#### 8.286 Lecture 15 October 31, 2013

# BLACK-BODY RADIATION AND THE EARLY HISTORY OF THE UNIVERSE

### Summary of Lecture 14: THE SPACETIME GEODESIC EQUATION

$$\frac{\mathrm{d}}{\mathrm{d}\tau} \left[ g_{\mu\nu} \, \frac{\mathrm{d}x^{\nu}}{\mathrm{d}\tau} \right] = \frac{1}{2} \frac{\partial g_{\lambda\sigma}}{\partial x^{\mu}} \, \frac{\mathrm{d}x^{\lambda}}{\mathrm{d}\tau} \frac{\mathrm{d}x^{\sigma}}{\mathrm{d}\tau} \; .$$

★ Use indices  $\mu$ ,  $\nu$ , etc., which are summed from 0 to 3, where  $x^0 \equiv t$ .

 $\checkmark$  Use  $\tau$  to parameterize the path, where  $\tau$  = proper time measured along the path.



#### THE SCHWARZSCHILD METRIC

$$ds^{2} = -c^{2}d\tau^{2} = -\left(1 - \frac{2GM}{rc^{2}}\right)c^{2}dt^{2} + \left(1 - \frac{2GM}{rc^{2}}\right)^{-1}dr^{2} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta d\phi^{2}.$$

- $\checkmark$  Describes the metric for any spherically symmetric mass distribution, for the region outside the mass distribution. M is the mass of the object, G is Newton's constant, and c is (of course) the speed of light.
- $\checkmark$  At  $r = R_S$ , where

$$R_{\rm S} = \frac{2GM}{c^2}$$

is called the Schwarzschild horizon, the metric is singular. But the singularity is not physical, and can be removed by a different choice of coordinates.  $R_{\rm S}$  is, however, a horizon: anything with  $r < R_{\rm S}$  can never get out.

## RADIAL GEODESICS

 $\checkmark$  For  $\mu = r$ , the geodesic equation gives

$$\frac{\mathrm{d}}{\mathrm{d}\tau} \left[ g_{rr} \frac{\mathrm{d}r}{\mathrm{d}\tau} \right] = \frac{1}{2} \partial_r g_{rr} \left( \frac{\mathrm{d}r}{\mathrm{d}\tau} \right)^2 + \frac{1}{2} \partial_r g_{tt} \left( \frac{\mathrm{d}t}{\mathrm{d}\tau} \right)^2 \; .$$

which simplifies to

$$\frac{\mathrm{d}^2 r}{\mathrm{d}\tau^2} = -\frac{GM}{r^2} \; .$$

It looks like Newton, but r is not really the distance from the origin, and  $\tau$  is the proper time measured along the trajectory.



## Solving the Radial Infall Equation

 $\checkmark$  The proper time  $\tau$  needed to reach radial variable r is

$$\tau(r) = \sqrt{\frac{r_0}{2GM}} \left\{ r_0 \tan^{-1} \left( \sqrt{\frac{r_0 - r}{r}} \right) + \sqrt{r(r_0 - r)} \right\} .$$

The infalling object will be ripped apart by the singularity at r = 0 in a finite amount of the object's proper time.

 $\bigstar$  But from the outside, it will take an infinite amount of coordinate time t before the object reaches the horizon.

