			Physics Department	Physics 8.286: The Early Universe Noven	Drof Alon Cuth
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November 6, 2007

QUIZ 2

Prof. Alan Guth

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BLANK PAGES AND A SUMMARY OF USEFUL INFORMATION ARE AT THE END OF THE EXAM.

	100	TOTAL
	30	4
	26	3
	20	2
	24	1
Score	Maximum	Problem

Your Name

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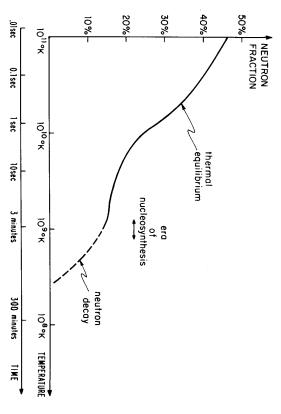
PROBLEM 1: DID YOU DO THE READING? (24 points)

- (a) (6 points) In 1948 Ralph A. Alpher and Robert Herman wrote a paper predicting a cosmic microwave background with a temperature of 5 K. The paper was based on a cosmological model that they had developed with George Gamow, in which the early universe was assumed to have been filled with hot neutrons. As the universe expanded and cooled the neutrons underwent beta decay into protons, electrons, and antineutrinos, until at some point the universe cooled enough for light elements to be synthesized. Alpher and Herman found that to account for the observed present abundances of light elements, the ratio of photons to nuclear particles must have been about 10⁹. Although the predicted from our present theory in two ways. Circle the two correct statements in the following list. (3 points for each right answer; circle at most 2.)
- (i) Gamow, Alpher, and Herman assumed that the neutron could decay, but now the neutron is thought to be absolutely stable.
- (ii) In the current theory, the universe started with nearly equal densities of protons and neutrons, not all neutrons as Gamow, Alpher, and Herman assumed.
- (iii) In the current theory, the universe started with mainly alpha particles, not all neutrons as Gamow, Alpher, and Herman assumed. (Note: an alpha particle is the nucleus of a helium atom, composed of two protons and two neutrons.)
- (iv) In the current theory, the conversion of neutrons into protons (and vice versa) took place mainly through collisions with electrons, positrons, neutrinos, and antineutrinos, not through the decay of the neutrons.
- (v) The ratio of photons to nuclear particles in the early universe is now believed to have been about 10^3 , not 10^9 as Alpher and Herman concluded.
- (b) (6 points) In Weinberg's "Recipe for a Hot Universe," he described the primordial composition of the universe in terms of three conserved quantities: electric charge, baryon number, and lepton number. If electric charge is measured in units of the electron charge, then all three quantities are integers for which the number density can be compared with the number density of photons. For each quantity, which choice most accurately describes the initial ratio of the number density of this quantity to the number density of photons:

	Lepton Number:	Baryon Number:	Electric Charge:
$(iv) \sim 10^{-6}$	(i) $\sim 10^9$	$\begin{array}{l} (\mathrm{i}) \sim 10^{-20} \\ (\mathrm{iv}) \sim 1 \end{array}$	(i) $\sim 10^9$ (iv) $\sim 10^{-6}$
(v) could be as high as ~ 1 , but is assumed to be very small	(ii) ~ 1000 (iii) ~ 1	(ii) $\sim 10^{-9}$ (iii) $\sim 10^{-6}$ (v) anywhere from 10^{-5} to 1	Electric Charge: (i) $\sim 10^9$ (ii) $\sim 1000(\text{iii}) \sim 1$ (iv) $\sim 10^{-6}$ (v) either zero or negligible

 (\circ)

 $(12 \ points)$ The figure below comes from Weinberg's Chapter 5, and is labeled The Shifting Neutron-Proton Balance.



- (i) (3 points) During the period labeled "thermal equilibrium," the neutron fraction is changing because (choose one):
- (A) The neutron is unstable, and decays into a proton, electron, and antineutrino with a lifetime of about 1 second.
- (B) The neutron is unstable, and decays into a proton, electron, and antineutrino with a lifetime of about 15 seconds.
- (C) The neutron is unstable, and decays into a proton, electron, and antineutrino with a lifetime of about 15 minutes.
- (D) Neutrons and protons can be converted from one into through reactions such as

antineutrino + proton \longleftrightarrow electron + neutron neutrino + neutron \longleftrightarrow positron + proton.

(E) Neutrons and protons can be converted from one into the other through reactions such as

antineutrino + proton \longleftrightarrow positron + neutron neutrino + neutron \longleftrightarrow electron + proton.

(F) Neutrons and protons can be created and destroyed by reactions such as

proton + neutrino \longleftrightarrow positron + antineutrino neutron + antineutrino \longleftrightarrow electron + positron.

- (ii) (3 points) During the period labeled "neutron decay," the neutron fraction is changing because (choose one):
- (A) The neutron is unstable, and decays into a proton, electron, and antineutrino with a lifetime of about 1 second.
- (B) The neutron is unstable, and decays into a proton, electron, and antineutrino with a lifetime of about 15 seconds.
- (C) The neutron is unstable, and decays into a proton, electron, and antineutrino with a lifetime of about 15 minutes.
- (D) Neutrons and protons can be converted from one into the other through reactions such as

antineutrino + proton \longleftrightarrow electron + neutron neutrino + neutron \longleftrightarrow positron + proton.

(E) Neutrons and protons can be converted from one into the other through reactions such as

antineutrino + proton \longleftrightarrow positron + neutron neutrino + neutron \longleftrightarrow electron + proton.

(F) Neutrons and protons can be created and destroyed by reactions such as

 $proton + neutrino \leftrightarrow positron + antineutrino$ neutron + antineutrino \leftrightarrow electron + positron.

- (iii) (3 points) The masses of the neutron and proton are not exactly equal, but instead
- (A) The neutron is more massive than a proton with a rest energy difference of 1.293 GeV (1 GeV = 10^9 eV).
- (B) The neutron is more massive than a proton with a rest energy difference of 1.293 MeV (1 MeV = 10^6 eV).
- (C) The neutron is more massive than a proton with a rest energy difference of 1.293 KeV (1 KeV = 10^3 eV).
- (D) The proton is more massive than a neutron with a rest energy difference of 1.293 GeV.
- (E) The proton is more massive than a neutron with a rest energy difference of 1.293 MeV.
- (F) The proton is more massive than a neutron with a rest energy difference of 1.293 KeV.

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- (iv) (3 points) During the period labeled "era of nucleosynthesis," (choose one:) **PROBLEM 2: THE EFFECT OF PRESSURE ON COSMOLOGICAL** EVOLUTION (20 points)
- This problem was Problem 2 of Problem Set 6.

A radiation-dominated universe behaves differently from a matter-dominated universe because the pressure of the radiation is significant. In this problem we explore the role of pressure for several fictitious forms of matter.

(a) (10 points) For the first fictitious form of matter, the mass density ρ decreases as the scale factor R(t) grows, with the relation

$$ho(t) \propto rac{1}{R^6(t)} \; .$$

(D

of the neutrons remain free.

About half the neutrons present combine with protons to form deu-

terium nuclei, which mostly survive until the present time, and the

 (\mathbf{F})

Essentially all the protons present combine with neutrons to form

deuterium nuclei, which mostly survive until the present time.

Essentially all the protons present combine with neutrons to form

helium nuclei, which mostly survive until the present time.

other half of the neutrons remain free.

(E)

0

About half the neutrons present combine with protons to form helium

nuclei, which mostly survive until the present time, and the other half

Essentially all the neutrons present combine with protons to form

deuterium nuclei, which mostly survive until the present time.

 (\mathbf{B})

(A) Essentially all the neutrons present combine with protons to form

helium nuclei, which mostly survive until the present time.

What is the pressure of this form of matter? [Hint: the answer is proportional to the mass density.]

- (b) (5 points) Find the behavior of the scale factor R(t) for a flat universe dominated by the form of matter described in part (a). You should be able to determine the function R(t) up to a constant factor.
- (c) (5 points) Now consider a universe dominated by a different form of fictitious matter, with a pressure given by

$$p = \frac{1}{2}\rho c^2 \; .$$

As the universe expands, the mass density of this form of matter behaves as

$$ho(t) \propto rac{1}{R^n(t)} \; .$$

Find the power n.

PROBLEM 3: A NEW SPECIES OF MESON (26 points)

charge Q = +e, the X⁻ has charge Q = -e, and the X⁰ has charge Q = 0, where e the particle is a boson, it has no spin, and it has three charge states: the X^+ has ing an additional, hypothetical meson, called an X. The X has roughly the same much lighter than the muon, which has a rest energy of 106 MeV. is much heavier than the electron, which has a rest energy of 0.511 MeV, and it is the X^+ is the antiparticle of the X^- , and the X^0 is its own antiparticle. The X is the magnitude of the charge of an electron. There are no additional antiparticles: properties as the pion, except that its mass is given by $mc^2 = 5$ MeV. That is, Suppose the calculations describing the early universe were modified by includ-

- (a) (6 points) What would be the number density of X's, in particles per cubic meter, when the temperature T was given by kT = 30 MeV? Include all the X's: the X^+ , the X^- , and the X^0 .
- (b) (5 points) Assuming (as in the standard picture) that the early universe is $g-cm^{-3}$. value of the mass density at $t = 10^{-3} \text{ sec}$? Express your answer in the units of accurately described by a flat, radiation-dominated model, what would be the
- $\widehat{\mathbf{c}}$ (5 points) What would be the value of kT, in MeV, at $t = 10^{-3}$ sec? You may positron pairs, and the three charge states of the X. icantly to the black-body radiation include the photons, neutrinos, electronassume that 5 MeV $\ll kT \ll 100$ MeV, so the particles contributing signif-
- d (10 points) When kT falls below 5 MeV, the X's will disappear from the thera temperature that that we will call T_{γ} . After the disappearance of the X's, energy to the neutrinos. After the disappearance of the X's, the photons and T_{γ}/T_{ν} , where T_{ν} is the temperature of the neutrinos? but before the disappearance of the electron-positron pairs, what is the ratio the electron positron pairs will be in thermal equilibrium with each other, at by the time kT falls near 5 MeV, so the disappearance of the X's transfers no In the imaginary world the neutrinos have decoupled from the rest of matter in which the neutrinos interact somewhat more weakly than in the real world. purpose of this problem we will discuss an imaginary world which has X's, and teracting significantly with the other particles at this temperature, but for the mal equilibrium mix. For realistic parameters the neutrinos would still be in-

PROBLEM 4: THE STABILITY OF SCHWARZSCHILD ORBITS (30 points

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self-contained answer to this question. allow the grader to assess your understanding, you are expected to present a logical is reproduced for your reference, but you should be aware that the solution to the present problem has important differences. You can copy from this material, but to the statement and the solution are reproduced at the end of this quiz. This material This problem is an elaboration of Problem 1 of Problem Set 5, for which both

of the orbits in a Schwarzschild geometry shows that the smallest *stable* circular orbit occurs for $r = 3R_S$. Circular orbits are possible for $\frac{3}{2}R_S < r < 3R_S$, but statement. they are not stable. In this problem we will explore the calculations behind this In the solution to that homework problem, it was stated that further analysis

be fixed, but all the other coordinates will vary as the body follows its orbit. orbit at $r(t) = r_0$, $\theta = \pi/2$, where r_0 is a constant. The coordinate θ will therefore We will consider a body which undergoes small oscillations about a circular

(a) (12 points) The first step, since $r(\tau)$ will not be a constant in this solution, will be to derive the equation of motion for $r(\tau)$. That is, for the Schwarzschild metric

$$ds^{2} = -c^{2}d\tau^{2} = -h(r)c^{2}dt^{2} + h(r)^{-1}dr^{2} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta \,d\phi^{2} , \quad (1)$$

where

$$h(r)\equiv 1-rac{R_S}{r}\;,$$

work out the explicit form of the geodesic equation

$$\frac{d}{d\tau} \left[g_{\mu\nu} \frac{dx^{\nu}}{d\tau} \right] = \frac{1}{2} \frac{\partial g_{\lambda\sigma}}{\partial x^{\mu}} \frac{dx^{\lambda}}{d\tau} \frac{dx^{\sigma}}{d\tau} , \qquad (2)$$

for the case $\mu = r$. You should use this result to find an explicit expression for

$$\frac{d^2r}{d\tau^2}$$

r, and the derivative with respect to τ of any coordinate, including $dt/d\tau$. You may allow your answer to contain h(r), its derivative h'(r) with respect to

Problem 4 Continues

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(b) (6 points) It is useful to consider r and ϕ to be the independent variables, while treating t as a dependent variable. Find an expression for

$$\left(\frac{dt}{d\tau}\right)^2$$

in terms of r, $dr/d\tau$, $d\phi/d\tau$, h(r), and c. Use this equation to simplify the expression for $d^2r/d\tau^2$ obtained in part (a). The goal is to obtain an expression of the form

$$\frac{d^2r}{d\tau^2} = f_0(r) + f_1(r) \left(\frac{d\phi}{d\tau}\right)^2 . \tag{3}$$

where the functions $f_0(r)$ and $f_1(r)$ might depend on R_S or c, and might be positive, negative, or zero. Note that the intermediate steps in the calculation involve a term proportional to $(dr/d\tau)^2$, but the net coefficient for this term vanishes.

(c) (7 points) To understand the orbit we will also need the equation of motion for ϕ . Evaluate the geodesic equation (2) for $\mu = \phi$, and write the result in terms of the quantity L, defined by

$$\equiv r^2 \frac{d\phi}{d\tau} \,. \tag{4}$$

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(d) (5 points) Finally, we come to the question of stability. Substituting Eq. (4) into Eq. (3), the equation of motion for r can be written as

$$\frac{d^2r}{d\tau^2} = f_0(r) + f_1(r)\frac{L^2}{r^4} \,.$$

Now consider a small perturbation about the circular orbit at $r = r_0$, and write an equation that determines the stability of the orbit. (That is, if some external force gives the orbiting body a small kick in the radial direction, how can you determine whether the perturbation will lead to stable oscillations, or whether it will start to grow?) You should express the stability requirement in terms of the unspecified functions $f_0(r)$ and $f_1(r)$. You are NOT asked to carry out the algebra of inserting the explicit forms that you have found for these functions.

USEFUL INFORMATION:

DOPPLER SHIFT (For motion along a line):

$$z = v/u$$
 (nonrelativistic, source moving)

$$z = \frac{v/u}{1 - v/u} \quad \text{(nonrelativistic, observer moving)}$$
$$z = \sqrt{\frac{1+\beta}{1-\alpha}} - 1 \quad \text{(special relativity, with } \beta = v/c\text{)}$$

COSMOLOGICAL REDSHIFT:

 $\bigvee 1 - \beta$

$$1 + z \equiv \frac{\lambda_{\rm observed}}{\lambda_{\rm emitted}} = \frac{R(t_{\rm observed})}{R(t_{\rm emitted})}$$

SPECIAL RELATIVITY:

Time Dilation Factor:

$$\gamma \equiv rac{1}{\sqrt{1-eta^2}} \ , \qquad \beta \equiv v/c$$

Lorentz-Fitzgerald Contraction Factor: γ

Relativity of Simultaneity:

Trailing clock reads later by an amount $\beta\ell_0/c$

Energy-Momentum Four-Vector:

$$p^{\mu} = \left(\frac{E}{c}, \vec{p}\right), \quad \vec{p} = \gamma m_0 \vec{v}, \quad E = \gamma m_0 c^2 = \sqrt{(m_0 c^2)^2 + |\vec{p}|^2 c^2},$$
$$p^2 \equiv |\vec{p}|^2 - (p^0)^2 = |\vec{p}|^2 - \frac{E^2}{c^2} = -(m_0 c)^2.$$

COSMOLOGICAL EVOLUTION:

$$H^2 = \left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi}{3}G\rho - \frac{kc^2}{R^2} \ , \quad \ddot{R} = -\frac{4\pi}{3}G\left(\rho + \frac{3p}{c^2}\right)R \ ,$$

$$\rho_m(t) = \frac{R^3(t_i)}{R^3(t)} \, \rho_m(t_i) \text{ (matter)}, \ \ \rho_r(t) = \frac{R^4(t_i)}{R^4(t)} \, \rho_r(t_i) \text{ (radiation)}.$$

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$$\dot{
ho} = -3rac{\dot{R}}{R}\left(
ho + rac{p}{c^2}
ight) \ , \ \ \Omega \equiv
ho/
ho_c \ , \ \ {
m where} \ \
ho_c = rac{3H^2}{8\pi G} \ .$$

Flat
$$(k = 0)$$
: $R(t) \propto t^{2/3}$ (matter-dominated),
 $R(t) \propto t^{1/2}$ (radiation-dominated),
 $\Omega = 1$.

EVOLUTION OF A MATTER-DOMINATED UNIVERSE:

$$\begin{array}{lll} \mbox{Closed } (k>0) \colon & ct = \alpha(\theta - \sin\theta) \ , & \frac{R}{\sqrt{k}} = \alpha(1 - \cos\theta) \ , \\ & \Omega = \frac{2}{1 + \cos\theta} > 1 \ , \\ & \mbox{where } \alpha \equiv \frac{4\pi}{3} \frac{G\rho}{c^2} \left(\frac{R}{\sqrt{k}}\right)^3 \ . \\ \mbox{Open } (k<0) \colon & ct = \alpha \left(\sinh\theta - \theta\right) \ , & \frac{R}{\sqrt{\kappa}} = \alpha \left(\cosh\theta - 1\right) \ , \\ & \Omega = \frac{2}{1 + \cosh\theta} < 1 \ , \\ & \mbox{where } \alpha \equiv \frac{4\pi}{3} \frac{G\rho}{c^2} \left(\frac{R}{\sqrt{\kappa}}\right)^3 \ , \\ & \kappa \equiv -k > 0 \ . \end{array}$$

ROBERTSON-WALKER METRIC:

$$ds^{2} = -c^{2} d\tau^{2} = -c^{2} dt^{2} + R^{2}(t) \left\{ \frac{dr^{2}}{1 - kr^{2}} + r^{2} \left(d\theta^{2} + \sin^{2} \theta \, d\phi^{2} \right) \right\}$$

SCHWARZSCHILD METRIC:

$$ds^{2} = -c^{2}d\tau^{2} = -\left(1 - \frac{2GM}{rc^{2}}\right)c^{2}dt^{2} + \left(1 - \frac{2GM}{rc^{2}}\right)^{-1}dr^{2} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta \,d\phi^{2} ,$$

GEODESIC EQUATION:

or:
$$\frac{d}{ds} \left\{ g_{ij} \frac{dx^j}{ds} \right\} = \frac{1}{2} \left(\partial_i g_{k\ell} \right) \frac{dx^k}{ds} \frac{dx^\ell}{ds}$$
$$\frac{d}{d\tau} \left\{ g_{\mu\nu} \frac{dx^\nu}{d\tau} \right\} = \frac{1}{2} \left(\partial_\mu g_{\lambda\sigma} \right) \frac{dx^\lambda}{d\tau} \frac{dx^\sigma}{d\tau}$$

 $g^*_{e^+e^-} = rac{3}{4} imes rac{1}{4} imes rac{1}{4} imes$ Fermion Species a

2

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Particle/ Spin states

$$\begin{split} & u = g \frac{\pi^2}{30} \frac{(kT)^4}{(\hbar c)^3} \qquad (\text{energy density}) \\ & p = \frac{1}{3} u \qquad \rho = u/c^2 \qquad (\text{pressure, mass density}) \\ & n = g^* \frac{\zeta(3)}{\pi^2} \frac{(kT)^3}{(\hbar c)^3} \qquad (\text{number density}) \\ & s = g \frac{2\pi^2}{45} \frac{k^4 T^3}{(\hbar c)^3} , \qquad (\text{entropy density}) \end{split}$$

where

$$g \equiv \begin{cases} 1 \text{ per spin state for bosons (integer spin)} \\ 7/8 \text{ per spin state for fermions (half-integer spin)} \end{cases}$$

$$g^* \equiv \begin{cases} 1 \text{ per spin state for bosons} \\ \frac{3/4}{2} \text{ per spin state for fermion} \end{cases}$$

$$(3/4 \text{ per spin state for fermions})$$

and

$$\begin{split} \zeta(3) &= \frac{1}{1^3} + \frac{1}{2^3} + \frac{1}{3^3} + \dots \approx 1.202 \ . \\ &= g_{\gamma}^* = 2 \ , \\ &= \frac{7}{2} \ \times \ 3 \ \times \ 2 \ \times \ 1 \ = \ 2 \end{split}$$

$$\begin{split} g_{\gamma} &= g_{\gamma}^{*} = 2 \ , \\ g_{\nu} &= \frac{7}{8} \times 3 \times 2 \times 1 = \frac{21}{4} \\ & \text{Fermion 3 species Particle/ Spin states} \\ g_{\nu}^{*} &= \frac{3}{4} \times 3 \times 2 \times 1 = \frac{9}{4} \\ g_{\nu}^{*} &= \frac{3}{4} \times 3 \text{ species Particle/ Spin states} \\ & \text{Fermion 3 species Particle/ Spin states} \\ & \text{Fermion 3 species Particle/ Spin states} \\ \end{array}$$

$$=\frac{3}{4}\times 3 \times 2 \times 1 = \frac{9}{2},$$

Fermion 3 species Particle/ Spin states
factor $\nu_{e},\nu_{\mu},\nu_{\tau}$ antiparticle $\sum_{e^{-}}^{+}=\frac{7}{8}\times 1 \times 2 \times 2 = \frac{7}{2},$

$$\mu_{e^-}=rac{7}{8} imes 1 imes 2 imes 2=rac{7}{2},$$

$$e^{-} = rac{7}{8} imes rac{1}{8} imes rac{1}{8} imes rac{2}{8} imes rac{2}{8} imes rac{1}{8} imes rac{2}{8} imes rac{2}{8} imes rac{7}{2} i$$

$$g_{e^+e^-} = \frac{7}{8} \times \frac{1}{8} \times \frac{2}{8} \times \frac{2}{8} = \frac{7}{2},$$

pecies
$$Particle/$$
 Spin states

EVOLUTION OF A FLAT RADIATION-DOMINATED

$$\begin{split} \rho &= \frac{3}{32\pi G t^2} \\ kT &= \left(\frac{45\hbar^3 c^5}{16\pi^3 g G}\right)^{1/4} \ \frac{1}{\sqrt{t}} \end{split}$$

For $m_{\mu} = 106 \text{ MeV} \gg kT \gg m_e = 0.511 \text{ MeV}, g = 10.75$ and then

$$kT = \frac{0.860 \text{ MeV}}{\sqrt{t \text{ (in sec)}}}$$

After the freeze-out of electron-positron pairs,

$$\frac{T_{\nu}}{T_{\gamma}} = \left(\frac{4}{11}\right)^{1/3} .$$

HORIZON DISTANCE:

$$\begin{split} \ell_{p,\text{horizon}}(t) &= R(t) \int_0^t \frac{c}{R(t')} dt' \\ &= \begin{cases} 3ct \quad \text{(flat, matter-dominated),} \\ 2ct \quad \text{(flat, radiation-dominated).} \end{cases} \end{split}$$

PHYSICAL CONSTANTS:

$$G = 6.673 \times 10^{-8} \text{ cm}^3 \cdot \text{g}^{-1} \cdot \text{s}^{-2}$$

k = Boltzmann's constant = 1.381×10^{-16} er

$$\begin{split} k = \text{Boltzmann's constant} &= 1.381 \times 10^{-16} \text{ erg/K} \\ &= 8.617 \times 10^{-5} \text{ eV/K} \\ \hbar &= \frac{h}{2\pi} = 1.055 \times 10^{-27} \text{ erg-s} \\ &= 6.582 \times 10^{-16} \text{ eV-s} \\ c &= 2.998 \times 10^{10} \text{ cm/s} \\ 1 \text{ yr} &= 3.156 \times 10^7 \text{ s} \\ 1 \text{ eV} &= 1.602 \times 10^{-12} \text{ erg} \\ 1 \text{ GeV} &= 10^9 \text{ eV} = 1.783 \times 10^{-24} \text{ gram (where } c \equiv 1) . \end{split}$$