

THE SCHWARZSCHILD METRIC

$$ds^{2} = -c^{2}d\tau^{2} = -\left(1 - \frac{2GM}{rc^{2}}\right)c^{2}dt^{2} + \left(1 - \frac{2GM}{rc^{2}}\right)^{-1}dr^{2} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta \,d\phi^{2} .$$

- ★ Describes the metric for any spherically symmetric mass distribution, for the region outside the mass distribution. M is the mass of the object, G is Newton's constant, and c is (of course) the speed of light.
- \checkmark At $r = R_S$, where

$$R_{\rm S} = \frac{2GM}{c^2}$$

is called the Schwarzschild horizon, the metric is singular. But the singularity is not physical, and can be removed by a different choice of coordinates. $R_{\rm S}$ is, however, a horizon: anything with $r < R_{\rm S}$ can never get out.

RADIAL GEODESICS

☆ For $\mu = r$, the geodesic equation gives

$$\frac{\mathrm{d}}{\mathrm{d}\tau} \left[g_{rr} \frac{\mathrm{d}r}{\mathrm{d}\tau} \right] = \frac{1}{2} \partial_r g_{rr} \left(\frac{\mathrm{d}r}{\mathrm{d}\tau} \right)^2 + \frac{1}{2} \partial_r g_{tt} \left(\frac{\mathrm{d}t}{\mathrm{d}\tau} \right)^2 \; .$$

which simplifies to

$$\frac{\mathrm{d}^2 r}{\mathrm{d}\tau^2} = -\frac{GM}{r^2} \; . \label{eq:delta_state}$$

It looks like Newton, but r is not really the distance from the origin, and τ is the proper time measured along the trajectory.

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Alan Guth, Black-Body Radiation and the Early History of the Universe, 8.286 Lecture 15, October 31, 2013, p. 2.

