

Relativistic Momentum

R5-1

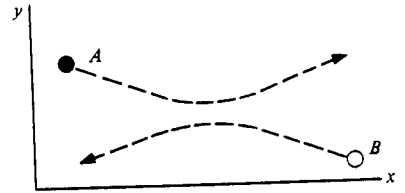
- What does Special Relativity do to Dynamics?
- What has to be changed to preserve 'Conservation of Momentum' - a pillar principle of Dynamics.

- Elastic collision of two identical particles A and B:

Frame-A : moves along x-axis fixed to A.
 Frame-B : moves along x-axis fixed to B.

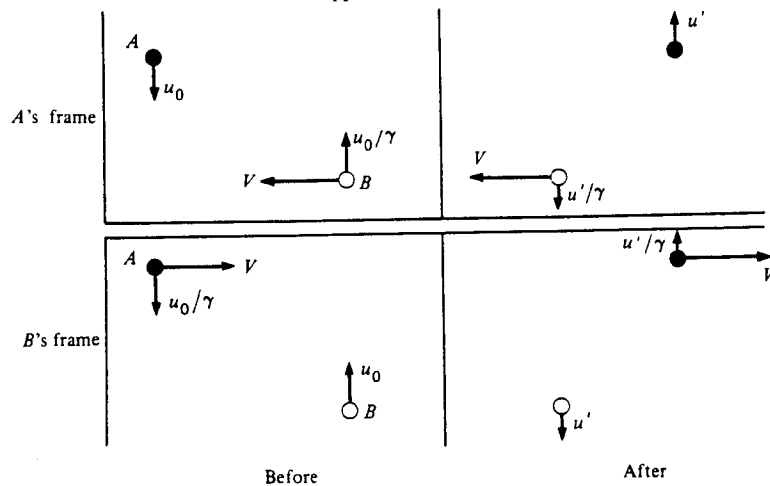
- Collision completely symmetrical:

- Each particle has the same speed u_0 in its own frame along the y-axis



- Collision alters y-velocities leaving x-motion constant.

- Relative x-velocity of the frames is V .



- The law of transformation of velocities gives the y -velocity of the opposite particle as

$$\frac{u_0}{\gamma} = u_0 \sqrt{1 - \frac{v^2}{c^2}}$$

- After the collision the y -velocities are both reversed.

Assume :

$$\vec{p} = m(\omega) \vec{\omega}$$

$m(\omega)$: Something that may depend on particle speed $\vec{\omega}$.

Frame - A :

x -momentum entirely due to particle - B

Before collision B's speed

$$\omega^2 = v^2 + \frac{u_0^2}{\gamma^2}$$

After collision

$$\omega'^2 = v^2 + \frac{u'^2}{\gamma^2}$$

writing conservation of momentum along -x

$$m(\omega) v = m(\omega') v$$

$$\therefore \omega = \omega'$$

$$\therefore u' = u_0$$

Write conservation of momentum in Frame-A along y :

$$-m(u_0) u_0 + m(\omega) \frac{u_0}{\gamma} = m(u_0) u_0 - m(\omega) \frac{u_0}{\gamma}$$

$$\therefore m(\omega) = \gamma m(u_0)$$

Let $m(u_0) = m_0$ as $u_0 \rightarrow 0$.

m_0 : rest mass. Mass measured in a frame where it is stationary.

In this limit $\omega = v$.

$$\therefore m(v) = \gamma m_0 = \frac{m_0}{\sqrt{1 - v^2/c^2}}$$

- Can be interpreted as the dependence of m on speed.

In general:

$$\vec{p} = \frac{m_0 \vec{u}}{\sqrt{1 - u^2/c^2}} = m \vec{u} = \gamma m_0 \vec{u}$$

$$m = \frac{m_0}{\sqrt{1 - u^2/c^2}} = \gamma m_0$$

Relativistic Energy

R5-4

- Generalize classical concept
- Preserve "Conservation of Energy"
- Classical Approach

We defined:

$$K_b - K_a = \int_a^b \frac{d\vec{p}}{dt} \cdot d\vec{r}$$

Consider a particle moving with velocity \vec{u}

$$\vec{p} = m \vec{u} \quad m = \text{constant here}$$

$$K_b - K_a = \int_a^b \frac{d}{dt} (m \vec{u}) \cdot d\vec{r}$$

$$= \int_a^b m \frac{d\vec{u}}{dt} \cdot \vec{u} dt$$

$$= \int_a^b m \vec{u} \cdot d\vec{u}$$

$$\text{Use } \vec{u} \cdot d\vec{u} = \frac{1}{2} d(\vec{u} \cdot \vec{u}) = \frac{1}{2} d(u^2) = u du$$

$$\therefore K_b - K_a = \frac{1}{2} m u_b^2 - \frac{1}{2} m u_a^2$$

Relativity:

R5-5

Try a similar approach

$$\vec{p} = \gamma m_0 \vec{u}$$

$$\begin{aligned} K_b - K_a &= \int_a^b \frac{d\vec{p}}{dt} \cdot d\vec{r} \\ &= \int_a^b \frac{d}{dt} \left[\frac{m_0 \vec{u}}{\sqrt{1 - u^2/c^2}} \right] \cdot \vec{u} dt \\ &= \int_a^b \vec{u} \cdot d \left[\frac{m_0 \vec{u}}{\sqrt{1 - u^2/c^2}} \right] \end{aligned}$$

The integrand is $\vec{u} \cdot d\vec{p} = d(\vec{u} \cdot \vec{p}) - \vec{p} \cdot d\vec{u}$

$$\begin{aligned} K_b - K_a &= \vec{u} \cdot \vec{p} \Big|_a^b - \int_a^b \vec{p} \cdot d\vec{u} \\ &= \frac{m_0 u^2}{\sqrt{1 - u^2/c^2}} \Big|_a^b - \int_a^b \frac{m_0 u du}{\sqrt{1 - u^2/c^2}} \end{aligned}$$

we used $\vec{u} \cdot d\vec{u} = u du$ [see earlier]

Integral is elementary,

$$K_b - K_a = \frac{m_0 u^2}{\sqrt{1 - u^2/c^2}} \Big|_a^b + m_0 c^2 \sqrt{1 - \frac{u^2}{c^2}} \Big|_a^b$$

b is an arbitrary point
 Assume at point a, $u_a = 0$

$$\begin{aligned}
 K &= \gamma m_0 u^2 + \frac{m_0 c^2}{\gamma} - m_0 c^2 \\
 &= \gamma m_0 \left[u^2 + c^2 \left(1 - \frac{u^2}{c^2} \right) \right] - m_0 c^2 \\
 &= \gamma m_0 c^2 - m_0 c^2
 \end{aligned}$$

$$K = mc^2 - m_0 c^2 \quad \text{where } m = \gamma m_0$$

What happens when $v \ll c$

$$\gamma \sim 1 + \frac{1}{2} \frac{u^2}{c^2} + \dots$$

$$K = \frac{m_0 c^2}{\sqrt{1 - \frac{u^2}{c^2}}} - m_0 c^2$$

$$\approx m_0 c^2 \left[1 + \frac{1}{2} \frac{u^2}{c^2} - 1 \right]$$

$$K = \frac{1}{2} m_0 u^2 \quad [\text{classical Result}]$$

RS-7

KE results from work done on particle to bring it from rest to speed u .

Rewrite:

$$mc^2 = K + m_0c^2$$

$$= \text{Work Done} + m_0c^2$$

Einstein:

$$mc^2 \equiv \text{Total Energy, } E, \text{ of particle}$$

$$= \text{External Work} + \text{"Rest" Energy}$$

$$\therefore E = mc^2 = \gamma m_0c^2$$

If an energy ΔE is added to an object, its mass changes by

$$\Delta m = \frac{\Delta E}{c^2}$$

Energy / Momentum

R5-8

Classically, $E = \frac{1}{2} m v^2 = \frac{p^2}{2m}$ $u=0$

$$\vec{p} = m\vec{u} = \frac{m_0 \vec{u}}{\sqrt{1-u^2/c^2}} = \gamma m_0 \vec{u}$$

$$E = mc^2 = \gamma m_0 c^2$$

$$p^2 = \gamma^2 m_0^2 u^2 = \frac{1}{1-u^2/c^2} m_0^2 u^2$$

$$\frac{u^2}{c^2} = \frac{p^2}{p^2 + m_0^2 c^2}$$

$$\gamma = \frac{1}{\sqrt{1-u^2/c^2}} = \sqrt{1 + \frac{p^2}{m_0^2 c^2}}$$

$$\therefore E = m_0 c^2 \sqrt{1 + \frac{p^2}{m_0^2 c^2}}$$

$$\therefore E^2 = (pc)^2 + (m_0 c^2)^2$$

$$E^2 - (pc)^2 = (m_0 c^2)^2$$

4-Vector
Invariant Quantity

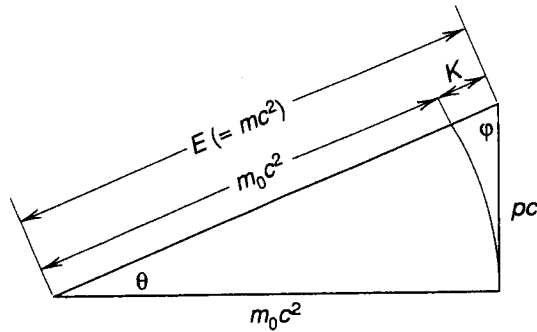


FIGURE 3-3. A mnemonic device, using a right triangle and the Pythagorean relation, to help in remembering the relations between total energy E , rest energy m_0c^2 , and momentum p ; see Eq. 3-13b. Shown also is the relation $E = m_0c^2 + K$ between total energy, rest energy, and kinetic energy. You can show that $\sin \theta = \beta$ and $\sin \phi = 1/\gamma$.

$$\sin \theta = \beta = \frac{u}{c}$$

$$\sin \phi = \frac{1}{\gamma}$$

Massless Particles

R5-10

$$E^2 = (pc)^2 + (m_0c^2)^2$$

If $m_0 = 0$

$$E = pc$$

$$\vec{p} = \gamma m_0 \vec{u} = \frac{m_0}{\sqrt{1 - u^2/c^2}} \vec{u}$$

As $m_0 \rightarrow 0$ \vec{p} must remain finite

Only possible if $u \rightarrow c$ as $m_0 \rightarrow 0$

\therefore massless particles must travel with c !

Photons have $m_0 = 0$

Neutrinos have $m_0 \sim 0$

[close/maybe?]

Forces and Relativity

R5-11

$$\begin{aligned}\vec{F} &= \frac{d\vec{p}}{dt} = \frac{d(m\vec{v})}{dt} \\ &= m \frac{d\vec{v}}{dt} + \vec{v} \frac{dm}{dt}\end{aligned}\quad (1)$$

$$m = E/c^2$$

$$\therefore \frac{dm}{dt} = \frac{1}{c^2} \frac{dE}{dt} = \frac{1}{c^2} \frac{d}{dt} (K + m_0 c^2) = \frac{1}{c^2} \frac{dK}{dt}$$

but

$$\frac{dK}{dt} = \frac{\vec{F} \cdot d\vec{r}}{dt} = \vec{F} \cdot \frac{d\vec{r}}{dt} = \vec{F} \cdot \vec{v}$$

so

$$\frac{dm}{dt} = \frac{1}{c^2} \vec{F} \cdot \vec{v}$$

Sub. into (1)

$$\vec{F} = m \frac{d\vec{v}}{dt} + \frac{\vec{v} (\vec{F} \cdot \vec{v})}{c^2}$$

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{1}{m} \left[\vec{F} - \frac{\vec{v}}{c^2} (\vec{F} \cdot \vec{v}) \right]$$

\vec{a} is no longer \parallel to \vec{F} ,
has a component \parallel to \vec{v} itself.

Relativistic Force / Acceleration

R5-12

