

# Weak Interaction Studies with



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in collaboration with  
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# Weak Currents and QCD

- One of the main goals of QCD is to understand nucleons and nuclei in terms of quarks and the forces of QCD.
- We believe that the spontaneous breaking of chiral symmetry of QCD is responsible for the large effective mass of quarks confined in nucleons and nuclei.
- The best way to study these effects is to measure the weak currents in nucleons and nuclei. The axial current is of primary importance.
- However, the axial form factor of light nuclei such as  $^3\text{He}$  has never been measured.
- Weak interaction measurements will also provide information about axial two-body currents.
- Weak capture of polarized electrons in  $^3\text{He}$  has great potential to fill this gap and also help test for 2<sup>nd</sup> class currents.
- Complementary to e-capture in proton discussed by A. Deur.



# Second Class Currents

Introduced by Weinberg Phys. Rev. **112**, 1375 (1958)

Represent the properties of strangeness conserving ( $\Delta S=0$ )  
Vector and Axial currents under G-parity transformation

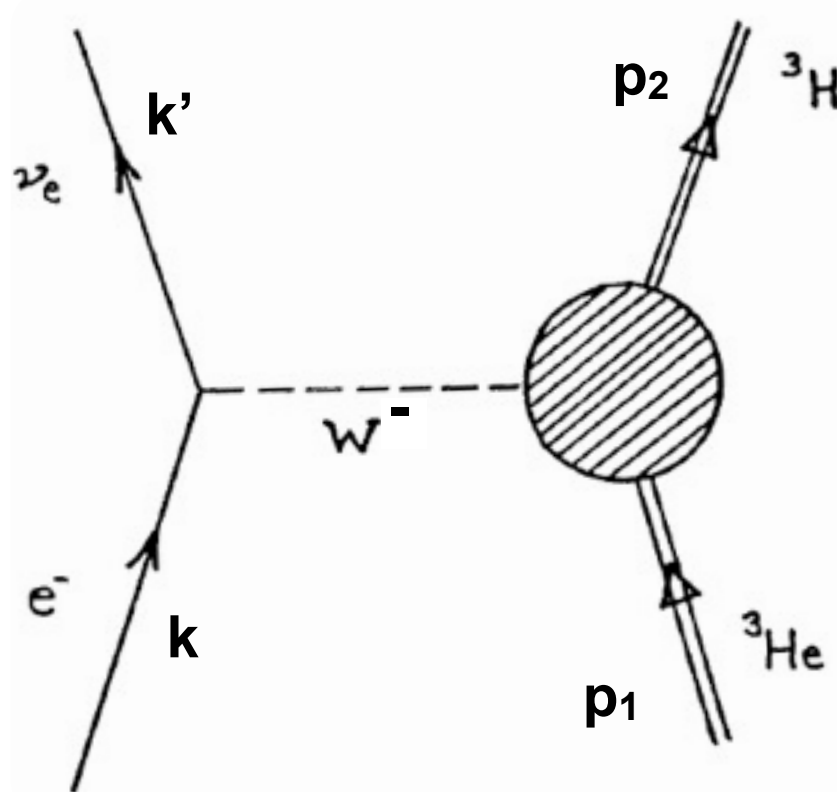
First class:  $GV_{\mu}^{(I)\pm}G^{-1} = V_{\mu}^{(I)\pm}, \quad GA_{\mu}^{(I)\pm}G^{-1} = -A_{\mu}^{(I)\pm}$

Second class:  $GV_{\mu}^{(II)\pm}G^{-1} = -V_{\mu}^{(II)\pm}, \quad GA_{\mu}^{(II)\pm}G^{-1} = A_{\mu}^{(II)\pm}$

**Experimentally, so far there is no evidence  
for or against the existence of second class currents**



# The $e + {}^3\text{He} \rightarrow {}^3\text{H} + \nu$ Reaction



$$\mathbf{q} = \mathbf{k} - \mathbf{k}', \quad \mathbf{P} = \mathbf{p}_1 - \mathbf{p}_2$$

Neglecting lepton masses

$$\frac{d\sigma}{d\Omega_n} = K(\mathbf{E}_e, \mathbf{E}_\nu, \theta_\nu) \sigma_0(f_E, f_M, f_A)$$

$f_E, f_M, f_A$  are the electric, magnetic and axial formfactors of  ${}^3\text{He}$

Neglecting electron mass, there are no contributions from 2<sup>nd</sup> class currents and no spin-momentum correlations.

Small cross section  $\sim 10^{-40} \text{ cm}^2/\text{sr}$



# Spin Dependent e-Capture

If the mass of the electron is not neglected  
(relevant for the low energies proposed here)

$$\frac{d\sigma}{d\Omega_n} = K(\mathbf{E}_e, \mathbf{E}_\nu, \theta_\nu) \sigma_0 \left[ 1 + \frac{m_e}{\sigma_0} [(\mathbf{k}' \cdot \mathbf{s})g_1(\mathbf{W}_1, \mathbf{W}_2, \mathbf{W}_5) + (\mathbf{P} \cdot \mathbf{s})g_2(\mathbf{W}_2, \mathbf{W}_5) + (\mathbf{q} \cdot \mathbf{s})g_3(\mathbf{W}_5, \mathbf{W}_6)] \right]$$

$\mathbf{q} = \mathbf{k} - \mathbf{k}'$ ;  $\mathbf{P} = \mathbf{p}_1 - \mathbf{p}_2$ ;  $\mathbf{s}$  = electron spin vector

$\mathbf{W}_1 - \mathbf{W}_6$  are functions of the nuclear form factors  
with  $\mathbf{W}_5$  and  $\mathbf{W}_6$  entirely due to second class currents



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Observation of a **non-zero correlation between  $\mathbf{q}$ -vector and electron spin direction** would be sufficient to **verify the existence of second-class currents**.



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Observation of a **non-zero correlation between  $\mathbf{q}$ -vector and electron spin direction** would be sufficient to **verify the existence of second-class currents**.

Also relevant for the  $\bar{e} + p \rightarrow \nu + n$  reactions discussed in the  
previous talk





# Spin Dependent e-Capture

**If the mass of the electron is not neglected**  
(relevant for the low energies proposed here)

Ignoring second class currents and for  $q^2 \ll M_{\text{Targ}}^2$

$$\frac{d\sigma}{d\Omega_n} = K(E_e, E_\nu, \theta_\nu) \sigma_0 [1 + A_1 + \gamma h m_e A_2 \cos \theta_n (1 + \frac{|f_A(q^2)|^2}{|f_V(q^2)|^2})],$$

$h$  is electron helicity,  $f_A$  and  $f_V$  are axial and vector formfactors of  $^3\text{He}$





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**The coefficient of the  $\cos\theta_n$  moment of the  
longitudinal asymmetry can give us the  
ratio of the axial to vector formfactor**



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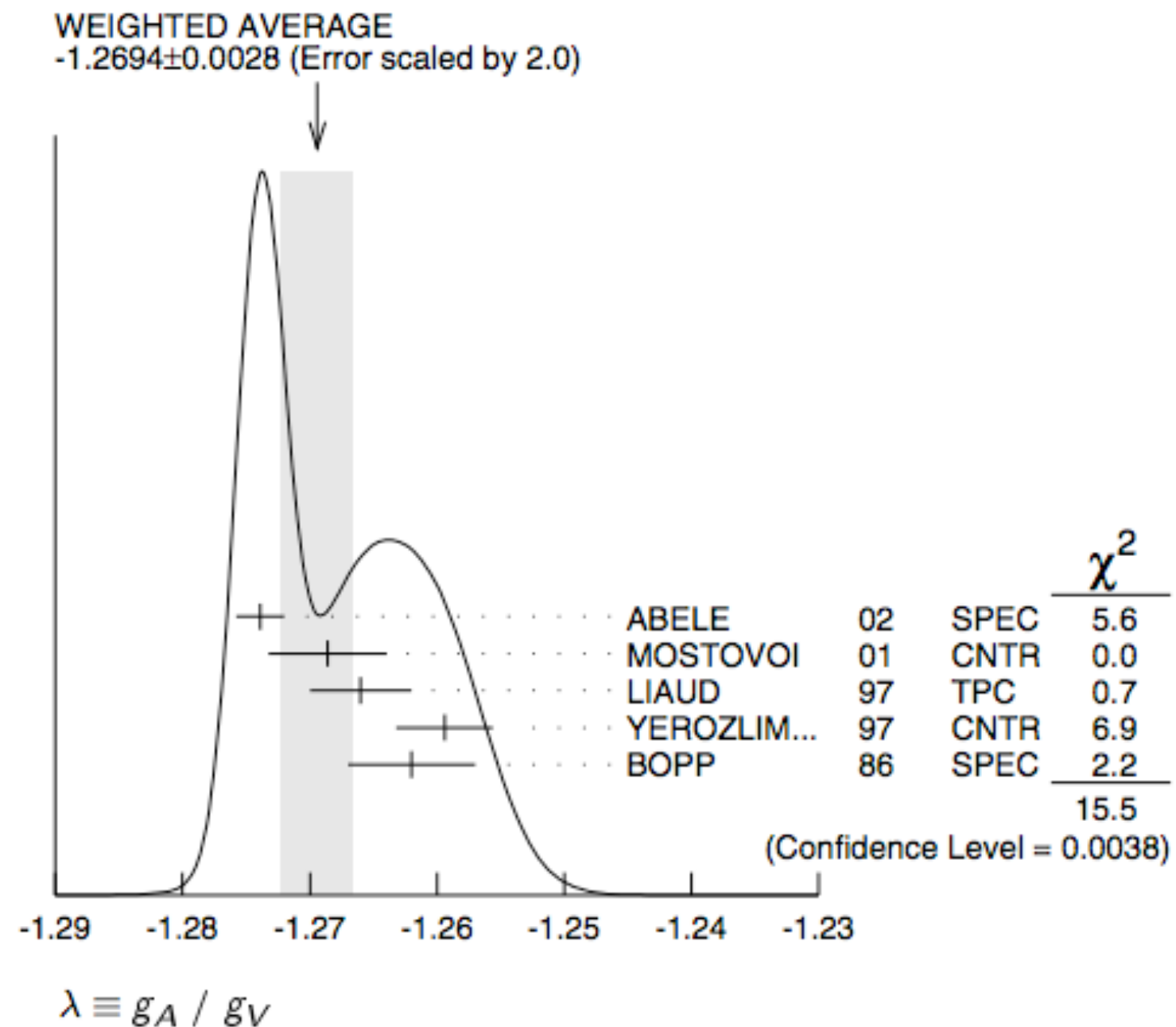
$h$  is electron helicity,  $f_A$  and  $f_V$  are axial and vector formfactors of  $^3\text{He}$

**The coefficient of the  $\cos\theta_n$  moment of the longitudinal asymmetry can give us the ratio of the axial to vector formfactor**

Implies that  $\vec{e} + p \rightarrow \nu + n$  reaction will allow access to  $|g_A(Q^2)/g_V(Q^2)|$  via a new method unlike the ones traditionally used to measure this ratio.



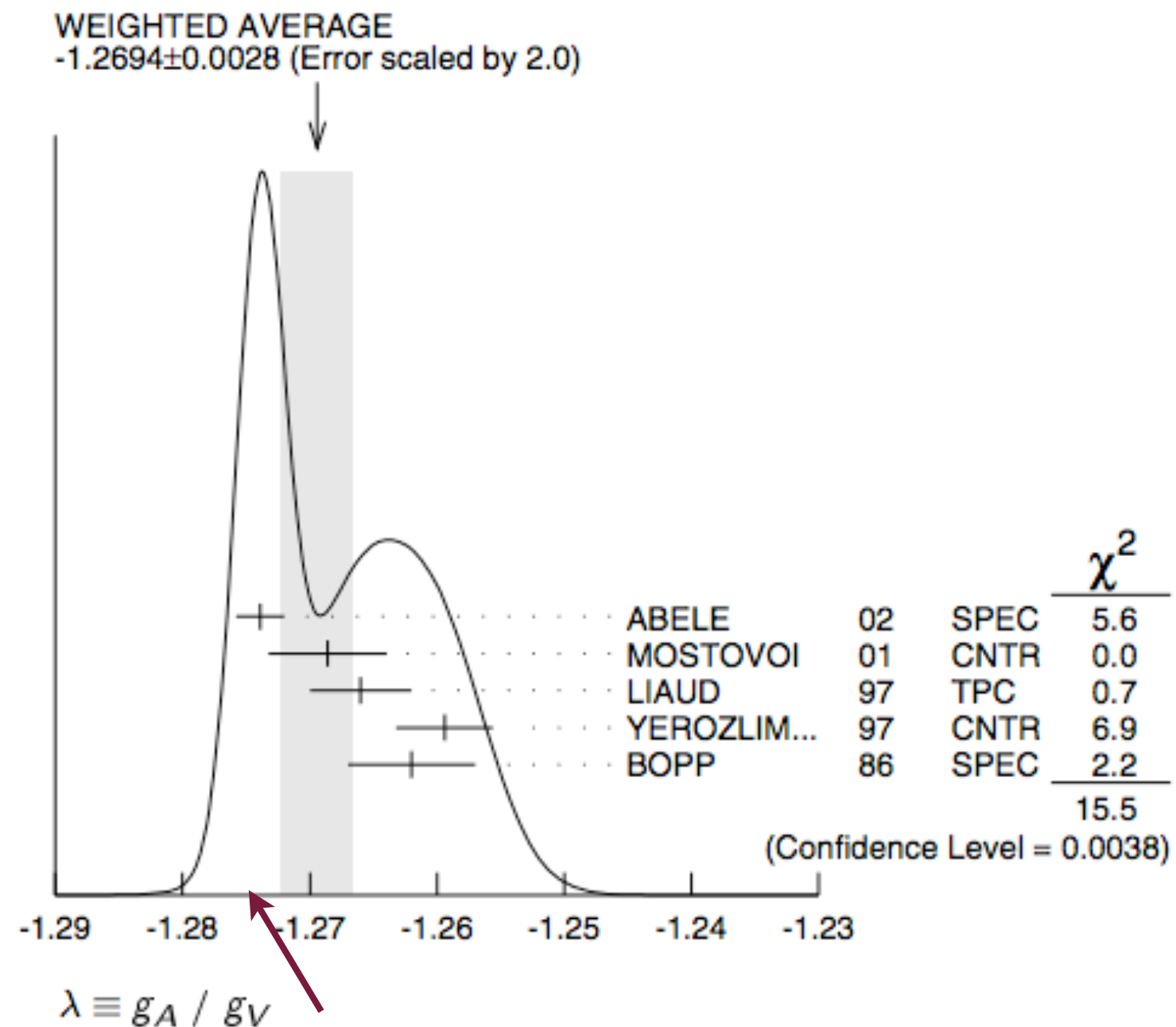
# PDG Compilation of $g_A/g_V$



The most precise measurements of  $g_A/g_V$  from neutron beta decay seems to be inconsistent with each other.



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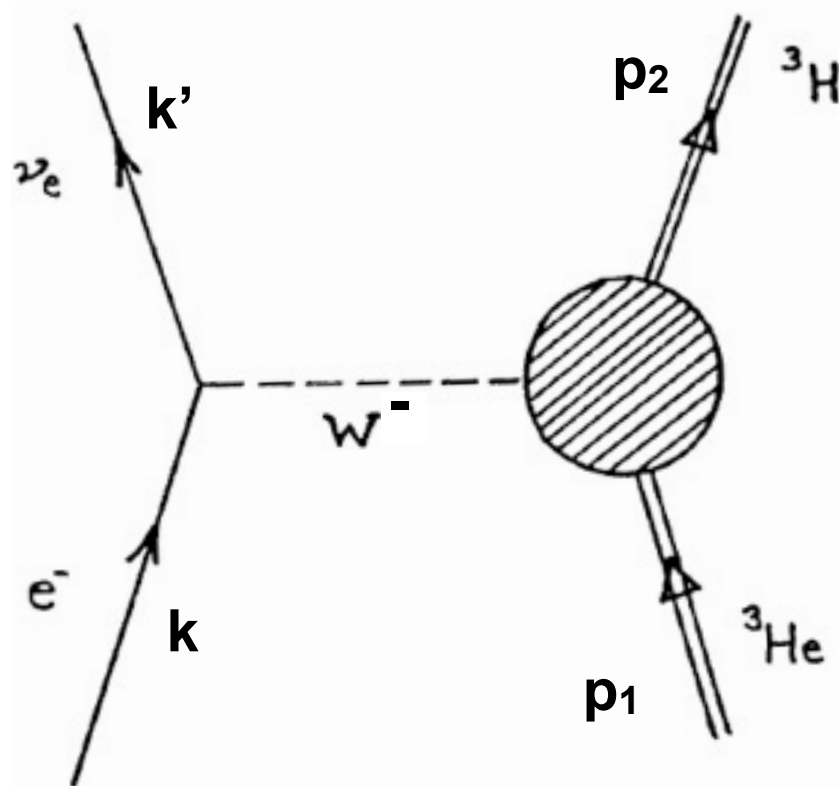


New LANCE results

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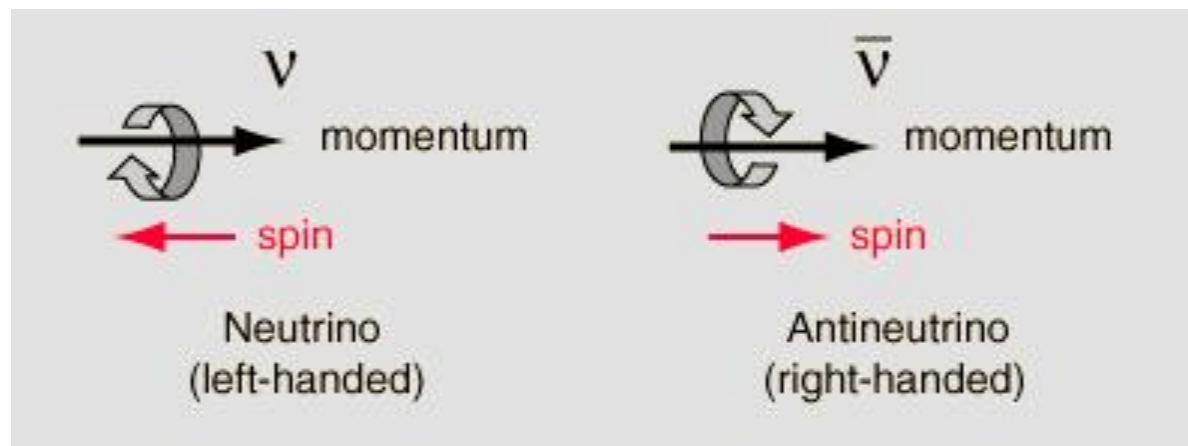
Want to measure a small cross section buried under electromagnetic background to extract the axial formfactor of  ${}^3\text{He}$

## Strategy:

- Detect  ${}^3\text{H}$  in backward lepton kinematics to enhance Weak/EM cross sections
- Detect elastic backward going electrons in coincidence to reduce backgrounds
- Use high intensity low energy beam with a thin target
- Use polarization degrees of freedom to cleanup parity conserving backgrounds



# Spin Dependent e-Capture



Only left-handed neutrinos exist in the Standard Model with massless neutrinos.

If polarized electrons are used



cross section is non-zero only for left handed electrons

For a symmetric detector any parity conserving (electromagnetic) background would cancel in:

$$\sigma (h=+1) - \sigma (h=-1)$$



# Backgrounds

- Elastic scattering from  $^3\text{He}$  and  $^3\text{H}$  (cross sections 8 - 9 orders of magnitude larger)  
Elastic rates  $\sim 10^5 - 10^6$  Hz compared to  $10^{-3}$  Hz for signal  
Target must be ultra-pure with  $^3\text{H}$  contamination at  $10^{-9}$  level  
Elastic background can be reduced by detecting the scattered  $e^-$  in coincidence
- Elastically scattered  $^3\text{He}$  undergoes charge exchange with cell wall to yield  $^3\text{H}$   
Detect scattered  $e^-$  in coincidence
- Elastically scattered  $^3\text{He}$  picking up an  $e^-$  to become indistinguishable from  $^3\text{H}$   
Charge sensitive detector  
Detect scattered  $e^-$  in coincidence

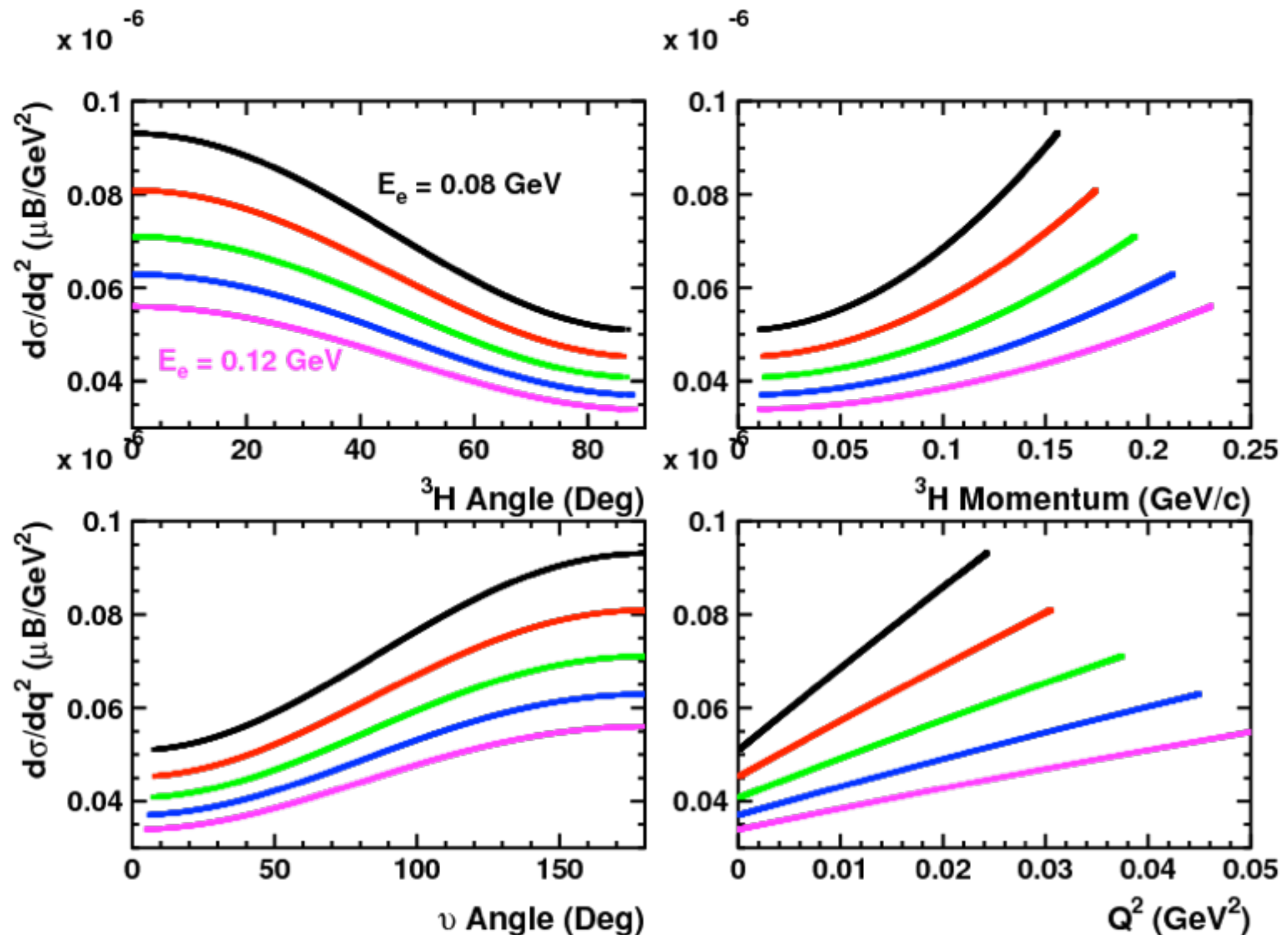
**Use longitudinally polarized electrons with high polarization**





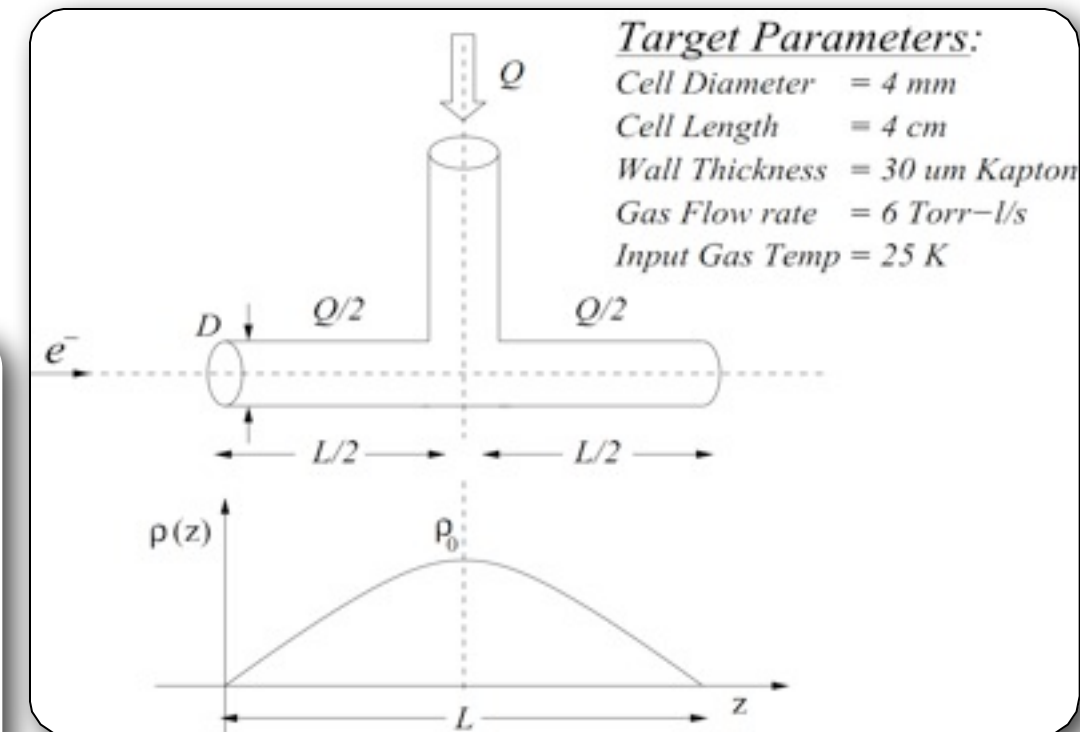
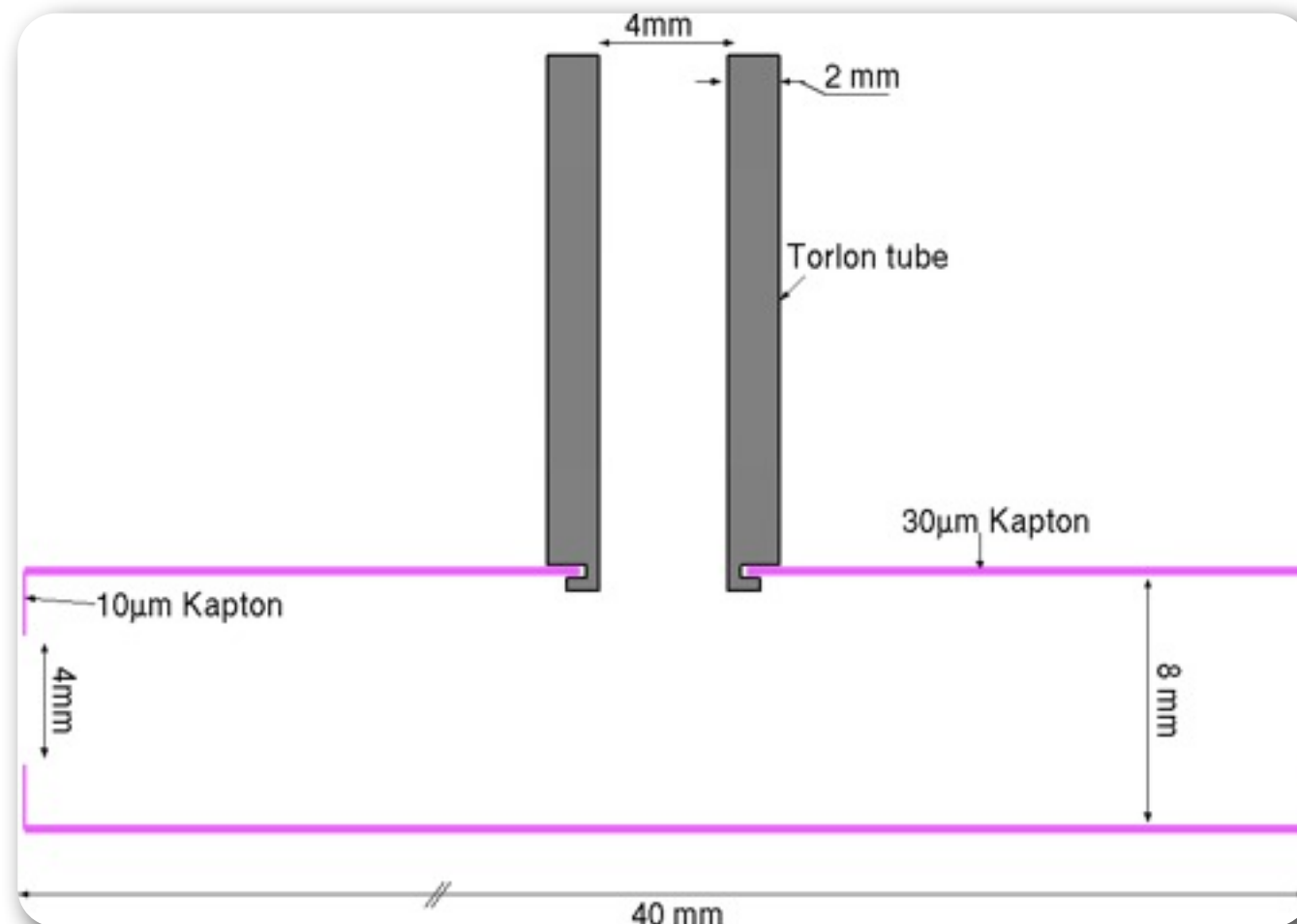
# Weak Charged Current Cross Sections

Using  $\langle r_c^{3\text{He}} \rangle = 1.87 \text{ fm}$ ;  $\langle r_m^{3\text{He}} \rangle = 1.74 \text{ fm}$   
 $\langle r_c^{3\text{H}} \rangle = 1.70 \text{ fm}$ ;  $\langle r_m^{3\text{H}} \rangle = 1.70 \text{ fm}$





# Windowless Target



under development at JLab for the  
**Proton Charge Radius Experiment (PRAD)**

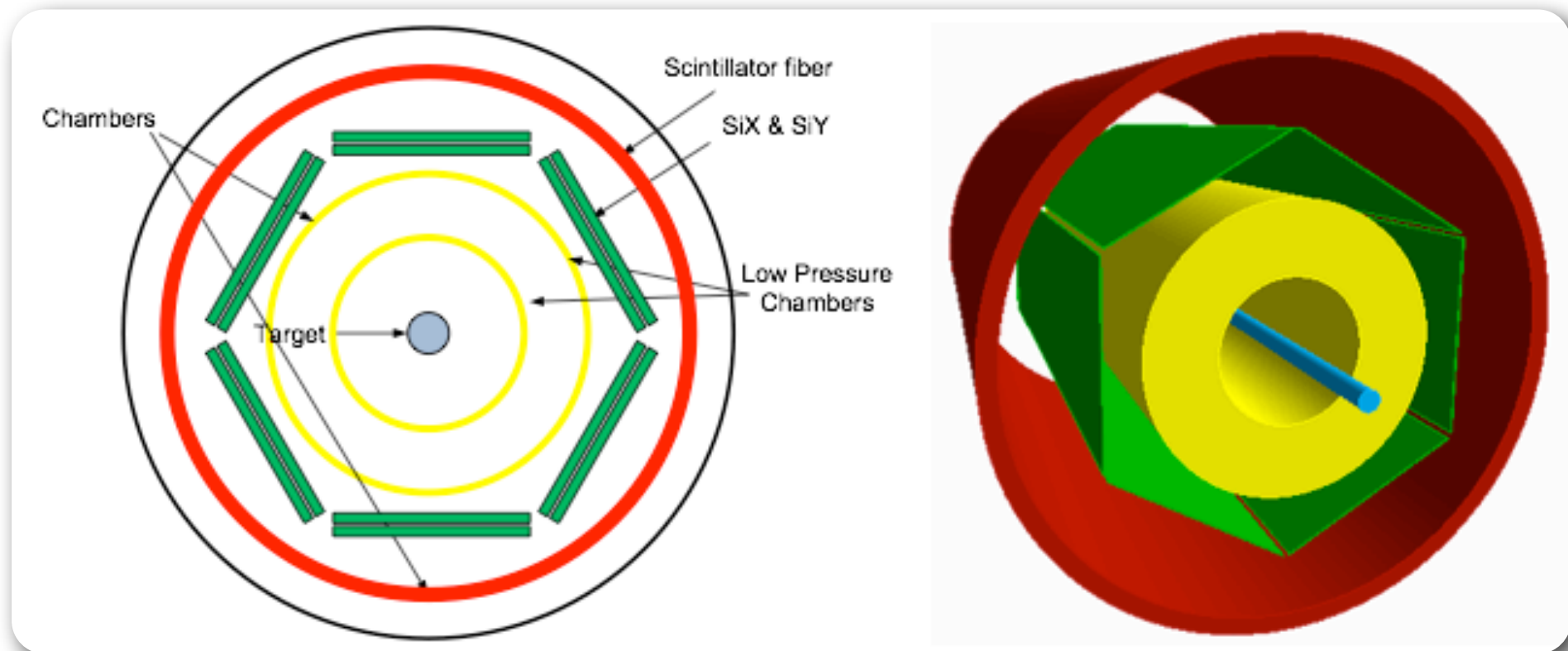
A. Gasparian (NCA&T), D. D(Miss. State), H. Gao (Duke) and M. Khandakar (Idaho State)

Similar target also being developed for the DarkLight experiment at the JLab FEL



# Detector Design

**Large acceptance ( $2\pi$  polar angle and  $20^\circ - 160^\circ$  azimuthal)  
with a solenoidal/toroidal field**

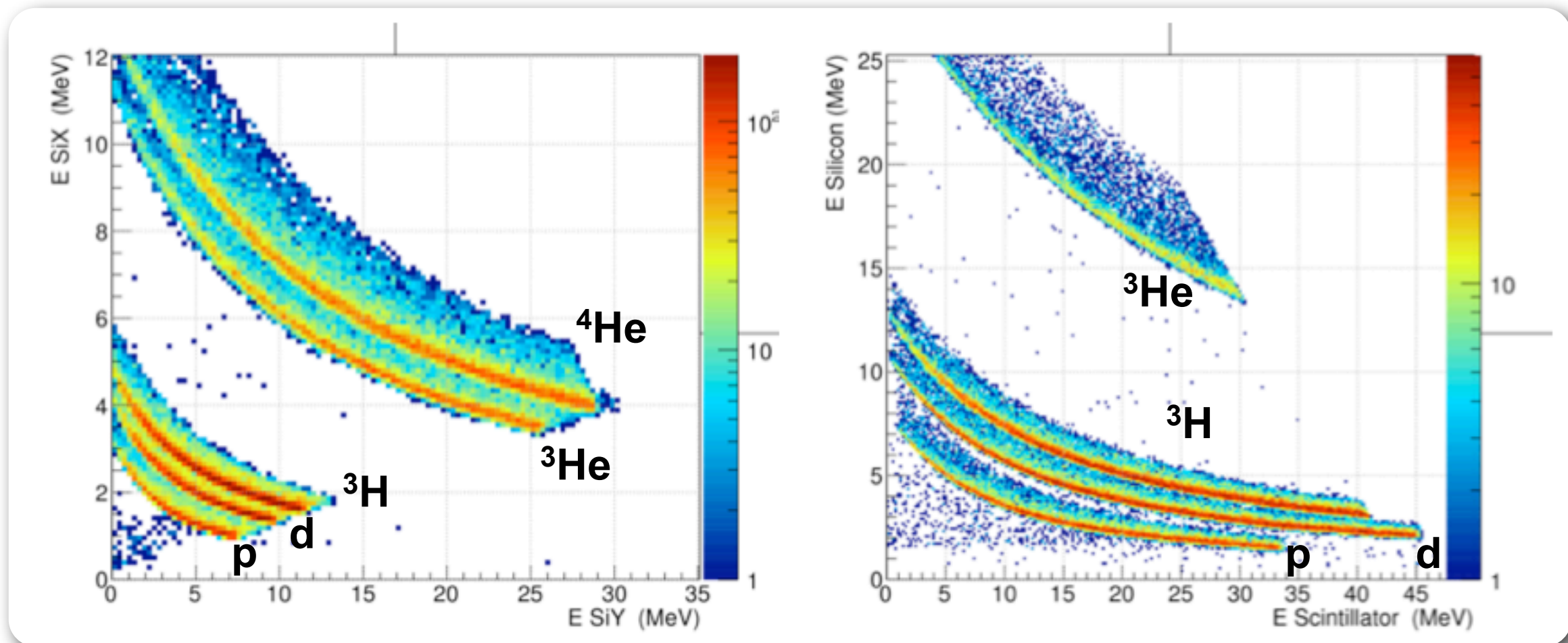


<b>Avalanche Chambers</b>	<b><math>r = 4 \text{ \&amp; } 8 \text{ cm}, \sigma \sim 100 \text{ ps/ } 200 \text{ }\mu\text{m}</math></b>
<b>Chamber Gas</b>	<b>2-10 Torr Isobutane</b>
<b>Chamber Windows</b>	<b><math>40\text{-}50 \text{ }\mu\text{g/cm}^2</math> Polypropylene</b>
<b>Silicon Micro-strips</b>	<b><math>r = 10 \text{ cm}, t \sim 100 \text{ \&amp; } 400 \text{ }\mu\text{m}, \sigma \sim 100\mu\text{m}</math></b>
<b>Scintillator</b>	<b><math>r = 15 \text{ cm}, t \sim 2 \text{ cm}</math></b>

Low Energy Recoil Detector under development  
by K. Hafidi (ANL), R. Dupré (Saclay) and S. Stepanyan (JLab)



# Detector Design



Momentum range  
of  $^3\text{H}$  80 - 550 MeV/c

× through tracks  
⊗ stopped tracks

Component				Particle Identification	Threshold (MeV/c)				
WC	SiX	SiY	SC		p	d	$^3\text{H}$	$^3\text{He}$	$^4\text{He}$
×	⊗			$\Delta E_{WC}, E_X, T$	55	75	95	150	180
×	×	⊗		$\Delta E_X, E_Y$	80	120	160	240	290
×	×	×	⊗	$\Delta E_{X+Y}, E_{SC}$	130	210	280	420	

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# Projected Rates

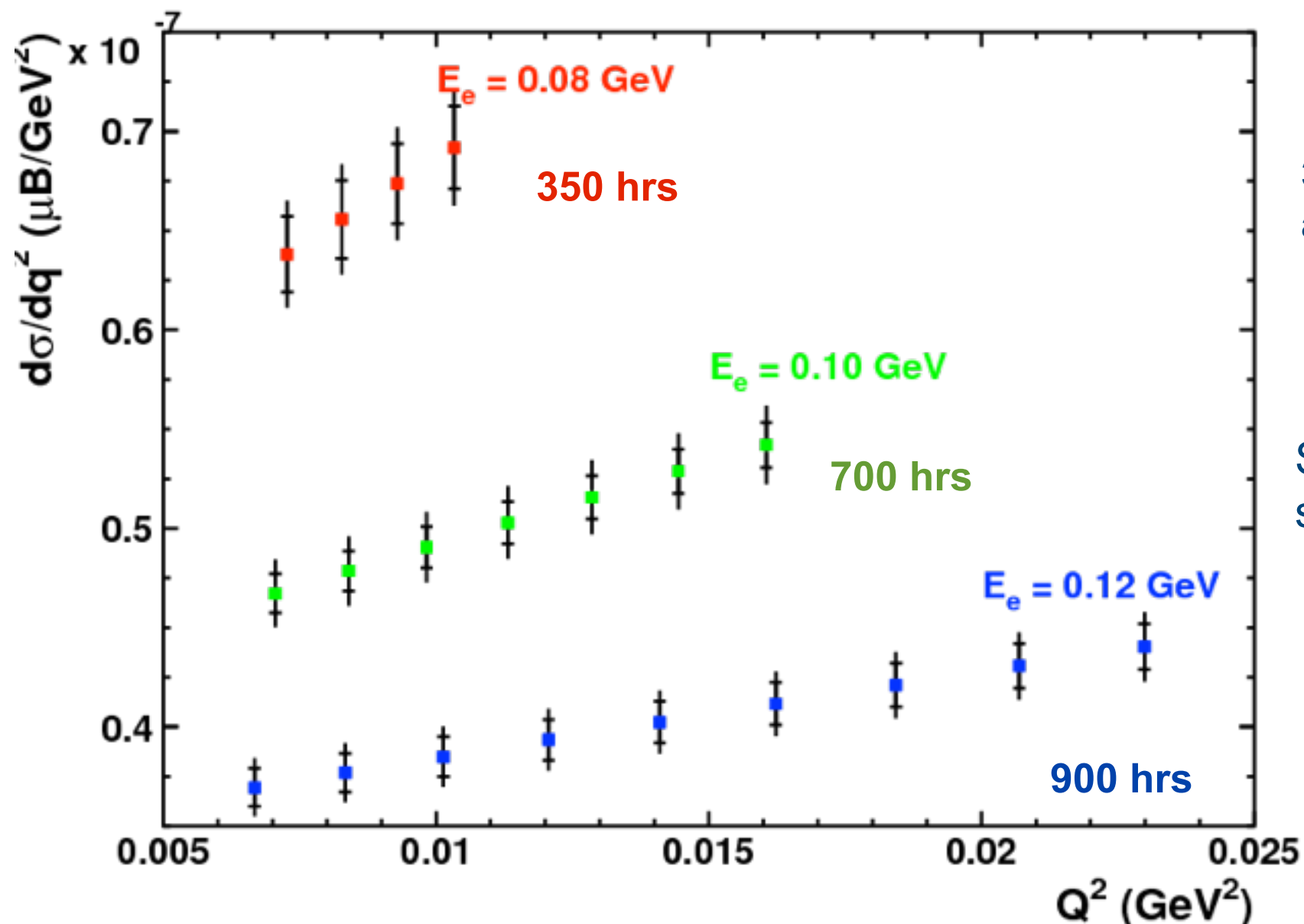
$I_e = 600 \mu\text{A}$ , polarized beam

Target density  $\sim 10^{19}$  atoms/cm<sup>2</sup>

Luminosity  $\sim 4 \times 10^{34}$  /cm<sup>2</sup>

Rates @  $E_e = 0.08$  GeV  $\sim 8 \times 10^{-4}$  /sec

$\Rightarrow \sim 3\%$  statistics in 350 hrs



3% systematic uncertainty  
assumed for all settings.

Stat. and total error  
shown for each point



# Summary of Experiment

1. **0.08 - 0.12 GeV**, highly polarized beam with **600  $\mu\text{A}$**  current
2. Windowless gas flow/jet target with ultra-pure  $^3\text{He}$
3. Solenoidal magnet (low energy/warm)
4. Large acceptance recoil detector with excellent PID  
detect recoil  $^3\text{H}$   
and elastically scattered electrons in coincidence (to reduce backgrounds)
5. Eliminate parity conserving EM backgrounds using  
(  $\sigma(h=+1) - \sigma(h=-1)$  )
6. Use electron scattering data on  $f_M^{^3\text{He}}(Q^2)$  to extract  $f_A^{^3\text{He}}(Q^2)$
7. Look for spin correlations in search of second class currents
8. Use spin correlations to measure  $f_A/f_V$  of  $^3\text{He}$



# Conclusions

- Weak capture of electrons on the proton and  $^3\text{He}$  can provide unique opportunities to measure the axial formfactor.
- Backgrounds can be controlled using polarized electrons and well designed detectors.
- Several ideal detectors and targets already under development for other experiments.
- At very low energies measurement of spin correlations can test for 2<sup>nd</sup> class currents.
- On protons one can measure  $|g_A/g_V|$  with completely different systematics compared to neutron beta decay.
- Weak capture of electrons would be complementary to the traditional neutrino scattering experiments and less resource intensive.
- Will help launch a new program to measure the charge changing weak current. The next reaction to attempt  $e + p \rightarrow \Lambda + \nu$   
(cross section 1/25 of capture cross section on p)