

③ Capillary rise:

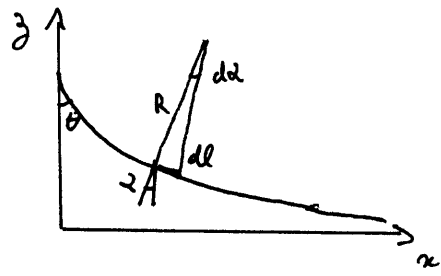
3.1 Menisci: (from greek 'me'ne' = moon).



non-dimensional analysis: $h = l_c f(\theta)$

$90^\circ > \theta > 0 \Rightarrow f(\theta) > 0$ (up)

$\theta > 90^\circ \Rightarrow f(\theta) < 0$ (down)



Laplace pressure = gravity.

$P_0 - \rho g z = P_0 - \frac{\sigma}{R}$

$\Rightarrow Rg = \frac{\sigma}{\rho g} = l_c^2$

sign: $\alpha > 0 \quad \frac{\pi}{2} - \theta \rightarrow 0$

$\Rightarrow dz < 0$

$\begin{cases} dz < 0 \\ dx > 0 \\ dl > 0 \end{cases}$



$dl = -R dz \Rightarrow R = -\frac{dl}{dz}$

$dz = -dl \cdot \sin \alpha \Rightarrow dl = -\frac{dz}{\sin \alpha}$

$\Rightarrow R = \frac{dz}{\sin \alpha dl}$

$\Rightarrow \frac{z dz}{\sin \alpha dl} = l_c^2 \rightarrow z dz = l_c^2 \sin \alpha dl$

far away $\begin{cases} z=0 \\ \alpha=0 \end{cases} \rightarrow \boxed{g^2 = l_c^2 (1 - \cos 2)}$

$\begin{cases} \alpha = \frac{\pi}{2} - \theta \\ z = R \end{cases} \rightarrow \boxed{h = l_c (1 - \sin \theta)^{1/2}}$ for $\theta < 90^\circ$

$h = -l_c (1 - \sin \theta)^{1/2}$ for $\theta > 90^\circ$

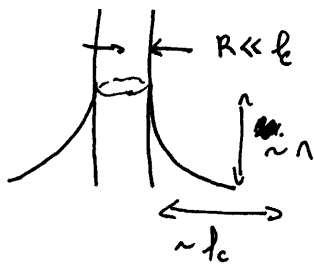
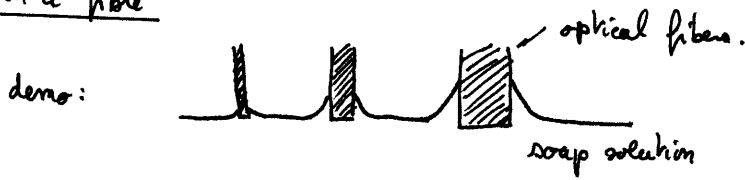
↳ conversion equations: messy

$$\frac{1}{R} = \frac{\ddot{z}}{(1+\dot{z}^2)^{3/2}}$$

shape: $x - x_0 = l_c \operatorname{Arctanh} \frac{2z}{l_c} - \sqrt{2} l_c \left(2 - \frac{z^2}{2l_c^2} \right)^{1/2}$

x_0 to get $z=h$ for $x=0$.

Meniscus on a fibre

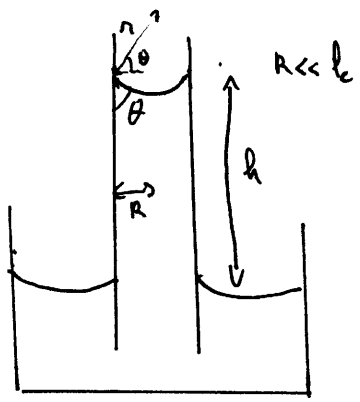
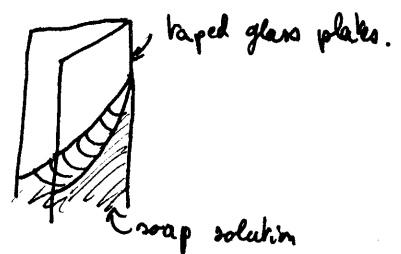
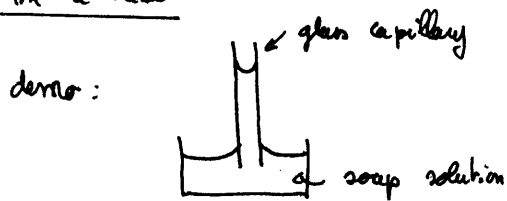


→ 2 curvatures ⇒ 2 radii R and l_c .

$$h \approx R \ln \left(\frac{2l_c}{R} \right) \quad \text{for } \theta=0 \text{ and } R \ll l_c$$

→ slide: experiment with etching liquid.

3.2 Rise in a tube



Laplace = hydrostatic

$$p_0 - \frac{2\sigma \cos \theta}{R} = p_0 - \rho g h$$

$$\Rightarrow h_{\uparrow} = \frac{2\sigma \cos \theta}{\rho g R} = 2 \frac{l_c^2}{R} \cos \theta \quad \text{Jurin's law.}$$

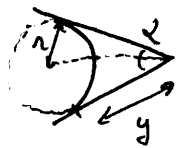
↳ rise if $\theta < 90^\circ$.

→ question: what happens if tube shorter than l_c ?

↳ profile in the edge experiment:

↳ ~~important~~ important for microfluidic devices → angular channels.

for $\theta = 0$

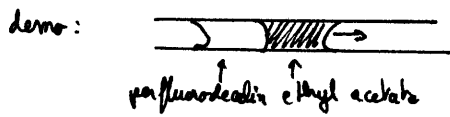


$$R = y \tan \frac{\theta}{2}$$

$$\Rightarrow \frac{\sigma}{r} = \rho g \Rightarrow \frac{\sigma}{y \tan \frac{\theta}{2}} = \rho g$$

$$y = \frac{l_c^2}{g \tan \frac{\theta}{2}}$$

3.3 Capillary forces:

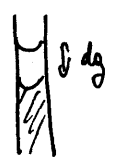


liquid train



conical tube.

force in a tube:



$$dE_s = (\sigma_{SL} - \sigma_{SV}) 2\pi R dy$$

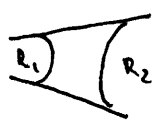
$$f = \frac{dE_s}{dy} = 2\pi R \sigma \cos \theta$$

perimeter of the contact line.

↳ if $\sigma_{SL} - \sigma_{SV} > \sigma_{SL}$ ($S > 0$, $\theta = 0$ and $\cos \theta = 1$).

↳ same force: the difference $\sigma_{SL} - \sigma_{SV} - \sigma_{LV}$ is balanced by viscous dissipation on the precursor film (de Gennes).

Note:



$R_2 > R_1$ → would be expected to move to the right.

⇒ misleading: work with pressures.

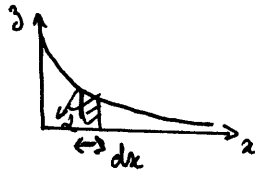
force on a plate:



moving contact line \rightarrow same argument

$$f = \frac{\sigma l \cos \theta}{\downarrow \text{perimeter}}$$

weight of the meniscus:



$$dm = \rho l g dx$$

$$dz = \frac{dl}{dx}$$

$$\tan \alpha = - \frac{dz}{dx}$$

$$\rightarrow dm = - \frac{\rho l g dz}{\tan \alpha}$$

$$\rightarrow \text{from shape: } g dz = l c^2 \sin^2 \alpha dx$$

$$\Rightarrow dm = - \rho l c^2 \frac{\sin^2 \alpha}{\tan \alpha} dx$$

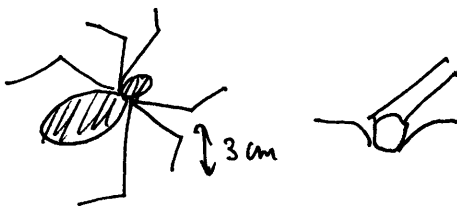
$$\rightarrow m = - \rho l c^2 \int_{\frac{\pi}{2}-\theta}^0 \sin^2 \alpha d\alpha = \rho l c^2 \cos \theta$$

$$\Rightarrow \underline{mg = f.}$$

\rightarrow if $\cos \alpha < 0$ force upwards \Rightarrow dense object can float.

\rightarrow slide = water sticks.

face it can hold:



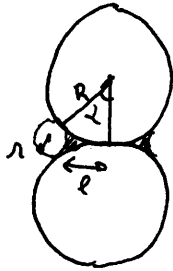
$$f \approx 2 \times 0.03 \times 0.07 \times 6 \approx 0.03 \text{ N} \approx 3g$$

\downarrow
2 sides

(with $\theta = 40^\circ$).

Contact force in between 2 grains:

↳ slide sand castle.



for $\theta = 0$.

shape of the meniscus assumed to be hemispherical. ($\lambda \ll 1$).

$\lambda \ll 1 \Rightarrow$ curvature.

$$R = (R+r) \cos \alpha \rightarrow r = R \frac{1 - \cos \alpha}{\cos \alpha}$$

$$\Delta P = \frac{\sigma}{r} = \frac{\sigma}{R} \frac{\cos \alpha}{1 - \cos \alpha}$$

$$f = \Delta P \cdot \pi r^2 = \frac{\sigma}{R} \frac{\cos \alpha}{1 - \cos \alpha} \pi R^2 V \cos^2 \alpha$$

$$\rightarrow f = \pi R \sigma \frac{1 + \cos \alpha}{\cos \alpha}$$

$$\lambda \ll 1 \Rightarrow \boxed{f \approx 2\pi R \sigma} \text{ independent of volume of liquid!}$$

\rightarrow for small $V \rightarrow$ roughness $\rightarrow f \downarrow$.

