

Basic Principles of Magnetic Resonance

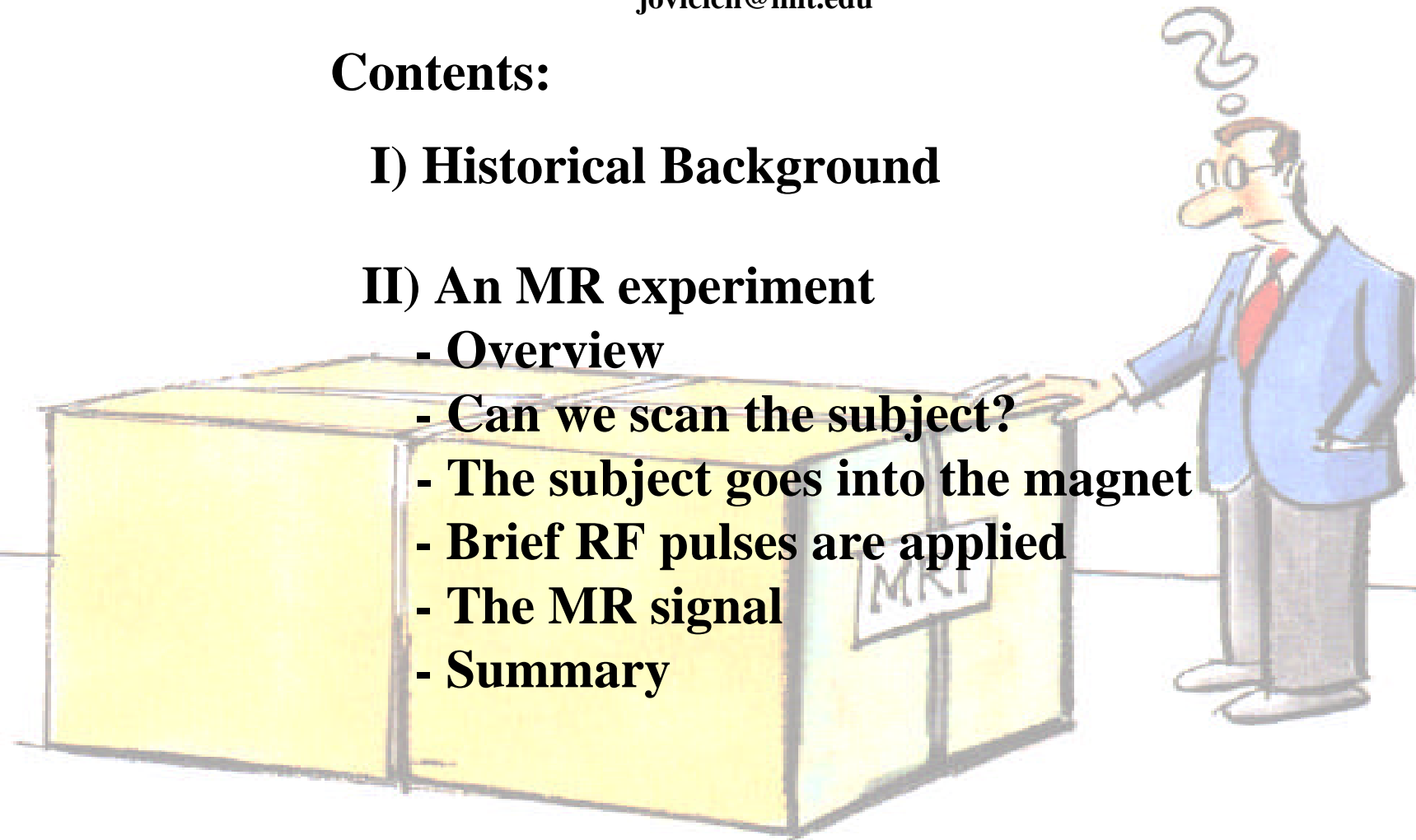
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Contents:

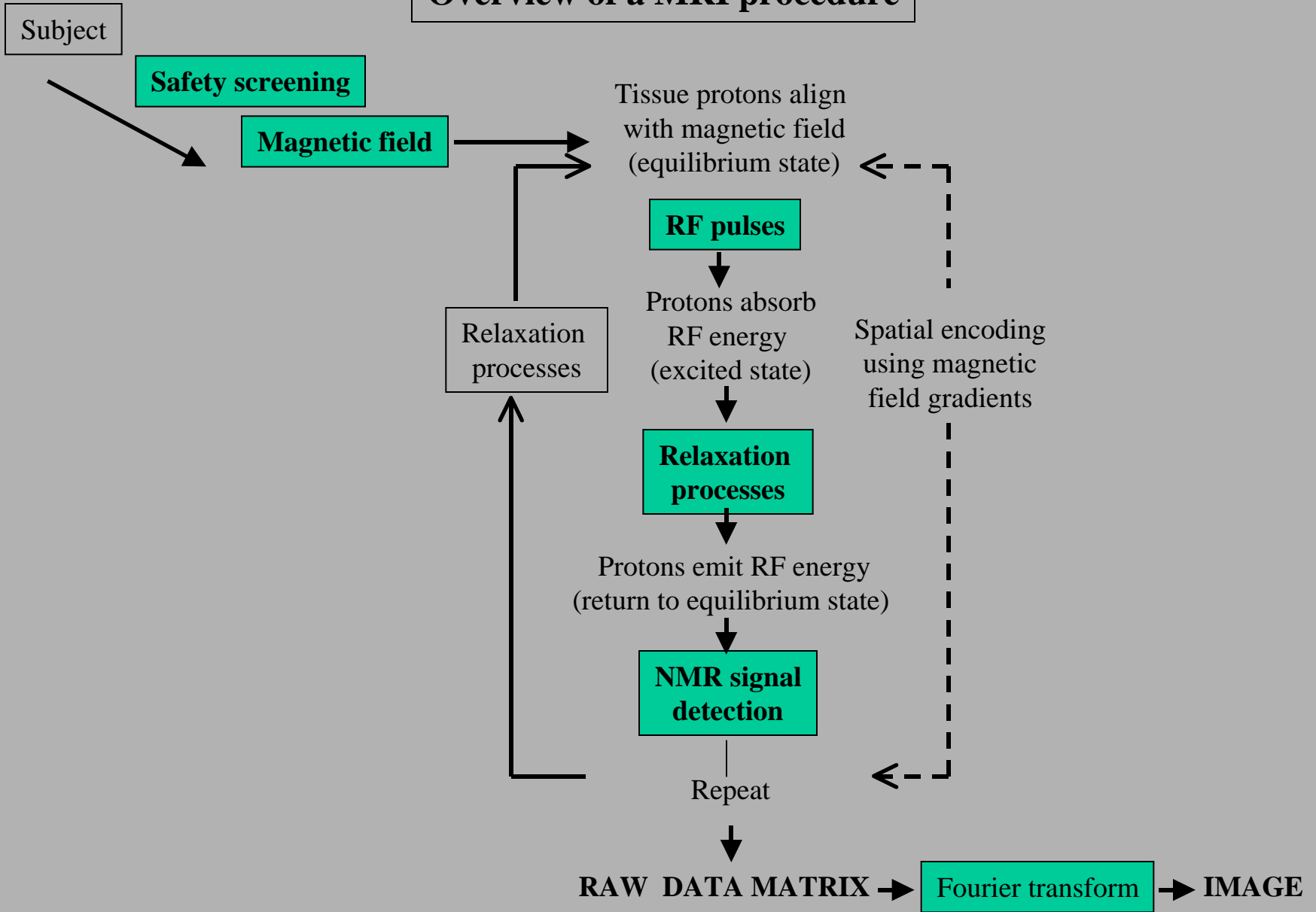
I) Historical Background

II) An MR experiment

- Overview
- Can we scan the subject?
- The subject goes into the magnet
- Brief RF pulses are applied
- The MR signal
- Summary



Overview of a MRI procedure



Can we scan the subject?

Safety Issues

- Not everybody can undergo a MR examination
- Screening is mandatory to ensure safety / quality during the MR examination

A typical MRI scanner



The main magnetic field:

- is **VERY** strong (3T ~ 30,000 times the earth's magnetic field)
- is **ALWAYS ON!!** (superconducting currents)
- it's strength decays $\sim 1 / r^3$ (r is the distance from the center)

Can we scan the subject?

Safety Issues

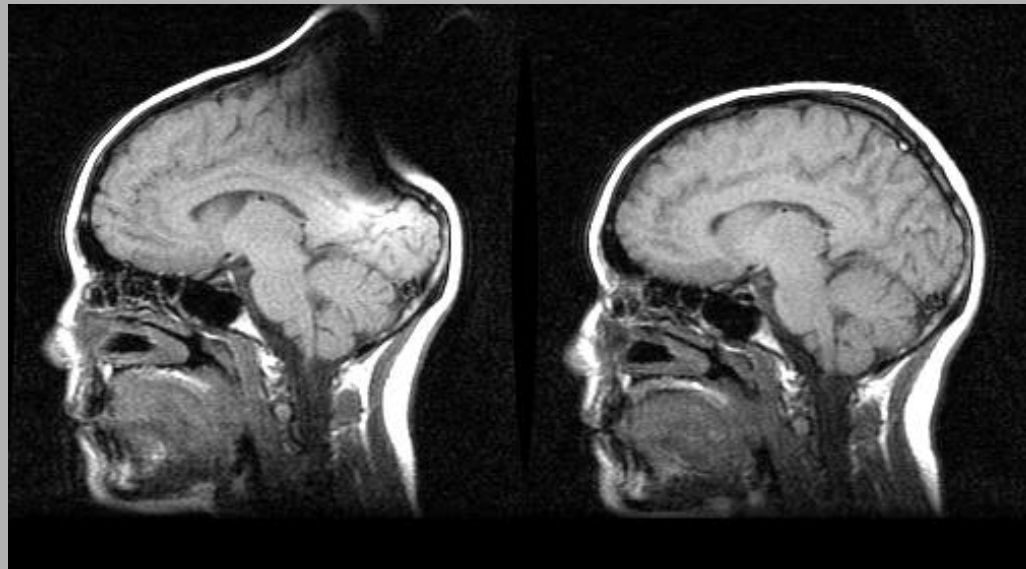
- Risks:
 - Ferromagnetic materials (unsafe)
 - will be subject to force / torque in the main magnetic field
 - risks associated with motion / displacements
 - active biomedical implants may not function properly



Can we scan the subject? Safety Issues

- Risks:
 - Ferromagnetic materials (unsafe)
 - Non ferromagnetic materials (safe but give image distortions)

Sagittal MRI
of a normal
female subject



With hair rubber band*

Without hair rubber band

J. Jovicich, MIT

* The two ends of the rubber band are joined by a ~ 2 mm cooper clamp.

Can we scan the subject?

Safety Issues

- Risks:
 - Ferromagnetic materials (unsafe)
 - Non ferromagnetic materials (safe but give image distortions)
 - Claustrophobia (subject unhappy image distortions)



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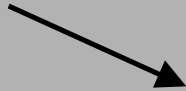
Can we scan the subject?

Safety Issues

- Risks:
 - Ferromagnetic materials (unsafe)
 - Non ferromagnetic materials (safe but give image distortions)
 - Claustrophobia (subject unhappy image distortions)
 - Movement during examination (potential problem for dynamic studies)
 - Acoustic protection (mandatory)

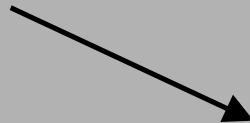
Overview of a MRI procedure

Subject



Safety screening

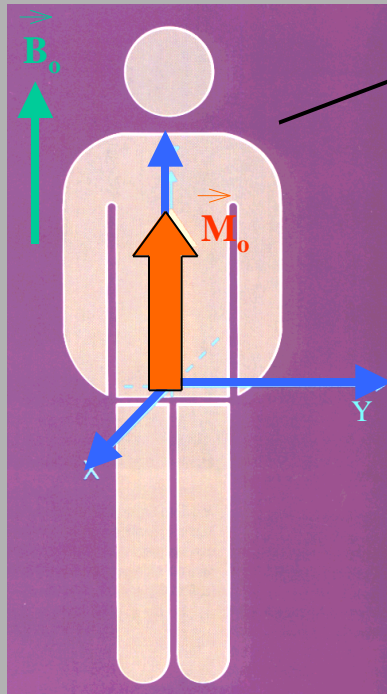
Magnetic field



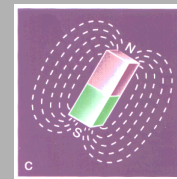
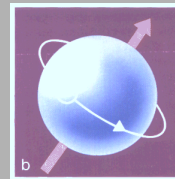
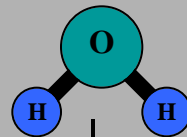
We will introduce the following concepts:

- equilibrium magnetization
- dynamics of the magnetization
- rotating coordinate system

The subject goes into the magnet...

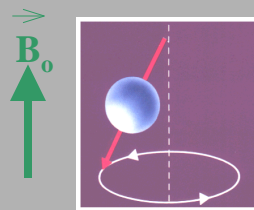


Single nucleus

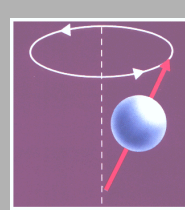


Anti-parallel

Parallel



High energy state
(E_1)



Low energy state
(E_0)

Water molecules

Hydrogen nucleus:
magnetic moment $\vec{\mu}$

Two energy states:

$$E = E_1 - E_0$$

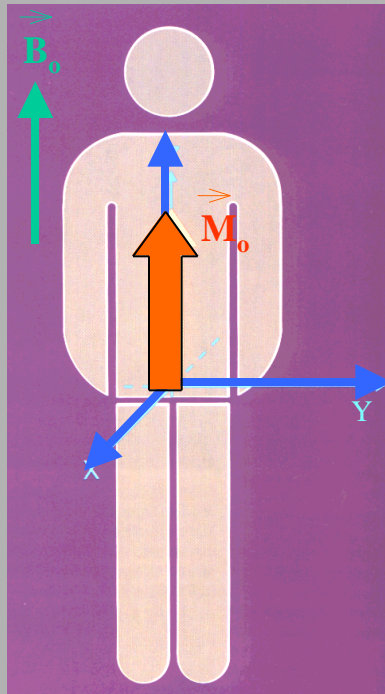
$$E = \hbar\omega_0$$

Precession frequency:

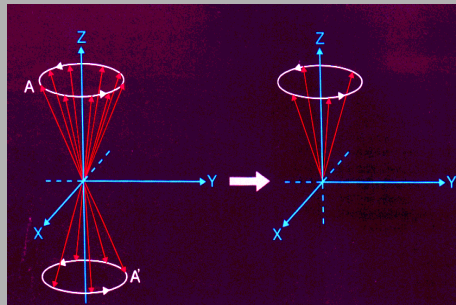
$$\omega_0 = \gamma \mathbf{B}_0$$

\vec{B}_0 : uniform static magnetic field
 \vec{M}_0 : static macroscopic magnetization

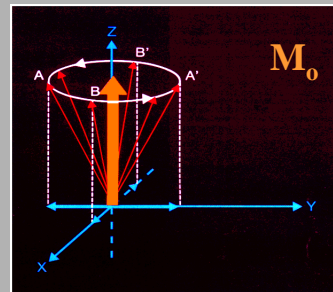
The subject goes into the magnet... (continued)



A group of protons



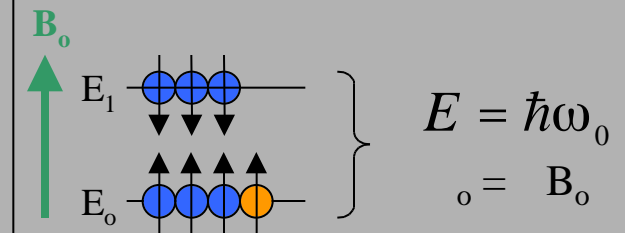
Net magnetization M_0



Equilibrium magnetization M_0 :

- M_z aligned with B_0
- $M_{xy} = 0$

Energy states



Spin states distribution

$$\frac{N_1}{N_0} = \exp \left(-\frac{\hbar\gamma B_0}{kT} \right)$$

N_1 = # of protons in state E_1

N_0 = # of protons in state E_0

k : Boltzmann's constant

T : temperature

Reminder of three main steps in MRI

0) Equilibrium (M_o along B_o)



1) RF excitation
(tip M_o away from equil.)



2) Precession of M_{xy} induces signal
(dephasing for a time TE)



3) Return to equilibrium
(recovery time TR)

Dynamics of
magnetization

Dynamics of the magnetization

- Equation of motion of \vec{M} in external \vec{B}_{ext}
- Classical mechanics formalism
- First, model a single nucleus
- Then, add up for sample
- Finally, consider relaxation (Bloch equations)

Equation of motion for a **single nucleus**:

- Mechanical moment angular momentum (spin)

$$\vec{T} = \frac{d\vec{L}}{dt}$$

- Magnetic moment spin

$$\vec{\mu} = \gamma\vec{L}$$

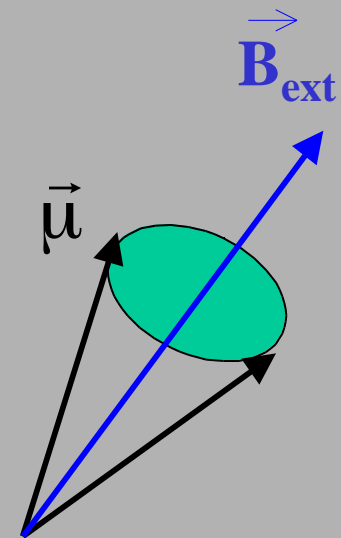
- Magnetic moment and magnetic field interaction

$$\vec{T} = \vec{\mu} \times \vec{B}_{ext}$$



$$\frac{d\vec{\mu}_{(t)}}{dt} = \vec{\mu}_{(t)} \times \gamma\vec{B}_{ext(t)}$$

Precession of
 $\vec{\mu}$ about \vec{B}_{ext}
with frequency
 $\omega = \gamma B_{ext}$



Equation of motion for the **magnetization vector**:

- Assuming no interaction between nuclei

$$\vec{M} = \vec{\mu}_0 + \vec{\mu}_1 + \vec{\mu}_2 + \dots = \vec{\mu}_i \quad (\text{spin excess})$$

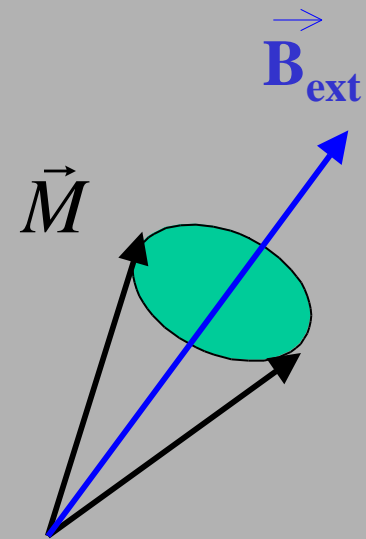
- And since for each nuclei

$$\frac{d\vec{\mu}_{i(t)}}{dt} = \vec{\mu}_{i(t)} \times \gamma \vec{B}_{ext(t)}$$



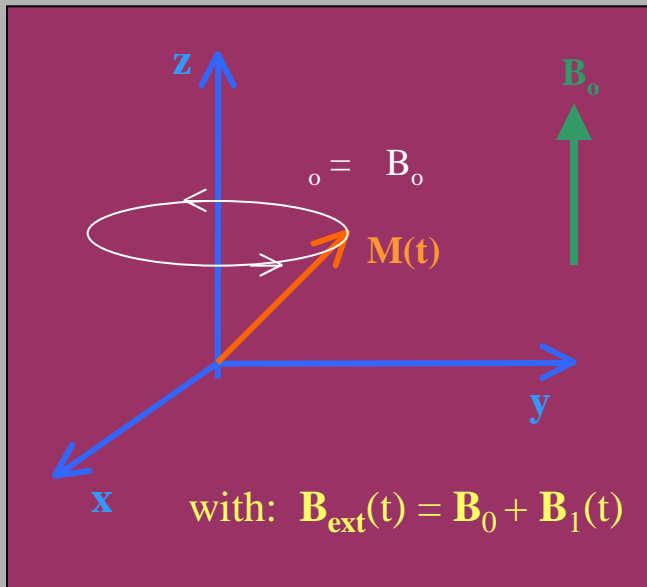
$$\frac{d\vec{M}_{(t)}}{dt} = \vec{M}_{(t)} \times \gamma \vec{B}_{ext(t)}$$

Precession of
 \vec{M} about \vec{B}_{ext}
with frequency
 $\omega = \gamma B_{ext}$



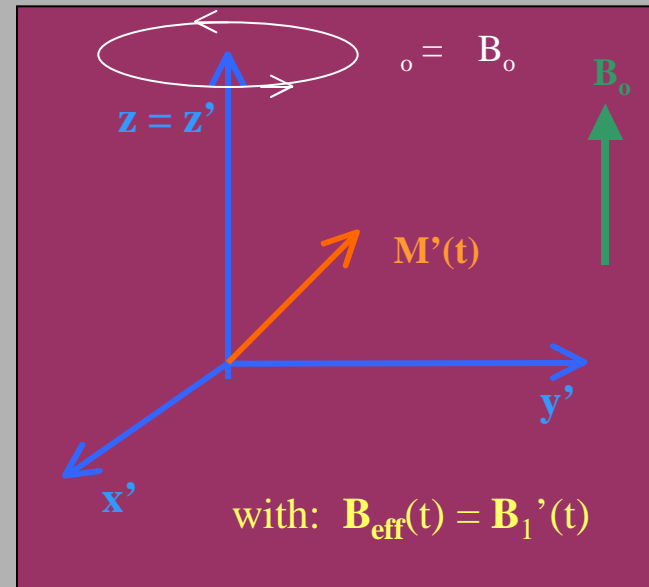
To describe the magnetization we need to choose a coordinate systems

Laboratory coordinate system



$$\frac{d\vec{M}_{(t)}}{dt} = \vec{M}_{(t)} \times \gamma \vec{B}_{ext(t)}$$

Rotating coordinate system



$$\frac{d\vec{M}'_{(t)}}{dt} = \vec{M}'_{(t)} \times \gamma \vec{B}_{eff(t)}$$

Example:

On-resonance spins:

Off-resonance spins:



o

o +



static

The NMR signal

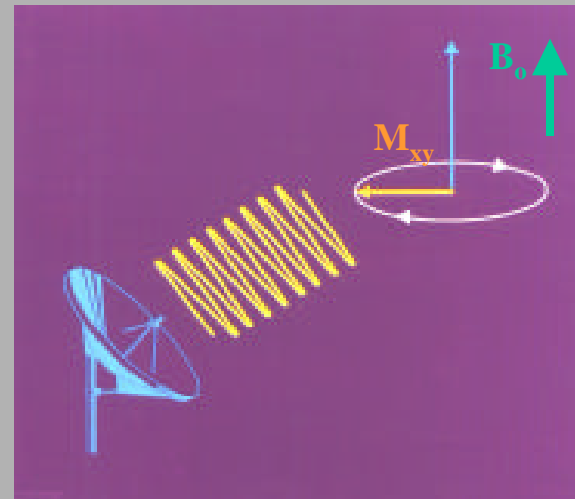
- **At equilibrium no signal:**

- static longitudinal magnetization M_0



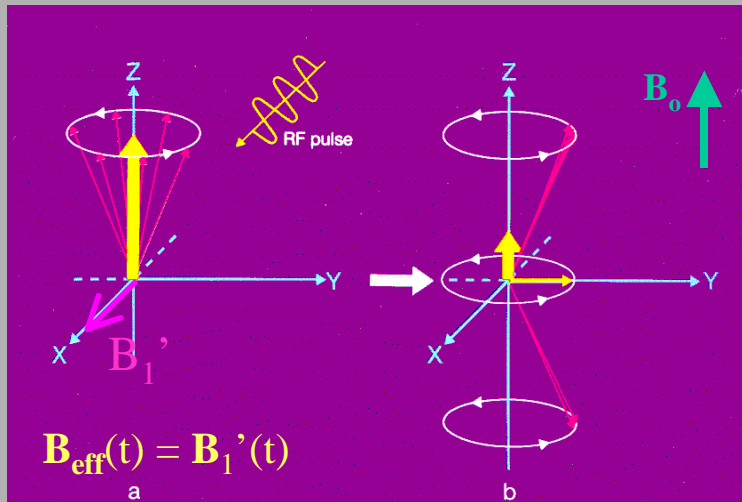
- **We need net transverse magnetization:**

- precession of M_{xy} about B_0
- rotating magnetic field
- induces current in a coil: MR signal



The NMR signal (continuation)

- With an RF pulse on resonance we can rotate M_0



Rotating frame:

- Dynamics: $\frac{d\vec{M}'_{(t)}}{dt} = \vec{M}'_{(t)} \times \gamma \vec{B}_{eff(t)}$
- $B_1'(t)$: B_1' constant and \perp to B_0
- M' precesses about B_1 with $\omega_1 = \gamma B_1'$
- Flip angle : $\theta = \gamma B_1' t$

- Signal relaxation, system goes back to equilibrium

$$M_z \rightarrow M_0 \quad (T_1 \text{ relaxation})$$

$$M_{xy} \rightarrow 0 \quad (\text{signal loss, } T_2 \text{ \& } T_2^* \text{ relaxation})$$

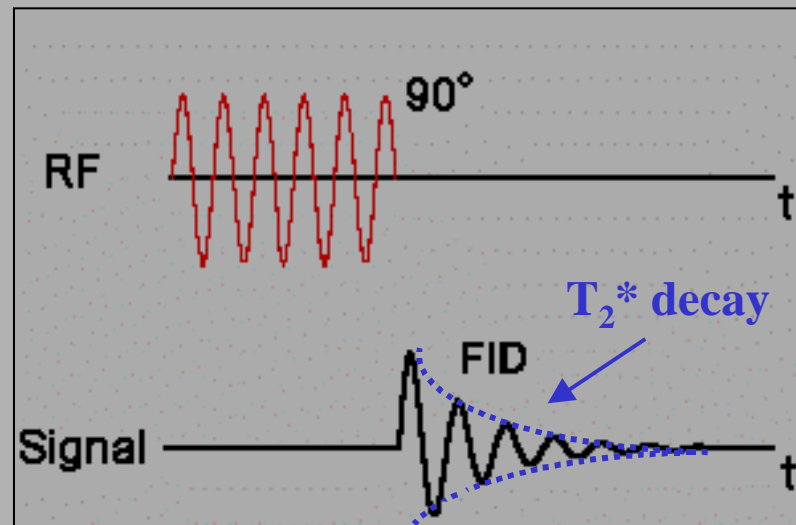
- Relaxation mechanisms give tissue contrast

The NMR signal (continuation)

Radio-frequency pulse
(oscillating $B_1(t)$ rotating at ω_0)

NMR signal
(transverse magnetization decay)
'Free Induction Decay'

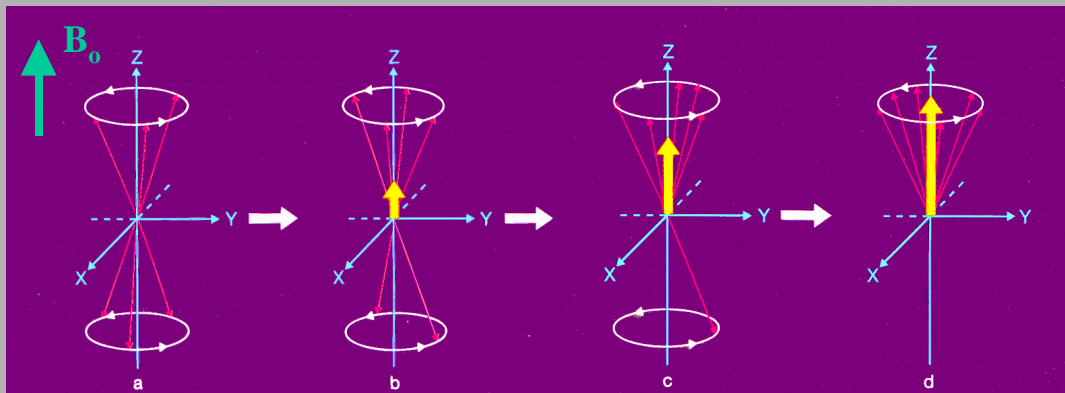
Schematic representation:



Relaxation mechanisms

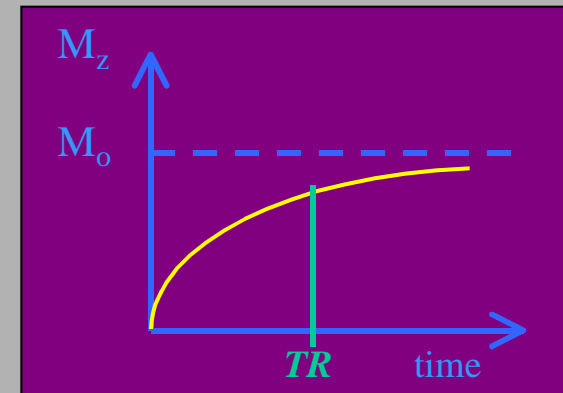
T_1 or spin-lattice relaxation (longitudinal magnetization)

- M_z defined by spin excess population between two energy states
- M_z recovery transitions between spin states
- transitions fluctuating transverse field on resonance molecular motion
- exponential recovery (T_1 100 - 3000 ms, longer for higher B_0)



$M_z=0$
after 90°
RF pulse

$M_z=M_0$
at equilibrium



$$\frac{dM_z}{dt} = \frac{M_0 - M_z}{T_1}$$

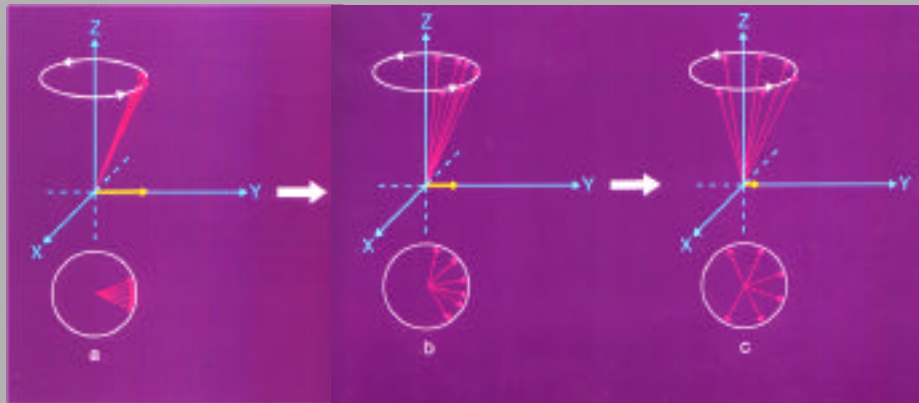
Repetition time (TR): time allowed for recovery, defines T_1 contrast

Relaxation mechanisms (continued)

T_2 or spin-spin relaxation (transverse magnetization)

- incoherent exchange of energy between spins
- molecular motion fluctuations in local B_z resonance frequency variations
- dephasing of transverse magnetization signal decay
- exponential decay (T_2 70 - 1000 ms)

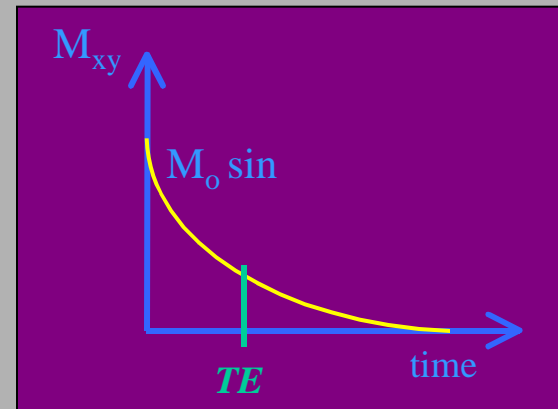
$M_{xy} \leftrightarrow$ NMR signal



M_{xy} max
after 90°
RF pulse

\rightarrow spins dephase \rightarrow

$M_{xy} = 0$
at equilibrium



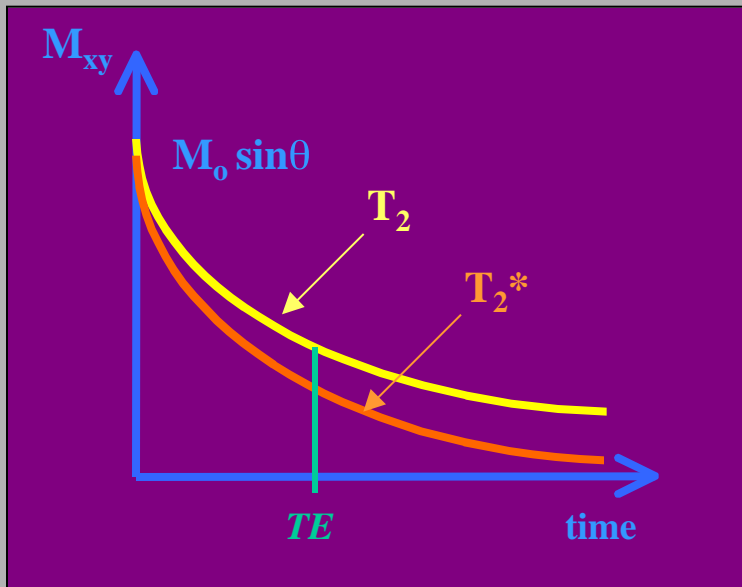
$$\frac{dM_{xy}}{dt} = -\frac{M_{xy}}{T_2}$$

Echo time (TE): time allowed for dephasing, defines T_2 contrast

Relaxation mechanisms (continued)

T_2^* relaxation

- dephasing of transverse magnetization due to both:
 - microscopic molecular interactions (T_2)
 - spatial variations of the external main field B (tissue/air, tissue/bone interfaces)
- exponential decay ($T_2^* \approx 30 - 100$ ms, shorter for higher B_0)



$$B = \mu \mu_0 H$$

Boundary conditions

Field distortions
(B)

μ_{tissue}

μ_{air}

$$\frac{1}{T_2^*} = \frac{1}{T_2} + \gamma \frac{\Delta B}{2}$$

Bloch Equations

- Dynamics of the magnetization + Relaxation effects:

$$\frac{d\vec{M}_{(t)}}{dt} = \vec{M}_{(t)} \times \gamma \vec{B}_{ext(t)} - \frac{(M_x \hat{i} + M_y \hat{j})}{T_2^*} - \frac{(M_z - M_0)}{T_1} \hat{k}$$

- Geometrical description: damped precession (SUMMARY)

