

9.5 Majorana Modes in a Very Simple Model

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9.5.1 The Model: A Quantum Wire

In 1 space dimension, spinless fermions on lattice points $j=1, \dots, L$.

$$H_1 = \sum_j -\omega (c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) - \mu (c_j^\dagger c_j - \frac{1}{2}) + \Delta c_j c_{j+1} + \Delta^* c_{j+1}^\dagger c_j$$

hopping
chemical potential
"condensate" interaction

$$\Delta = |\Delta| e^{i\theta}$$

Write - introducing Majorana operators -

$$\gamma_{2j-1} = e^{i\frac{\theta}{2}} c_j + e^{-i\frac{\theta}{2}} c_j^\dagger \quad \gamma_{2j} = -i(e^{i\frac{\theta}{2}} c_j - e^{-i\frac{\theta}{2}} c_j^\dagger)$$

These satisfy the Clifford algebra $\{\gamma_\ell, \gamma_m\} = 2\delta_{\ell m}$ and $\gamma_\ell^* = \gamma_\ell$. Now

$$H_1 = \frac{i}{2} \sum_j \left[-\mu \gamma_{2j-1} \gamma_{2j} + (\omega + |\Delta|) \gamma_{2j} \gamma_{2j+1} + (-\omega + |\Delta|) \gamma_{2j-1} \gamma_{2j+2} \right]$$

so $c_j = e^{-i\frac{\theta}{2}} \left(\frac{\gamma_{2j-1} + i\gamma_{2j}}{2} \right) \quad c_j^\dagger = e^{i\frac{\theta}{2}} \left(\frac{\gamma_{2j-1} - i\gamma_{2j}}{2} \right)$

9.5.2 Two simple phases

a) If $|\Delta| = \omega = 0, \mu < 0$ (totally trivial!)

$$H_1 = \sum -\mu (c_j^\dagger c_j - \frac{1}{2})$$

$$= -\frac{i}{2} \mu \sum \delta_{2j-1} \delta_{2j}$$

ground state: totally unoccupied

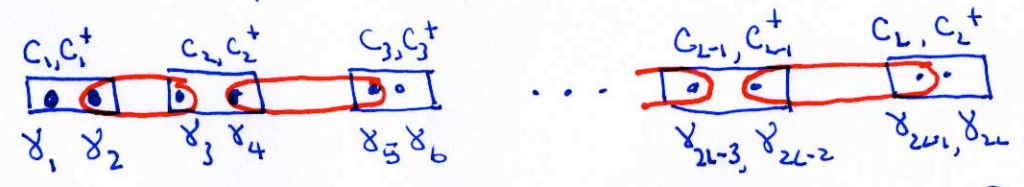
the δ_{2j-1} and δ_{2j} pair off $\delta_{2j-1} \delta_{2j} |g\rangle = -i |g\rangle$

$$H_1 = \frac{i}{2} \sum_j \left[-\mu \delta_{2j-1} \delta_{2j} + (\omega + |\Delta|) \delta_{2j} \delta_{2j+1} + (-\omega + |\Delta|) \delta_{2j-1} \delta_{2j+2} \right]$$

b) If $|\Delta| = \omega > 0, \mu = 0$

$$H_1 = \frac{i}{2} \omega \sum \delta_{2j} \delta_{2j+1}$$

so δ_{2j} pairs with δ_{2j+1} . Note that these come from different sites!



□ lattice site pairing
○ correlation pairing

$$\delta_{2j} \delta_{2j+1} |g\rangle = i |g\rangle$$

γ_1 and γ_2 decouple from the Hamiltonian!

We have 2 ~~degenerate~~ degenerate ground states $|\psi_0\rangle, |\psi_1\rangle$ with

$$\gamma_1 \gamma_{2L} |\psi_0\rangle = i |\psi_0\rangle \quad \gamma_1 \gamma_{2L} |\psi_1\rangle = -i |\psi_1\rangle$$

(Note: We are simultaneously diagonalizing commuting operators ...)

The \mathbb{Z}_2 parity operator

$$P = \prod_j (1 - 2c_j^\dagger c_j) = \prod_j (-i \gamma_{2j-1} \gamma_{2j})$$

commutes with all our bilinears. We see

$$P |\psi_0\rangle = |\psi_0\rangle \quad P |\psi_1\rangle = -|\psi_1\rangle$$

so $|\psi_0\rangle$ contains ~~an~~ an even # of fermions, $|\psi_1\rangle$ an odd number.

9.5.3 Pumping

~~The~~ An adiabatic change in the ^{phase of the} order parameter θ through 2π changes $\gamma_i \rightarrow -\gamma_i$.

[Exercise: Make this a good argument, using Berry's phase]

This same transformation, $\delta_i \rightarrow -\delta_i$, is implemented by P , regarded as a unitary transformation (Pulling δ_i through gives an odd # of - signs: one for every index $1, \dots, 2L$ except i .)

This same transformation is implemented by passing an electron from one side to the other!

Interpretation: We go to the same ground state, but an electron is transported from one end to the other.

[Project: really demonstrate that!
- follow the flow]