Finite Difference Solutions to the Lees-Dorodnitsyn Compressible Boundary Layer Equations



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Motivation

- Compressible, and specifically hypersonic, boundary layers can have complex velocity and temperature profiles
- Lees-Doroditsyn transformation results in decreased need for boundary layer scaling
- Boundary layer analysis is necessary to assess surface heating, skin friction, and external flow displacement effects
- Developing robust methods for modeling hypersonic boundary layers



TSL L-D Coordinate Transformation

$$\xi = \int_{0}^{x} \rho_{e} u_{e} \mu_{e} \, dx \qquad \eta = \frac{u_{e}}{\sqrt{2\xi}} \int_{0}^{y} \rho \, dy$$

$$(Cf'')' + ff'' = \frac{2\xi}{u_{e}} \left[(f')^{2} - \frac{\rho_{e}}{\rho} \right] \frac{du_{e}}{d\xi} + 2\xi \left(f' \frac{\partial f'}{\partial \xi} - \frac{\partial f}{\partial \xi} f'' \right)$$

$$\frac{\partial p}{\partial \eta} = 0$$

$$\left(\frac{C}{Pr} g' \right)' + fg' = 2\xi \left[f' \frac{\partial g}{\partial \xi} + \frac{f'g}{h_{e}} \frac{\partial h_{e}}{\partial \xi} - g' \frac{\partial f}{\partial \xi} + \frac{\rho_{e} u_{e}}{\rho h_{e}} f' \frac{du_{e}}{d\xi} \right] - C \frac{u_{e}^{2}}{h_{e}} (f'')^{2}$$

$$C = \frac{\rho\mu}{\rho_e\mu_e}$$
, $g = \frac{h}{h_e}$, and, $f' = \frac{u}{u_e}$

Lees-Dorodnitsyn Equations

- For 2D, Laminar, Compressible Boundary Layers
- Effectively Parabolic
- 5th Order system of coupled equations
- Introduce F, U, S, H, and Q to simplify into 5 1st order equations

•
$$F = f = \frac{\varphi}{\sqrt{2\xi}}$$
 (stream function relation)
• $U = \frac{\partial f}{\partial \eta} = \frac{u}{u_e}$ (velocity relation)
• $S = \frac{\partial U}{\partial \eta} = \frac{\partial^2 f}{\partial \eta^2}$ (shear stress relation)
• $H = g = \frac{h}{h_e} = \frac{T}{T_e} = \frac{\rho_e}{\rho}$ (enthalpy/temperature relation)
• $Q = \frac{\partial g}{\partial \eta} = \frac{\partial H}{\partial \eta}$ (heat transfer relation)

Flow Parameters

FLUID MODELS

- Viscosity Sutherland's Law • $\mu = \frac{\mu_{ref} \left(\frac{T}{T_{ref}}\right)^{\frac{3}{2}} (T_{ref} + S_{ref})}{T + S_{ref}}$
- Calorically Perfect Gas
 - $h = c_p T$
 - $\circ \gamma = 1.4$
- Constant Pr
 - 0.71 (Van Driest), 0.75, or 1
- Isothermal wall

SPECIFIED PARAMETERS

ξmax
 η_e
 T_e → h_e
 Ma_e(ξ) → u_e
 $\frac{h_w}{h_e}$ → h_w
 P_e = ρ_e

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$$\frac{h_w}{h_e}$$
 → h_w
 P_e = ρ_e

Boundary Conditions

- 5 Required for a 5th Order system:
- Wall:
 - $\circ F_{wall} = 0$
 - $U_{wall} = 0$
 - $\circ G_{wall} = G_{specified}$
- Edge:
 - $U_{edge} = 1$ • $G_{edge} = 1$
- Initial condition @ $\xi = 0$
 - 1D Flat plate similarity solution (calculated with Newton-Raphson method)

Residual Form

F' - U = 0U' - S = 0

$$(CS)' + FS + \frac{2\xi}{u_e} [H - (U)^2] \frac{du_e}{d\xi} + 2\xi \left(\frac{\partial F}{\partial \xi}S - U\frac{\partial U}{\partial \xi}\right) = 0$$

$$\left(\frac{C}{Pr}Q\right)' + FQ - 2\xi \left[U\frac{\partial H}{\partial\xi} + \frac{UH}{h_e}\frac{\partial h_e}{\partial\xi} - Q\frac{\partial F}{\partial\xi} + \frac{Hu_e}{h_e}U\frac{du_e}{d\xi}\right] + C\frac{u_e^2}{h_e}(S)^2 = 0$$
$$H' - Q = 0$$

F, *U*, *S*, *H*, *Q* are unknowns

Finite Difference Stencil

BOX SCHEME

- Space-march in ξ direction
- 2nd Order
- Centered Difference in both directions
- Similar to Crank Nicholson



3-POINT BACKWARD DIFFERENCE

- Space-march in ξ direction
- 2nd Order
- Centered Difference in η
- 3-pt Backward Difference in ξ



Box Stencil

$$X = X_{i+\frac{1}{2}}^{j+\frac{1}{2}} = \frac{1}{4} \left(X_i^j + X_i^{j+1} + X_{i+1}^j + X_{i+1}^{j+1} \right)$$

$$X' = {X'}_{i+\frac{1}{2}}^{j+\frac{1}{2}} = \frac{\left(X_{i+\frac{1}{2}}^{j+1} - X_{i+\frac{1}{2}}^{j}\right)}{\Delta\eta} = \frac{\left(\frac{1}{2}\left(X_{i+1}^{j+1} + X_{i}^{j+1}\right) - \frac{1}{2}\left(X_{i+1}^{j} + X_{i}^{j}\right)\right)}{\Delta\eta}$$

$$\frac{\partial X}{\partial \xi} = \frac{\partial X}{\partial \xi} \Big|_{i+\frac{1}{2}}^{j+\frac{1}{2}} = \frac{\left(X_{i+1}^{j+\frac{1}{2}} - X_{i}^{j+\frac{1}{2}}\right)}{\Delta \xi} = \frac{\left(\frac{1}{2}\left(X_{i+1}^{j+1} + X_{i+1}^{j}\right) - \frac{1}{2}\left(X_{i}^{j+1} + X_{i}^{j}\right)\right)}{\Delta \xi}$$





3pt - Backward Stencil

$$\begin{aligned} & \text{Spt - Backward Stencil} \\ & x = x_i^{j+\frac{1}{2}} = \frac{1}{2}(x_i^j + x_i^{j+1}) \\ & \text{Finite difference} \\ & \text{approximations:} \quad x' = x'_i^{j+\frac{1}{2}} = \frac{(X_i^{j+1} - X_i^j)}{\Delta \eta} \\ & \frac{\partial X}{\partial \xi} = \frac{\partial X}{\partial \xi} \Big|_i^{j+\frac{1}{2}} = \frac{(\frac{3}{2}X_i^{j+\frac{1}{2}} - 2X_{i-1}^{j+\frac{1}{2}} + \frac{1}{2}X_{i-2}^{j+\frac{1}{2}})}{2\Delta \xi} = \frac{\left(\frac{3}{2}(X_i^{j+1} + X_i^j) - 2(X_{i-1}^{j+1} + X_{i-1}^j) + \frac{1}{2}(X_{i-1}^{j+1} + X_{i-1}^j)\right)}{4\Delta \xi} \end{aligned}$$
Analytical Evaluations:
$$u_e = u_{e_i} = u_e(\xi_i) \\ h_e = h_{e_i} = h_e(\xi_i) \\ \frac{\partial u_e}{\partial \xi} = \frac{\partial u_e}{\partial \xi} \Big|_i = \frac{\partial h_e}{\partial \xi} (\xi_i) \\ \frac{\partial u_e}{\partial \xi} = \frac{\partial u_e}{\partial \xi} \Big|_i = \frac{\partial u_e}{\partial \xi} (\xi_i) \end{aligned}$$

Newton-Raphson Solver

- Necessary for Non-linear system
- 10 Diagonal Sparse Matrix
- Analytically calculated jacobian
- Utilized MATLAB built in matrix solver
- Quadratic convergence
- Convergence criteria based on magnitude of max residual



Newton vs. MATLAB FSolve

- Compared the Newton-Raphson Method against MATLAB's nonlinear solver
- Fsolve is a quasi-Newton method
- Assumed both solution methods had similar accuracies

	Newton-Raphson	MATLAB FSolve
Runtime	6-7 sec	116 sec
# of Iterations	3-4	2
Implementation Complexity	~650 lines	~125 lines
# of Function Evaluations	525	243

(Solved for 30 x 30 system)

Selected Results



Error Analysis

• L2 Error Analysis of $\delta^*(\xi)$, displacement thickness

$$\delta^* = \int_0^{\eta_e} \left(1 - \frac{\rho}{\rho_e} U \right) \frac{\partial y}{\partial \eta} \, d\eta = \int_0^{\eta_e} \left(1 - \frac{\rho}{\rho_e} U \right) \frac{H\sqrt{2\xi}}{u_e \rho_e} \, d\eta$$

- Compared to solution with fine resolution
- Roughly 2nd order convergence, as expected
- Comparable performance of Box Scheme and 3pt scheme



Conclusions

- Newton-Raphson (in this case) is significantly preferable to FSolve
- Stability is not a significant concern for the backwards difference scheme
- Oscillations are not present in the box scheme
- There is no clear advantage to using the box scheme instead of the 3-pt backwards difference
- Further sampling of the solution space may be necessary to determine overall behavior of solution methods

Thank you

Questions?

Appendix

Analytical Evaluations:

$$\begin{split} \xi &= \xi_{i+\frac{1}{2}} & \frac{\partial h_e}{\partial \xi} = \frac{\partial h_e}{\partial \xi} \bigg|_{i+\frac{1}{2}} = \frac{\partial h_e}{\partial \xi} \left(\xi_{i+\frac{1}{2}}\right) \\ h_e &= h_{e_{i+\frac{1}{2}}} = h_e \left(\xi_{i+\frac{1}{2}}\right) & \frac{\partial u_e}{\partial \xi} = \frac{\partial u_e}{\partial \xi} \bigg|_{i+\frac{1}{2}} = \frac{\partial u_e}{\partial \xi} \left(\xi_{i+\frac{1}{2}}\right) \end{split}$$

Analytical Evaluations:

$$\begin{split} \xi &= \xi_i \\ u_e &= u_{e_i} = u_e(\xi_i) \\ h_e &= h_{e_i} = h_e(\xi_i) \end{split}$$

$$\frac{\partial h_e}{\partial \xi} = \frac{\partial h_e}{\partial \xi} \bigg|_i = \frac{\partial h_e}{\partial \xi} (\xi_i)$$
$$\frac{\partial u_e}{\partial \xi} = \frac{\partial u_e}{\partial \xi} \bigg|_i = \frac{\partial u_e}{\partial \xi} (\xi_i)$$

Chapman Rubesin factor

• C = C(g)

•
$$C_{i+\frac{1}{2}}^{j+\frac{1}{2}} = C\left(g_{i+\frac{1}{2}}^{j+\frac{1}{2}}\right)$$

•
$$C_{i+\frac{1}{2}}^{j+\frac{1}{2}} = \frac{1}{4} \left[C(g_i^j) + C(g_{i+1}^j) + C(g_i^{j+1}) + C(g_{i+1}^{j+1}) \right]$$

Similarity Equations

2 - Coupled 2nd and 3rd Order Diff. Eqns

$$(Cf'')' + ff'' = 0$$
$$\left(\frac{C}{Pr}g'\right)' + fg' + C\frac{u_e^2}{h_e}(f'')^2 = 0$$



5 - Coupled 1st Order ODEs (Drela Notation)

F' = UU' = S(CS)' + FS = 0H' = Q $\left(\frac{C}{Pr}Q\right)' + FQ + C\frac{u_e^2}{h_e}(S)^2 = 0$

Where
$$\phi' = rac{\partial \phi}{\partial \eta}$$