

Novel Features in eRHIC

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e-RHIC
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Beam Dynamics
(V. Ptitsyn), GL
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(E. Pozdeyev)
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(N. Tsoupas)
SRF
(I. Ben Zvi), GL
(A. Burrill)
(H. Hahn)
(D. Naik)
(L. Hammons)
Polarization
(M. Bai), GL
(H. Huang)
(J. Kewish)
(A. Luccio)
(A. Zelenski)
IRs
(C. Montag), GL
(A. Drees)
(J. Beebe-Wang)

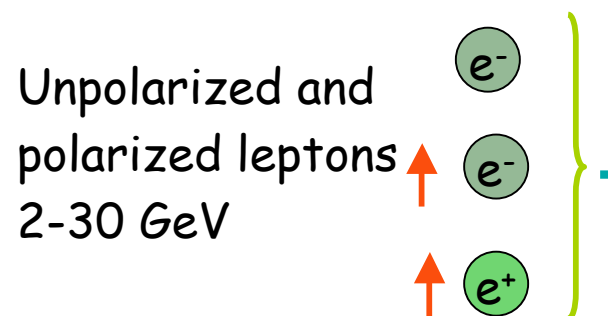
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Novel features

- *Coherent electron cooling - the key for many novel features in eRHIC*
- *Choosing the focus: ERL for electrons*
 - Advantages and challenges of ERL driver
 - spin transparency
 - R&D items for ERL-based eRHIC
- *eRHIC is the future of RHIC: eRHIC staging*
 - Energy challenge
 - $20 \text{ GeV } e^- \times 325 \text{ GeV } p^+$ and $30 \text{ GeV } e^- \times 125 \text{ GeV/n heavy ions}$
 - Loss on synchrotron radiation
 - Polarized beam current
- *Luminosity challenge:*
 - Can eRHIC deliver $10^{35} \text{ cm}^{-2} \text{ sec}^{-1}$ luminosity?
 - High rep-rate, crab cavities, coating RHIC arc vacuum chambers and more
- *Other novelties and oldies*
 - Low (350 MHz) RF frequency, no 3rd harmonic, higher real estate gradient
 - Small magnets for re-circulating passes, resistive-wall losses
 - e-lens or fast a quads for matching ERL beam
 - compact and flexible separators and combiners
 - Possibility of eRHIC II up-grade

eRHIC Scope - QCD Factory

Electron accelerator

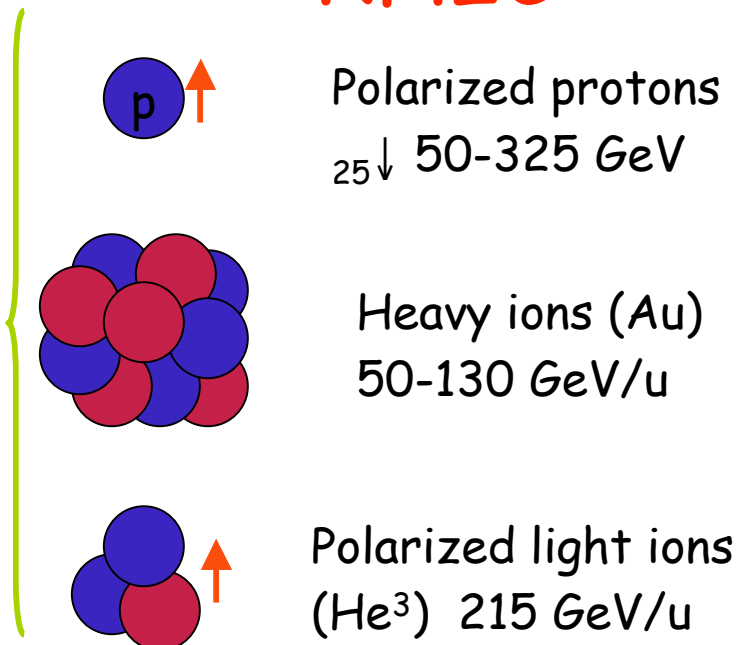


70% beam polarization goal
Positrons at low intensities

Center mass energy range: 15-200 GeV

New requirements: eA program for eRHIC needs as high as possible energies of electron beams even with a trade-off for the luminosity. 20 GeV is absolutely essential and 30 GeV is strongly desirable.

RHIC



ERL spin transparency at all energies 4

Bargman, Mitchel, Telegdi equation

$$\frac{d\hat{s}}{dt} = \frac{e}{mc} \hat{s} \times \left[\left(\frac{g}{2} - 1 + \frac{1}{\gamma} \right) \vec{B} - \frac{\gamma}{\gamma+1} \left(\frac{g}{2} - 1 \right) \hat{\beta} (\hat{\beta} \cdot \vec{B}) - \left(\frac{g}{2} - \frac{\gamma}{\gamma+1} \right) [\vec{\beta} \times \vec{E}] \right]$$

$$a = g/2 - 1 = 1.1596521884 \cdot 10^{-3}$$

$$\hat{\mu} = \frac{g}{2} \frac{e}{m_o} \hat{s} = (1+a) \frac{e}{m_o} \hat{s}; \quad \nu_{spin} = a \cdot \gamma = \frac{E_e}{0.44065 [GeV]}$$

$$\Delta\varphi = a \cdot \gamma\theta$$

Total angle

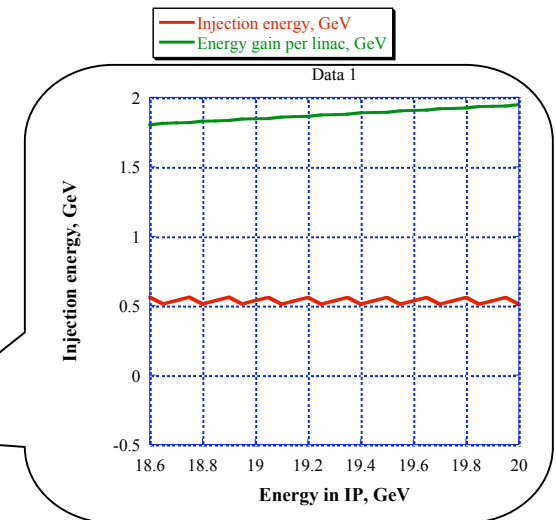
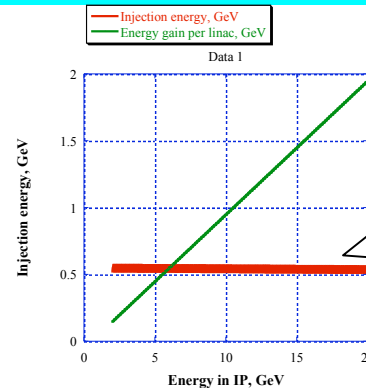
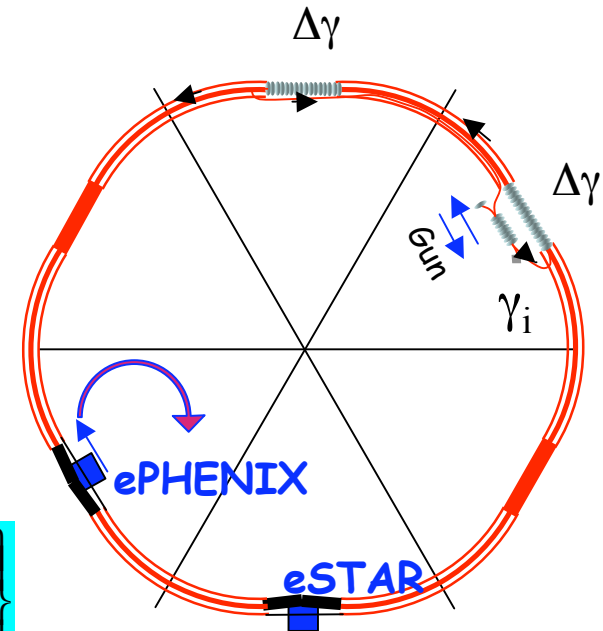
$$\varphi = 2\pi a \cdot \left((n - 1/2) \gamma_i + \{2n(n - 2) - 1/6\} \Delta\gamma \right) + \varphi_i$$

Has solution
for all
energies!

$$\begin{cases} \Delta\gamma = (\gamma_f - \gamma_i) / 2n \\ 2\pi a \cdot \left((n - 1/2) \gamma_i + \{n(n - 2) - 1/3\} \Delta\gamma \right) + \varphi_i = \theta + N\pi \end{cases}$$

$$E_i = \frac{0.44065 [GeV]}{n + 1 + 1/3n} \bmod \left(\varphi_f - \varphi_i - \left(n - 2 - \frac{1}{3n} \right) \frac{E_f}{0.44065 [GeV]}, \pi \right)$$

$$\delta E_{i \max} = \pm 37 \text{ MeV} \quad \forall n = 5$$



Main advantages of ERL + cooling

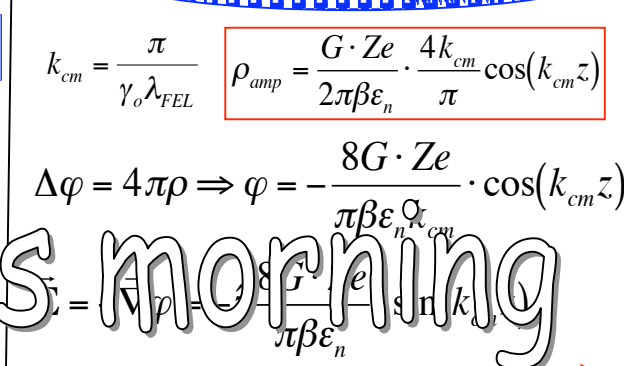
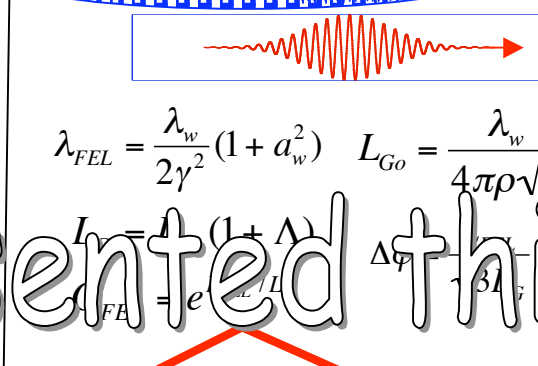
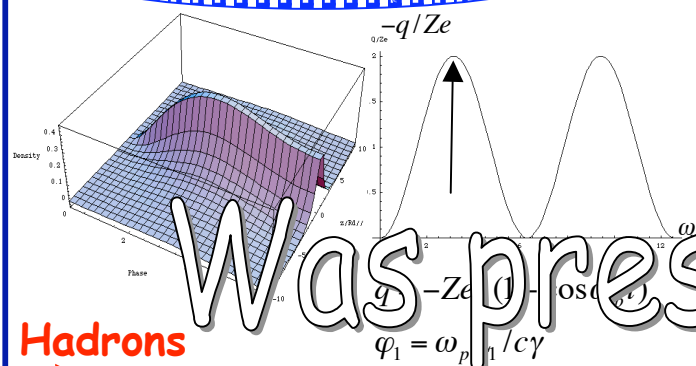
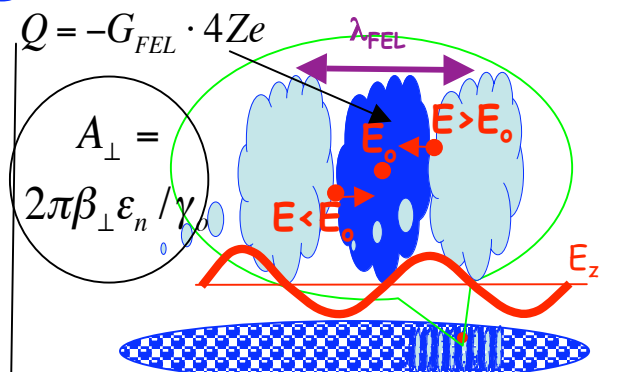
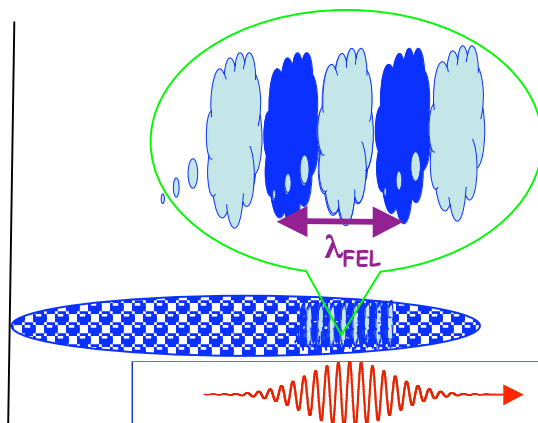
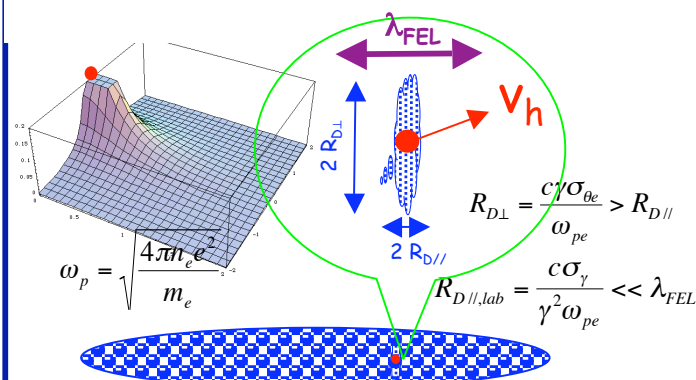
$$L = \gamma_p \frac{f_{col} N_p}{\beta_p^* r_p} \xi_p \quad \xi_p = \frac{r_p}{4\pi} \cdot \frac{N_e}{\varepsilon_{p \text{ norm}}};$$

$$\frac{N_e}{\varepsilon_{p \text{ norm}}} = \text{const} \Rightarrow \xi_p = \text{const}; \quad L = \text{const}$$

$$N_e \propto \varepsilon_{p \text{ norm}} \Rightarrow I_e \propto \varepsilon_{p \text{ norm}} \Rightarrow P_{SR} \propto \varepsilon_{p \text{ norm}}!$$

- Main point is very simple: if one cools the emittance of a hadron beam in electron-hadron collider, the intensity of the electron beam can be reduced proportionally without any loss in luminosity or increase in the beam-beam parameter for hadrons
- Hadron beam size is reduced in the IR triplets - hence it opens possibility of further β^* squeeze and increase in luminosity
- Electron beam current goes down -> relaxed gun!, losses for synchrotron radiation going down, X-ray background in the detectors goes down....

Coherent electron cooling



Was presented this morning

Hadrons

Electrons

$$Q_{\lambda_{FEL}} \approx \int_0^{\lambda_{FEL}} \rho(z) \cos(k_{FEL} z) dz$$

$$Q_{\lambda_{FEL}} (\max) \approx -2Ze; \rho_k = -Ze \frac{4k}{\pi A_{\perp}}$$

Modulator: region 1
a quarter to a half
of plasma oscillation

Longitudinal dispersion for
hadrons

$$\Delta t = -D \cdot \frac{\gamma - \gamma_o}{\gamma_o}; D = D_{free} + D_{chicane};$$

$$D_{free} = \frac{L}{\gamma^2}; D_{chicane} = l_{chicane} \cdot \theta^2$$

Amplifier of the e-beam modulation
via FEL with gain $G_{FEL} \sim 10^2 - 10^3$

Most versatile option

$$\Delta E_i = -\frac{8G \cdot Z^2 e^2}{\pi \beta \epsilon_n} L_2 \cdot \sin\left(k_{FEL} D \frac{E - E_o}{E_o}\right) \cdot \left(\frac{\sin \varphi_2}{\varphi_{p2}}\right) \cdot \left(\frac{\sin \varphi_1}{2}\right)^2$$

Kicker: region 2,
less than a quarter of
plasma oscillation

Coherent e-Cooling for eRHIC

(protons are the main challenge)

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Main Parameters	CeC	
Modulator Length	15	m
Kicker length	5	m
Peak current, e	100.0	A
Amplification	200.00	
Wavelength	500	nm
λ_w	5	cm
FEL bandwidth	0.1	

Cooling time		
Emittance, Full bunch	0.086	hrs
Ampl, Full bunch	0.171	hrs
Local	50.23	sec
Length of the system	32.49	m
FEL length	12.49	m
FEL gain length	0.99	m

Hadrons		
Z	1	
A	1	
Energy per nucleon	325	GeV
Energy per nucleon	3.250000E+11	eV
γ	346.38	=
N, part/bunch	2.00E+11	
Charge	32.04	nC
Bunch length	0.433	nsec
Bunch lengt, RMS	0.130	m
Peak current	29.50	A
Emittance, norm	2	mm mrad
Emittance, m rad	5.77398E-09	
$\sigma_{E/E}$	4.00E-04	

Electrons		
	1	
Energy	0.177	GeV
Energy	1.770E+08	eV
γ	346.38	
N, part/bunch	3.12E+10	
Charge	5.0	nC
Bunch length	0.050	nsec
Bunch lengt, full	0.015	m
Peak current	100.0	A
Emittance, norm, RMS	5	mm mrad
Emittance, RMS	1.443E-08	m rad
$\sigma_{E/E}$	2.26E-04	
σ_E	4.00E+04	eV
Long emittance	2.000E-06	eV sec

Stationary state: IBS vs. CeC

$$\frac{\sigma_\epsilon^2}{\tau_{IBS//}} = \frac{Nr_c^2 c}{2^5 \pi \gamma^3 \epsilon_x^{3/2} \sigma_s} \left\langle \frac{f(\chi_m)}{\beta_y v} \right\rangle; \quad \frac{\epsilon_x}{\tau_{IBS\perp}} = \frac{Nr_c^2 c}{2^5 \pi \gamma^3 \epsilon_x^{3/2} \sigma_s} \left\langle \frac{H}{\beta_y^{1/2}} f(\chi_m) \right\rangle; \kappa = 1$$

$$f(\chi_m) = \int_{\chi_m}^{\infty} \frac{d\chi}{\chi} \ln \left(\frac{\chi}{\chi_m} \right) e^{-\chi}; \quad \chi_m = \frac{r_c m^2 c^4}{b_{\max} \sigma_E^2}; b_{\max} \cong n^{-1/3}; \quad r_c = \frac{e^2}{mc^2}; \quad (e^- \rightarrow Ze; m^- \rightarrow Am)$$

J. LeDuff, "Single and Multiple Touschek effects",
 Proceedings of CERN Accelerator School,
 Rhodes, Greece, 20 September - 1 October, 1993,
 Editor: S. Turner, CERN 95-06, 22 November 1995,
 Vol. II, p. 57

$$X = \frac{\epsilon_x}{\epsilon_{xo}}; S = \left(\frac{\sigma_s}{\sigma_{so}} \right)^2 = \left(\frac{\sigma_E}{\sigma_{sE}} \right)^2;$$

$$\frac{dX}{dt} = \frac{1}{\tau_{IBS\perp}} \frac{1}{X^{3/2} S^{1/2}} - \frac{\xi_{\perp}}{\tau_{CeC}} \frac{1}{S};$$

$$\frac{dS}{dt} = \frac{1}{\tau_{IBS//}} \frac{1}{X^{3/2} Y} - \frac{1-2\xi_{\perp}}{\tau_{CeC}} \frac{1}{X};$$

$$\epsilon_{xn0} = 2 \mu m; \quad \sigma_{s0} = 13 \text{ cm}; \quad \sigma_{\delta 0} = 4 \cdot 10^{-4}$$

$$\tau_{IBS\perp} = 4.6 \text{ hrs}; \quad \tau_{IBS//} = 1.6 \text{ hrs};$$

*IBS in RHIC for
 eRHIC, 250 GeV, $N_p = 2 \cdot 10^{11}$
 Beta-cool, © A. Fedotov*

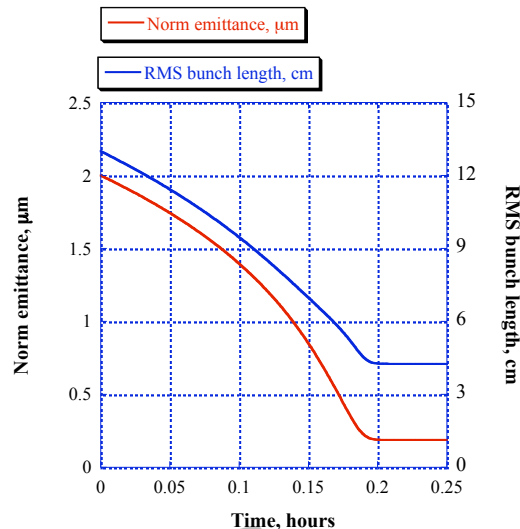
Stationary solution:

$$X = \frac{\tau_{CeC}}{\sqrt{\tau_{IBS//} \tau_{IBS\perp}}} \frac{1}{\sqrt{\xi_{\perp} (1-2\xi_{\perp})}}; \quad S = \frac{\tau_{CeC}}{\tau_{IBS//}} \cdot \sqrt{\frac{\tau_{IBS\perp}}{\tau_{IBS//}}} \cdot \sqrt{\frac{\xi_{\perp}}{(1-2\xi_{\perp})^3}}$$

$$\epsilon_{xn} = 0.2 \mu m; \quad \sigma_s = 4.9 \text{ cm}$$

This allows

- a) keep the luminosity as it is
- b) reduce polarized beam current down to 25 mA (5 mA for e-I)
- c) increase electron beam energy to 20 GeV (30 GeV for e-I)
- d) increase luminosity by reducing β^* from 25 cm down to 5 cm



Staging of eRHIC: Energy Reach and Luminosity

- **MEIC: Medium Energy Electron-Ion Collider**
 - Located at IP2 (with a modest detector)
 - 2 GeV e^- x 250 GeV p (45 GeV c.m.), $L \sim 10^{32} \text{ cm}^{-2} \text{ sec}^{-1}$
- **eRHIC - Full energy, nominal luminosity , inside RHIC tunnel**
 - Polarized 20 GeV e^- x 325 GeV p (160 GeV c.m.), $L \sim 4 \cdot 10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$
 - 30 GeV e^- x 120 GeV/n Au (120 GeV c.m.), $L \sim 10^{31} \text{ cm}^{-2} \text{ sec}^{-1}$
 - 20 GeV e^- x 120 GeV/n Au (120 GeV c.m.), $L \sim 5 \cdot 10^{31} \text{ cm}^{-2} \text{ sec}^{-1}$
- **eRHIC - High luminosity at reduced energy, inside RHIC tunnel**
 - Polarized 10 GeV e^- x 325 GeV p, $L \sim 10^{35} \text{ cm}^{-2} \text{ sec}^{-1}$
 - Smaller improvements (3-4 fold) in e-Ion collisions

More detail during discussion on staged eRHIC,
Today 5:25 p.m.

Staging of eRHIC:

Cost, Re-use, Beams and Energetics

- **MEIC: Medium Energy Electron-Ion Collider**
 - **Cost estimate - \$150M (in 2007 \$)**
 - 90% of ERL hardware will be use in the phase I (and will reduce cost of eRHIC)
 - Possible use of the detector components for eRHIC detectors
- **eRHIC - phase I**
 - **Based on present RHIC beam intensities**
 - With coherent electron cooling requirements on the electron beam current is 25 mA
 - 20 GeV, 25 mA electron beam losses **1.92 MW total for synchrotron radiation***.
 - 30 GeV, 5 mA electron beam loses 1.98 MW for synchrotron radiation
 - Power density is 1 kW/meter and is well within B-factory limits (8 kW/m)
- **eRHIC - phase II**
 - **Requires crab cavities, new injections, Cu-coating of RHIC vacuum chambers, new level of intensities in RHIC**
 - Polarized electron source current of 400 mA
 - 10 GeV, 400 mA electron beam losses 1.96 MW total for synchrotron radiation, power density is 1 kW/meter

**Compare it with 15 MW power loos for 10 GeV electrons in ELIC!*

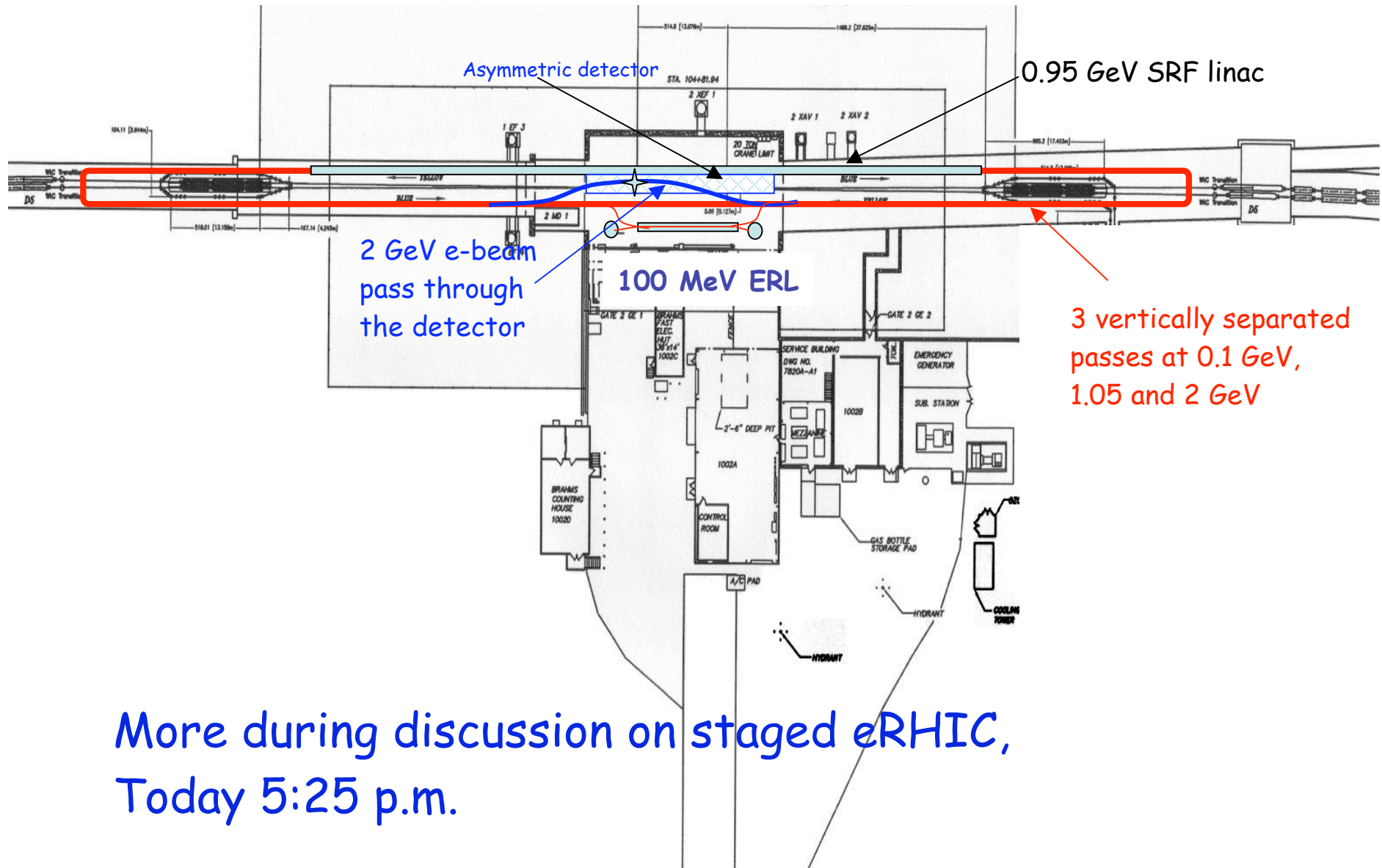
More detail during discussion on staged eRHIC, today 5:25 p.m.

Possible future up-grade - eRHIC II

c.m. Energy of HERA with 100x Luminosity

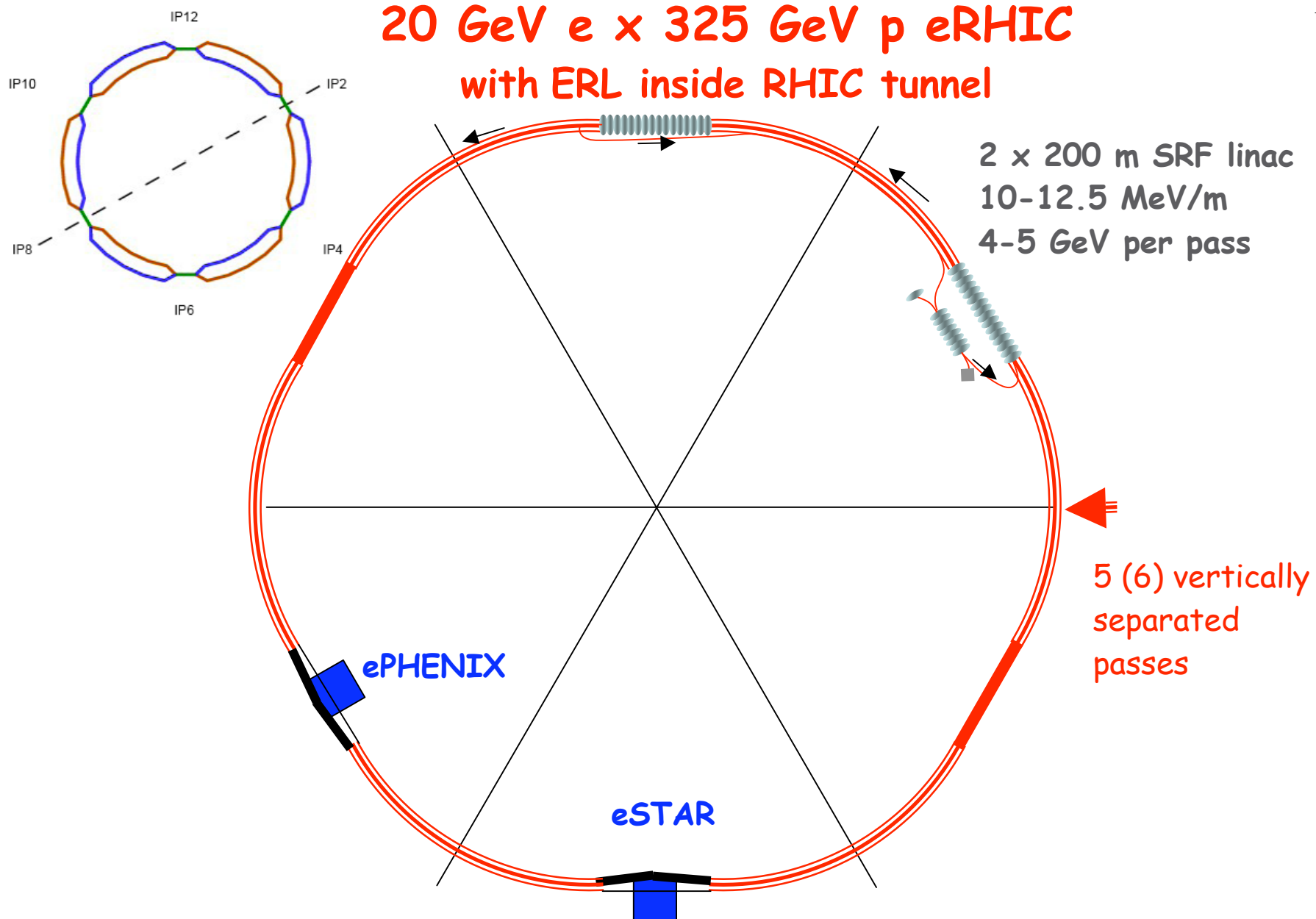
- **eRHIC II:** replacing RHIC-ring magnets by 8 T
 - proton energy in RHIC to ~ 800 GeV
 - will require more snakes for polarized proton operation
 - heavy ions with ~ 300 GeV/n
- **eRHIC II - Full energy, nominal luminosity**
 - inside RHIC tunnel
 - Polarized 20 GeV e^- x 800 GeV p (~ 300 GeV c.m.), $L \sim 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$
 - 30 GeV e x 300 GeV/n Au (~ 200 GeV c.m.), $L \sim 10^{32} \text{ cm}^{-2} \text{ sec}^{-1}$

MEIC with 2 GeV ERL @ IP2

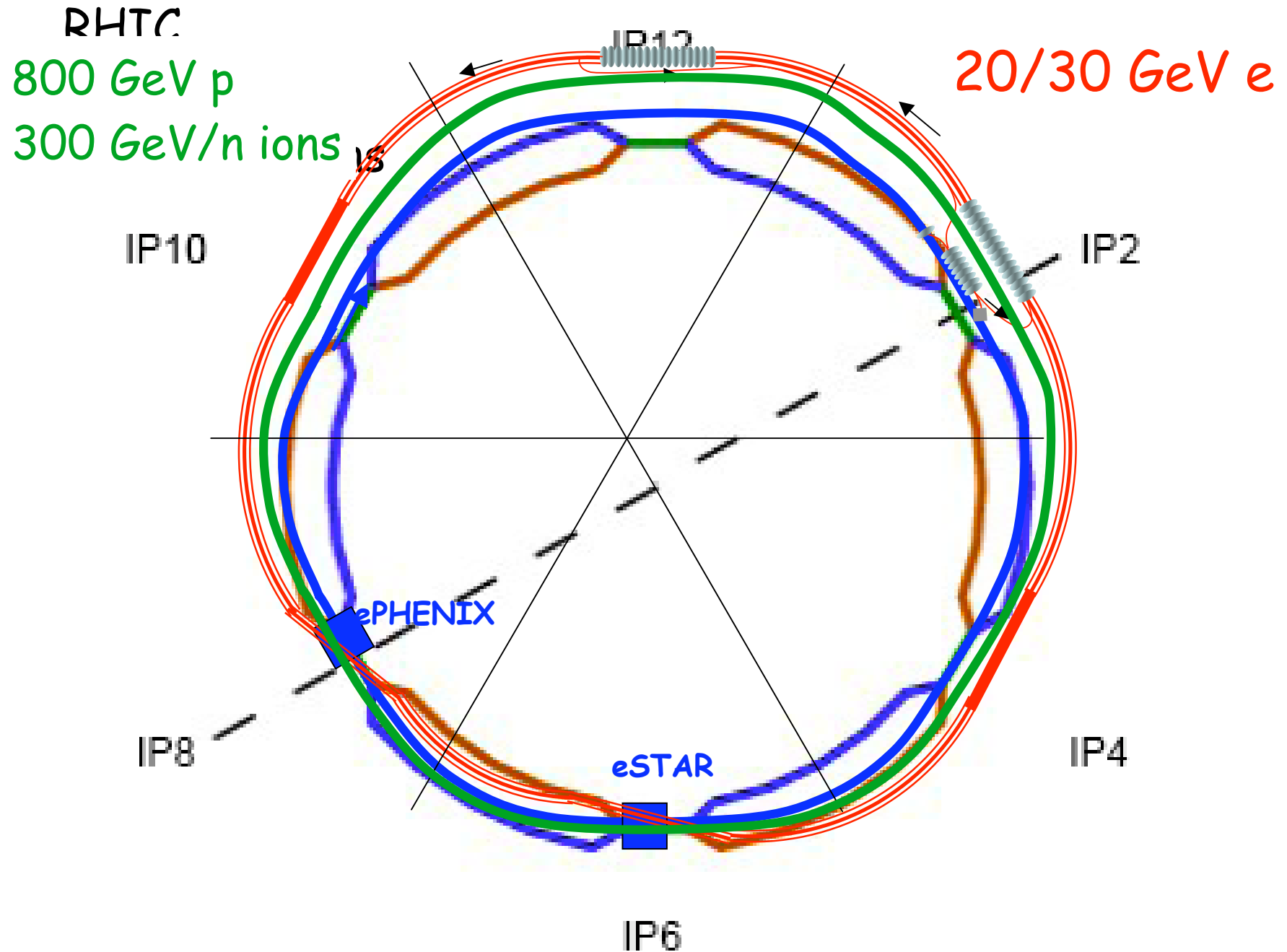


More during discussion on staged eRHIC,
Today 5:25 p.m.

20 GeV e x 325 GeV p eRHIC with ERL inside RHIC tunnel



Staging of eRHIC with ERL inside RHIC tunnel

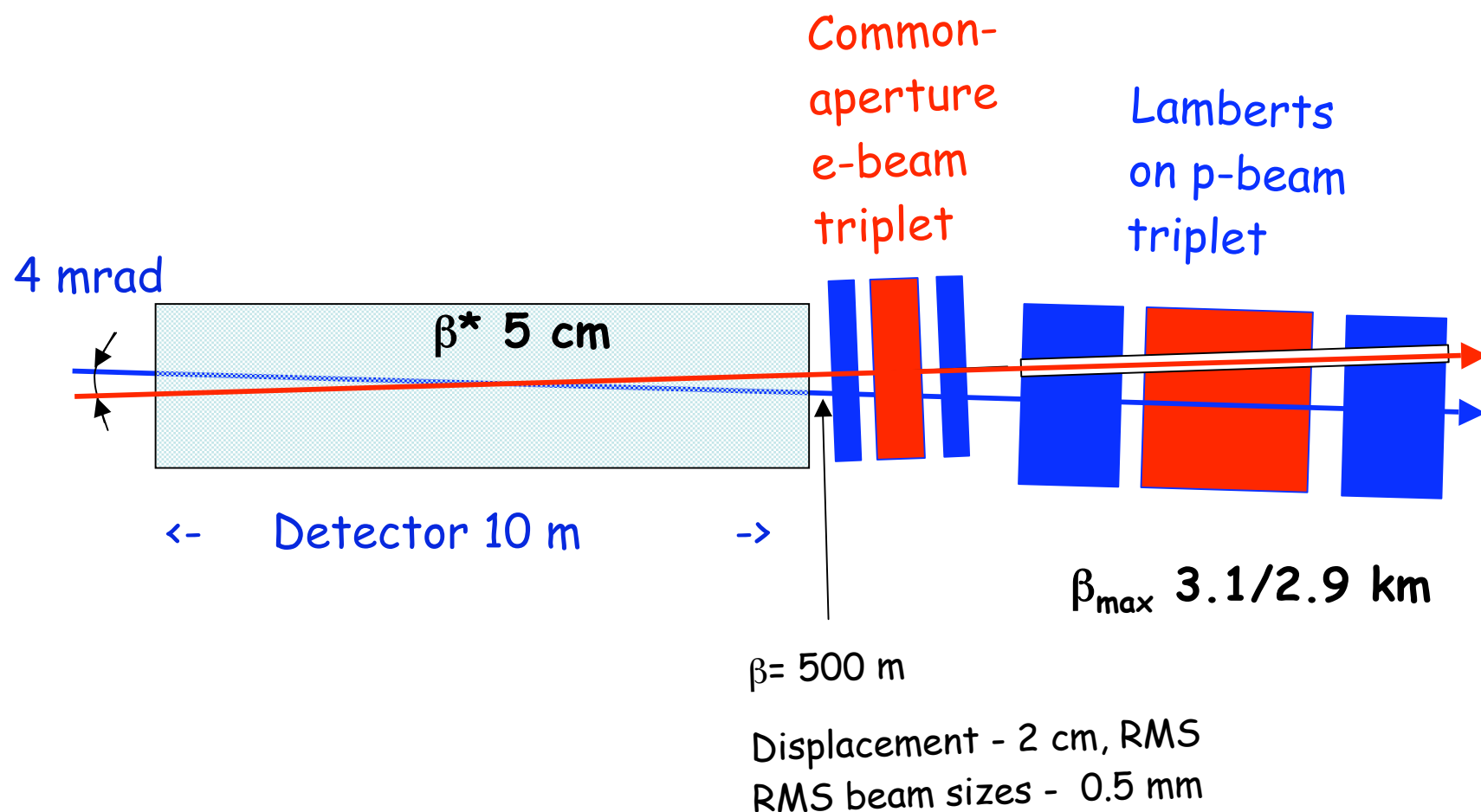


Borrowing ideas: B-factory, KEK, JLab's ELIC and LHeC ¹⁵

- We are considering possibility of IP's with crossing angle and crab cavities
 - **MAIN REASON** - There will be no synchrotron radiation background problem in the detector and we can afford 10 m + of element-free IR for the detectors
 - β^* 5 cm for both protons and electrons
 - Hadron bunch length < 5 cm
 - Emittance ~ 0.8 nm (normalized ~0.2 μm for protons and 30 μm for electrons)
 - RMS angular spread 0.1 mrad
 - Crossing angle ~ 2 mrad (per beam - 40 RMS sizes of hadron beam) - angle mostly required for separating beams in triplets
- We are considering possibility of using rather small-aperture Lambertson-quads for such a scheme
 - At 5 m, electron and proton beam will be separated ~ 2 cm and beam-sizes will be only 0.5 mm RMS

Straw-man IR lay out

no bending of electron trajectory in IR -> no X-rays

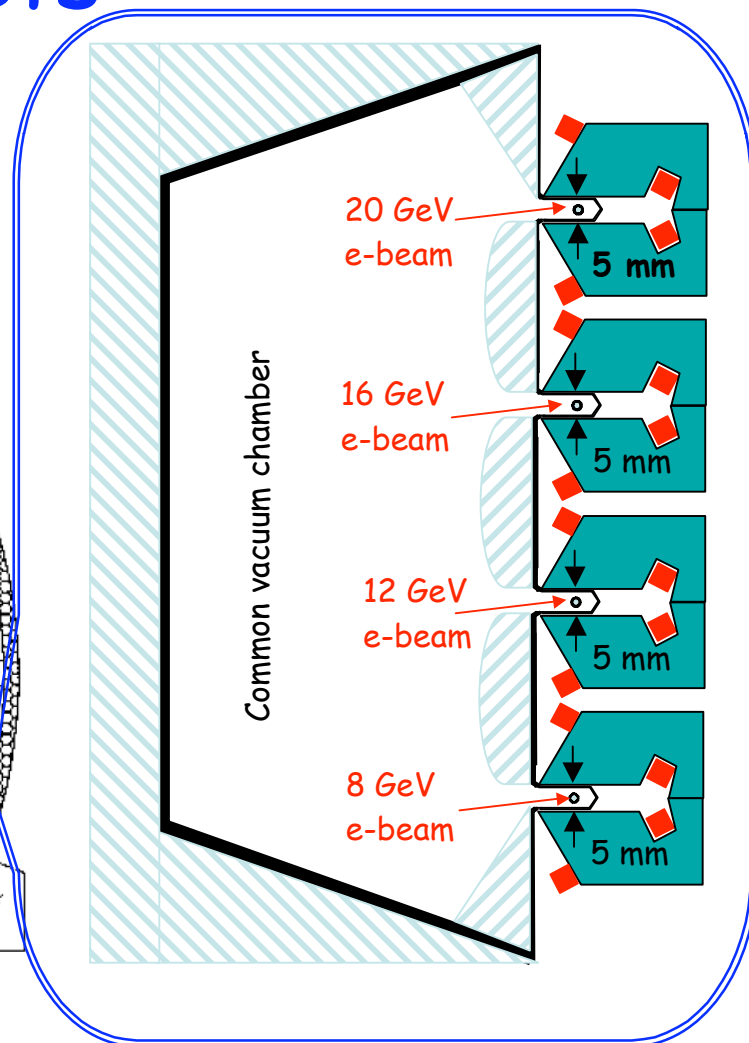
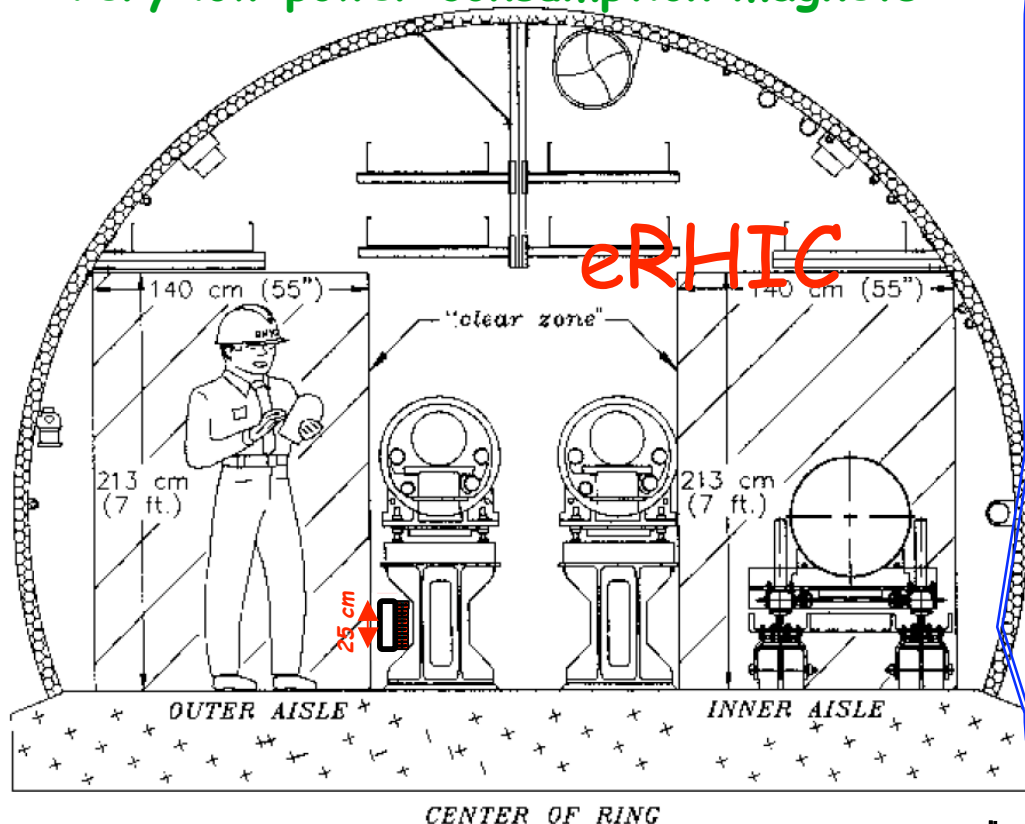


Can eRHIC deliver luminosity $\sim 10^{35} \text{ cm}^{-2} \text{ sec}^{-1}$?

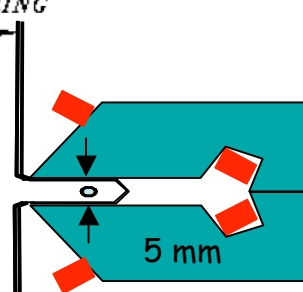
- The answer is Yes. With coherent electron cooling eRHIC it can reach luminosity of $0.2 \cdot 10^{35} \text{ cm}^{-2} \text{ sec}^{-1}$ with $\beta^* = 5 \text{ cm}$ with presently designed proton intensities
 - The question is what will be compromises?
 - Another question is what additional modification of RHIC it will require → →
- **Compromises**
 - Lower electron beam energy ($\sim 10 \text{ GeV}$) to keep power bill (for loss of synchrotron radiation)
 - 5-10 times higher collision rate ($\sim 100 \text{ MHz}$)
- **Additional developments**
 - New injection system supporting higher rep-rate
 - Coating RHIC arc's vacuum chamber
 - Crossing angle and crab cavity

eRHIC loop magnets

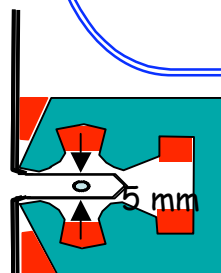
- Small gap provides for low current
- Very low power consumption magnets



C- Dipole



C-Quad

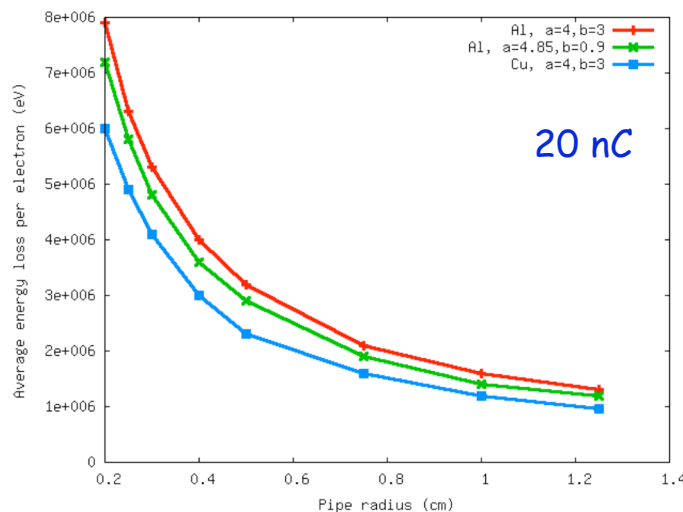
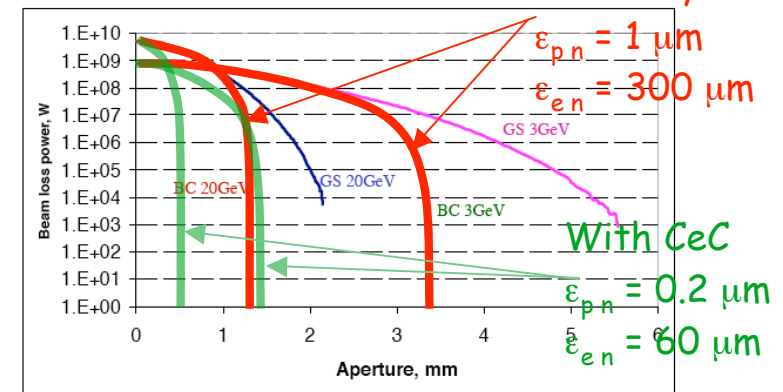


Limitations on the aperture for electron beam

- Magnetic field quality
- Alignment accuracy
- e-beam loss

resistive-wall induced energy spread and energy loss

Power loss and magnet aperture
(based on $\beta=50\text{m}$)

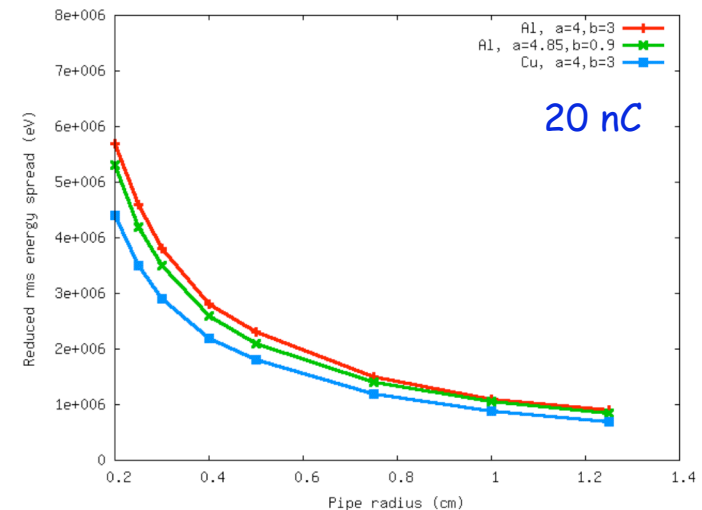


©E.Pozdeyev

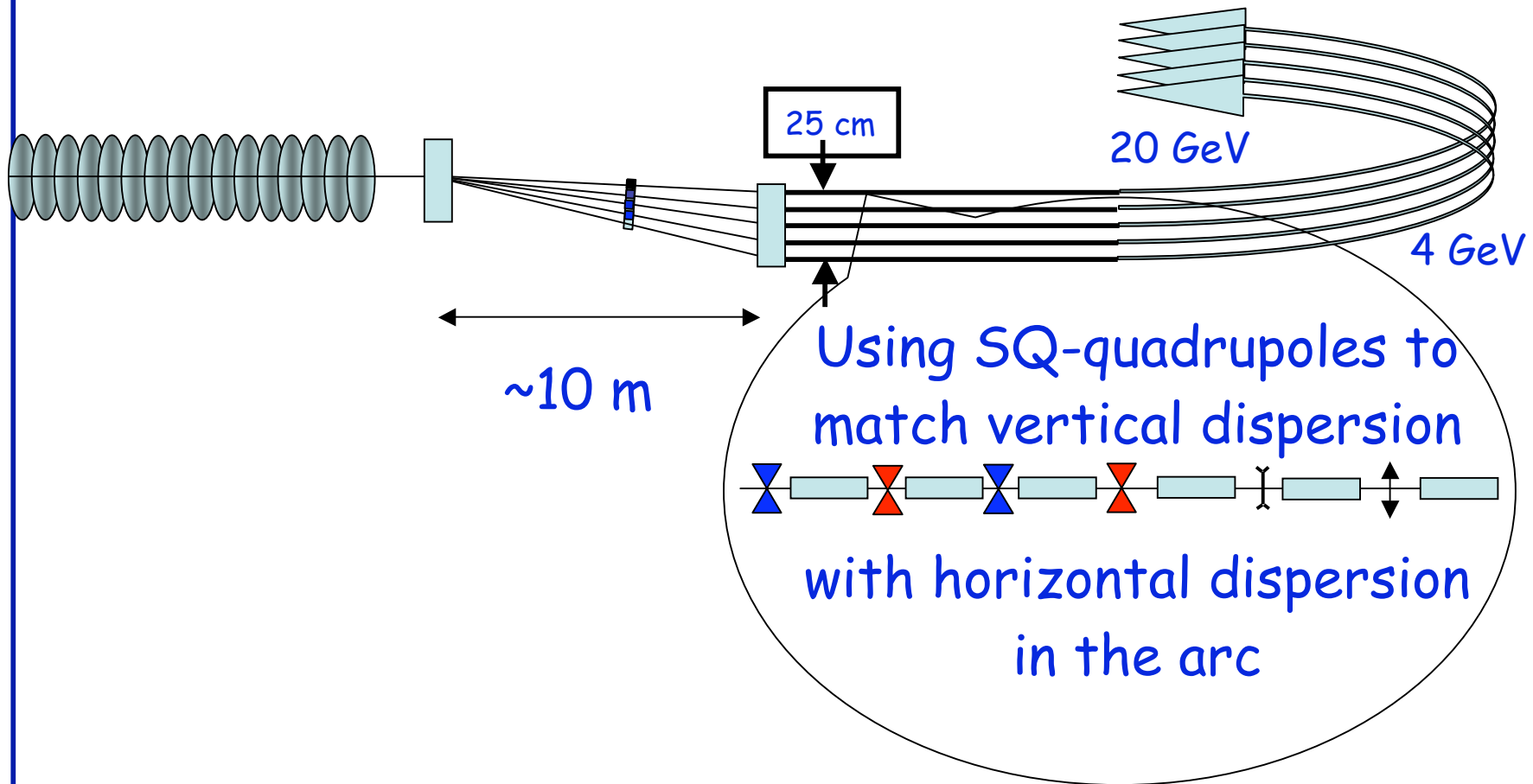
$N_e = 20 \text{ nC/bunch/e}$
Loss $\sim 1\text{MW}$ with 5 mm aperture

With CeC -

$N_e \rightarrow 2 \text{ nC/bunch/e}$
Loss $\sim 10\text{kW}$ with 5 mm aperture



Compact spreaders/combiners



This concept allows to use most of the RHIC straight sections for SF linacs and to use part of the arcs for matching

Conclusions

- High energy, high luminosity ERL-based electron-ion and polarized electron-proton collider is the most promising approach for eRHIC
- Presently there is no show-stoppers and a significant amount of R&D
- There is a clear possibility for eRHIC staging (will be discussed later today)

Back-up slides

Advantages & Challenges of ERL based eRHIC²³

$$L = \left(\frac{4\pi\gamma_i\gamma_e}{r_i r_e} \right) (\xi_i \xi_e) (\sigma'_i \sigma'_e) f \quad \longrightarrow \quad L = \gamma_i f N_i \frac{\xi_i Z_i}{\beta_i^* r_i}$$

- This scheme takes full advantage of cooling of the hadron beams
- Allows use of RHIC tunnel for the return passes and thus allow much higher energy of electrons compared with the storage ring.
- High luminosity up to $10^{34} - 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$
- Allows multiple IPs
- Allows higher range of CM-energies with high luminosities
- Full spin transparency at all energies
- No machine elements inside detector(s)
- No significant limitation on the lengths of detectors
- Energy of ERL is simply upgradeable
- Relatively novel technology
- Needs R&D on polarized gun
- Needs completion of e-cooling R&D (CeC and conventional)

In eRHIC luminosity is determined by the hadron beam!

$$L = f_c \frac{N_e N_h}{4\pi\beta_h^* \epsilon_h} \cdot h \left(\frac{\sigma_s}{\beta_h^*} \right)$$

Round beams

$$\beta_e^* \epsilon_e = \beta_h^* \epsilon_h$$

$$L = \gamma_h \cdot (f_c \cdot N_h) \cdot \frac{\xi_h \cdot Z_h}{\beta_h^* \cdot r_h} \cdot h \left(\frac{\sigma_s}{\beta_h^*} \right)$$

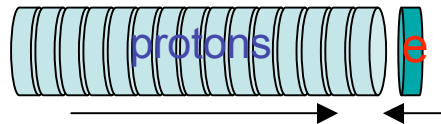
$$\xi_h = \frac{N_e}{\gamma_h} \frac{r_h}{4\pi Z \epsilon_h}$$

$$\xi_h \rightarrow 0.02 \quad \Leftrightarrow \quad L_{p\ e} \rightarrow 0.3 \cdot 10^{34}$$

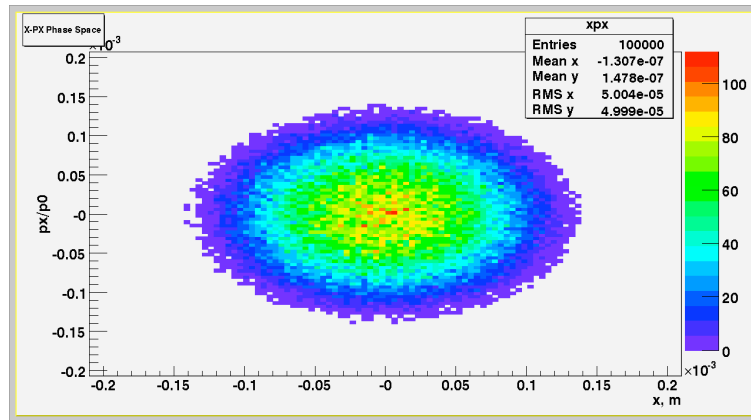
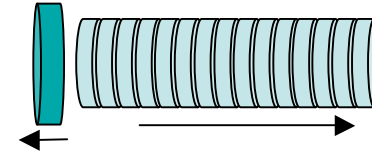
Thus, reducing (cooling) emittance of hadron beam, ϵ_h , allows to proportionally reduce electron beam current ($N_e \sim \epsilon_h$). This in return reduces strain on photocathode, loss on synchrotron radiation -> means higher energy!, X-ray back-ground in detectors.... In combination with reduction of the bunch length, this also allows reduction of β^* and an increase of the luminosity.

Thus, strong cooling makes eRHIC a perfect EIC!

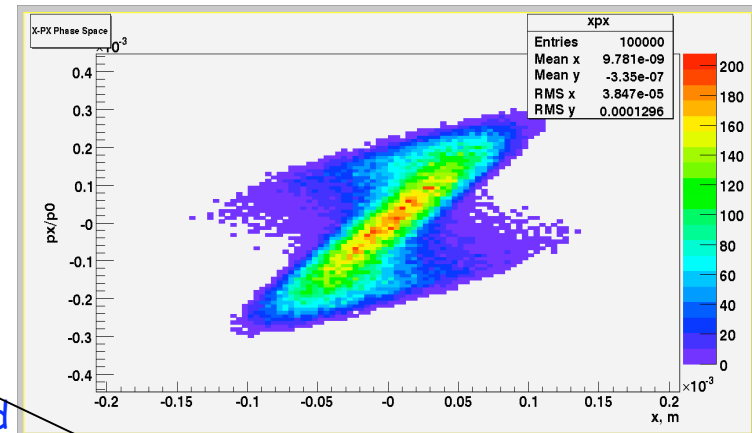
Beam mismatch - e-lens or ferrite lens for compensation



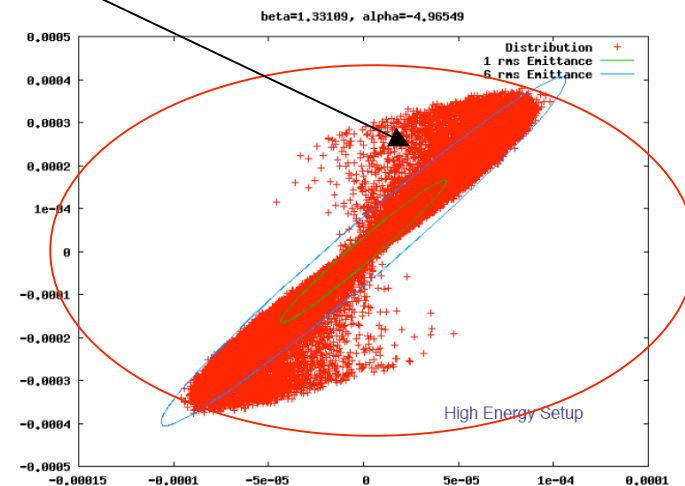
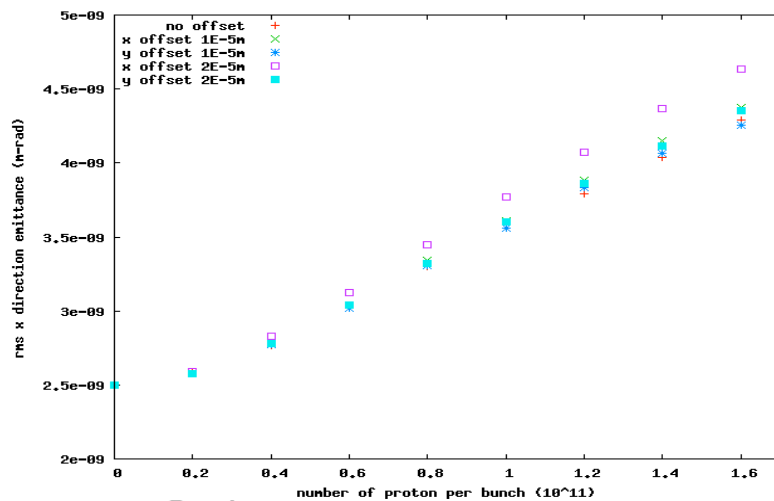
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Interaction



Optimized



Main advantages of ERL + cooling (cont..)

- Where is the limit?

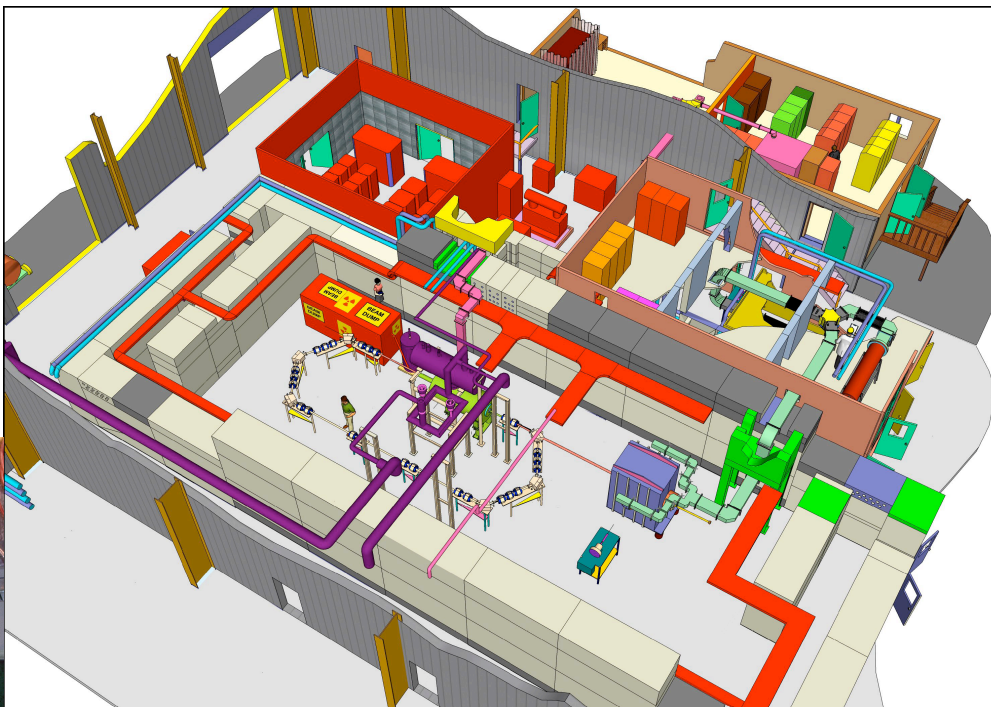
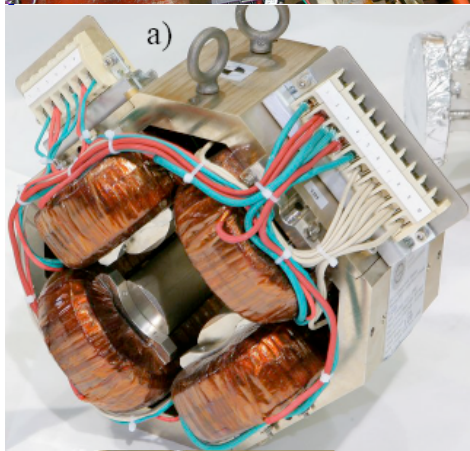
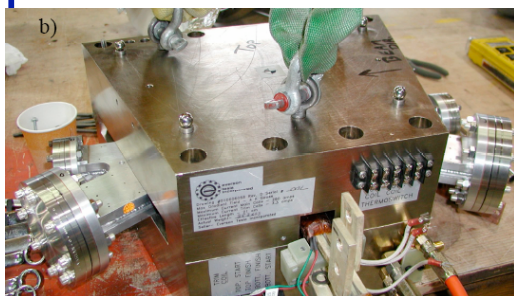
$$D = \frac{Z_h N_h r_e}{\gamma_e \beta_h^* \epsilon_h} \sigma_{s \ h}$$

- Electron beam disruption (which better describes affect on electron beam in linac case) can cause emittance growth and kink instability of the hadron beam

$$\Lambda = D \cdot \xi_h / Q_{s \ h}$$

R&D ERL

Commissioning start 2/09



PoP of coherent electron cooling

- Use existing R&D ERL
- Design & simulations - 2008-2010
- RHIC modification for PoP - 2011
- Moving R&D ERL and installing it at RHIC - 2012
 - ? - should we speed it up to be ahead of NP LRP ?
- Total budget - \$9M-\$10M

Topics of active research for eRHIC

- High charge / high average current, normal and polarized e guns
- High current ERLs
- High energy electron cooling of protons/ions
 - Electron cooling requires SRF-ERL technology
- Integration of interaction region design with detector geometry
- Detailed studies of disruption of the electron beam and kink instability
- Study possibility of shortening hadron bunches in RHIC or of suppressing kink instability by feedback

Major R&D issues

- **Ring-ring:**
 - The accommodation of synchrotron radiation power load on vacuum chamber. (To go beyond $5 \cdot 10^{32} \text{ cm}^{-2}\text{s}^{-1}$ luminosity).
- **Linac-ring:**
 - High current polarized electron source
 - Energy recovery technology for high energy and high current beams
- **Ion ring:**
 - Beam cooling techniques development (electron, stochastic).
 - Increasing total current (ions per bunch and number of bunches).
 - Polarized He^3 production (EBIS) and acceleration

Beam parameters

RHIC	main case
Ring circumference [m]	3834
Number of bunches	360
Beam rep-rate [MHz]	28.15
Protons: number of bunches	180
Beam energy [GeV]	26 - 250
Protons per bunch (max)	$2.0 \cdot 10^{11}$
Normalized 96% emittance [μm]	14.5
β^* [m]	0.26
RMS Bunch length [m]	0.2
Beam-beam tune shift in eRHIC	0.005
Synchrotron tune, Q_s	0.0028
Gold ions: number of bunches	180
Beam energy [GeV/u]	50 - 100
Ions per bunch (max)	$2.0 \cdot 10^9$
Normalized 96% emittance [μm]	6
β^* [m]	0.25
RMS Bunch length [m]	0.2
Beam-beam tune shift	0.005
Synchrotron tune, Q_s	0.0026
Electrons:	
Beam rep-rate [MHz]	14
Beam energy [GeV]	2 - 20
RMS normalized emittance [μm]	5- 50 <i>for $N_e = 10^{10} / 10^{11}$ e^- per bunch</i>
β^*	<i>$\sim 1\text{m}$, to fit beam-size of hadron beam</i>
RMS Bunch length [m]	0.01
Electrons per bunch	$0.1 - 1.0 \cdot 10^{11}$
Charge per bunch [nC]	1.6 - 16
Average e-beam current [A]	0.045 - 0.22