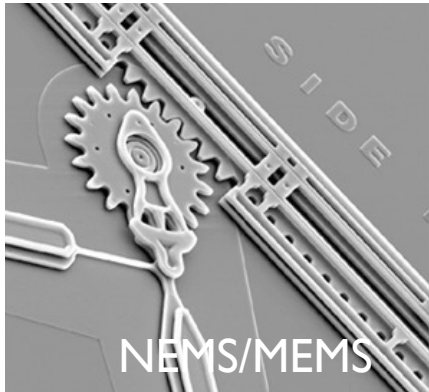


SIMULATING AND COMPARING ALTERNATIVE FORMULATIONS OF THE BOLTZMANN EQUATION

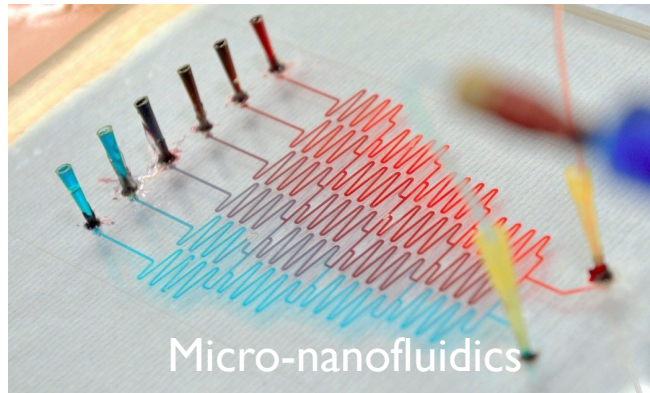
EVAN MASSARO



APPLICATIONS



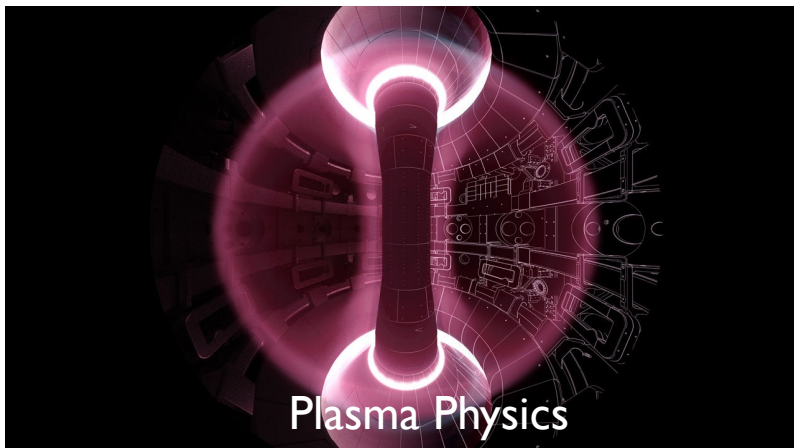
NEMS/MEMS



Micro-nanofluidics



Hypersonics



Plasma Physics



Galactic Dynamics

BACKGROUND

- The Navier-Stokes equations rely on the continuum hypothesis^[1], such that the fluid intensive property functions are well behaved (often C^∞)
- Knudsen number $Kn = \lambda/L$ is the ratio of the mean free path of particle over the characteristic length
 - measures the degree to which the state locally departs from equilibrium
 - $Kn = 0 \Rightarrow$ local equilibrium \Rightarrow Euler Equations
 - $Kn \lesssim 0.01 \Rightarrow$ sufficiently continuous \Rightarrow Navier-Stokes
- For non-continuum flows and any Kn , need a new governing equation derived from first principles: e.g. the Boltzmann Equation

[1] Pierre Lermusiaux, *Recitation Navier-Stokes*, MIT 2.29 (2021)

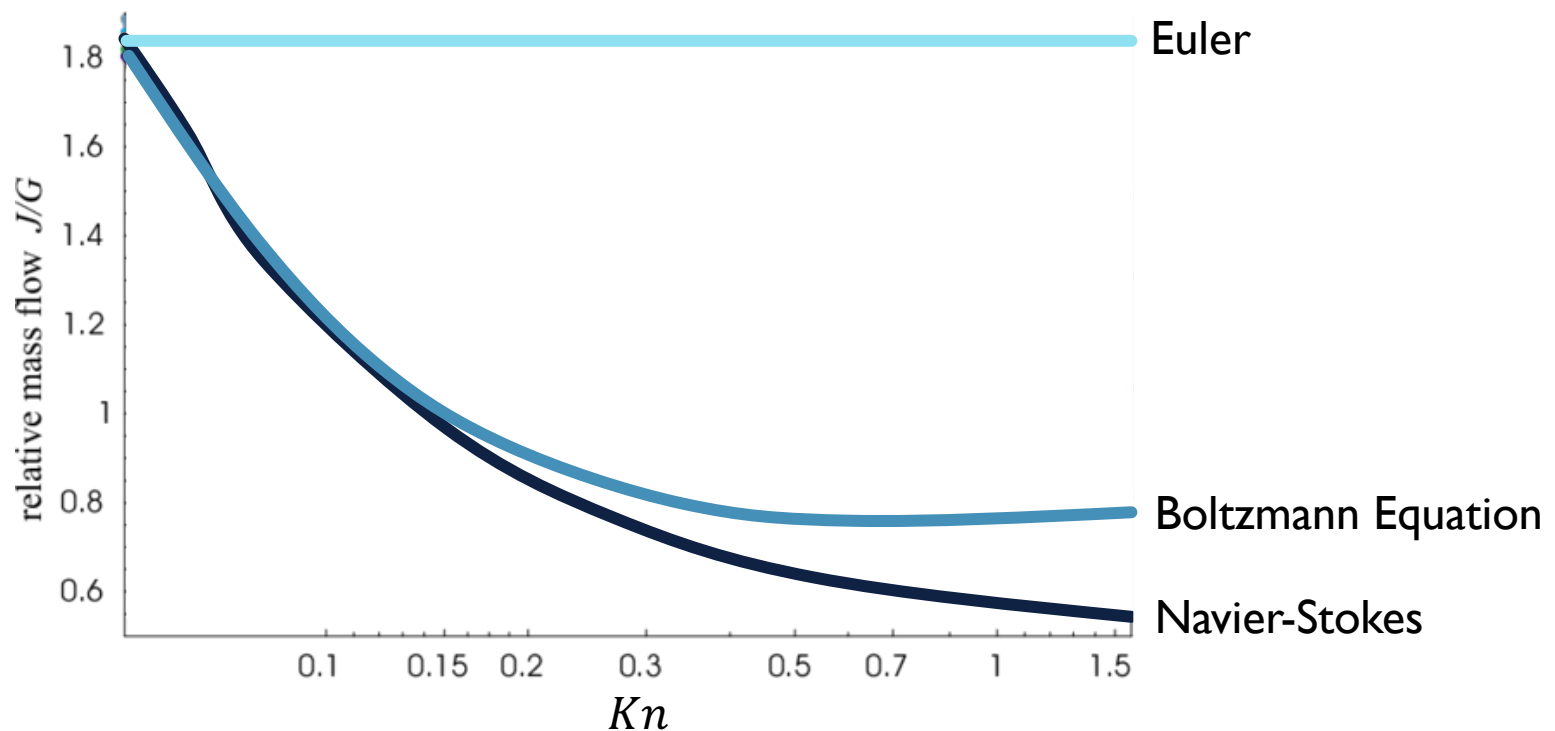
BOLTZMANN EQUATION

- The Boltzmann Equation is a (6+1)–dimensional PDE that describes the evolution of the single particle distribution function (pdf) $f(\mathbf{x}, \mathbf{c}, t)$ of a gas

$$\frac{\partial f}{\partial t} + \mathbf{c} \cdot \frac{\partial f}{\partial \mathbf{x}} + \frac{\mathbf{F}}{m} \cdot \frac{\partial f}{\partial \mathbf{c}} = \frac{1}{Kn} Q(f, f) \quad (\text{Eq I})$$

- In the ballistic regime $Kn \gg 1$, Eq I approximately only depends on LHS
- In the diffusive regime $Kn \ll 1$, Eq I is collision dominated and thus close to local equilibrium
 - Chapman-Enskog derives the Euler and Navier-Stokes as the zeroth and first order correction in a perturbative expansion $f = f^{(0)} + Kn f^{(1)} + \mathcal{O}(Kn^2)$
- For $0.1 \lesssim Kn \lesssim 10$ only the Boltzmann Equation (or Molecular Dynamics) is appropriate

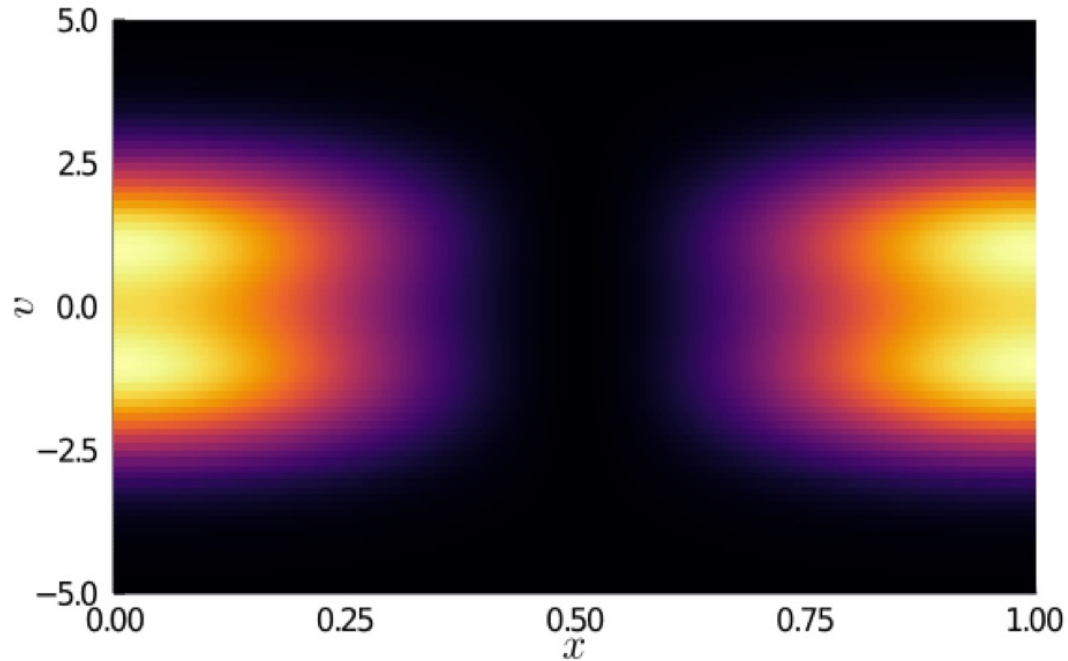
EXAMPLE I



[2] Adapted from Henning Struchtrup and Manuel Torrilhon, *Higher-Order Effects in Rarefied Channel Flows*, [10.1103/PhysRevE.78.046301](https://doi.org/10.1103/PhysRevE.78.046301) (2008)

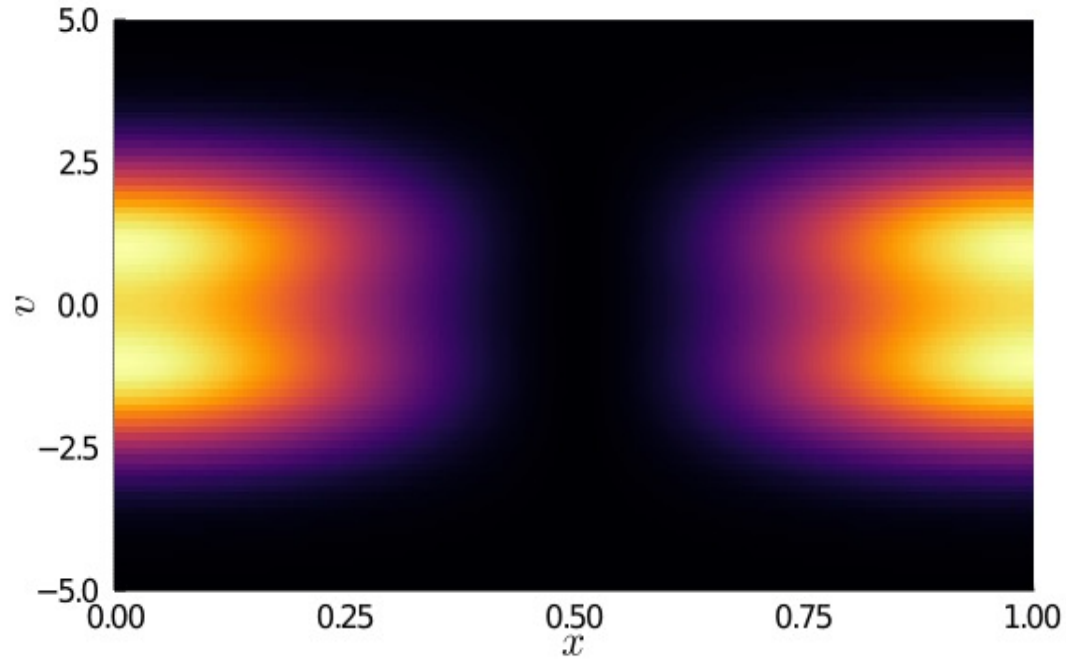
EXAMPLE II

$t = 0.0$



EXAMPLE II

$t = 0.0$



NUMERICAL METHODS I

- Two body collision integral

$$Q(f, f) = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} [f(\mathbf{v}') f(\mathbf{v}'_*) - f(\mathbf{v}) f(\mathbf{v}_*)] \mathcal{B}(\cos \theta, g) d\Omega d\mathbf{v}_*$$

- Evaluating this integral is difficult due to the high dimensionality and nonlinearity
- Fast spectral method is used to discretize space and the polar angles
- Flux reconstruction in space computed within the Lagrange polynomial basis

$$l_p = \prod_{q=0, q \neq p}^m \left(\frac{r - r_q}{r_p - r_q} \right) \quad \hat{f}^\delta = \sum_{p=0}^m \hat{f}_p^\delta l_p \quad \hat{F}^{\delta D} = \sum_{p=0}^m \hat{F}_p^{\delta D} l_p$$

[4] Tianbai Xiao, A Flux Reconstruction Kinetic Scheme for the Boltzmann Equation, arxiv.org/pdf/2103.10371.pdf (2021)

NUMERICAL METHODS II

- Spectral representation

$$\hat{f}^\delta(t, \mathbf{r}, \mathbf{v}) = \sum_{k=-N/2}^{N/2-1} \hat{f}_k^\delta(t, \mathbf{r}) \exp(i\xi_k \cdot \mathbf{v})$$

$$\hat{Q}_k^\delta = \sum_{l, m=-N/2, (l+m=k)}^{N/2-1} \hat{f}_l^\delta \hat{f}_m^\delta [\beta(l, m) - \beta(m, m)]$$

$$\beta(l, m) = 4 \int \left[\int_0^R \rho \cos(\rho \xi_l \cdot \mathbf{e}) d\rho \right] \left[\int \delta(\mathbf{e} \cdot \mathbf{e}') \int_0^R \rho' \Theta \cos(\rho' \xi_m \cdot \mathbf{e}') d\rho' d\mathbf{e}' \right] d\mathbf{e}$$

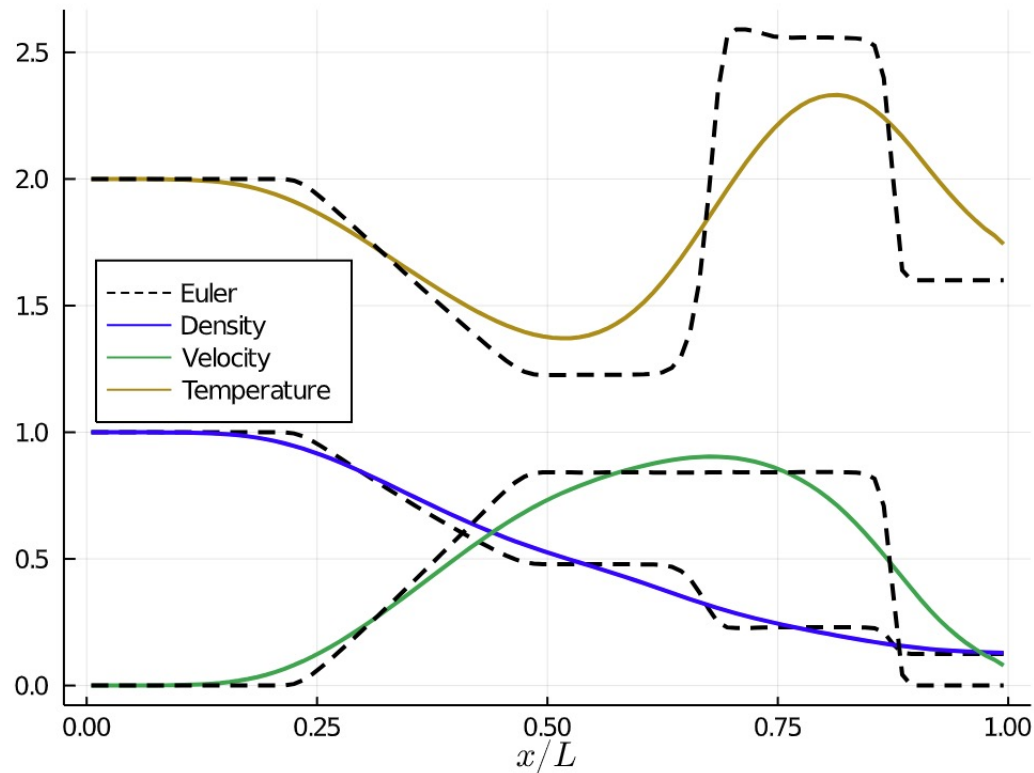
$$\frac{\partial \hat{f}_{i,j}^\delta}{\partial t} = -\nabla_{\mathbf{r}} \cdot \hat{\mathbf{F}}_{i,j}^\delta + \hat{Q}_{i,j}^\delta = \hat{\mathcal{C}}_{i,j}^\delta$$

- Use singly diagonally implicit Runge–Kutta (SDIRK) for time integration

[4] Tianbai Xiao, A Flux Reconstruction Kinetic Scheme for the Boltzmann Equation, arxiv.org/pdf/2103.10371.pdf (2021)

SOD SHOCK TUBE

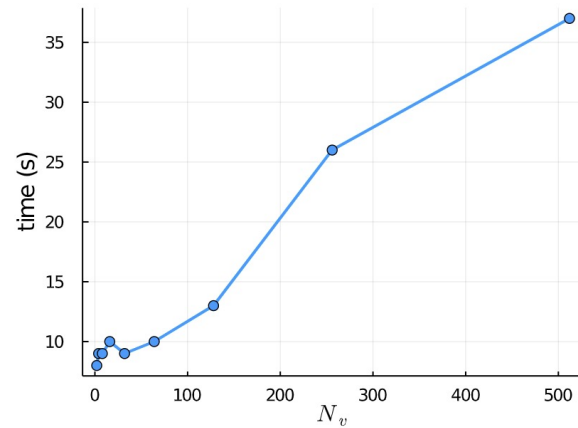
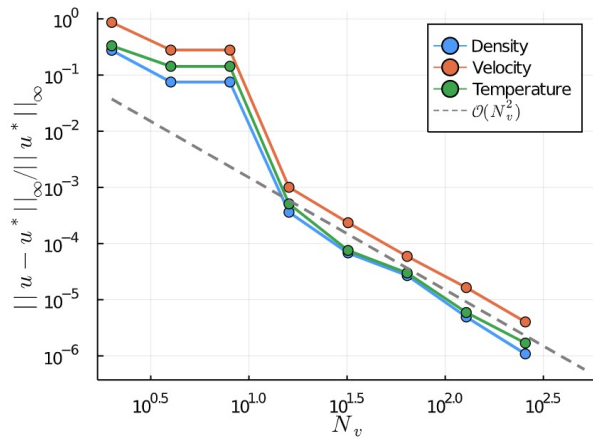
$$Kn = 0.1$$



[4] Code based in Kinetic.jl, <https://xiaotianbai.com/Kinetic.jl/dev> (2021)

ERROR ANALYSIS

order = 2



order = 3

