

Supplementary Information

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After substituting our analytical solutions (Eqs. [14], [15], [16], and [17]) into the nondimensionalized, modified sickle cell model (Eqs. [10], [11], [12], and [13]) and dropping negligible terms (Eq. [19]), we are able to extract a system of autonomous nonlinear ordinary differential equations (Eqs. [20], [21], and [22]) and a system of nonlinear algebraic equations (Eq. [25]). The variables ψ_1, ψ_2 , and ψ_3 are undefined functions of τ . The system of nonlinear ordinary differential equations that we will obtain expresses a dynamic relationship among the variables ψ_1, ψ_2, ψ_3 ; and the system of nonlinear algebraic equations expresses a relationship among the free parameters (see [18] and [24]). In the following analysis, we will separately examine each equation in the nondimensionalized, sickle cell model (see [9], [10], [11] and [12]) and then combine our results in the last section.

Equation [9].

Substituting [13] and [14] into [9] yields

$$\psi_1 e^{k_1 r' + k_2 z'} (A_{11} r'^2 + A_{12} r'^3) = 0 \quad [25]$$

where $A_{1i} \ 1 \leq i \leq 2$ are nonlinear algebraic equations. These nonlinear algebraic equations are functions of the free parameters in [18] such that

$$A_{1i} = A_{1i}(k_1, k_2, a_1, a_2, a_3, b_1, b_2, b_3, b_4, b_5, b_6).$$

For [25] to hold true, then

$$A_{11} = 0 \text{ and } A_{12} = 0, \quad [26]$$

because $\psi_1 = 0$ only admits trivial solutions to v'_r and v'_z .

Equation [10].

Substituting [13] and [14] into [10] yields

$$\begin{aligned} \frac{d\psi_1}{d\tau} & \left[e^{k_1 r' + k_2 z'} (B_{11} r' + B_{12} r'^2 + B_{13} r'^3) \right] \\ & + \psi_1^2 e^{k_1 r' + k_2 z'} (B_{21} r' + B_{22} r'^2 + B_{23} r'^3) \\ & + \psi_1^2 e^{2k_1 r' + 2k_2 z'} (B_{32} r'^2 + B_{33} r'^3) \\ & = l_1 \psi_1 e^{k_1 r' + k_2 z'} (B_{41} r' + B_{42} r'^2 + B_{43} r'^3) \end{aligned} \quad [27]$$

where $l_1 = \frac{\mu T}{\rho L^2}$ and $B_{ij} \ 1 \leq i, j \leq 4$ are nonlinear algebraic equations. These nonlinear algebraic equations are functions of the free parameters in [18] such that

$$B_{ij} = B_{ij}(k_1, k_2, a_1, a_2, a_3, b_1, b_2, b_3, b_4, b_5, b_6).$$

According to the Galerkin method, we can arbitrarily assume relations between the B_{ij} 's. Lets set

$$\begin{aligned} B_{2i} &= B_{1i} \quad \text{for } 1 \leq i \leq 3, \\ B_{3i} &= 0 \quad \text{for } 2 \leq i \leq 3, \\ B_{4i} &= B_{1i} \quad \text{for } 1 \leq i \leq 3. \end{aligned} \quad [28]$$

Substituting [28] into [27] yields

$$\begin{aligned} \frac{d\psi_1}{d\tau} & \left[e^{k_1 r' + k_2 z'} (B_{11} r' + B_{12} r'^2 + B_{13} r'^3) \right] \\ & + \psi_1^2 e^{k_1 r' + k_2 z'} (B_{11} r' + B_{12} r'^2 + B_{13} r'^3) \\ & + \psi_1^2 e^{2k_1 r' + 2k_2 z'} (0 \cdot r'^2 + 0 \cdot r'^3) \\ & = l_1 \psi_1 e^{k_1 r' + k_2 z'} (B_{11} r' + B_{12} r'^2 + B_{13} r'^3). \end{aligned} \quad [29]$$

[29] can be factored and then rewritten as

$$\left(\frac{d\psi_1}{d\tau} + \psi_1^2 - l_1 \psi_1 \right) e^{k_1 r' + k_2 z'} (B_{11} r' + B_{12} r'^2 + B_{13} r'^3) = 0. \quad [30]$$

The above relation implies that

$$\frac{d\psi_1}{d\tau} + \psi_1^2 - l_1 \psi_1 = 0, \quad [31]$$

because having B_{11}, B_{12} , and B_{13} equal to zero would yield trivial solutions for v'_r and v'_z . [31] can then be rewritten as

$$\frac{d\psi_1}{d\tau} = l_1 \psi_1 - \psi_1^2 \quad [32]$$

Equation [11].

Substituting [13], [14], [15] and [16] into [11] yields

$$\begin{aligned} \frac{d\psi_2}{d\tau} & \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ & + \psi_1 \psi_2 \left[e^{k_1 r' + k_2 z'} (C_{213} r'^3) \right] \\ & + \psi_1 \left[e^{k_1 r' + k_2 z'} (C_{313} r'^3 + C_{312} r'^2) + C_{303} r'^3 + C_{302} r'^2 + C_{301} r' + C_{300} \right] \\ & = \psi_2 \left[e^{k_1 r' + k_2 z'} (C_{411} r' + C_{412} r'^2 + C_{413} r'^3) \right. \\ & \quad \left. + (C_{403} r'^3 + C_{402} r'^2 + C_{401} r' + C_{400}) \right] \\ & + \psi_2^2 \left[e^{k_1 r' + k_2 z'} (C_{513} r'^3) + C_{503} r'^3 + C_{500} \right] \\ & + \psi_3 \left[e^{k_1 r' + k_2 z'} (C_{613} r'^3) + C_{603} r'^3 + C_{600} \right] \\ & + \psi_2 \psi_3 \left[e^{k_1 r' + k_2 z'} (C_{713} r'^3) + C_{703} r'^3 + C_{700} \right] \\ & \quad + l_2 \psi_2^2 \psi_3 \left[e^{k_1 r' + k_2 z'} (C_{813} r'^3) + C_{803} r'^3 + C_{800} \right] \\ & \quad + C_{003} r'^3 + C_{000} \end{aligned} \quad [33]$$

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where $l_2 = \frac{k^{Hb} [Hb] P_e T}{\alpha P_{O_2, 50\%}^{Hb}}$ and C_{ijk} $_{0 \leq i \leq 12, 0 \leq j \leq 1, 0 \leq k \leq 3}$ are nonlinear algebraic equations. These nonlinear algebraic equations are functions of the free parameters in [18] such that

$$C_{ijk} = C_{ijk}(k_1, k_2, a_1, a_2, a_3, b_1, \dots, b_6, c_1, \dots, c_{23}, d_1, \dots, d_{23}).$$

According to the Galerkin method, we can assume arbitrary relations between the C_{ijk} 's. Lets set

$$C_{ijk} = \begin{cases} \tilde{e}_{i-2} \cdot C_{1jk} & \text{for } i = 3, 0 \leq j \leq 1, k = \{0, 3\} \\ -\tilde{e}_{i-2} \cdot C_{1jk} & \text{for } 4 \leq i \leq 5, 0 \leq j \leq 1, k = \{0, 3\} \\ \tilde{e}_{i-2} \cdot C_{1jk} & \text{for } 6 \leq i \leq 8, 0 \leq j \leq 1, k = \{0, 3\} \end{cases} \quad [34]$$

and also let

$$\begin{aligned} C_{213} = 0, C_{312} = 0, C_{302} = 0, C_{301} = 0, C_{412} = 0, \\ C_{411} = 0, C_{402} = 0, C_{401} = 0, C_{003} = 0, C_{000} = 0. \end{aligned} \quad [35]$$

where \tilde{e}_i $_{1 \leq i \leq 6}$ are undefined, proportionality constants. Substituting [34] and [35] into [33] yields

$$\begin{aligned} & \frac{d\psi_2}{d\tau} \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ & + \left[\psi_1 \psi_2 e^{k_1 r' + k_2 z'} (0 \cdot r'^3) \right] \\ & + \psi_1 \left[e^{k_1 r' + k_2 z'} (\tilde{e}_1 C_{113} r'^3 + 0 \cdot r'^2) \right. \\ & \quad \left. + \tilde{e}_1 C_{103} r'^3 + 0 \cdot r'^2 + 0 \cdot r' + \tilde{e}_1 C_{100} \right] \\ = & \psi_2 \left[e^{k_1 r' + k_2 z'} (0 \cdot r' + 0 \cdot r'^2 - \tilde{e}_2 C_{113} r'^3) \right. \\ & \quad \left. + (-\tilde{e}_2 C_{103} r'^3 + 0 \cdot r'^2 + 0 \cdot r' - \tilde{e}_2 C_{100}) \right] \\ & + \psi_2^2 \left[-e^{k_1 r' + k_2 z'} (\tilde{e}_3 C_{113} r'^3) - \tilde{e}_3 C_{103} r'^3 - \tilde{e}_3 C_{100} \right] \\ & + \psi_3 \left[e^{k_1 r' + k_2 z'} (\tilde{e}_4 C_{113} r'^3) + \tilde{e}_4 C_{103} r'^3 + \tilde{e}_4 C_{100} \right] \\ & + \psi_2 \psi_3 \left[e^{k_1 r' + k_2 z'} (\tilde{e}_5 C_{113} r'^3) + \tilde{e}_5 C_{103} r'^3 + \tilde{e}_5 C_{100} \right] \\ & + l_2 \psi_2^2 \psi_3 \left[e^{k_1 r' + k_2 z'} (\tilde{e}_6 C_{113} r'^3) + \tilde{e}_6 C_{103} r'^3 + \tilde{e}_6 C_{100} \right] \\ & \quad + 0 \cdot r'^3 + 0 \quad [36] \end{aligned}$$

and, after simplifying [36], we obtain

$$\begin{aligned} & \frac{d\psi_2}{d\tau} \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ & + \tilde{e}_1 \psi_1 \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ = & -\tilde{e}_2 \psi_2 \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ & - \tilde{e}_3 \psi_2^2 \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ & + \tilde{e}_4 \psi_3 \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ & + \tilde{e}_5 \psi_2 \psi_3 \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] \\ & + \tilde{e}_6 l_2 \psi_2^2 \psi_3 \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right]. \quad [37] \end{aligned}$$

Factoring [37] yields

$$\begin{aligned} & \left(\frac{d\psi_2}{d\tau} + \tilde{e}_1 \psi_1 + \tilde{e}_2 \psi_2 + \tilde{e}_3 \psi_2^2 - \psi_3 (\tilde{e}_4 + \tilde{e}_5 \psi_2 \right. \\ & \quad \left. + \tilde{e}_6 l_2 \psi_2^2) \right) \cdot \left[e^{k_1 r' + k_2 z'} (C_{113} r'^3) + C_{103} r'^3 + C_{100} \right] = 0 \end{aligned}$$

where \tilde{e}_i $_{7 \leq i \leq 12}$ are undefined, arbitrary constants. Substituting [42] and [43] into [41] yields

The above relation implies that

$$\frac{d\psi_2}{d\tau} + \tilde{e}_1 \psi_1 + \tilde{e}_2 \psi_2 + \tilde{e}_3 \psi_2^2 - \psi_3 (\tilde{e}_4 + \tilde{e}_5 \psi_2 + \tilde{e}_6 l_2 \psi_2^2) = 0, \quad [39]$$

because having C_{113}, C_{103} , and C_{100} equal to zero would yield trivial solutions for v'_r, v'_z, S' and P' . Equation [39] can then be rewritten as

$$\frac{d\psi_2}{d\tau} = -\tilde{e}_1 \psi_1 - (\tilde{e}_2 \psi_2 + \tilde{e}_3 \psi_2^2) + \psi_3 (\tilde{e}_4 + \tilde{e}_5 \psi_2 + \tilde{e}_6 l_2 \psi_2^2). \quad [40]$$

Equation [12].

Substituting [13],[14],[15], and [16] into [12] yields

$$\begin{aligned} & \frac{d\psi_3}{d\tau} \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\ & + \left[\psi_1 \psi_3 e^{k_1 r' + k_2 z'} (D_{213} r'^3) \right] \\ & + \psi_1 \left[e^{k_1 r' + k_2 z'} (D_{313} r'^3 + D_{312} r'^2) + D_{303} r'^3 + D_{302} r'^2 + D_{301} r' + D_{300} \right] \\ = & \psi_2 \left[e^{k_1 r' + k_2 z'} (D_{411} r' + D_{412} r'^2 + D_{413} r'^3) \right. \\ & \quad \left. + (D_{403} r'^3 + D_{402} r'^2 + D_{401} r' + D_{400}) \right] \\ & + \psi_2^2 \left[e^{k_1 r' + k_2 z'} (D_{513} r'^3) + D_{503} r'^3 + D_{500} \right] \\ & + \psi_3 \left[e^{k_1 r' + k_2 z'} (D_{613} r'^3) + D_{603} r'^3 + D_{600} \right] \\ & + \psi_2 \psi_3 \left[e^{k_1 r' + k_2 z'} (D_{713} r'^3) + D_{703} r'^3 + D_{700} \right] \\ & + l_3 \psi_2^2 \psi_3 \left[e^{k_1 r' + k_2 z'} (D_{813} r'^3) + D_{803} r'^3 + D_{800} \right] \\ & \quad + D_{003} r'^3 + D_{000} \quad [41] \end{aligned}$$

where $l_3 = \frac{k^{Hb} P_e^2 T}{P_{O_2, 50\%}^{Hb}}$ and D_{ijk} $_{0 \leq i \leq 8, 0 \leq j \leq 1, 0 \leq k \leq 3}$ are nonlinear algebraic equations. These nonlinear algebraic equations are functions of the free parameters in [18] such that

$$D_{ijk} = D_{ijk}(k_1, k_2, a_1, a_2, a_3, b_1, \dots, b_6, c_1, \dots, c_{23}, d_1, \dots, d_{23}).$$

According to the Galerkin method, we can assume arbitrary relations between the D_{ijk} 's. Lets set

$$D_{ijk} = \begin{cases} \tilde{e}_{i+4} \cdot D_{1jk} & \text{for } 3 \leq i \leq 5, 0 \leq j \leq 1, k = \{0, 3\} \\ -\tilde{e}_{i+4} \cdot D_{1jk} & \text{for } 6 \leq i \leq 8, 0 \leq j \leq 1, k = \{0, 3\} \end{cases} \quad [42]$$

and

$$\begin{aligned} D_{213} = 0, D_{312} = 0, D_{302} = 0, D_{301} = 0, D_{412} = 0, \\ D_{411} = 0, D_{402} = 0, D_{401} = 0, D_{003} = 0, D_{000} = 0. \end{aligned} \quad [43]$$

$$\begin{aligned}
 & \frac{d\psi_3}{d\tau} \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\
 & + \left[\psi_1 \psi_3 e^{k_1 r' + k_2 z'} (0 \cdot r'^3) \right] \\
 & + \psi_1 \left[e^{k_1 r' + k_2 z'} (\tilde{e}_7 D_{113} r'^3 + 0 \cdot r'^2) \right. \\
 & \quad \left. + \tilde{e}_7 D_{103} r'^3 + 0 \cdot r'^2 + 0 \cdot r' + \tilde{e}_7 D_{100} \right] \\
 & = \psi_2 \left[e^{k_1 r' + k_2 z'} (0 \cdot r' + 0 \cdot r'^2 + \tilde{e}_8 D_{113} r'^3) \right. \\
 & \quad \left. + (\tilde{e}_8 D_{103} r'^3 + 0 \cdot r'^2 + 0 \cdot r' + \tilde{e}_8 D_{100}) \right] \\
 & + \psi_2^2 \left[e^{k_1 r' + k_2 z'} (\tilde{e}_9 D_{113} r'^3) + \tilde{e}_9 D_{103} r'^3 + \tilde{e}_9 D_{100} \right] \\
 & + \psi_3 \left[-e^{k_1 r' + k_2 z'} (\tilde{e}_{10} D_{113} r'^3) - \tilde{e}_{10} D_{103} r'^3 - \tilde{e}_{10} D_{100} \right] \\
 & + \psi_2 \psi_3 \left[-e^{k_1 r' + k_2 z'} (\tilde{e}_{11} D_{113} r'^3) - \tilde{e}_{11} D_{103} r'^3 - \tilde{e}_{11} D_{100} \right] \\
 & + l_3 \psi_2^2 \psi_3 \left[-e^{k_1 r' + k_2 z'} (\tilde{e}_{12} D_{113} r'^3) - \tilde{e}_{12} D_{103} r'^3 - \tilde{e}_{12} D_{100} \right] \\
 & \quad + 0 \cdot r'^3 + 0 \quad [44]
 \end{aligned}$$

and, after simplifying [44], we obtain

$$\begin{aligned}
 & \frac{d\psi_3}{d\tau} \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\
 & + \tilde{e}_7 \psi_1 \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\
 & = \tilde{e}_8 \psi_2 \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\
 & \quad + \tilde{e}_9 \psi_2^2 \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\
 & \quad - \tilde{e}_{10} \psi_3 \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\
 & \quad - \tilde{e}_{11} \psi_2 \psi_3 \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] \\
 & \quad - \tilde{e}_{12} l_3 \psi_2^2 \psi_3 \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right]. \quad [45]
 \end{aligned}$$

Factoring [45] yields

$$\begin{aligned}
 & \left(\frac{d\psi_3}{d\tau} + \tilde{e}_7 \psi_1 - \tilde{e}_8 \psi_2 - \tilde{e}_9 \psi_2^2 + \psi_3 (\tilde{e}_{10} + \tilde{e}_{11} \psi_2 \right. \\
 & \quad \left. + \tilde{e}_{12} l_3 \psi_2^2) \right) \cdot \left[e^{k_1 r' + k_2 z'} (D_{113} r'^3) + D_{103} r'^3 + D_{100} \right] = 0.
 \end{aligned}$$

The above relation implies that

$$\frac{d\psi_3}{d\tau} + \tilde{e}_7 \psi_1 - \tilde{e}_8 \psi_2 - \tilde{e}_9 \psi_2^2 + \psi_3 (\tilde{e}_{10} + \tilde{e}_{11} \psi_2 + \tilde{e}_{12} l_3 \psi_2^2) = 0, \quad [46]$$

because having D_{113}, D_{103} and D_{100} equal to zero yields trivial solutions for v'_r, v'_z, S' and P' . Equation [46] can then be rewritten as

$$\frac{d\psi_3}{d\tau} = -\tilde{e}_7 \psi_1 + \tilde{e}_8 \psi_2 + \tilde{e}_9 \psi_2^2 - \psi_3 (\tilde{e}_{10} + \tilde{e}_{11} \psi_2 + \tilde{e}_{12} l_3 \psi_2^2). \quad [47]$$

Results

We can combine the results from each subsection, which yields a system of nonlinear ordinary differential equations (Eqs. [32], [40], [47]) and a system of nonlinear algebraic equations (Eqs. [26], [28], [34], [35], [42], [43]) which are shown below. The system of nonlinear algebraic equations is written in standard form such that the right hand side of each algebraic equation is zero. In [25], the system of F_i $1 \leq i \leq 78$ (such that $F_i = 0$) represent the system of nonlinear algebraic equations in [51],[52], [53],[54], and [55].

System of Ordinary Differential Equations

$$\frac{d\psi_1}{d\tau} = l_1 \psi_1 - \psi_1^2 \quad [48]$$

$$\frac{d\psi_2}{d\tau} = -\tilde{e}_1 \psi_1 - (\tilde{e}_2 \psi_2 + \tilde{e}_3 \psi_2^2) + \psi_3 (\tilde{e}_4 + \tilde{e}_5 \psi_2 + \tilde{e}_6 l_2 \psi_2^2) \quad [49]$$

$$\frac{d\psi_3}{d\tau} = -\tilde{e}_7 \psi_1 + \tilde{e}_8 \psi_2 + \tilde{e}_9 \psi_2^2 - \psi_3 (\tilde{e}_{10} + \tilde{e}_{11} \psi_2 + \tilde{e}_{12} l_3 \psi_2^2) \quad [50]$$

where

$$l_1 = \frac{\mu T}{\rho L^2}, l_2 = \frac{k^{Hb} [Hb] P_c T}{\alpha \cdot P_{O_2, 50\%}^{Hb}}, l_3 = \frac{k^{Hb} P_c^2 T}{P_{O_2, 50\%}^{Hb}},$$

and \tilde{e}_i $1 \leq i \leq 12$ are free parameters.

System of Nonlinear Algebraic Equations

$$A_{11} = 0 \text{ and } A_{12} = 0 \quad [51]$$

$$\begin{aligned}
 B_{2i} - B_{1i} &= 0 & \text{for } 1 \leq i \leq 3 \\
 B_{3i} &= 0 & \text{for } 2 \leq i \leq 3 \\
 B_{4i} - B_{1i} &= 0 & \text{for } 1 \leq i \leq 3
 \end{aligned} \quad [52]$$

$$\begin{aligned}
 C_{ijk} - \tilde{e}_{i-2} \cdot C_{1jk} &= 0 & \text{for } i = 3, 0 \leq j \leq 1, k = \{0, 3\} \\
 C_{ijk} + \tilde{e}_{i-2} \cdot C_{1jk} &= 0 & \text{for } 4 \leq i \leq 5, 0 \leq j \leq 1, k = \{0, 3\} \\
 C_{ijk} - \tilde{e}_{i-2} \cdot C_{1jk} &= 0 & \text{for } 6 \leq i \leq 8, 0 \leq j \leq 1, k = \{0, 3\}
 \end{aligned} \quad [53]$$

$$\begin{aligned}
 C_{213} = 0, C_{312} = 0, C_{302} = 0, C_{301} = 0, C_{412} = 0, \\
 C_{411} = 0, C_{402} = 0, C_{401} = 0, C_{003} = 0, C_{000} = 0.
 \end{aligned} \quad [54]$$

$$\begin{aligned}
 D_{ijk} - \tilde{e}_{i+4} \cdot D_{1jk} &= 0 & \text{for } 3 \leq i \leq 5, 0 \leq j \leq 1, k = \{0, 3\} \\
 D_{ijk} + \tilde{e}_{i+4} \cdot D_{1jk} &= 0 & \text{for } 6 \leq i \leq 8, 0 \leq j \leq 1, k = \{0, 3\}
 \end{aligned} \quad [55]$$

$$\begin{aligned}
 D_{213} = 0, D_{312} = 0, D_{302} = 0, D_{301} = 0, D_{412} = 0, \\
 D_{411} = 0, D_{402} = 0, D_{401} = 0, D_{003} = 0, D_{000} = 0.
 \end{aligned} \quad [56]$$