

Fundamental transport mechanisms, fabrication and potential applications of nanoporous atomically thin membranes

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Graphene and other two-dimensional materials offer a new approach to controlling mass transport at the nanoscale. These materials can sustain nanoscale pores in their rigid lattices and due to their minimum possible material thickness, high mechanical strength and chemical robustness, they could be used to address persistent challenges in membrane separations. Here we discuss theoretical and experimental developments in the emerging field of nanoporous atomically thin membranes, focusing on the fundamental mechanisms of gas- and liquid-phase transport, membrane fabrication techniques and advances towards practical application. We highlight potential functional characteristics of the membranes and discuss applications where they are expected to offer advantages. Finally, we outline the major scientific questions and technological challenges that need to be addressed to bridge the gap from theoretical simulations and proof-of-concept experiments to real-world applications.

Membranes are thin barriers that permit selective mass transport, and form the basis of a number of separation processes typically driven by gradients in pressure, electric potential, concentration or temperature^{1,2}. With the advantages of modularity, scalability, compactness and high energy efficiency, membrane separations have become pervasive³ in applications related to energy⁴, water⁵, food⁶, biotechnology⁷ and chemical processing³. Major applications include water desalination^{8–10}, natural gas purification¹¹, production of nitrogen from air¹, haemodialysis¹², bioprocessing⁷, solvent- and petrochemical-based separations^{13,14}, and production of ultrapure water¹⁵. Beyond separations, membranes find use in fuel cells⁴, drug delivery¹², bio/chemical sensors¹⁶ and energy harvesting from mixing processes⁵.

The performance of membrane-based systems (for example, energy efficiency, productivity or product recovery) depends on the constituent membrane characteristics, which place constraints on and define trade-offs in membrane process design. In addition to being economical and manufacturable, an ideal membrane should easily allow flow of the desired species (high permeance), reject undesired species (high selectivity) and exhibit robustness in operation (high chemical, mechanical and thermal stability, low fouling) (see Box 1).

Although remarkable progress has been achieved in membrane technology, persistent challenges remain, specifically (1) overcoming the trade-off between selectivity and permeability of the membrane material, (2) mitigation of fouling and (3) robust operation under harsh conditions. Specific examples include the need for chlorine-resistant membranes for water desalination^{5,10,17} and high-permeance membranes for production of oxygen from air¹. To address these challenges, the past few decades have seen the exploration of various membrane structures and materials, including novel polymers¹, inorganic membranes (zeolites¹⁸, silica, carbon¹⁹, ceramics), nanomaterials (metal-organic frameworks²⁰, carbon nanotubes²¹ (CNTs), carbon nanomembranes²²) and polymer–inorganic

‘mixed-matrix’ membranes^{1,19}, among others. These developments seek to achieve improved membrane performance by (1) the use of favourable membrane structures, for example, a thin selective layer, or (2) the use of materials with favourable properties, for example, antifouling properties and frictionless transport in CNTs²³, or high permeability and selectivity of materials with intrinsic porosity²⁴.

The rise of graphene and other atomically thin materials in the past decade²⁵ has opened new possibilities in membrane technology. The atomic thickness of these materials makes them the thinnest possible barrier²⁶, which, combined with their remarkable mechanical strength²⁷, chemical robustness²⁸ and ability to sustain selective, nanometre-scale pores²⁹, evokes the possibility — at least in principle — of realizing an ideal nanoporous atomically thin membrane that could be tailored to a range of applications while offering high permeance, high selectivity and high chemical resistance. This review focuses on nanoporous atomically thin membranes (NATMs)^{30–37} where transport is governed by flow across rigid pores in a continuous, atomically thin layer; for a discussion of membranes where other mechanisms influence transport, such as graphene oxide or carbon nanomembranes, the reader is directed to several excellent reviews^{22,32,38–41}.

Transport mechanisms

Membranes achieve selective transport through a variety of mechanisms operative over different length scales (Fig. 1a). At the smallest scale, dense polymeric membranes (without defined pores), such as reverse osmosis membranes for water desalination and many gas separation membranes, operate by a solution-diffusion mechanism^{1,42}. Here, selectivity results from differences in species solubilities and diffusivities in the membrane material; solubility depends on the molecular structure, membrane porosity and chemical affinity, whereas diffusivity is governed by thermally activated rearrangements of the polymer chains that strongly favour size-dependent diffusion of smaller molecules. When selectivity

Box 1 | Definitions, transport equations and length scales for NATMs.

Definitions of transport properties

$$\text{Permeance} \equiv \frac{\text{Mass, volume or molar flow rate across membrane}}{\text{Membrane area} \times \text{Pressure difference}}$$

where ‘pressure difference’ may also refer to a difference in chemical potential, concentration, electric potential and so on.

$$\text{Selectivity} \equiv \frac{\text{Permeance of species A}}{\text{Permeance of species B}}$$

$$\text{Rejection} \equiv 1 - \frac{\text{Solute concentration in permeate solution}}{\text{Solute concentration in feed solution}}$$

$$\text{Pore permeation coefficient} \equiv \frac{\text{Mass, volume or molar flow rate across pore}}{\text{Pressure difference}}$$

$$\text{Pore effective area} \equiv \frac{\text{Pore permeation coefficient}}{\text{Ideal gas flux (incident on a surface)}}$$

Some basic transport equations

Ideal gas flux.

$$J_{\text{ideal gas}} = \frac{\Delta P}{\sqrt{2\pi MR_u T}}$$

ΔP is the gas pressure, M is the molecular weight of the gas, R_u is the universal gas constant and T is the absolute temperature.

Permeation coefficient for viscous flow in pore.

$\Pi = \pi R^4 / \mu (3\pi R + 8L_p)$, where R is the pore radius, μ is the dynamic viscosity and L_p (0.535 nm for graphene) is the pore length (from ref. 102). Sampson’s model: $\Pi = R^3 / 3\mu$; obtained for $L_p = 0$.

Ion conductance of pore.

$$G = \sigma \left(\frac{L_p}{\pi R^2} + \frac{1}{2R} \right)^{-1}$$

σ is the bulk ionic conductivity. This relation is valid for neutral pores.

Length scales

Graphene. Carbon atom van der Waals diameter: $D_{\text{vdw}} = 3.4 \text{ \AA}$. Carbon bond length in graphene: $a = 1.42 \text{ \AA}$. Area of a hexagonal ring in graphene is 5.24 \AA^2 .

Pore geometry. D_c is pore diameter defined by centres of edge atoms. $D_p \equiv D_c - D_{\text{vdw}}/\sqrt{2}$ is the pore diameter used for gas transport. $D_{p,\text{vdw}} = D_c - D_{\text{vdw}}$ is the pore diameter adjusted for van der Waals size of pore edge atoms, used for water/ion transport.

Gases. Kinetic diameter (\AA): He (2.6), H_2O (2.65), H_2 (2.89), CO_2 (3.3), Ar (3.4), O_2 (3.46), N_2 (3.64), CH_4 (3.8), SF_6 (5.5).

Water molecule. Mean van der Waals diameter: 2.8 \AA .

Ions. Diameter of hydrated ion (\AA): Li^+ (7.64), Na^+ (7.16), K^+ (6.62), Mg^{2+} (8.56), Ca^{2+} (8.24), F^- (7.04), Cl^- (6.64), Br^- (6.60).

See Supplementary Section I for additional data, units and source references.

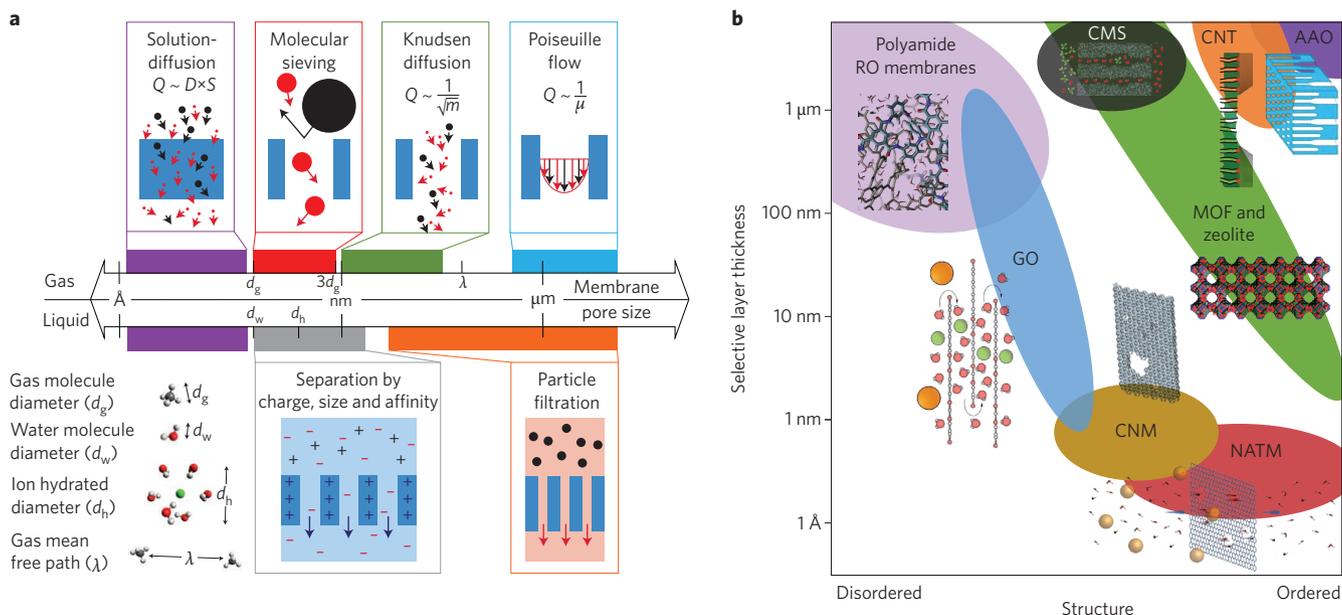


Figure 1 | Membrane characteristics and length scales. a, Length-scale dependence of membrane transport mechanisms. Relative scales of gas and water molecules, hydrated ions and gas mean free path are depicted on bottom left. Q , flux; D , diffusivity; S , sorption coefficient; m , molecular mass; μ , viscosity. **b**, Membrane structure–thickness map with some illustrative examples. From left to right, the structures change from disordered to ordered. Ideal NATMs are the thinnest with a high degree of order. Panel **b** images adapted from: ref. 193, Elsevier (polyamide reverse osmosis (RO) membranes); ref. 194, RSC (carbon molecular sieves (CMS)); ref. 21, AAAS (CNT); ref. 195, IOP (anodized aluminum oxide (AAO)); ref. 40, AAAS (graphene oxide (GO)); ref. 196, RSC (metal organic framework (MOF) and zeolite); ref. 197, American Chemical Society (carbon nanomembrane (CNM)); ref. 198, Macmillan Publishers Ltd (NATM).

is governed by diffusion alone (as in many polymeric gas separation membranes)⁴³, a more permeable material typically provides less selectivity and results in a trade-off between permeability and selectivity known as the Robeson limit in the context of gas separations^{43,44}. Overcoming this trade-off requires incorporation of additional mechanisms such as chemical affinity or molecular sieving, where smaller molecules pass through while larger ones are sterically impeded (for example, in high free-volume polymers or carbon molecular sieves)^{5,24,45}.

Due to the small free volume available, gas transport in membranes with pore diameters just beyond the molecular size is governed by phenomena such as diffusion, surface adsorption and condensation of gas molecules in the membrane pores^{24,43}. In pores that are much larger than molecular size but smaller than the gas mean free path, gas transport is governed by Knudsen diffusion^{43,46}, where molecules with lower molecular mass travel faster and have higher permeance. In liquid environments, transport in pores that are larger than molecules or ions is influenced by differences in species diffusivity, steric effects, chemical affinity and electrostatic interactions, including surface charge and dielectric effects⁴⁷.

Although transport in membranes is complex, a thin selective layer with precisely controlled pores and chemical functionality — which is potentially realizable with NATMs — is ideal for achieving high permeance and high selectivity (Fig. 1b). Realization of NATMs requires an understanding of gas- and liquid-phase transport across pores in atomically thin materials, which we now discuss.

Gas transport across atomically thin nanopores. In 2008, Bunch *et al.*²⁶ demonstrated that graphene is impermeable to helium and other gases, which opened the possibility of creating selectively permeable pores in graphene to realize selective membranes^{26,48}. In their experiment, exfoliated pristine graphene ‘nanoballoons’ suspended over pressurized microcavities retained gases to within detection limits (Fig. 2a). This impermeability was attributed to the fact that pristine graphene does not have gas-permeable defects, the space between the carbon atoms in graphene is too small to allow for transport of gases and the contribution of quantum tunnelling is negligible²⁶.

A large number of theoretical studies^{48–51} have explored the transport of different gases (for example, CH₄, CO₂, H₂, N₂, CO, H₂S, O₂, noble gases, alkanes and isotopes) through graphene and other NATMs with pore edges terminated with atoms including C, H, N, F and O (Fig. 2b) (see Supplementary Section II for a detailed list of references). The majority of these studies have focused on graphene and similar materials with inherent nanoporosity, such as porous graphene with various modifications^{50,52–56}, graphdiyne^{57,58}, graphyne⁵⁹ and two-dimensional (2D) polymers based on polyphenylene^{60–62}, porphyrin⁶³ and cyclohexa-*m*-phenylene⁶¹; exploration of other materials including hexagonal boron nitride (hBN)⁶⁴ has also begun.

These studies use one of two approaches to compute transport properties. (1) The energy barrier for a gas molecule to cross a pore is calculated using density functional theory (DFT)⁵¹ or other quantum mechanical methods^{50,62}; permeance is then estimated using a transition state approach given the gas molecule kinetic energy distribution⁵⁰, whereas selectivity is often estimated as the ratio of Arrhenius factors⁵¹. (2) The rate of molecules crossing the pore is calculated using classical (or first principles⁴⁸) molecular dynamics simulations, which directly yields the permeance^{49,56,65}. These studies have shed light on transport mechanisms^{66,67}, and revealed that permeance and selectivity can depend on differences in molecule size, mass, surface adsorption⁶⁸, interaction with functional groups on the pore rim, conformational entropy⁵⁴ and tunnelling rates^{62,69,70}.

To a good approximation, gas transport across atomically thin pores can be understood by comparing the diameters of the gas molecule (D_m) and the pore (D_p), though one needs to exercise caution as diameter definitions vary (Fig. 2d). When the pore is slightly larger

than the gas molecule, gas transport is similar to effusion, that is, molecular flow across a thin aperture smaller than the gas mean free path, where the pore presents an effective area $A_{\text{eff}} \approx \pi/4(D_p - D_m)^2$ (see Box 1) that is smaller than the pore area $A_{\text{pore}} = (\pi/4)D_p^2$. Here, transport is dominated by steric considerations. When the pore size is approximately equal to or slightly smaller than the size of the gas molecule, the pore presents an energetic barrier to transport that depends on the molecule–pore repulsive interactions and any compliance due to bond stretching or flexing. We refer to this as the activated regime. If we assume a rigid pore with fixed atoms and approximate the molecule–pore interaction by the Lennard-Jones potential, the barrier height is given by

$$E \approx \frac{\pi D_c}{a} 4\epsilon \left[\left(\frac{\sigma}{D_c/2} \right)^{12} - \left(\frac{\sigma}{D_c/2} \right)^6 \right] \approx \frac{4\pi\epsilon D_c}{a} \left(\frac{\sigma}{D_c/2} \right)^{12}$$

where ϵ and σ are the Lennard-Jones parameters, $D_c/2$ is the distance from the pore centre to the centre of the atoms constituting the pore, a is the distance between adjacent atoms on the pore rim and $\pi D_c/a$ gives the average number of atoms on the pore rim (Fig. 2d). Assuming that all gas molecules with kinetic energies exceeding this value cross the pore, we obtain an Arrhenius-like expression for the effective pore size

$$\frac{A_{\text{eff}}}{A_{\text{pore}}} \approx \frac{1}{2} \operatorname{erfc} \left(\sqrt{\frac{E}{k_B T}} \right)$$

where $k_B T$ is the Boltzmann factor (see Box 1 and Supplementary Section IV).

The effective pore areas predicted from simulations for different gas molecules across pores with different structures and functional groups, compiled from 33 literature reports, agree with this model and indeed collapse into two regimes (Fig. 2e)^{48–59,61,63–65,68,69,71–85}. The steric regime has a high A_{eff} and a high permeation coefficient, with $0.0025 \leq A_{\text{eff}}/A_{\text{pore}} < 1$ and $\pi/4(D_p - D_m)^2 \geq 0.24 \text{ \AA}^2$. A_{eff} and the permeation coefficient decrease rapidly in the activated regime as the pore size is reduced below the molecule size. For the majority (~75%) of simulated pores, effective pore areas and the corresponding permeation coefficients (see Box 1) of the most permeable molecule range from 10^{-4} – 10 \AA^2 and 10^{-23} – $10^{-19} \text{ mol s}^{-1} \text{ Pa}^{-1}$, respectively, which, for a pore density of 10^{12} cm^{-2} , corresponds to a membrane permeance of 10^{-7} – $10^{-3} \text{ mol m}^{-2} \text{ s}^{-1} \text{ Pa}^{-1}$ (10^2 – 10^6 GPU). For example, graphene pores with diameters of 3.6 and 4.8 Å were reported to present effective areas of 9.9×10^{-3} and 5.2 \AA^2 (effective diameters of 0.11 and 2.57 Å), permeation coefficients of 1.8×10^{-23} and $9.3 \times 10^{-21} \text{ mol s}^{-1} \text{ Pa}^{-1}$ and permeance of 5.3×10^2 and $2.8 \times 10^5 \text{ GPU}$, respectively, to hydrogen^{75,81} (kinetic diameter 2.89 Å) (see Supplementary Section I for equations relating these parameters). However, some simulations involve pore densities of up to $8.5 \times 10^{14} \text{ cm}^{-2}$, resulting in permeances up to $10^{-1} \text{ mol m}^{-2} \text{ s}^{-1} \text{ Pa}^{-1}$ (10^8 GPU).

Given two different gas molecules, there are three possible regimes depending on pore size (Fig. 2d). (1) For large pores, transport is in the steric regime for both molecules and high selectivity is not possible. (2) Selectivity by molecular sieving accompanied by high permeance occurs when the pore size lies in between the sizes of the two gas molecules; transport of the smaller molecule is in the steric regime while that of the larger one is in the activated regime. (3) High selectivity at very low permeance occurs for small pores when transport of both molecules is in the activated regime. Interestingly, this simple steric exclusion and activated transport model predicts a trade-off between permeance and selectivity, analogous to the Robeson limit, and is consistent with published reports from theoretical studies (Figs 2f and 5f); smaller pore sizes lead to higher selectivity and lower permeance due to the highly nonlinear dependence of the energy barrier on pore size.

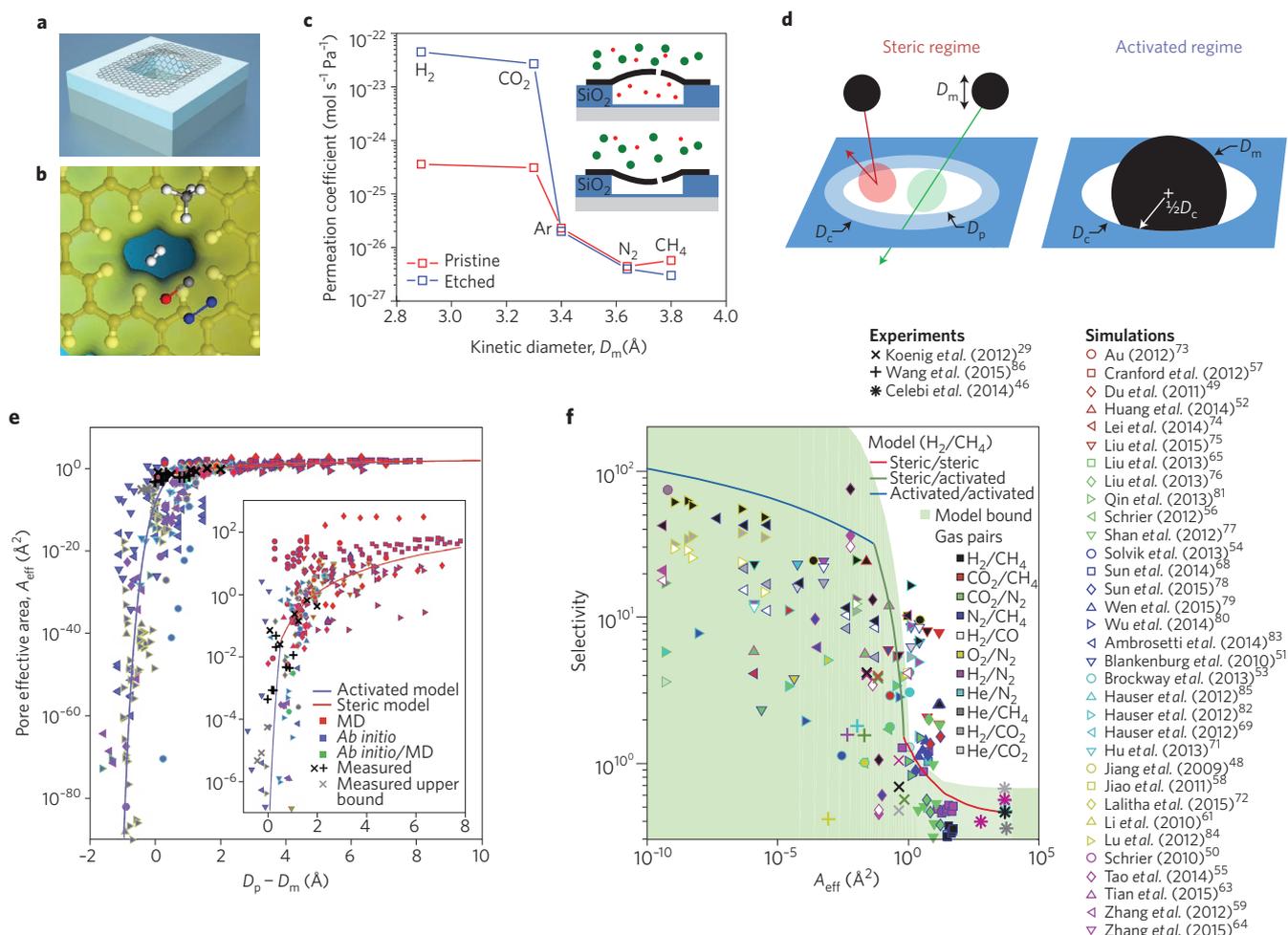


Figure 2 | Gas transport across NATM pores. **a**, Schematic of a graphene nanoballoon experiment²⁶ in which graphene suspended over a cavity retained pressurized helium. **b**, Visualization from simulations⁵⁵ of a hydrogen-functionalized graphene pore along with H₂, N₂, CO and CH₄ molecules. H, C, N and O atoms are shown in white, grey, blue and red, respectively, and the 0.02 Å⁻³ electron density isosurface is shown in yellow. **c**, Permeation coefficient of pristine and ultraviolet/ozone etched bilayer graphene measured in ref. 29 showing molecular sieving through a graphene nanopore, with H₂ and CO₂ preferentially effusing through the etched graphene. Inset: small gas molecules (in red) escape through the selective pore while big molecules (in green) cannot enter, leading to a measurable membrane deflection. **d**, Illustration of transport model and mechanism of steric/activated gas transport. Carbon centres located at D_c with pore diameter defined as D_p = D_c - D_{vdw}/√2, where D_{vdw} is the van der Waals diameter of a carbon atom. **e, f**, Compilation of simulation and experimental data of gas flow through NATM pores for different gas molecules and pore sizes qualitatively agrees with scaling model and illustrates the steric/activated transport regimes (see Supplementary Section III for data extraction and Supplementary Section IV for scaling model). Marker border colour indicates the source. **e**, Pore effective area decreases sharply when the pore size falls below the molecule size. Marker fill colour indicates the simulation method used. ‘Measured upper bound’ indicates an experimentally measured upper bound on permeance, where the flow rate through the graphene pore was below the measurement resolution. ‘Ab initio / MD’ indicates that the reported effective area is based on a permeance (for gas B), which was computed as the ratio of the permeance of a more permeable gas species (gas A), determined by classical molecular dynamics (MD), and the selectivity (A/B), determined by *ab initio* calculations. **f**, Selectivity versus effective pore area. For clarity, the curve predicted by the model is shown for H₂/CH₄ since curves for other gas pairs (shown in Fig. S3) are qualitatively similar. Marker fill colour identifies the gas pair. Shaded area bounds the model predictions for all gas pairs. Note the double-log scale on the y axis. Panels adapted from: **a**, ref. 26, American Chemical Society; **b**, ref. 55, American Chemical Society; **c**, ref. 29, Macmillan Publishers Ltd.

Molecular dynamics tends to better resolve high permeance, but has difficulty quantifying high selectivities ($\geq 10^3$) because of the limited number of observable molecule crossings imposed by a limited simulation time (≤ 100 ns). DFT-based approaches can resolve higher selectivities with reported values ranging up to 10^{78} , but some of these correspond to impractically low permeance.

Although the above model captures the basic dependence of permeance and selectivity on molecule size, other phenomena can determine selectivity when differences in size are not sufficiently large to be the determining factor. For example, permeance of gas molecules may be enhanced by adsorption and surface diffusion of the gas molecules towards the pore⁶⁷ or by interactions with

partial charges on the pore edge. In porous graphene, permeance is enhanced in the order (1) SO₂ (ref. 60) > CO₂ (refs 56,60,77) > CH₄ (refs 56,60,68) > H₂O (ref. 60) > N₂ (refs 49,56,60,68,77,79) / O₂ (refs 56,60) > H₂ (refs 49,68,79) > He (ref. 68) (permeance enhancements in the case of N₂ and O₂ are comparable); (2) H₂S (ref. 60) > CH₄; and (3) paraffins > olefins⁵⁴. Similar effects have been reported for variations of porous graphene^{54,56} and 2D polyphenylene⁶⁰, and are also implicated in enhanced CO₂ transport in graphene oxide membranes³². Entropic barriers arising from constrained conformations of molecules passing through the pore can lead to selective transport of shorter hydrocarbon molecules compared with longer ones⁵⁴. Quantum tunnelling can dominate

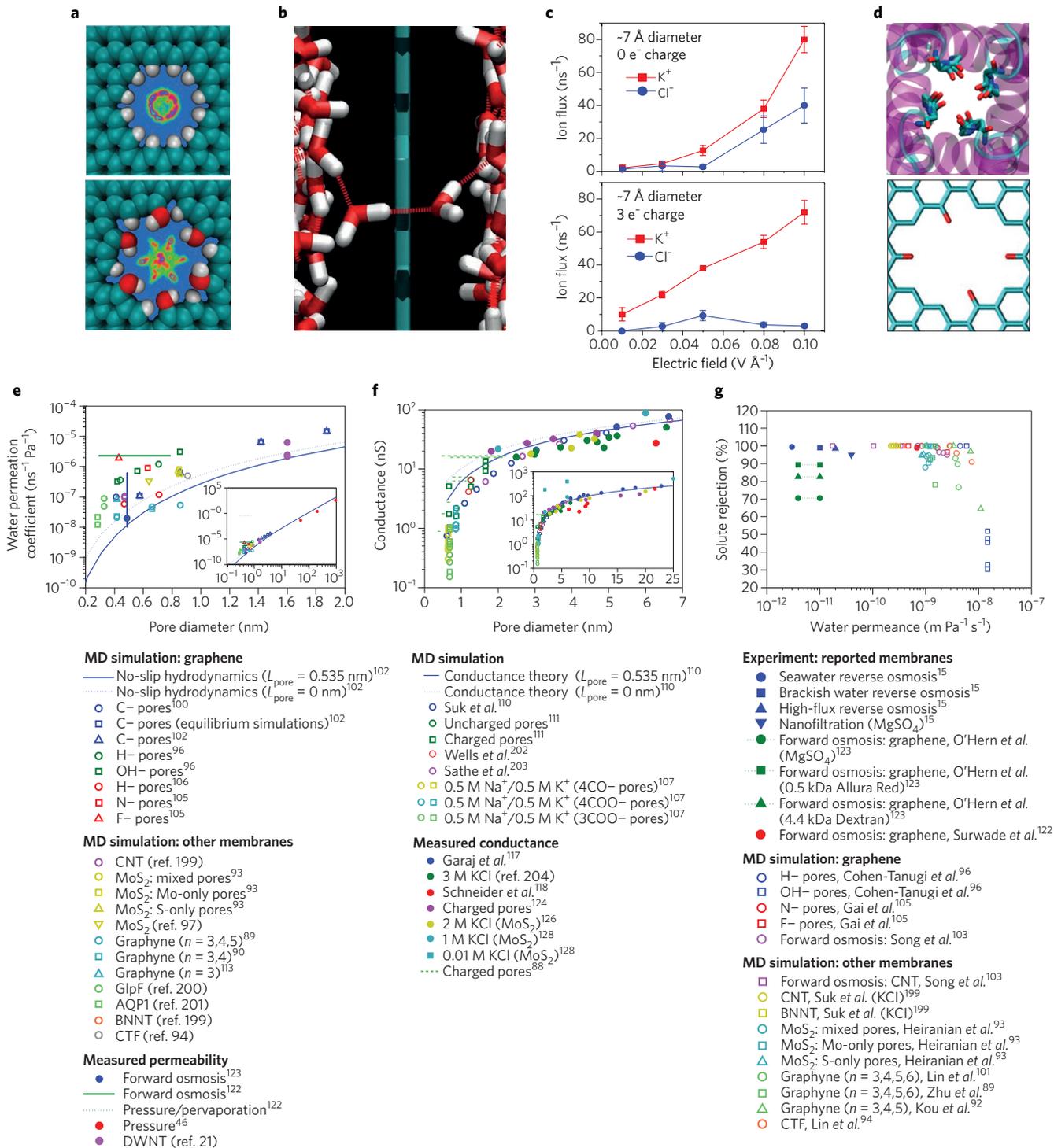


Figure 3 | Water and ion transport across NATM pores. **a**, Preferential configurations adopted by water molecules traversing hydrogen or hydroxyl terminated graphene pores⁹⁶ are illustrated by water-oxygen density maps. Red denotes high density. **b**, Single-file water transport in a graphene pore¹⁰⁰. **c**, Nonlinear increase in ion flux with electric field illustrates the ion dehydration barrier. Addition of a negative pore charge enhances the flux of K⁺ while suppressing Cl⁻ flux (from ref. 111). **d**, Illustration of a C=O terminated graphene pore designed to mimic a KcsA potassium ion channel (from ref. 107). **e**, Compiled data of simulated and experimentally measured water permeation coefficients show increase with pore diameter consistent with continuum theory (see Box 1). 3CO, 3COO, and 4COO denote pores terminated by 3 carbonyl and 3 or 4 carboxylate groups, respectively. **f**, Compiled data of simulated and experimentally measured ionic conductance for electrically driven transport of KCl compare well with continuum theory (see Box 1). Conductance is scaled to 1 M KCl (see Supplementary Section V). Estimated pore diameters are reported in refs 88,124. **g**, Compiled data of simulated and experimentally measured solute rejection and water permeance with respect to reported NATM membrane area, compared with that of commercial membranes from ref. 15. Solute is NaCl unless otherwise noted. In **e-g**, open markers are from molecular dynamics (MD) simulations, filled markers are for experimentally measured values, and horizontal solid or dashed lines indicate the estimated range for experimentally measured values. GlpF, aquaglyceroporin; AQP1, aquaporin 1; BNNT, boron nitride nanotube; CTF, covalent triazine framework; DWNT, double-walled carbon nanotube. Panels adapted from: **a**, ref. 96, American Chemical Society; **b**, ref. 100, American Chemical Society; **c**, ref. 111, AIP; **d**, ref. 107, American Chemical Society.

transport^{62,69,70} at low temperatures where the kinetic energy of gas molecules is much smaller than the energy barrier, which has been proposed for separating helium isotopes⁶⁹.

Experimental evidence of molecular sieving through NATMs was first reported in 2012 by Koenig *et al.*²⁹ (Fig. 2c). Gas permeance of initially impermeable bilayer graphene nanoballoon membranes was monitored as nanopores were created in the graphene by ultraviolet/ozone etching, where a single pore is likely to dominate transport⁸⁶. One graphene membrane permitted selective transport of H₂ and CO₂ over the larger Ar, N₂ and CH₄ molecules (~3.4 Å pore) (Fig. 2c) and another membrane permitted transport of the above molecules but not that of larger SF₆ (~4.9 Å pore), with observed selectivities exceeding 10,000 (ref. 29). In subsequent studies with a single-layer graphene membrane, permeance was observed to decrease with increasing kinetic diameter for He, Ne, H₂ and Ar, but was anomalously high for N₂O and CO₂ (ref. 86), consistent with flux enhancement due to polar interactions^{56,60,77}. The permeation coefficients in these experiments were in the range of 10⁻²³–10⁻²¹ mol s⁻¹ Pa⁻¹ (pore effective area 5 × 10⁻³–0.4 Å²) in reasonable agreement with theoretical studies (Fig. 2e,f).

Although experimentally measured gas permeance generally agreed with simulations, transport through single-layer graphene pores was observed to fluctuate on a timescale of minutes, which was postulated to originate from thermally activated rearrangement of molecular bonds at the pore^{86,87}. Similar observations in multiple pores⁸⁶ and in liquid environments⁸⁸ suggest that dynamic fluctuations in transport across NATM pores may be common, although understanding their origin remains elusive. This also highlights the need for more controlled experiments and more realistic simulations that complement those on idealized model pores. For example, the majority of simulations fix the 2D lattice or the position and orientation of the gas molecules in calculating the energy barrier, although some studies take these into account^{45,52,56,64,68,73,78,79,82,83}. DFT calculations that permit lattice deformation⁸² or molecular dynamics simulations using the adaptive intermolecular reactive empirical bond order (AIREBO) potential that allows for deformation and thermal fluctuations^{56,68,73,78,79} indicate that assuming a fixed lattice can greatly over-predict selectivities^{79,82}.

Gas transport in larger pores is less complicated, but permeance is still very high owing to the atomic thickness of the membrane. Celebi *et al.*⁴⁶ showed that graphene bilayer membranes with an array of 7.6-nm-diameter pores machined by a focused ion beam exhibited H₂/CO₂ selectivity consistent with effusion, but had more than three orders of magnitude higher permeance (~10⁻² mol m⁻² s⁻¹ Pa⁻¹, based on graphene area) than existing gas separation membranes with similar selectivity. Gas flow followed free-molecular effusion theory for small pores (≤50 nm) and a modified Sampson's model (Box 1) for large pores.

Water and ion transport across atomically thin nanopores.

Motivated by applications in water desalination and purification, theoretical studies using classical molecular dynamics (or other methods^{89–91}) have focused on transport of water and ions driven by pressure^{89,90,92–102}, osmosis^{103–106} and electric fields^{107–111} across graphene with different pore terminations (C (refs 95,100,102,110,111), H (refs 96,98,99,103,106,109,112), hydroxyl^{95,96,103,106}, carbonyl¹⁰⁷, carboxyl^{95,106,107}, amine⁹⁵, F (refs 99,103,105,109), N (refs 105,109), O (ref. 108)), variants of graphyne^{89–92,101,104,113,114}, MoS₂ (refs 93,97) and covalent organic frameworks⁹⁴.

Transport of water molecules (~2.8 Å van der Waals diameter) through NATM pores is largely determined by pore size. Graphyne-2 pores with ~2.9 Å diameter present a small barrier to water transport⁹¹; pores larger than ~3 Å (which corresponds to atom centre-to-centre diameters, D_c , of ~6 Å) easily permit transport of water^{95,103,105}. When the pore size is below ~2 nm, transport of water is also influenced by hydrogen bonding and

structuring of water molecules, which depends on the pore geometry and functional groups^{95,96,98,100,102,103,105,112,115}. The smallest pores in graphene^{95,100,102,103}, graphyne^{89,90,113} and MoS₂⁹³ that permit water transport can accommodate only a single water molecule in their cross-section and therefore exhibit single-file movement of water molecules. In pores with diameter ≤1.5 nm (ref. 112), water molecules adopt certain preferential configurations as they pass through the pores^{95,101,112} (Fig. 3a,b); similar effects are also observed in other nanoscale conduits such as CNTs¹¹⁶. In NATM pores, a significant fraction of the pressure drop occurs in the entrance and exit regions outside the pore¹⁰². As, in many cases, entrance and exit resistance dominates transport in CNTs, flow rates through NATM pores are comparable to those in CNTs of similar diameter^{100,102,113} (Fig. 3e). These flow rates are, in fact, predicted reasonably accurately (within an order of magnitude) by continuum no-slip hydrodynamics (see Box 1) for NATM pores¹⁰². Flow enhancement is observed in molybdenum-terminated MoS₂ pores that mimic the conical shape of biological aquaporin channels and reduce the flow resistance at the entrance and exit⁹³. The effect of pore functional groups on water flow tends to be modest, with hydrophilic groups (for example, -OH, Mo, -N) reported to enhance permeation of water by up to twofold^{93,94,96,99,105,106} by attracting water molecules to the pore, compared with hydrophobic groups (for example, -H), which can present an entropic barrier to transport⁹⁶.

Transport of hydrated ions across atomically thin pores is influenced by electrostatic interactions, coordination with functional groups or charges at pore edges, and steric exclusion of the ion hydration shell^{137,96,105,109–111} (Fig. 3c,d,f). The diameters of the first hydration shell for Na⁺, K⁺ and Cl⁻ are ~6.6–7.2 Å, which require pore diameters exceeding ~7 Å (D_c exceeding ~10 Å) for ions to pass through with an intact hydration shell⁹⁵. Neutrally charged pores smaller than the ion hydration size therefore present a barrier to ion transport that depends on the energy required for dehydration and the electric polarizability of the NATM material⁴⁷. For example, Zhao *et al.*¹¹¹ observed that 4-Å-diameter neutrally charged graphene pores completely excluded both K⁺ and Cl⁻, but 8 Å pores permitted passage of both ions. Ionic transport across larger, neutral pores is dominated by the 'access resistance' associated with the pore entrance and exit that is approximated well by continuum theory¹¹⁰ (Box 1 and Fig. 3f).

Pore functionalization can radically alter ionic transport, especially if the pore is smaller than the size of the hydrated ion. Charged or partially charged functional groups along the pore edge can lower the energy barrier for ions of opposite charge and increase the barrier for ions of like charge⁹⁵, leading to cation/anion selectivity. For example, Sint *et al.*¹⁰⁹ showed that a ~5-Å-diameter graphene pore could be made cation selective by functionalizing the pore with electronegative F or N atoms. Similarly, Zhao *et al.*¹¹¹ observed that negatively charged, 4-Å- and 8-Å-diameter graphene pores permitted selective transport of K⁺, but the ionic flux increased nonlinearly with voltage as higher voltages facilitated partial dehydration of the K⁺ ion (Fig. 3c). Counterintuitively, by partially shedding their hydration shells, larger ions with lower hydration energies traverse pores more easily than smaller ions with tightly bound hydration shells¹⁰⁹. Ion transport can also drive fluid flow under an applied potential difference¹¹¹, a phenomenon called electroosmosis that finds use in fluidic pumping. Functionalized atomically thin pores that structurally resemble biological ion channels (Fig. 3d) show some similar behaviours, such as voltage gating and cation/cation selectivity^{107,108}, although the pores are much thinner than their biological counterparts. This raises the intriguing question as to what extent atomically thin pores can be tailored to achieve high ion selectivity, voltage sensitivity and other functionalities of biological channels.

The smaller size of water molecules compared with that of hydrated ions enables their separation when NATM pores are large enough to pass water, but small enough to block hydrated ions (Fig. 3g). Cohen-Tanugi and Grossman⁹⁶ demonstrated rejection of Na⁺ and

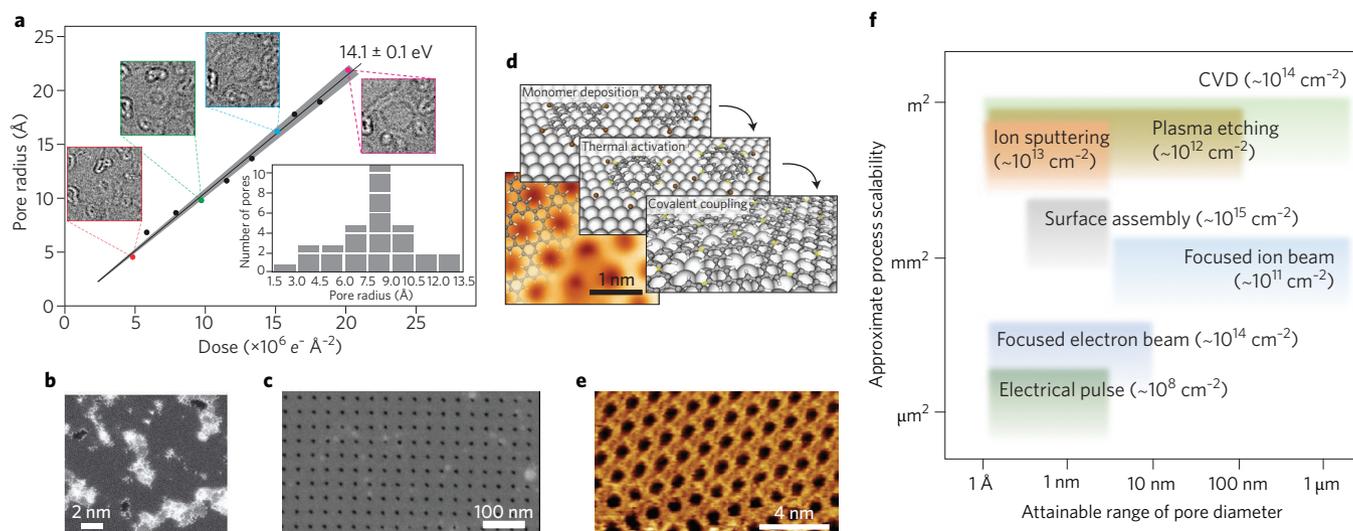


Figure 4 | NATM pore creation methods. **a**, Electron irradiation-mediated growth of pores in graphene nucleated by argon ion bombardment. Inset depicts a pore size distribution¹⁴⁵. **b, c**, Electron microscopy images of pores in single-layer graphene created by oxygen plasma¹²² and focused ion beam machining⁴⁶, respectively. **d, e**, Scanning transmission electron microscopy images and schematics of bottom-up synthesis of porous 2D polymers by self-assembly on a Ag(111) surface⁵¹ and at an air/water interface¹⁷⁰, respectively. **f**, Pore diameter versus approximate upper bound of membrane area that can be fabricated using different pore creation methods. Numbers in parentheses represent estimates of the highest possible pore density at the smallest possible pore size for each method. Panels adapted from: **a**, ref. 145, PNAS; **b**, ref. 122, Macmillan Publishers Ltd; **c**, ref. 46, AAAS; **e**, ref. 170, American Chemical Society. Panel **d** reproduced from ref. 51, Wiley.

Cl⁻ ions for hydrogen- and hydroxyl-terminated graphene pores with diameters below 5.5 Å, even at realistic operating pressures for reverse osmosis (10–100 bar)⁹⁸. Heiranian *et al.*⁹³ showed that MoS₂ pores below 6 Å could reject NaCl while supporting high water flux. Pressure-driven separation has also been shown for trihalomethanes in 6 Å graphene pores⁹⁹, NaCl in graphyne-3 (~4 Å)^{89,90,92,101}, NaCl in graphyne-4 (~6 Å) with a fixed lattice⁸⁹ (<100% rejection^{90,92,101,104} with a non-fixed lattice), CuSO₄, benzene and CCl₄ in graphyne-3 (ref. 101), and Mg²⁺, Ca²⁺, K⁺, Na⁺ and Cl⁻ in graphyne-3 and variants⁹⁰. Rejection of NaCl has also been demonstrated in forward osmosis in graphene^{103,105,106} and graphyne-3¹⁰⁴, consistent with selectivity governed by size effects. Compiled results indicate the theoretically predicted potential for high solute rejection and permeance of NATMs compared with existing membranes, although this performance remains to be experimentally realized (Fig. 3g).

A few experimental studies have probed transport across atomically thin pores^{46,117–129}, with particular focus on electrically driven ionic transport across pores in graphene^{117–119,124}, hBN¹²⁵ and MoS₂ (ref. 126), motivated by applications in DNA analysis. Ionic conductance of large pores (>2 nm) in graphene and MoS₂ agrees reasonably well with continuum descriptions of ion transport^{110,117,124,126,127}, although a surprisingly high K⁺/Cl⁻ selectivity was recently observed in large (20 nm) graphene pores, which was attributed to the surface charge on graphene¹²⁴. Experimental data on the behaviour of smaller pores (≤2 nm) is limited (Fig. 3f). Hints of ion transport through subnanometre pore defects^{109,110} smaller than the ion hydration shell were first seen in a study by Garaj *et al.*¹¹⁷, who observed that salts of larger monovalent ions with lower hydration energies displayed higher conductance than salts of smaller ions. Jain *et al.*⁸⁸ attempted to isolate pore defects by suspending graphene across a silicon nitride nanopore. They observed linear, rectified and nonlinear current–voltage characteristics consistent with ion dehydration and electrostatic effects in subnanometre pores. Their observation of cation/cation selectivity and voltage-activated fluctuations⁸⁸ also hints at parallels between atomically thin nanopores and biological ion channels^{107,108}. Nonlinear current–voltage characteristics, effects of charge quantization and high ionic current density driven by salt concentration

difference were recently observed in MoS₂ nanopores^{128,129}, which may find use in energy harvesting from salinity gradients.

Ionic selectivity in macroscale single-layer graphene membranes was reported by O’Hern *et al.*¹²¹, who introduced subnanometre pores in graphene by ion bombardment followed by chemical etching. Although there was significant leakage through defects, the membranes exhibited some K⁺/Cl⁻ selectivity consistent with electrostatic exclusion of Cl⁻ due to negatively charged pore functional groups¹¹¹ expected in oxidative etching; with further etching, the membrane permitted selective transport of KCl over a larger organic molecule (~1.0 nm size). Osmosis-driven water flux measured across similar graphene membranes with ~0.5 nm pores¹²³ was consistent with molecular dynamics predictions by Suk and Aluru¹⁰². Water/ion selectivity was demonstrated by Surwade *et al.*¹²² using oxygen plasma to introduce ~1 nm pores at a density of ~10¹² cm⁻² in single-layer graphene placed on a micrometre-scale aperture. For a certain plasma treatment, all NaCl ions were rejected and water permeated at high rates, though only one side of the membrane was wetted and evaporation may have played a role. Water transport has also been measured across larger (50–1,000 nm) pores in bilayer graphene, with results consistent with continuum theory⁴⁶. Although a picture of water and ion transport across atomically thin pores is beginning to emerge, experimental data are sparse, there are some discrepancies among the simulation results and ion selectivity remains to be understood (Fig. 3f).

Compared with water or ions, protons show distinct transport behaviours. Geim and co-workers discovered that monolayers of (nonporous) pristine graphene and hBN are permeable to protons¹³⁰ and have high selectivity (~10) over deuterons¹³¹, opening new possibilities for deuterium/hydrogen separation. The observation that small, otherwise impermeable defects in graphene could also permit selective proton transport^{132,133} suggests the possibility of creating high-permeance proton-selective NATMs from materials such as graphyne-1.

Fabrication of NATMs

Practical realization of NATMs requires synthesis of a continuous layer of atomically thin material with controlled porosity as well as

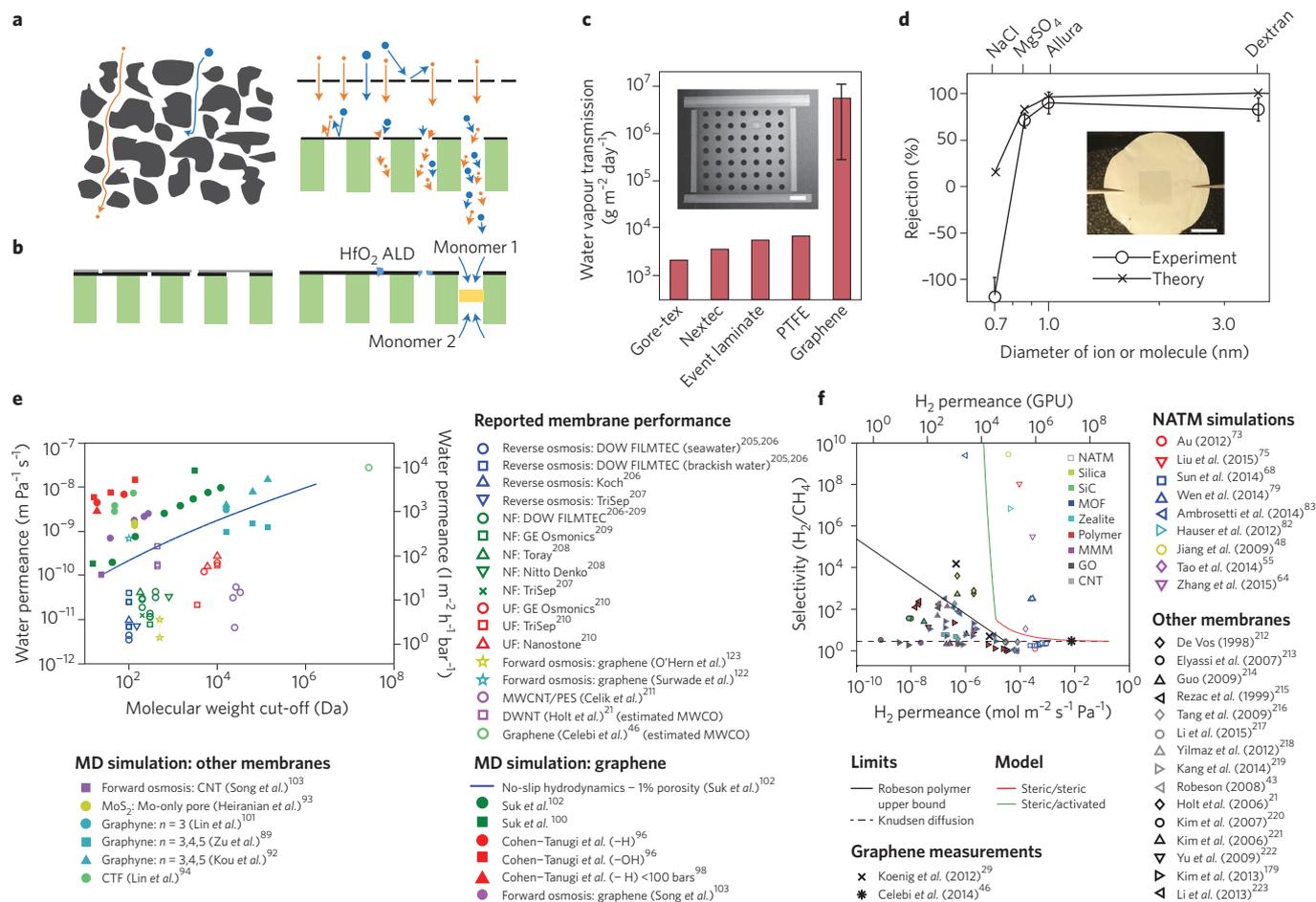


Figure 5 | Transport characteristics of NATMs. **a**, Pathways for transport of molecules of different sizes across membranes with torturous pores (left) and unsupported and supported NATMs (right). **b**, Schematics of defect sealing in NATMs. Left: stacking of multiple layers. Right: sealing of small defects using atomic layer deposition (ALD) of hafnia, and of large defects by reaction of monomers that form plugs only at defects¹²³. **c**, Fast water vapour permeation across NATMs with 400 nm pores formed by focused ion beam milling in two-layer graphene⁴⁶. PTFE, Polytetrafluoroethylene. The inset shows a membrane comprising freestanding circular areas of graphene into which pores are drilled. Scale bar, 10 μm. **d**, Forward osmosis-driven nanofiltration across single-layer graphene membranes (inset) with sealed defects. Results are compared with continuum theory assuming only size exclusion¹²³. Scale bar, 0.5 cm. **e**, Compilation of simulation and experimental water permeance data versus estimated molecular weight cut-off across NATMs compared with a few other illustrative membranes. **f**, Compilation of simulation and experimental selectivity–permeance data for H₂/CH₄ separation across NATMs compared with a few other membranes. In **e,f**, permeability data for membrane materials were converted to permeance assuming a 100-nm-thick selective layer. NATMs without inherent porosity, such as graphene, were assumed to have a pore density of 10¹² cm⁻². NATMs with inherent porosity had pore densities of up to 8.5 × 10¹⁴ cm⁻². CTF, covalent triazine framework; NF, nanofiltration; UF, ultrafiltration; MWCNT/PES, multi-walled carbon nanotube/polyethersulfone; DWNT, double-walled carbon nanotube; MWCO, molecular weight cut-off; MMM, mixed-matrix membrane; GO, graphene oxide. Panel **c** reproduced from ref. 46, AAAS. Panel **d** adapted from ref. 123, American Chemical Society.

the ability to handle the materials using suitable porous supports. NATMs can be fabricated ‘top-down’ by creating pores in an initially nonporous atomically thin material, or by ‘bottom-up’ synthesis of an intrinsically porous material.

Top-down fabrication. Chemical vapour deposition (CVD) has emerged as a tunable and versatile method for producing continuous layers of 2D materials over large areas¹³⁴ and lends itself easily to fabrication of NATMs. Within a year of the first demonstration of large-area CVD growth of monolayer graphene on copper by Li *et al.*¹³⁵ in 2009, Samsung demonstrated roll-to-roll transfer of 30-inch sheets of CVD graphene¹³⁶, and SONY later demonstrated roll-to-roll synthesis and transfer to produce 100-m-long graphene-coated films¹³⁷. Several methods are now available to handle and transfer graphene and other 2D materials to various substrates¹³⁸, and membranes have been made without transfer by introducing pores in the catalytic substrate

before¹³⁹ or after¹⁴⁰ CVD growth. Besides CVD methods, mechanical exfoliation of 2D layered materials produces thin, pristine flakes that are ideally suited for physics experiments or microscale membranes²⁹, whereas liquid-phase exfoliated 2D materials¹³⁴ are less suitable due to the need to assemble and fill in gaps between flakes. Experimental work on NATMs has therefore largely focused on exfoliated and CVD graphene, for which technical know-how is most advanced.

Subnanometre pores can be introduced in the initially nonporous material by ‘top-down’ approaches using ion irradiation^{121,141–147} or chemical/plasma etching^{29,86,122,148–152} (Fig. 4a–c). Etchant chemistry dictates the pore functional groups and can stabilize the pores¹⁵³, whereas pore density and size depend on the interplay between nucleation of defects and their growth into larger pores¹⁵⁴. For example, thermal oxidation of graphene or graphite in oxygen produces a low (<10¹⁰ cm⁻²) density of large (>10 nm) holes¹⁵⁵, while ozone tends to easily nucleate defects to yield a high density of pores^{154,156}

and was used for creating gas-selective pores in graphene nanoballoons^{29,86}. Oxygen plasma can also induce^{148,149} and grow nanometre-scale pores in suspended single-layer graphene^{122,151} at a high density ($\sim 10^{12}$ cm⁻², Fig. 4b). Nucleation of defects followed by growth can produce more uniform pores with controllable size and density¹⁴⁶. This was realized in graphene using oxygen and hydrogen plasma, respectively, to nucleate and grow pores¹⁴⁹. Ion irradiation can also controllably nucleate defects, which can be followed by electron-induced sputtering or chemical etching to create a high density ($\sim 10^{12}$ – 10^{13} cm⁻²) of subnanometre pores^{121,123,145} (Fig. 4a). Ion angle of incidence, energy and type determine the kind of defect, whereas the ion fluence determines their density^{141,142}. An energy threshold of ~ 30 eV is required to produce defects in graphene¹⁵⁷, indicating that ion impact could also play a role in plasma treatment¹⁴⁶. Whereas these methods produce a distribution of pore sizes^{121,145} and perhaps work well only in single-layer materials, focused ion beams can precisely machine a large number of individual pores down to ~ 8 nm in diameter⁴⁶ (Fig. 4c). Electron beams focused below 1 nm can create precise subnanometre to few-nanometre pores^{117,158,159}. Other methods to create pores include the use of block copolymers¹⁶⁰ or substrates¹⁶¹ as templates, reduction of graphene oxide¹⁶², catalyst-induced etching^{163,164}, electrochemical machining by atomic force microscopy⁸⁶, or the use of electrical pulses to create pores with subnanometre accuracy in graphene¹²⁷ or MoS₂ (ref. 165).

Bottom-up fabrication. In contrast to the creation of pores in an initially nonporous material, recent advances in the synthesis of intrinsically porous covalently cross-linked single or few-layer 2D polymers^{166,167}, including graphyne and its variants¹⁶⁸, have potential for creating a high density of atomically precise pores tailor-made for specific applications, directly in a covalently bonded single layer^{51,59,90,94,101,113}. For example, surface-assisted synthesis of nanoporous graphene by aryl–aryl coupling of polyphenylene-based precursors or cyclodehydrogenation using polycyclic aromatic hydrocarbons has been proposed⁵¹ (Fig. 4d). Graphdiyne has been synthesized at the millimetre scale by cross-linking of hexaethynylbenzene on copper under nitrogen atmosphere¹⁶⁹, which may permit selective transport of hydrogen⁵⁹, and possibly water⁹¹ or protons. Recently, a 2D polymer with ~ 0.8 nm pores was synthesized by ultraviolet cross-linking of triptycene-based amphiphilic monomers self-assembled at the water/air interface in a Langmuir–Blodgett trough¹⁷⁰ (Fig. 4e). The high quality of the resulting polymer and the ability to transfer centimetre-scale areas to porous supports makes it a promising route to realize NATMs.

Another potential route to NATM fabrication is by modifying synthesis processes that normally result in nonporous materials. Pores could potentially be created in a modified CVD process by controlling the ratio of the growth precursor and etchant (for example, CH₄ and H₂ for graphene)¹⁷¹, introducing etchants during or post synthesis, templating of the catalytic substrate¹⁶¹ or substitutional doping to generate precisely defined pores¹⁷². However, atomic precision and high porosity may be difficult to realize with these approaches.

Scalability, quality and stability. Among the top-down methods, pore creation in CVD-grown materials by chemical and plasma routes is amenable to fabrication of large-area NATMs (Fig. 4f). Machining using focused ion beams is less scalable and appropriate for smaller membranes, whereas electrochemical methods and machining using atomic force microscopy or tightly focused electron beams are suitable for creation of a few pores for microscale membranes. Bottom-up approaches remain relatively unexplored for membrane applications, and key challenges are to engineer practical processes for synthesis of continuous layers over sufficiently large areas¹⁶⁷. As of now, roll-to-roll processed 2D materials are not commercially available, and CVD-grown materials remain inherently polycrystalline with intrinsic vacancy defects, grain boundaries and wrinkles that contribute to

leakage pathways^{120,173}. Compared with applications in electronics, the quality requirements for atomically thin materials for membrane applications are more stringent in some respects, as small pinhole defects can severely compromise selectivity. To address these issues, synthesis processes need to be tailored to minimize pinhole defects, and methods to assess quality for membrane applications need to be developed. In addition, polymer-based methods to transfer atomically thin materials tend to leave surface contamination, the effects of which on transport across the membrane are poorly understood. Techniques for clean transfer and control of surface contamination are also essential to ensure quality and reproducibility.

Mechanical stability of NATMs is essential to withstand handling, abrasion and high pressures. Graphene has a remarkably high fracture strength of 130 GPa (~ 44 N m⁻¹) and tolerates strains of $\sim 25\%$ (ref. 27), which would facilitate its handling on a support. Although pores and defects can decrease the fracture strength by about an order of magnitude^{151,174}, it is more than sufficient to withstand high pressures when the atomically thin layer is properly supported. The ability to withstand pressure scales inversely with the support pore diameter, and molecular dynamics simulations indicate that nanoporous graphene can withstand a pressure of 570 bar when suspended over 1 μ m pores¹⁷⁵. Sufficiently high strengths are also expected of other atomically thin materials including MoS₂ (ref. 176), graphyne^{89,101} and covalent organic frameworks⁹⁴. Adhesion and abrasion resistance are other important considerations related to slippage, delamination and wear of NATMs. Graphene exhibits a high adhesion energy¹⁷⁷ and good abrasion resistance¹⁷⁸, but abrasion resistance in the context of membranes is not well understood and protective coatings will be required if the atomically thin layer is damaged easily. Similarly, although NATMs are expected to be chemically stable^{28,152}, studies that specifically investigate this aspect are needed.

From pores to membranes. Provided that transport across a pore is not influenced by its neighbours, the net flux across NATMs is the sum of those across the membrane pores. Using this assumption, the predicted permeance ranges from 10⁴ to 10⁸ GPU and exceeds 1,000 l m⁻² h⁻¹ bar⁻¹ for gas separations and water desalination, respectively, significantly surpassing the permeance of polymeric membranes (Fig. 5e,f). These estimates constitute upper bounds, since real membranes will inevitably have a distribution of pore sizes and defects that are situated in parallel with each other (Fig. 5a). NATMs are therefore extremely sensitive to leakage through defects, and their size selectivity is determined by the largest pores¹²³ — presenting a major engineering challenge that may well turn out to be the Achilles heel for some applications. It is for this reason that only a few studies have demonstrated selective transport across non-microscopic NATMs, which requires strategies to minimize the impact of defects by independent stacking of layers^{46,173}, sealing of defects¹²³, appropriate choice of the porous support¹⁷³ or other mechanisms¹³¹ (Fig. 5b).

Celebi *et al.*⁴⁶ stacked two layers of graphene on a silicon support to form a nearly impermeable layer, into which ordered arrays of pores with diameters of 8–1,000 nm were machined using a focused ion beam. The measured water flow rates (per unit graphene area) were five- to sevenfold greater than commercial ultrafiltration membranes, whereas vapour transport rates were two to three orders of magnitude higher than those in commercial membranes (Fig. 5c). Boutilier *et al.*¹⁷³ showed that gas transport through stacked graphene layers can be explained by random alignment of defects, resulting in exponential decrease in gas permeance with selectivity approaching the Knudsen limit. Kim *et al.*¹⁷⁹ stacked five layers of graphene on a polymeric membrane to enhance its barrier properties. The O₂/N₂ selectivity exceeded the Knudsen limit, and was attributed to defects and inter-layer transport. Although there are no reports of creating aligned subnanometre pores in multilayer materials, the stacking approach may be useful for materials like graphyne that have high intrinsic porosity.

O'Hern *et al.*¹²³ created centimetre-scale single-layer graphene NATMs by multiscale sealing of defects. Nanoscale defects were sealed by atomic layer deposition, which tends to deposit preferentially on defects and wrinkles, whereas the impermeability of graphene was exploited to seal larger defects using interfacial polymerization, where monomers introduced on opposite sides of the graphene layer react to form polymer plugs only where graphene is missing (Fig. 5b). After introducing subnanometre pores by ion irradiation and etching, the resulting membrane exhibited water permeance based on graphene area comparable to polyamide reverse osmosis membranes, and rejected organic molecules (≥ 1 nm) and divalent MgSO_4 (but not NaCl) under osmotically driven flow (Fig. 5d).

Beyond these considerations, an appropriate support layer is essential to achieve high selectivity in the presence of imperfections and defects¹⁷³. An ideal support — or coating — will add a transport resistance in series to both selective pores and defects, thereby limiting leakage without adversely impacting the permeance. This resistive layer must provide parallel transport pathways or be thinner than the spacing between non-selective defects in the atomically thin layer¹⁷³. Such an approach was critical to early commercialization of gas separation membranes¹¹. Beyond its role in minimizing leakage, the support must have high surface porosity to make use of most of the NATM area, have high mechanical strength, good chemical resistance, higher permeance than the atomically thin layer, and provide stable adhesion and facilitate transfer or coating of the atomically thin layer. Most experimental studies have used specialized supports (for example, polycarbonate track-etched membranes¹²³ or microfabricated silicon supports⁴⁶) with low surface porosity, which results in high permeance with respect to graphene area but low permeance based on total membrane area, and are suboptimal for practical separations. Although graphene NATMs have been fabricated on more practical supports such as poly(1-methylsilyl-1-propyne)¹⁷⁹, polypropylene¹⁸⁰ and polyvinylidene difluoride¹⁸⁰, meeting all of the above requirements is challenging, especially for pressure-driven separations.

Potential applications of NATMs

The path for new membrane technologies to advance beyond the laboratory is tortuous¹. Membranes that are difficult to scale up have found use in research and analytical applications¹⁸¹, whereas polyamide reverse osmosis membranes have grown to become the dominant desalination technology today. Although NATMs are still in the early stages of development, they present characteristics that are potentially advantageous for addressing persistent challenges in membrane separations.

A key advantage of NATMs is their high permeance that could enable higher energy efficiency and compact, high-productivity membrane separation systems. For seawater and brackish water, respectively, increasing the permeance of desalination membranes threefold could reduce the membrane area by 44% and 63% or reduce the energy consumption by 15% and 46% (ref. 182); the smaller gains for seawater are due to its higher osmotic pressure and build-up of the rejected salt at the membrane. More substantial gains in energy consumption and productivity are expected for applications involving low solute concentrations (for example, reverse osmosis treatment for drinking water, ultrapure water, and nanofiltration for water, pharmaceutical, food and beverage, and biotechnology industries^{13,183}), provided that fouling is not enhanced at higher fluxes¹⁸⁴. The combination of high selectivity and permeance is also desirable for a number of gas separations¹⁹, such as O_2 or N_2 from air, CO_2/CH_4 for natural gas sweetening and $\text{H}_2\text{O}/\text{air}$ for dehumidification¹⁸⁵.

A second advantage of NATMs is their potential for increased robustness due to high chemical resistance and mechanical strength, which could prove useful for operation under harsh conditions. For example, NATMs may be able to better withstand chlorine and permit desalination of high-salinity water by high-pressure reverse osmosis, or allow for more aggressive cleaning procedures¹⁸⁴ to maintain

membrane performance and extend their useful life. NATMs may help to address the challenges of aging, compaction and influence of solvents in organic solvent nanofiltration¹³, and potentially open new possibilities in the separation of liquefied gases and fuels.

Third, the molecular-sieving mechanism of transport is expected to lead to rejection of all species that significantly exceed the pore size^{99,101}. If NATMs could provide universal size-based rejection regardless of the solute or the fluid in which separation is performed, it would benefit applications involving a diversity of solutes such as in removal of persistent organic pollutants^{8,10} for water reclamation, removal of boron from seawater¹⁷ or in chemical processing^{3,13}.

However, actual application is contingent on commensurate advances in technology and practical membranes meeting these expectations⁴¹. The predicted performance and most of the proposed applications have yet to be realized even at the lab scale. Significant advances including ensuring a tight pore size distribution, design of appropriate supports and minimizing leakage are prerequisite for proof-of-concept studies. Manufacture at appropriate scale and cost, packaging into modules that minimize concentration polarization¹⁸⁶, characterization of fouling and chemical resistance, and demonstration of long-term performance under realistic conditions are essential⁴¹ for large-scale applications such as water desalination and natural gas separations. Small-scale applications such as laboratory and analytical separations or nanofiltration — which have less stringent demands on selectivity but would benefit from the high permeance and chemical stability of NATMs — are likely to emerge first.

Beyond separations, NATMs could find use in a variety of other applications such as in fuel cells^{130,132,133}, microfluidics¹¹¹, sensing¹⁶, surface micropatterning¹⁸⁷, energy harvesting¹²⁹ and a variety of biomedical applications, such as drug release, biosensing and immunoisolation¹². Nanopores in NATMs may find use in detection and analysis of single molecules, ions and binding events, or interrogating discrete steps in chemical reactions¹⁶⁵. In these applications, NATMs have the potential to offer higher selectivity, sensitivity, faster time-to-result or ease of fabrication. The more exotic capabilities of NATMs, such as isotope separation by quantum effects¹³¹ and perhaps undiscovered phenomena that do not have parallels in other membranes, could potentially open new applications. Finally, NATMs can provide fundamental insights into nanofluidic transport mechanisms that are relevant to atomically thin materials for barrier applications¹⁸⁸, for electron microscopy¹⁸⁹ and in graphene oxide membranes³².

Future perspectives

Regardless of the extent and timeframe in which NATMs will have an impact, they represent a unique class of membrane with tremendous potential for advancing membrane technology and are likely to persist far into the future.

Although theoretical studies have shed light on transport mechanisms in idealized pores, an understanding of the structure and behaviour of pores in real materials has only just started to emerge. Controlled experiments at the single pore level, complemented by theoretical studies on realistic pores, are required to probe the effects of different pore structures, functional groups, dynamic behaviours and the role of surface contaminants. Simulations are needed to advance pore creation techniques and explore different applications of NATMs, which can guide experimental developments. For example, are there pore structures that enable selective transport of water over boron, urea and other species that are difficult to reject using state-of-the-art membranes? Can NATMs separate water from ethanol? How effective are NATMs for separations in organic solvents?

Realization of practical membranes requires advances on four fronts. First, improved methods are needed to create a high density of uniform pores that will enable separations such as water desalination. Here, key challenges are achieving sufficiently precise pores using top-down techniques and scale-up of bottom-up synthesis methods. Second, better understanding of the origin of defects that

may arise during synthesis, handling or creation of selective pores, methods to limit leakage through defects, and quality control during synthesis are indispensable for realization of functional membranes. Theoretical studies are needed that shed light on how the distribution of pore size and defects as well as the membrane support structure determine the actual permeance and selectivity. Third, design or selection of an appropriate porous support layer and development of methods to promote adhesion and to coat the atomically thin layer on the support are essential for full utilization of the properties of NATMs. Fourth, packaging of membranes, design of modules tailored to NATMs, experimental and theoretical studies to understand fouling, chemical resistance, and potential degradation mechanisms, and cost-effective scale-up of membrane manufacturing are critical for advancing NATMs towards real-world applications.

Finally, it is likely that there will be opportunities to use NATMs in novel formats or in novel devices. For example, NATMs could act as porous scaffolds to realize novel membranes with improved performance by coating¹⁹⁰ or functionalizing¹⁹¹ them with ultrathin layers of materials¹⁹², perhaps by grafting polymers or materials that are difficult to cast into thin membranes. The fact that it is possible to make NATMs from conducting as well as insulating atomically thin materials raises the possibility of electrically actuated membranes that can actively modulate transport or pump ions⁹. Other intriguing possibilities arise due to the atomically thin nature that enables reactions across pores in NATMs — which could find use in selectively sealing leaky pores, functionalizing the membranes or performing chemical reactions across NATMs where pores control the stereochemistry.

Although significant technological challenges remain for practical applications, NATMs have provided us with the fascinating opportunity to explore fundamental questions in fluid flow at the smallest possible length scale. It is quite likely that these fundamental studies will continue to provide new — and perhaps surprising — insights into mass transport at the nanoscale. On the practical side, NATMs have tremendous potential, but also face considerable challenges that need to be recognized and addressed if the field is to move ahead. Only time will reveal whether NATMs live up to their promise.

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Additional information

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Competing financial interests

R.K. is a co-founder and has equity in a start-up company aimed at commercializing graphene membranes.