MASSACHUSETTS INSTITUTE OF TECHNOLOGY Physics Department November 5, 2011

Physics 8.286: The Early Universe Prof. Alan Guth

## PROBLEM SET 6

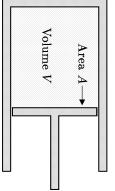
DUE DATE: Thursday, November 10, 2011

**READING ASSIGNMENT:** Steven Weinberg, The First Three Minutes, Chapter 7. Barbara Ryden, Introduction to Cosmology, Chapter 10. We have skipped Chapters 7–9 for now, but we will come back to at least some of them. Chapter 10, about Nucleosynthesis and the Early Universe, makes good parallel reading to Weinberg's book, and really has no dependence on the chapters that we are skipping.

**UPCOMING QUIZ:** Thursday, December 8, 2011.

## **PROBLEM 1: GAS PRESSURE AND ENERGY CONSERVATION** (10 points)

In this problem we will pursue the implications of the conservation of energy. Consider first a gas contained in a chamber with a movable piston, as shown below:



Let U denote the total energy of the gas, and let p denote the pressure. Suppose that the piston is moved a distance dx to the right. (We suppose that the motion is slow, so that the gas particles have time to respond and to maintain a uniform pressure throughout the volume.) The gas exerts a force pA on the piston, so the gas does work dW = pAdx as the piston is moved. Note that the volume increases by an amount dV = Adx, so dW = pdV. The energy of the gas decreases by this

$$dU = -pdV . (P1.1)$$

It turns out that this relation is valid whenever the volume of a gas is changed, regardless of the shape of the volume.

Now consider a homogeneous, isotropic, expanding universe, described by a scale factor a(t). Let u denote the energy density of the gas that fills it. (Remember

8.286 PROBLEM SET 6, FALL 2011

that  $u = \rho c^2$ , where  $\rho$  is the mass density of the gas.) We will consider a fixed coordinate volume  $V_{coord}$ , so the physical volume will vary as

$$y_{\rm s}(t) = a^3(t) V_{\rm coord} \ . \tag{P1.2}$$

 $V_{\rm ph}$ 

The energy of the gas in this region is then given by

$$U = V_{\rm phys} u . \tag{P1.3}$$

(a) Using these relations, show that

$$\frac{d}{dt}\left(a^{3}\rho c^{2}\right) = -p\frac{d}{dt}\left(a^{3}\right),\qquad(P1.4)$$

and then that

$$\dot{\rho} = -3\frac{a}{a}\left(\rho + \frac{p}{c^2}\right) , \qquad (P1.5)$$

where the dot denotes differentiation with respect to t.

(b) The scale factor evolves according to the relation

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3}G\rho - \frac{kc^2}{a^2}.$$
(P1.6)

Using Eqs. (P1.5) and (P1.6), show that

$$= -\frac{4\pi}{3}G\left(\rho + \frac{3p}{c^2}\right)a \ . \tag{P1.7}$$

a:

This equation describes directly the deceleration of the cosmic expansion. Note that there are contributions from the mass density  $\rho$ , but also from the pressure p.

(c) So far our equations have been valid for any sort of a gas, but let us now specialize to the case of black-body radiation. For this case we know that  $\rho = bT^4$ , where b is a constant and T is the temperature. We also know that as the universe expands, aT remains constant. Using these facts and Eq. (P1.5), find an expression for p in terms of  $\rho$ .

8.286 PROBLEM SET 6, FALL 2011 p. 3	8.286 PROBLEM SET 6, FALL 2011 p. 4
PROBLEM 2: THE EFFECT OF PRESSURE ON COSMOLOGICAL EVOLUTION (10 points)	PROBLEM 4: EFFECT OF AN EXTRA NEUTRINO SPECIES (5 points)
<ul> <li>A radiation-dominated universe behaves differently from a matter-dominated universe because the pressure of the radiation is significant. In this problem we explore the role of pressure for several fictitious forms of matter.</li> <li>(a) For the first fictitious form of matter, the mass density ρ decreases as the scale factor a(t) grows, with the relation</li> </ul>	According to the standard assumptions (which were used in the lecture notes), there are three species of effectively massless neutrinos. In the temperature range of 1 MeV $< kT < 100$ MeV, the mass density of the universe is believed to have been dominated by the black-body radiation of photons, electron-positron pairs, and these neutrinos, all of which were in thermal equilibrium.
$ ho(t) \propto rac{1}{a^6(t)}$ .	<ul> <li>(a) Under these assumptions, how long did it take (starting from the instant of the big bang) for the temperature to fall to the value such that kT = 1 MeV?</li> <li>(b) How much time would it have taken if there were one other species of massless</li> </ul>
<ul> <li>What is the pressure of this form of matter? [Hint: the answer is proportional to the mass density.]</li> <li>(b) Find the behavior of the scale factor a(t) for a flat universe dominated by the form of matter described in part (a). You should be able to determine the function a(t) up to a constant factor.</li> </ul>	<ul> <li>(b) How much time would it have taken if there were one other species of massless neutrino, in addition to the three which we are currently assuming?</li> <li>(c) What would be the mass density of the universe when kT = 1 MeV under the standard assumptions, and what would it be if there were one other species of massless neutrino?</li> </ul>
(c) Now consider a universe dominated by a different form of fictitious matter, with a pressure given by $p = \frac{1}{2}\rho c^2 .$	Total points for Problem Set 6: 30.
$ ho(t) \propto rac{1}{a^n(t)}  .$	
Find the power $n$ .	
PROBLEM 3: TIME EVOLUTION OF A UNIVERSE WITH MYSTE- RIOUS STUFF (5 points)	
Suppose that a model universe is filled with a peculiar form of matter for which	
$ ho \propto rac{1}{a^5(t)} .$	
Assuming that the model universe is flat, calculate (a) The behavior of the scale factor, $a(t)$ . You should be able to find $a(t)$ up to an arbitrary constant of proportionality.	
(b) The value of the Hubble parameter $H(t)$ , as a function of t.	
(c) The physical horizon distance, $\ell_{p,\text{horizon}}(t)$ .	

(d) The mass density  $\rho(t)$ .