## 2.2 Multiple variables

## 2.2.1 Many coupled ODEs

The results of the previous section can be generalized to multiple variables indexed by  $i = 1, 2, \dots, n$ . The set of coordinates  $\{x_i\}$  can be regraded as a point in n dimensional space, and can also be represented as a vector  $\vec{x}$  extending from the origin to this point. The generalized equations of motion can now be represented as

$$\dot{x}_i = F_i(\{x_i\})$$
 for  $i = 1, 2, \dots n$ , or equivalently as  $\dot{\vec{x}} = \vec{F}(\vec{x})$ . (2.2.1)

The linearized equations take the form

$$\dot{x}_i = \sum_{j=1}^n F_{ij} x_j \text{ for } i = 1, 2, \dots n, \quad \text{ or equivalently as } \dot{\vec{x}} = \mathbf{F} \vec{x}, \qquad (2.2.2)$$

in terms of the  $n \times n$  matrix formed from  $n^2$  elements  $\{F_{ij}\}$ .

A particular class of linear equations is obtained from gradient descent in a quadratic potential, which can be written as

$$V(\{x_i\}) = \frac{1}{2} \sum_{i=1}^{n} \sum_{j=1}^{n} M_{ij} x_i x_j.$$
(2.2.3)

It may appear that  $n^2$  elements are needed to specify the potential. This is in fact not the case since after summation over both *i* and *j*, only the symmetric part  $(M_{ij} + M_{ji})$ contributes as the coefficient of the term  $x_i x_j$ , while the antisymmetric part  $(M_{ij} - M_{ji})$ vanishes. Thus a general quadratic potential can be represented by n(n + 1)/2 elements forming a symmetric matrix, in which case  $F_{ij} = M_{ij}$  in Eq. (2.2.2).