Dephasing time 0.1 μsec*

Observation of Coherent Superposition of Macroscopic States

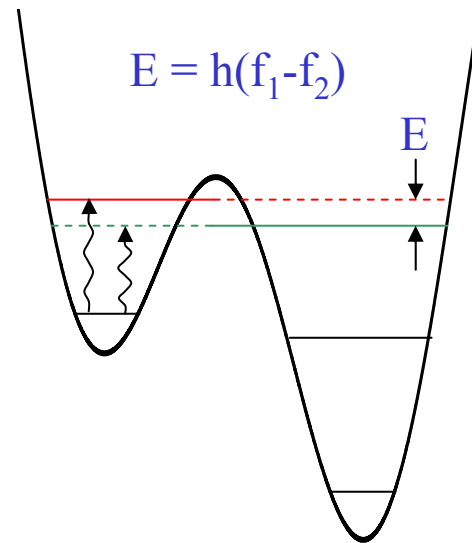
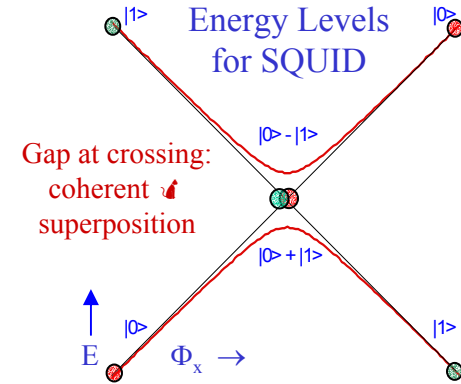
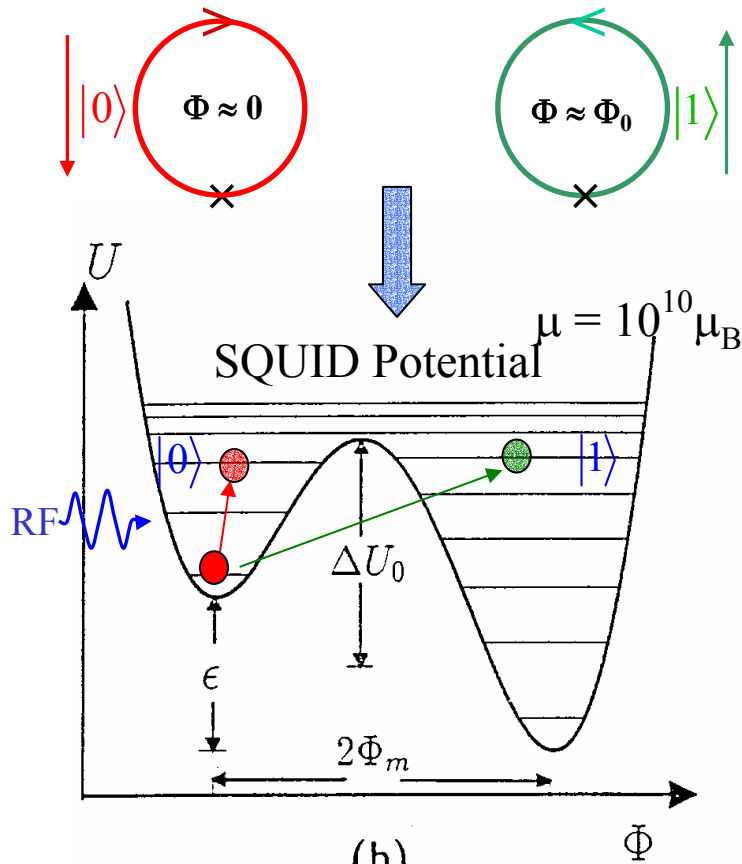
Jonathan Friedman, Vijay Patel, Wei Chen, Sergey Tolpygo and James Lukens

The SQUID

(Superconducting QUantum Interference Device)

- SQUIDs like integer flux quanta, Φ_0

$$\Phi_x \approx \frac{1}{2} \Phi_0$$



(b)

Φ

J. Friedman, *et al.*, Nature, **406**, 43 (July, 2000).

C. Van der Wal, *et al.*, Science, 290, 773 (Oct. 2000)

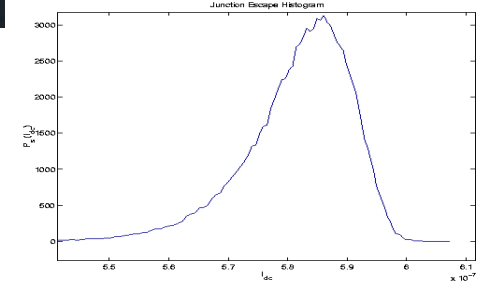
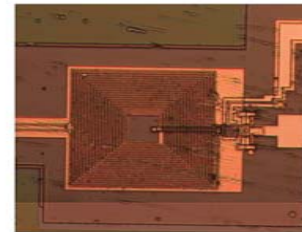
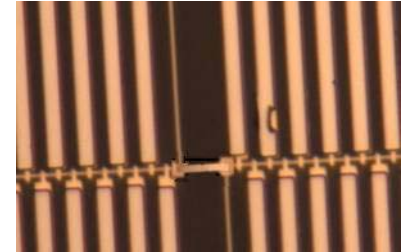
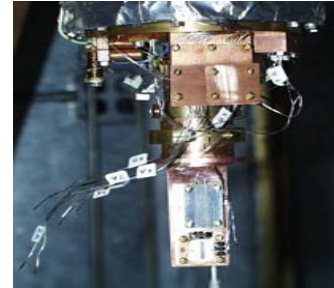
Quantum Computing with Superconducting Devices

F.C. Wellstood, C.J. Lobb, J.R. Anderson, and A.J. Dragt, Univ. of Md.

lobb@squid.umd.edu / <http://www.physics.umd.edu/sqc/>

Objective

- Measure energy levels and decoherence rates in single Josephson junctions and SQUIDs
- Manipulate states of the systems
- Perform 1-qubit operations
- Design and test 2-qubit systems of junctions and SQUIDs



Objective Approach

- Build Josephson junctions that are well isolated from measurement leads to achieve low dissipation and long coherence times at milliKelvin temperatures.
- Measure macroscopic quantum tunneling, energy levels and decoherence rate.
- Use microwaves to pump from $|0\rangle$ to $|1\rangle$
- Model SQUID qubits to guide experimental program.

Status

- Built resistively isolated Al/AlO_x/Al junctions and measured switching distributions with and without microwave excitation
- Assembled SQUID detection scheme for measuring junctions and rf SQUIDs
- Measured switching distributions for SQUID at mK temperatures, $\Delta\Phi = 10^{-3}\Phi_0$

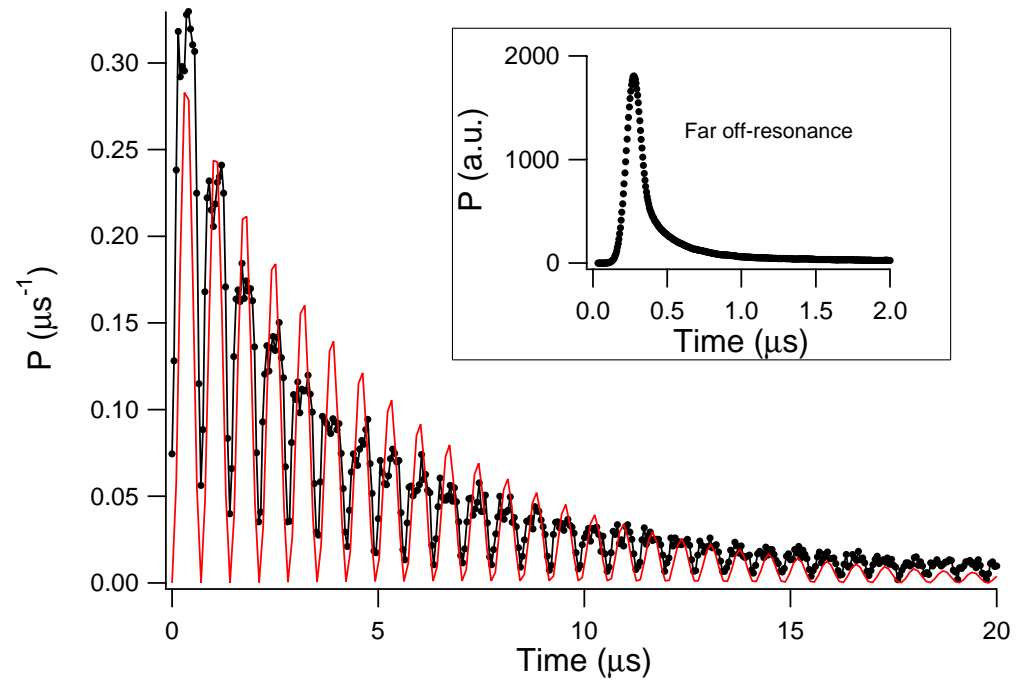
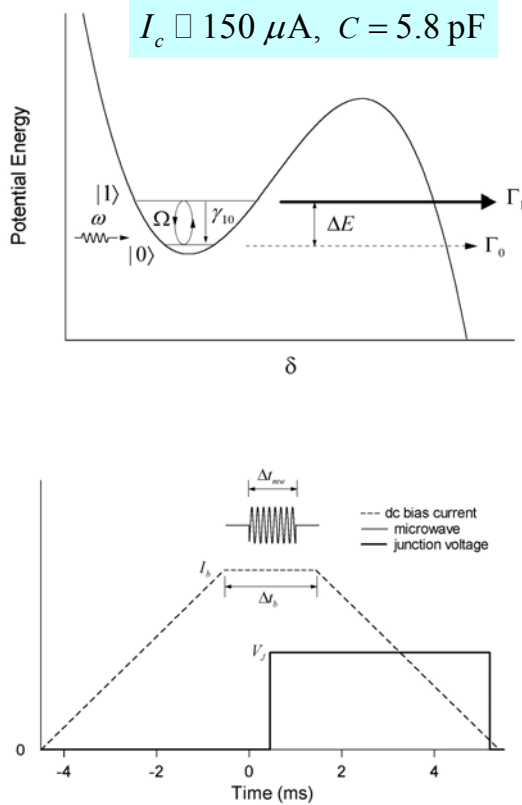


Recent Results from DURINT Quantum Computing Project

Siyuan Han, Yang Yu et al., University of Kansas



Observation of Rabi oscillations in a Josephson Tunnel Junction



$$\Gamma = \frac{1}{2}(\Gamma_0 + \Gamma_1 + \gamma_{10}) + \gamma_\phi$$



decoherence time $> 4.9 \mu\text{s}$

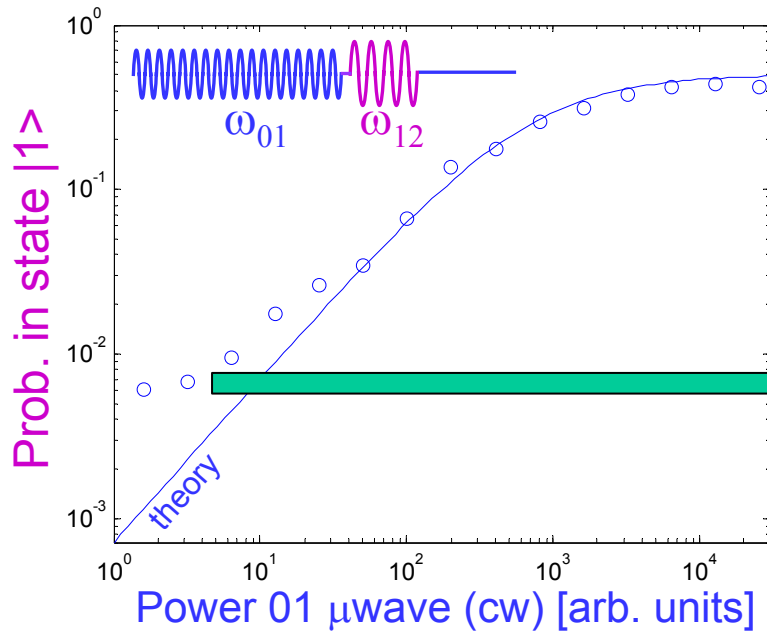
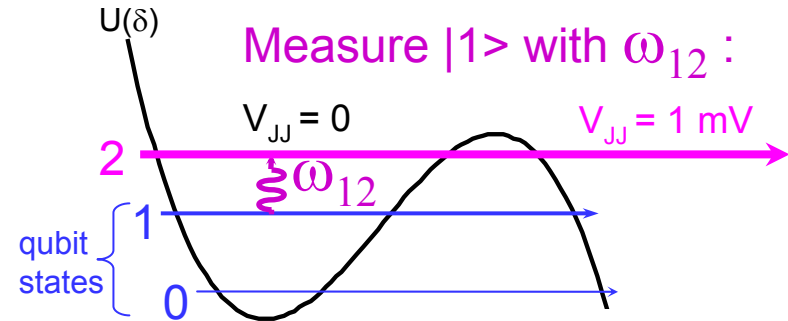
Generation and detection of RO in JJ

Quantum Computing with Current-Biased Josephson Junctions

John Martinis, S.W. Nam, J. Aumentado, C. Urbina; NIST Boulder

Key Advantages:

- Large area junctions - reliably fabricated using optical lithography
- Large capacitance allows coupling to many other qubits



$|1\rangle$ state meas : $>80\%$

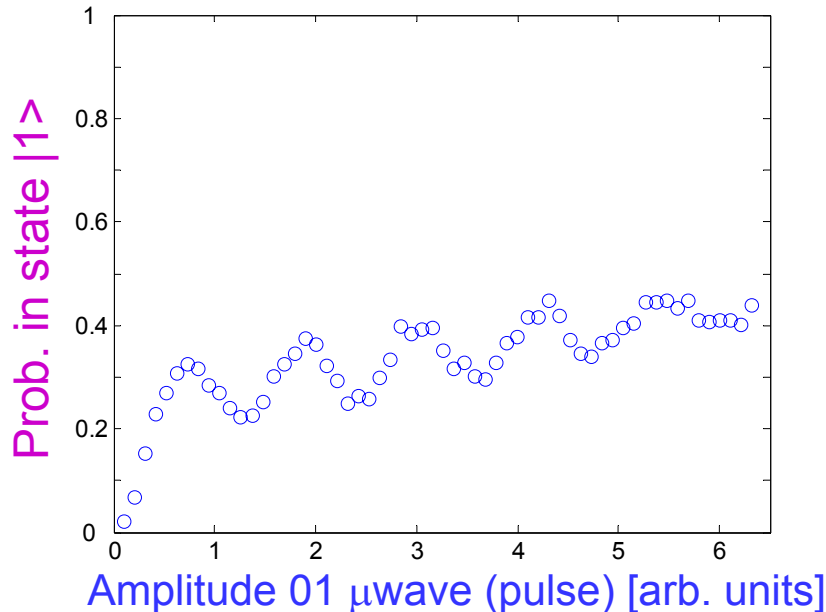
High Fidelity
state preparation
& measurement

$|0\rangle$ state prep : $>99\%$

$|0\rangle$ state meas : $>99\%$

Quantum Computing with Current-Biased Josephson Junctions

John Martinis, S.W. Nam, J. Aumentado, C. Urbina; NIST Boulder



Rabi Oscillations

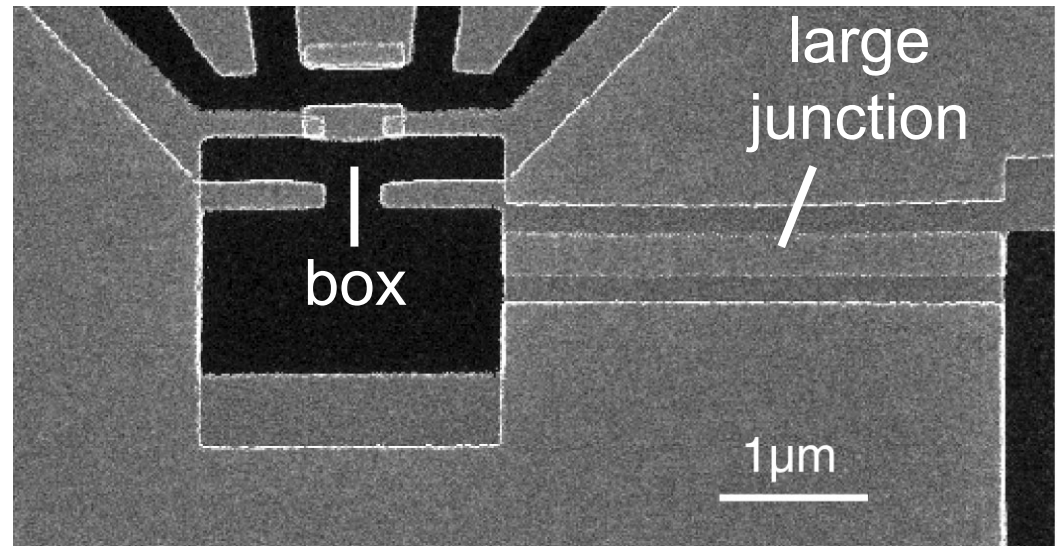
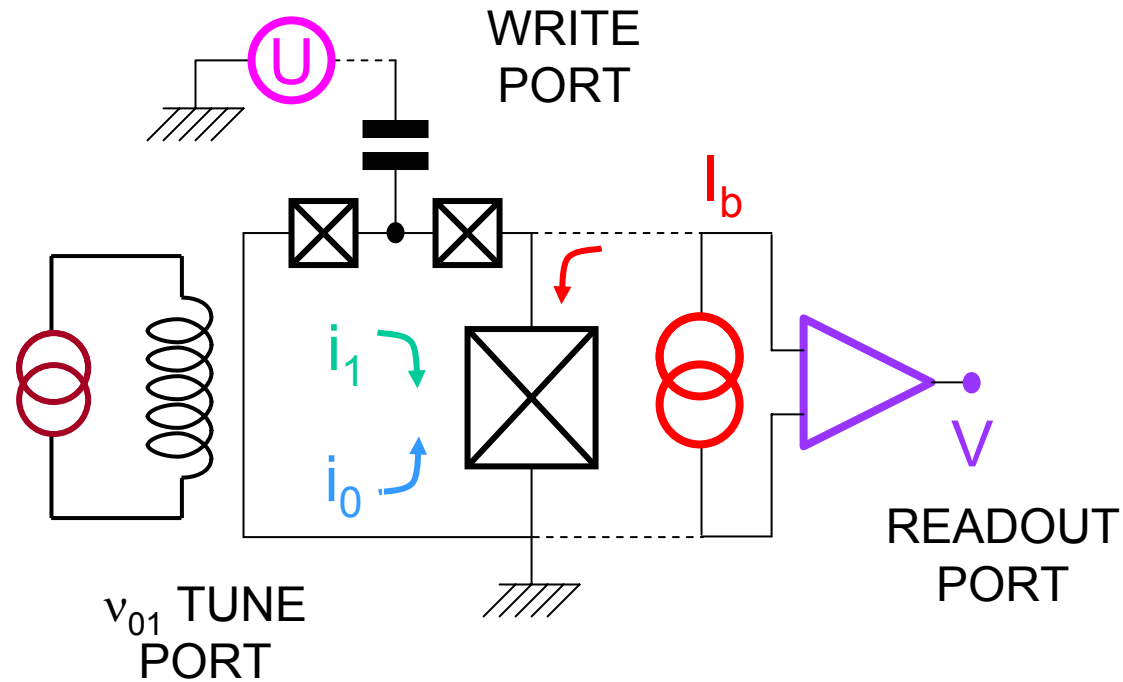
$$\tau_{\text{coh}} \sim 40 \text{ ns}$$

CHARGE-FLUX QUBIT

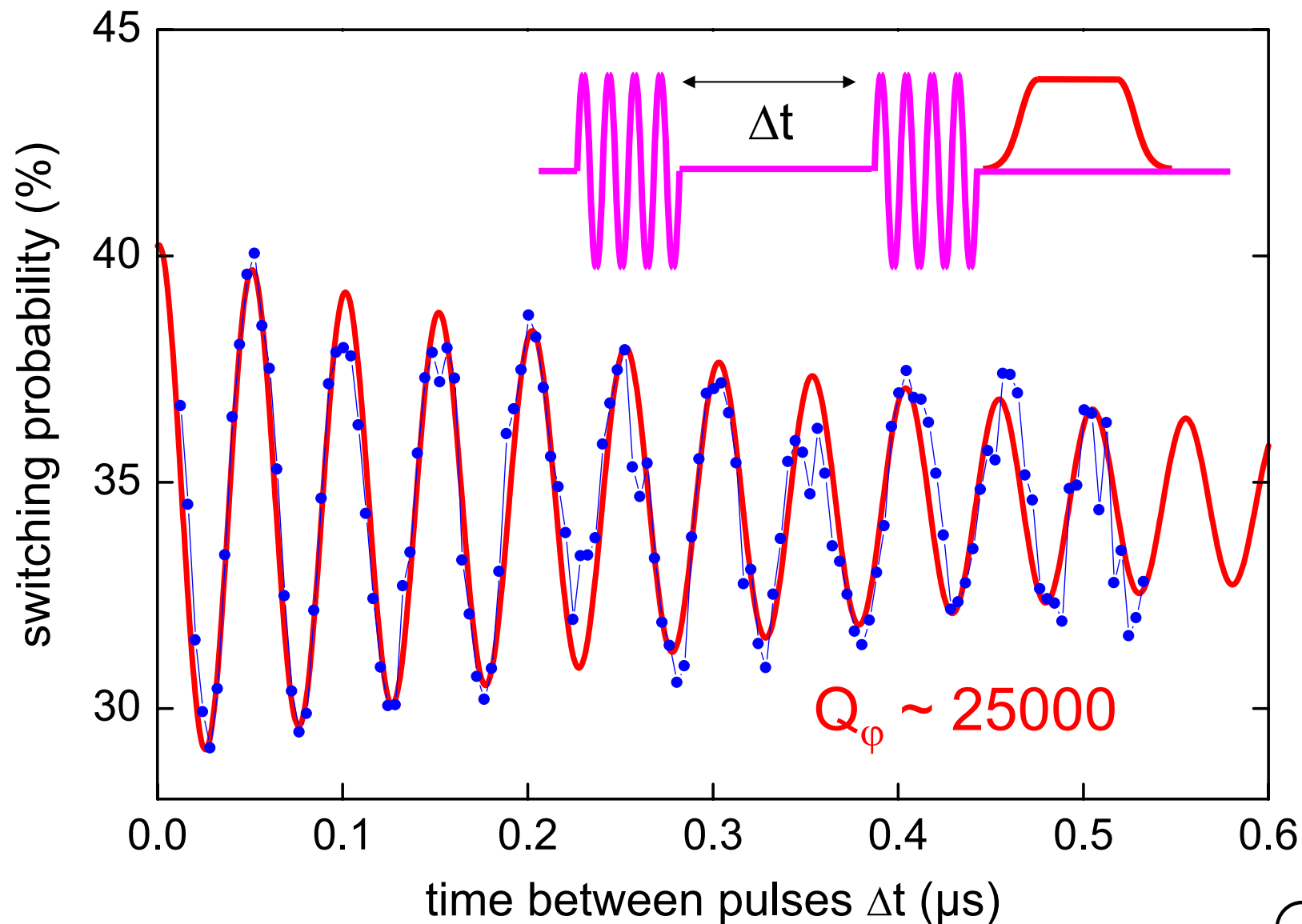
Quantronics Group
CEA-Saclay
France

M. Devoret (now at Yale)
D. Esteve, C. Urbina
D. Vion, H. Pothier
P. Joyez, A. Cottet

Coherence time measured
by Ramsey fringes : 500ns
Qubit transition frequency:
16.5 GHz; coherence quality
factor: 25 000



RAMSEY FRINGES



Outline

1. Introduction to Quantum Computation

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2. Background on Superconductors

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- b. Uses for Superconductors

3. Quantum Circuits

4. Building a Quantum Computer with Superconductors

- a. Types of Superconducting Qubits
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5. Outline of Class



Advantages of Superconductors for Quantum Computing

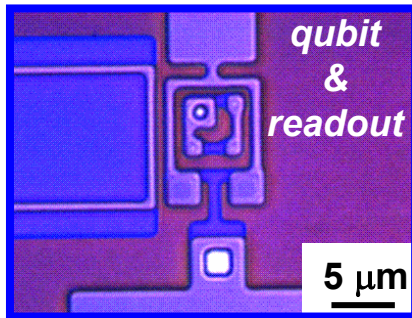
- Employs lithographic technology
- Scalable to large circuits
- Combined with on-chip, ultra-fast control electronics
 - Microwave Oscillators
 - Single Flux Quantum classical electronics





Circuits Fabricated at Lincoln Laboratory

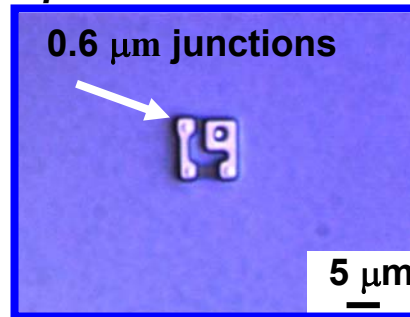
- Persistent-current qubit



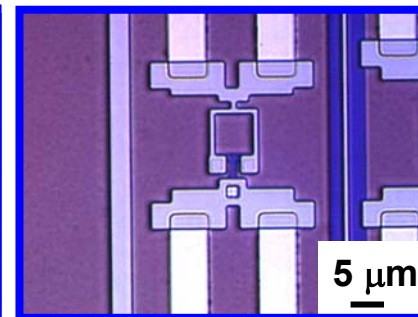
monolithic

flip-chip

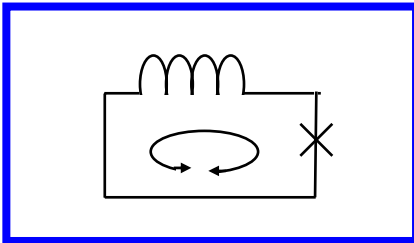
qubit



readout

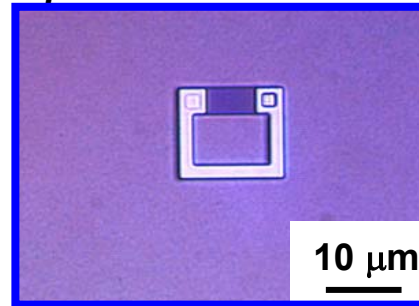


- RF SQUID qubit

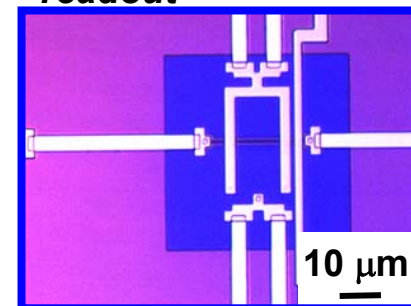


Friedman et al.

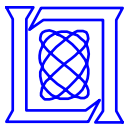
qubit



readout

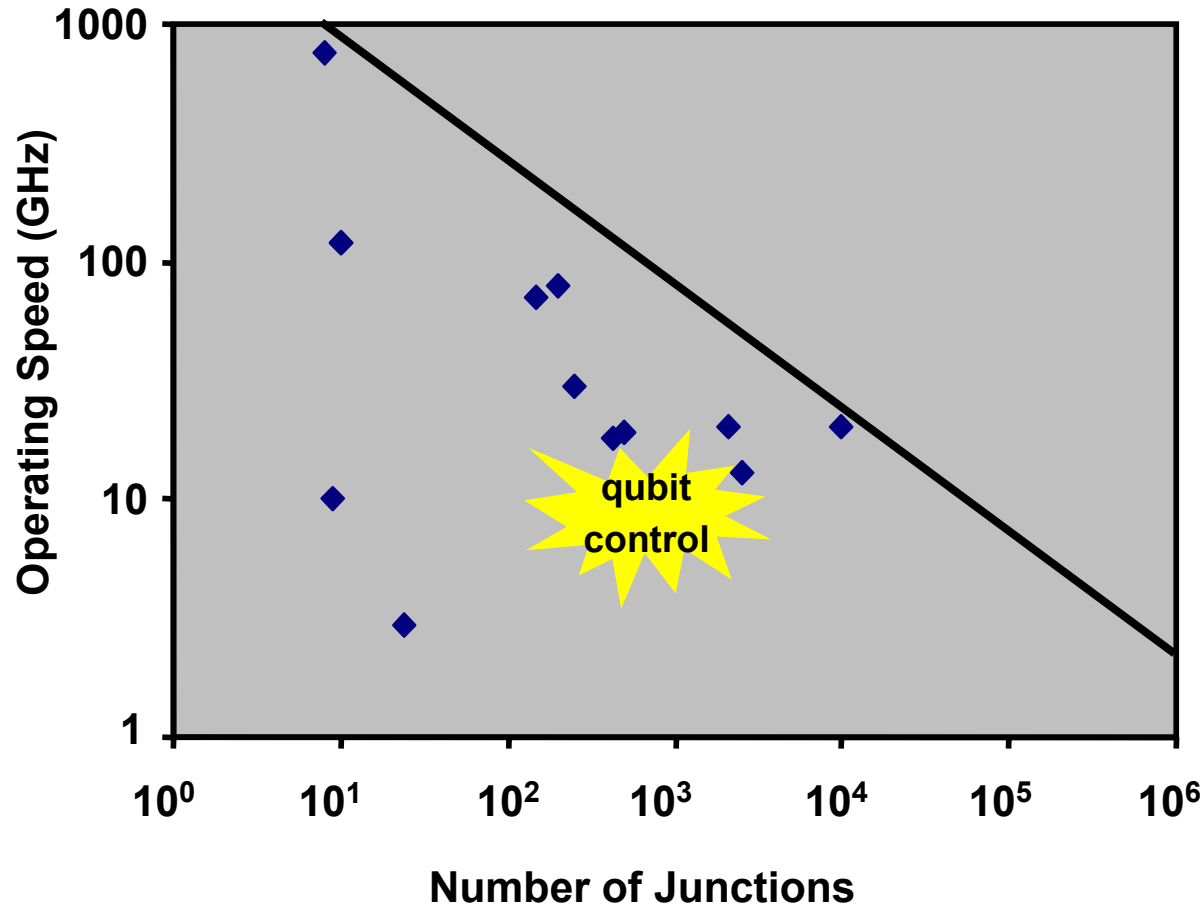


Flip chip process has no hard-wire connections, magnetic field is the only coupling between circuits



Feasibility of Superconductive Control Electronics Fabrication

Industry-Wide Demonstrations of Josephson
Junction Circuits (at 4.2 K)



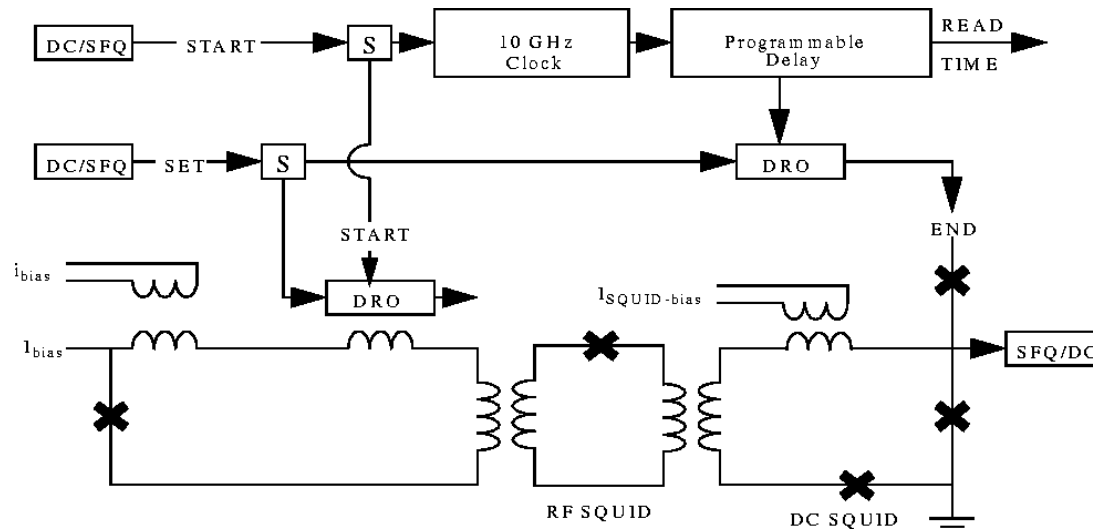
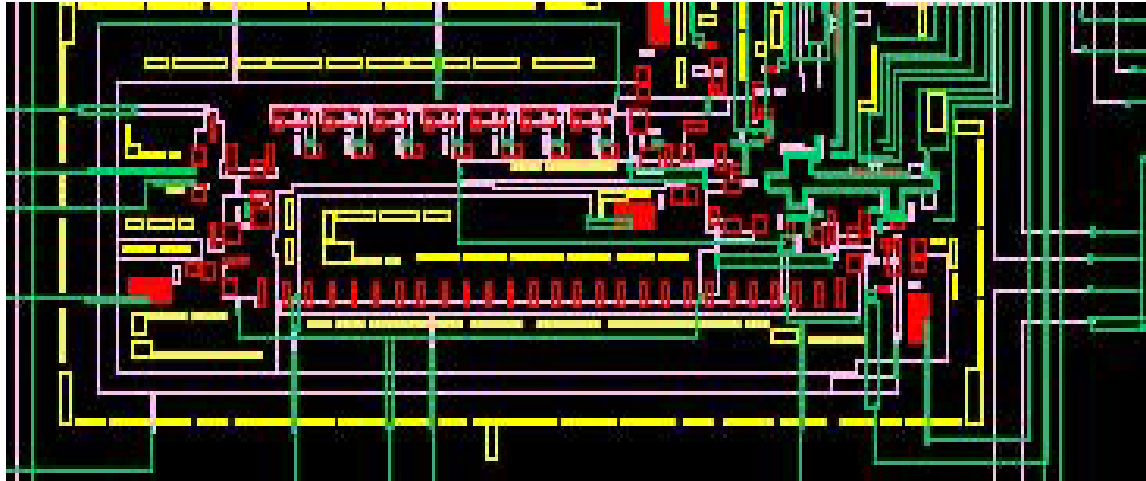
SOURCES
Berkeley
Hitachi
Hypres
Lincoln Laboratory
NEC
Northrop-Grumman
Phys.-Tech.-Bundesanstalt
State University of Moscow
U. of Stonybrook
TRW
others

On-chip Control for an RF-SQUID

M.J. Feldman, M.F. Bocko, Univ. of Rochester

feldman@ece.rochester.edu

www.ece.rochester.edu/~sde/





Quantum Computation with Superconducting Quantum Devices

T.P. Orlando, Yang Yu, D.Nakada, B. Singh, J. Lee, D. Berns,
Ken Segall, D. Crankshaw, B Cord- MIT

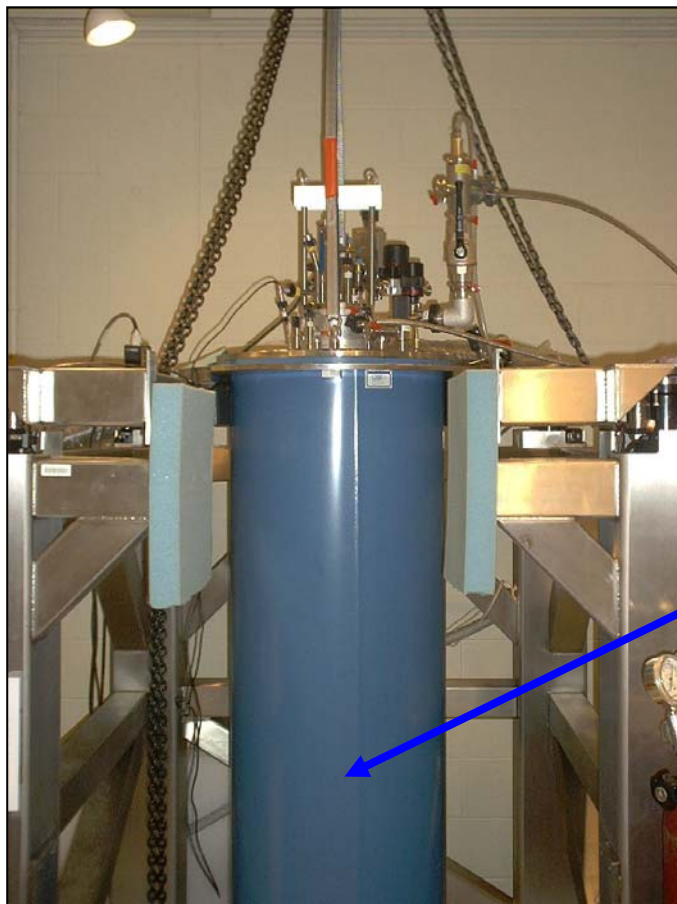


ARDA

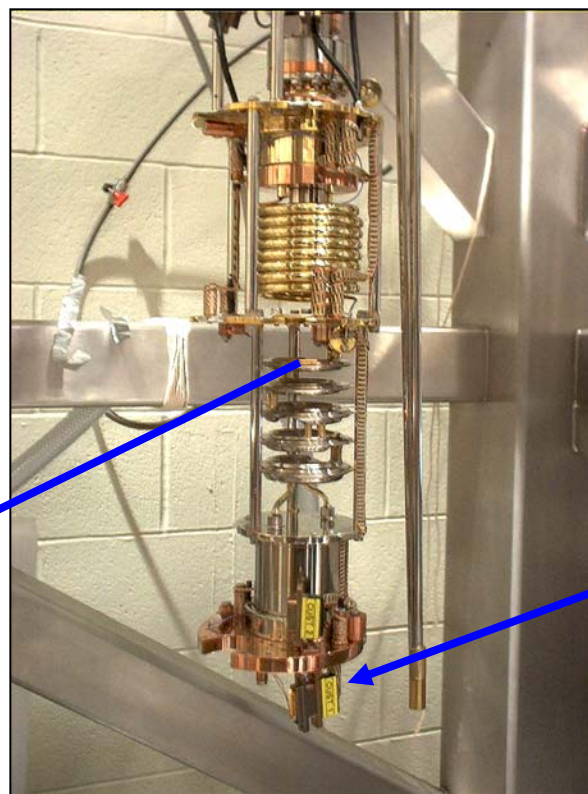


1/27/03

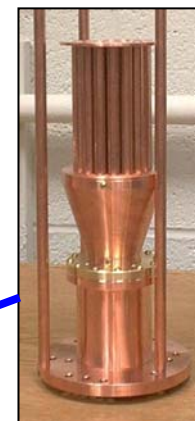
Dilution Refrigerator



Insert



Sample Holder



Massachusetts Institute of Technology
Installed and to begin dc data taking in February and ac data taking in April



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5. **Outline of Class**



Outline of Class 6.975

1. **General Introduction and Overview of Quantized Circuits and Quantum Computing**
2. **Simple Quantized Circuits**
 - a. Resistanceless circuits
 - b. Superconducting Josephson circuit elements
 - c. Energy storage in resistanceless circuits
 - d. Quantization of LC Oscillator Circuit
 - e. Quantization of Josephson junctions
3. **Superconducting Quantum Circuits and Qubits**
 - a. Charge and phase descriptions
 - b. Solutions to Schrödinger's Equation for circuits
 - c. Single Josephson junction
 - d. Cooper pair box
 - e. Quantum RF SQUID
 - f. Persistent Current Qubits
 - g. Hybrid circuits



Outline of Class 6.975

4. Quantum Computing

- a. Two state system model
- b. Qubits and coupled qubits
- c. Manipulations of qubits
- d. Fast control circuitry: SFQ electronics

5. Dissipation in Quantum Circuits

- a. How to model a resistor
- b. Decoherence: relaxation and dephasing
- c. Spin-Boson model



Outline of Class 6.975

6. Quantization of General Circuits

- a. General network theory for classical circuits
 - i. Branch, nodal, and mesh formulations
 - ii. Inclusion of sources
 - iii. Lagrangian and Hamiltonian formulations
- b. Network theory for quantum circuits
 - i. General canonical variables: charge, phase, and mesh currents
 - ii. Quantization without sources
 - iii. Transformations among canonical variables
- c. Quantization with voltage and current sources

7. Assessment and Future Research in Quantum Circuits

