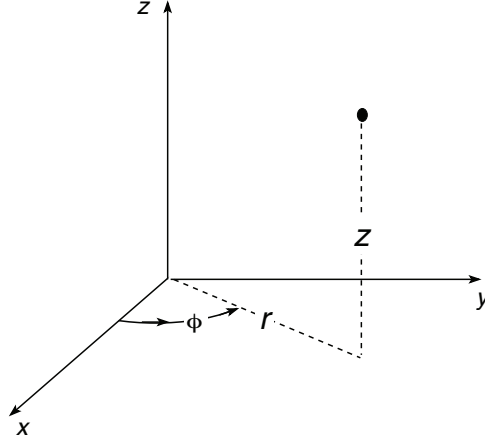


3.6.3 Cylindrical and spherical coordinates

Cylindrical coordinates simply extend the 2D polar coordinates, Eq. (3.6.11), by adding a third coordinate, z , pointing out of the plane.



From the Jacobian

$$\mathbf{J}(r, \phi, z) = \begin{bmatrix} \cos \phi & -r \sin \phi & 0 \\ \sin \phi & r \cos \phi & 0 \\ 0 & 0 & 1 \end{bmatrix}, \quad (3.6.26)$$

we obtain the metric factors

$$h_r = 1, \quad h_\phi = r, \quad h_z = 1, \quad \text{and} \quad |\det \mathbf{J}| = r. \quad (3.6.27)$$

For a scalar field $\Phi(r, \phi, z)$, we then obtain

$$\nabla \Phi = \left(\frac{\partial \Phi}{\partial r}, \frac{1}{r} \frac{\partial \Phi}{\partial \phi}, \frac{\partial \Phi}{\partial z} \right), \quad (3.6.28)$$

and

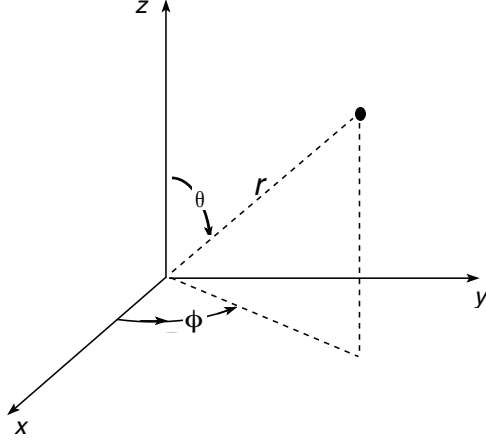
$$\nabla^2 \Phi = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \Phi}{\partial \phi^2} + \frac{\partial^2 \Phi}{\partial z^2}. \quad (3.6.29)$$

A vector field can also be presented in these coordinates with components (v_r, v_ϕ, v_z) that are functions of (r, ϕ, z) . The divergence of the vector is then obtained as

$$\vec{\nabla} \cdot \vec{v} = \frac{1}{r} \frac{\partial}{\partial r} (r v_r) + \frac{1}{r} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}. \quad (3.6.30)$$

Spherical coordinates provide another representation of three dimensional space, replacing the axis z with the angle θ to the z axis, such that $z = r \cos(\theta)$. With r now indicating the distance to the origin, its projection onto the 2D plane has length $r \sin(\theta)$, such that

$$x = r \sin(\theta) \cos(\phi) \quad y = r \sin(\theta) \sin(\phi) \quad z = r \cos(\theta), \quad \text{with} \quad 0 \leq \theta \leq \pi \quad \text{and} \quad 0 \leq \phi < 2\pi. \quad (3.6.31)$$



The Jacobian associated with this transformation is

$$\mathbf{J}(r, \theta, \phi) = \begin{bmatrix} \frac{\partial x}{\partial r} & \frac{\partial x}{\partial \theta} & \frac{\partial x}{\partial \phi} \\ \frac{\partial y}{\partial r} & \frac{\partial y}{\partial \theta} & \frac{\partial y}{\partial \phi} \\ \frac{\partial z}{\partial r} & \frac{\partial z}{\partial \theta} & \frac{\partial z}{\partial \phi} \end{bmatrix} = \begin{bmatrix} \sin \theta \cos \phi & r \cos \theta \cos \phi & -r \sin \theta \sin \phi \\ \sin \theta \sin \phi & r \cos \theta \sin \phi & r \sin \theta \cos \phi \\ \cos \theta & -r \sin \theta & 0 \end{bmatrix}. \quad (3.6.32)$$

From the magnitudes of the column vectors, we find

$$h_r = 1 \quad h_\theta = r, \quad h_\phi = \sin \theta, \quad \text{while} \quad |\det \mathbf{J}| = r^2 \sin \theta. \quad (3.6.33)$$

For a scalar field $\Phi(r, \theta, \phi)$, we then have

$$\vec{\nabla} \Phi = \left(\frac{\partial \Phi}{\partial r}, \frac{1}{r} \frac{\partial \Phi}{\partial \theta}, \frac{1}{r \sin \theta} \frac{\partial \Phi}{\partial \phi} \right), \quad (3.6.34)$$

and

$$\nabla^2 \Phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \Phi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Phi}{\partial \phi^2}. \quad (3.6.35)$$

For a vector field $\vec{v} = (v_r, v_\theta, v_\phi)$,

$$\vec{\nabla} \cdot \vec{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}. \quad (3.6.36)$$