

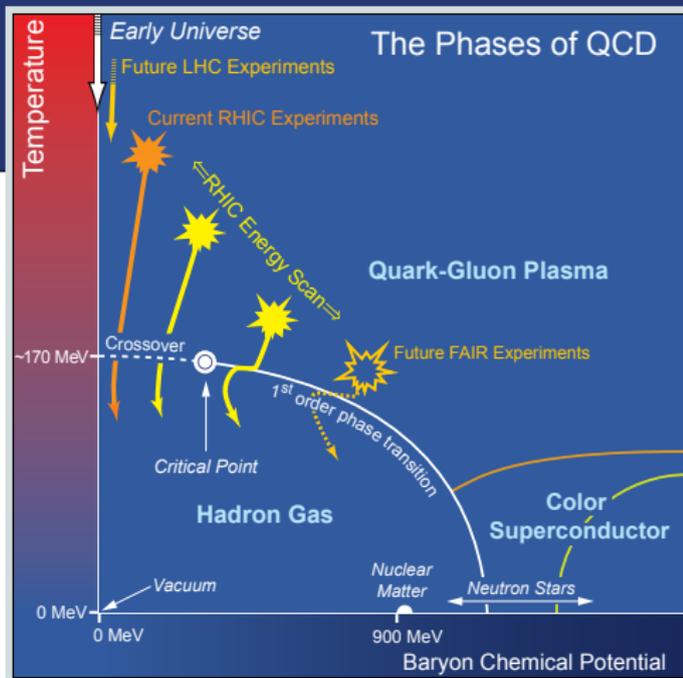
Transport Properties of Strong-Interaction Matter

PANIC11

July 28, 2011, Cambridge, MA

Jochen Wambach

- ▶ ideal fluid @ RHIC and LHC?
- ▶ shear viscosity @ finite T
- ▶ pion dynamics @ low T
- ▶ effects of χ S restoration
- ▶ use NJL model



Nambu-Jona-Lasinio model

bulk thermodynamics

NJL Lagrangian ($N_f = 2$)

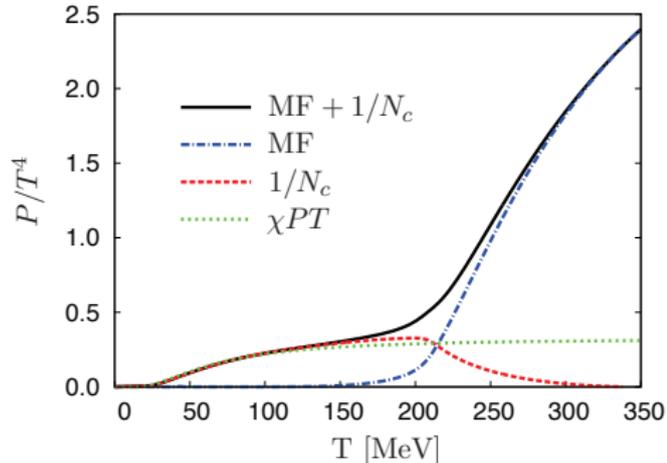
$$\mathcal{L} = \bar{\psi}(i\partial - m_0)\psi + g[(\bar{\psi}\psi)^2 + (\bar{\psi}i\gamma_5\vec{\tau}\psi)^2]$$

Nambu, Jona-Lasinio, Phys. Rev. 1961

thermodynamic potential
NLO in $1/N_c$ -expansion

$$\Omega = \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \dots$$

The diagrams represent Feynman diagrams for the thermodynamic potential. The first diagram is a single loop. The second diagram is a two-loop diagram with a wavy line. The third diagram is a two-loop diagram with a wavy line and a four-point vertex.



Radzhabov et al., 2011, arXiv:1102.0664

Viscous hydrodynamics

(1st order)



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assumptions:

system described locally by $u_\mu(x)$ local thermal equilibrium (EoS $\rightarrow p(\epsilon)$)

basic quantities:

energy-momentum tensor: $T^{\mu\nu}(x)$ 4-current: $J^\mu(x)$

conservation laws:

energy-momentum conservation: $\partial_\mu T^{\mu\nu}(x) = 0$ number conservation: $\partial_\mu J^\mu(x) = 0$

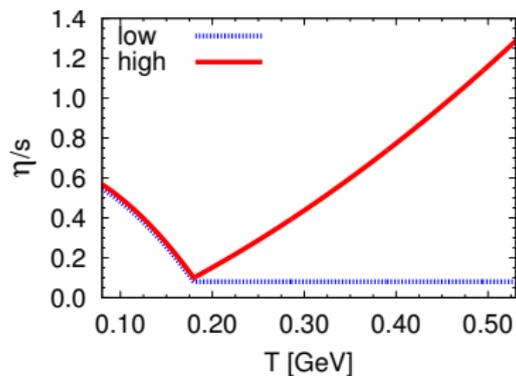
$$T^{\mu\nu} = (\epsilon + p)u^\mu u^\nu - g^{\mu\nu}p + \pi^{\mu\nu}; \quad J^\mu = nu^\mu + \nu^\mu$$

$$\pi^{\mu\nu} = \eta \left(\partial^\mu u^\nu + \partial^\nu u^\mu - u^\mu u_\lambda \partial^\lambda u^\nu - u^\nu u_\lambda \partial^\lambda u^\mu \right)$$

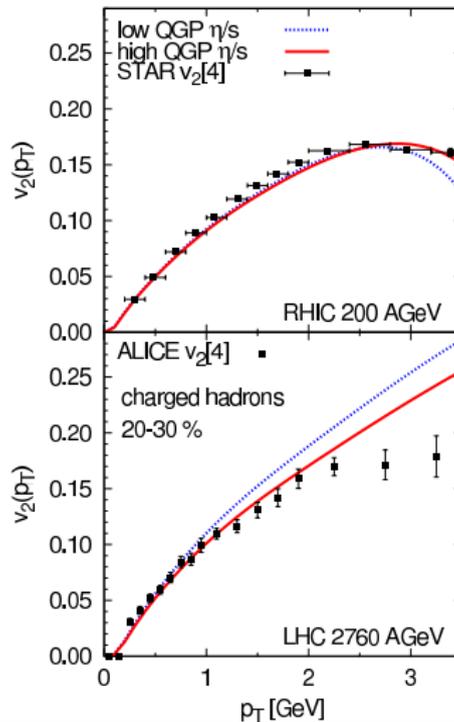
$$+ \left(\zeta - \frac{2}{3}\eta \right) (g^{\mu\nu} - u^\mu u^\nu) \partial_\lambda u^\lambda$$

$$\nu^\mu = \kappa \left(\frac{nT}{\epsilon + p} \right)^2 (\partial^\mu - u^\mu u^\nu \partial_\nu) \left(\frac{\mu}{T} \right)$$

Hydro Elliptic Flow



H. Niemi et al., arXiv:1101.2442



Transport coefficients

linear response theory



linear response theory describes how a system near equilibrium
responds to a weak perturbation

$$\delta \langle O(x) \rangle_f = i \int_{t>0} d^4 x' \underbrace{\langle [O(x), O(x')] \rangle}_{2p\text{-correlator}} \underbrace{f(x')}_{\text{weak pert.}}$$

viscosities:

$$\begin{aligned} & \frac{1}{\omega} \lim_{\vec{q} \rightarrow 0} \int d^3 r \int_0^\infty dt e^{i(\omega t - \vec{q} \cdot \vec{r})} \langle [T_{ik}(x), T_{lm}(0)] \rangle \\ & = \eta(\omega) \left(\delta_{ik} \delta_{km} + \delta_{im} \delta_{kl} - \frac{2}{3} \delta_{ik} \delta_{lm} \right) + \zeta(\omega) \delta_{ik} \delta_{lm} \end{aligned}$$

static limit:

$$\eta = \frac{1}{20} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int d^3 r \int_0^\infty dt e^{i\omega t} \langle [T_{ij}(x), T_{ij}(0)] \rangle, \quad i \neq j$$

$$\zeta = \frac{1}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int d^3 r \int_0^\infty dt e^{i\omega t} \langle [T_{ii}(x), T_{ii}(0)] \rangle$$

$$\kappa = \frac{1}{6T} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int d^3 r \int_0^\infty dt e^{i\omega t} \langle [l_i(x), l_i(0)] \rangle; \quad l_i = T_{0i} - \frac{\epsilon + p}{n} J_i$$

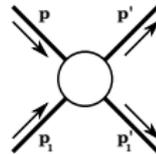
Quantum relativistic kinetic theory

equivalent to Green-Kubo in the dilute gas limit

Boltzmann-Ueling-Uhlenbeck equation:

mean free path $\lambda = 1/n\sigma$ much larger than interaction range, c.f. $d = 1/m$ or $d = \sqrt{\sigma/\pi}$

$$\frac{d}{dt} f_a(\vec{x}, \vec{p}, t) = \sum_b C[f_a, f_b]$$



$$C[f_a, f_b] = \frac{1}{1 + \delta_{ab}} \frac{g_b}{2E_b} \int \frac{d^3 p'}{(2\pi)^3} \int \frac{d^3 p_1}{(2\pi)^3} \int \frac{d^3 p'_1}{(2\pi)^3} (2\pi)^4 \delta(\mathbf{p} + \mathbf{p}_1 - \mathbf{p}' - \mathbf{p}'_1) |\mathcal{M}_{ab}|^2 \times \\ \times [f'_a f'_{1b} (1 + f_a)(1 + f_{1b}) - f_a f_{1b} (1 + f'_a)(1 + f'_{1b})]$$

- ▶ linearization of the BUU-equation (2^{nd} -order Chapman-Enskog expansion)

$$f_a = f_a^{(0)} + \epsilon f_a^{(1)} + \dots; \quad f_a^{(1)} = f_a^{(0)} (1 + f_a^{(0)}) \phi_a$$

- ▶ shear viscosity

$$\eta = \sum_a \frac{g_a}{15} \frac{4\pi}{(2\pi)^3} \int_0^\infty dp \frac{p^4}{E_p} f_a^{(0)} (1 + f_a^{(0)}) \mathcal{B}_a(p)$$

- ▶ expansion of $\mathcal{B}_a(p)$ in orthogonal polynomials ('generalized Sonine functions')

Nambu-Jona-Lasinio model

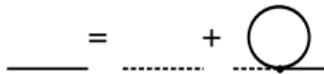
quark and meson masses

NJL Lagranian ($N_f = 2$)

$$\mathcal{L} = \bar{\psi}(i\partial - m_0)\psi + g[(\bar{\psi}\psi)^2 + (\bar{\psi}i\gamma_5\vec{\tau}\psi)^2]$$

Nambu, Jona-Lasinio, Phys. Rev. 1961

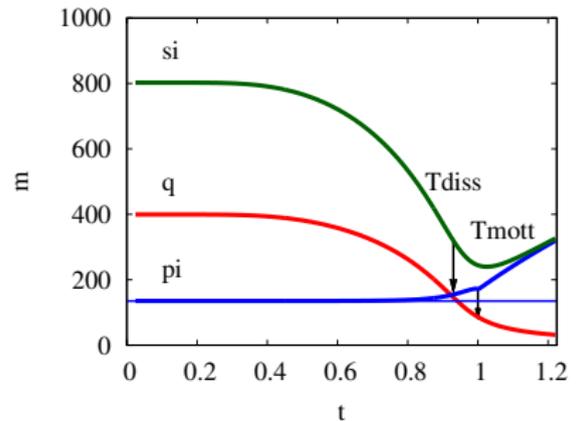
Hartree quark propagator



RPA meson propagators



parameters: Buballa, Phys. Rept. 2005



'dissociation'
temperature

$$m_\sigma = 2m_\pi$$

'Mott'
temperature

$$m_\pi = 2m_q$$

$$T_{diss} \approx 93\% T_{Mott}$$

Pion-Pion scattering

$\pi\pi$ scattering amplitude

$$i\mathcal{M}_{\pi\pi} = \text{[Box Diagram]} + \text{[Triangle Diagram with } \sigma \text{]}$$

Quack, et al., Phys. Lett. B 348 1 (1994)

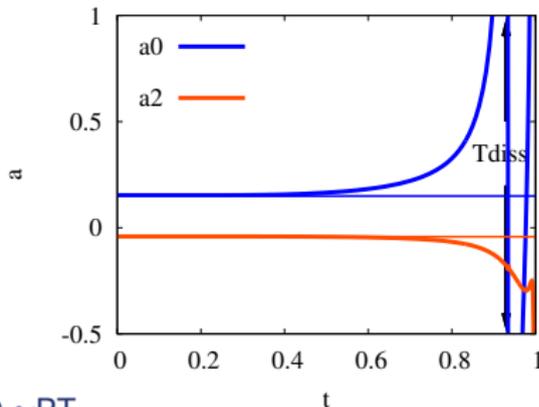
$\pi\pi$ scattering length

$$a^l = \frac{1}{32\pi m_\pi} \mathcal{M}_{\pi\pi}^l(s = 4m_\pi^2)$$

$\pi\pi$ cross section

$$\frac{d\sigma}{d\Omega} = \frac{|\mathcal{M}_{\pi\pi}|^2}{64\pi^2 E_{cm}^2}$$

$$|\mathcal{M}_{\pi\pi}|^2 = \frac{1}{9} \sum_l (2l+1) |\mathcal{M}_{\pi\pi}^l|^2$$



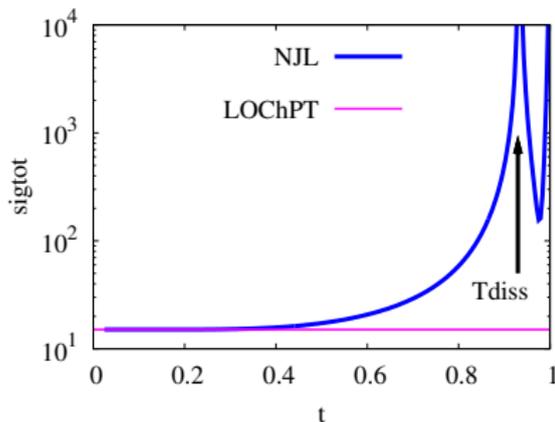
► LO- χ PT

$$a_W^0 = \frac{7m_\pi}{32\pi f_\pi^2}; \quad a_W^2 = -\frac{2m_\pi}{32\pi f_\pi^2}$$

Weinberg, Phys. Rev. Lett. 1966

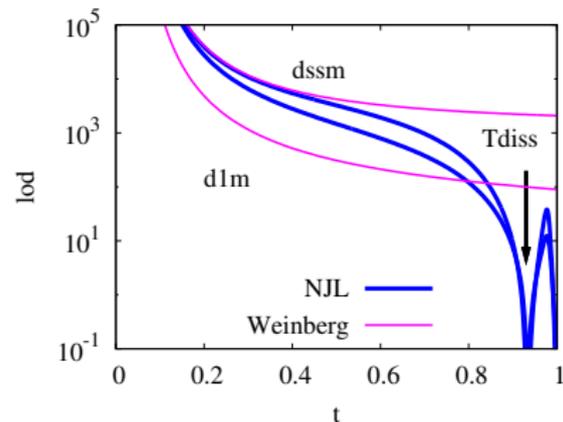
- $a^{0,2}$ large at T_{diss} , T_{Mott}
 \leftrightarrow ultra-cold atomic gases near 'Feshbach res.'

σ_{tot} at threshold



- ▶ in vacuum: $NJL \approx LO\text{-}\chi PT$
- ▶ strong coupling at transition temperatures

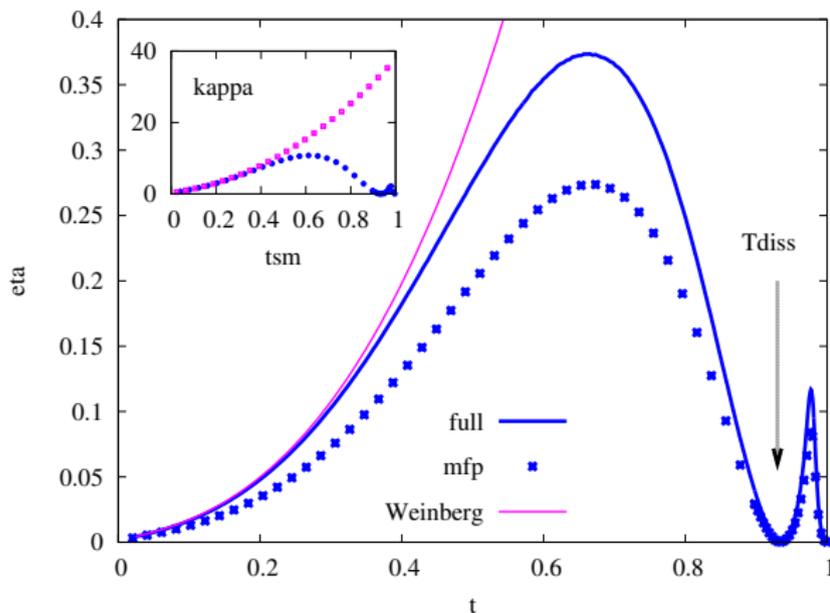
validity of relativistic kinetic theory



- ▶ Weinberg: valid in the whole T -range considered
- ▶ NJL: valid up to $\sim 88\% T_{Mott}$

Validity check according to Itakura, et al., Phys. Rev. D 2008

Shear viscosity and thermal conductivity



- ▶ η and κ similar behavior
- ▶ full RKT vs. mfp-estimate
 $\eta = \bar{\rho}/3\sigma$
- ▶ χ PT & NJL agree for small T
- ▶ important modifications at higher T
- ▶ divergence of a^0 at T_{diss} visible
- ▶ strong coupling at T_{diss} and T_{Mott}

Shear viscosity

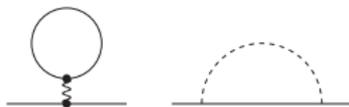
Green-Kubo formalism

static shear viscosity using Green-Kubo formula

$$\eta = - \lim_{\omega \rightarrow 0} \lim_{\vec{q} \rightarrow 0} \frac{d}{d\omega} \text{Im} \left(\text{quark self-energy diagram} \right)$$

- ▶ quark spectral function
- ▶ quark vertex

quark self energy: (NLO)



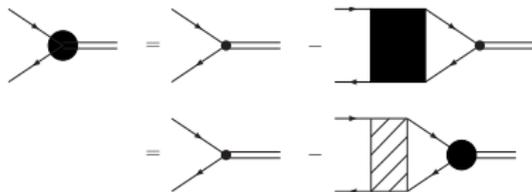
entropy density: (NLO)

$$s = - \frac{\partial \Omega}{\partial T}; \quad \Omega = \text{quark loop} + \text{ghost loop} + \text{gluon loop} + \text{quark-gluon loop} + \dots$$

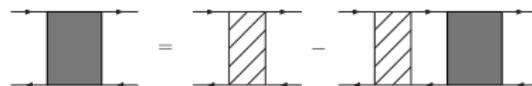
Shear viscosity

$1/N_c$ - expansion

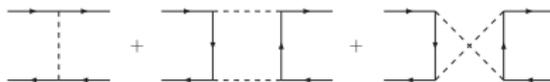
3-point vertex:



4-point vertex:



NLO Kernel:



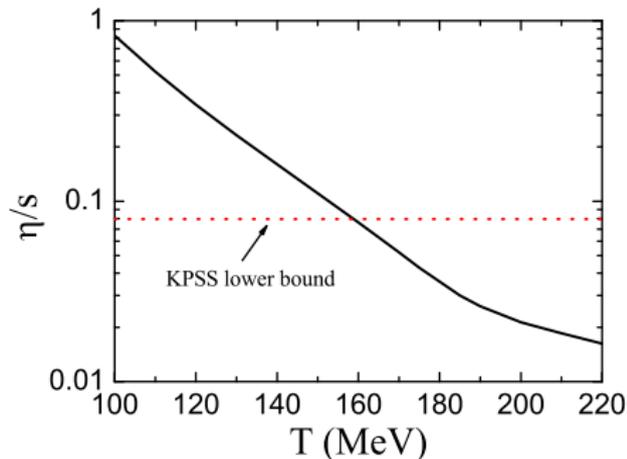
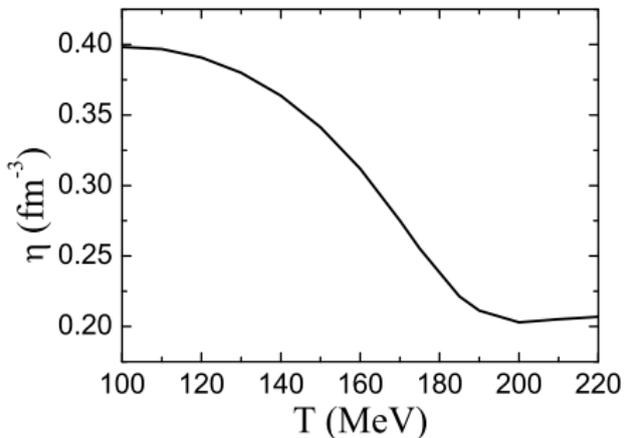
difficult to solve!

for scalar $O(N)$ theories

Aarts and Resco, JHEP 2004

Shear viscosity

Kubo results



Wei-jie Fu, J. Wambach and Yue-liang Wu, in progress

shear viscosity in the NJL model ($1/N_c$ - expansion)

- ▶ kinetic theory
 - ▶ in-medium $\pi\pi$ - scattering
 - ▶ scattering amplitude \mathcal{M} from scattering length
 - ▶ agreement with χPT at low temperatures
 - ▶ σ large at T_{diss} and T_{Mott}
 - ▶ minima of η/s at transition temperatures
- ▶ Green-Kubo formalism
 - ▶ η strongly decreases near chiral restoration
 - ▶ η/s falls below ADS/CFT bound

future directions

- ▶ improve cross sections to include higher partial waves and more species
- ▶ PNJL calculation within Kubo formalism
- ▶ study finite chemical potential (CEP)