

Probing Fundamental Physics with Molecules

Carsten Zülch, Konstantin Gaul, Robert Berger

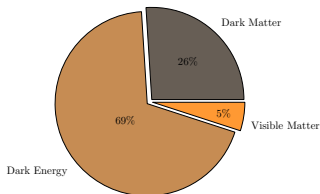
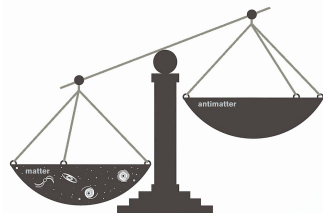
Fachbereich Chemie
Philipps-Universität Marburg, Germany

@New Opportunities for Fundamental Physics Research with
Radioactive Molecules

July 01, 2021

Discrete Symmetry Violation — Towards New Physics

- Standard Model leaves important phenomena unexplained:



- Baryon asymmetry
 - Dark matter and dark energy
- ⇒ Connected to symmetry violation

Discrete Symmetry Violation:

- Charge conjugation $C: \psi \rightarrow \bar{\psi}$
- Parity $\mathcal{P}: x, y, z \rightarrow -x, -y, -z$
- Time-Reversal $\mathcal{T}: t \rightarrow -t$

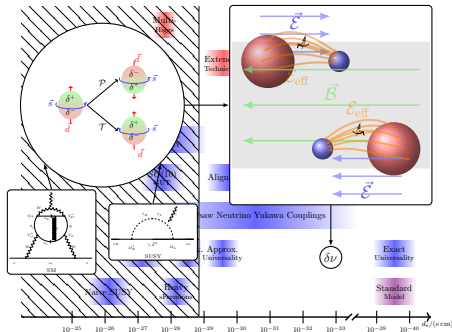
CPT -theorem $\Rightarrow \cancel{\mathcal{T}} \leftrightarrow \cancel{CP}$

Discrete Symmetry Violation — Towards New Physics

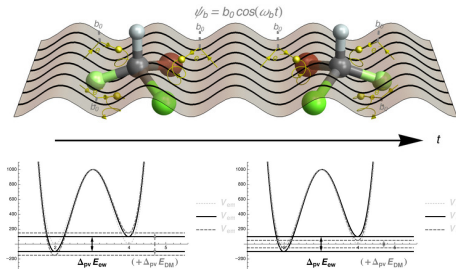
Molecules as versatile probes of New Physics

$\mathcal{P}$, $\mathcal{T}$ -Violation in molecules

Present Experimental Limit:
 1.1×10^{-29} e cm



Sensing dark matter with chiral molecules

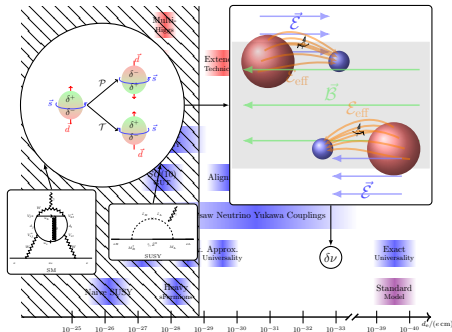


Discrete Symmetry Violation — Towards New Physics

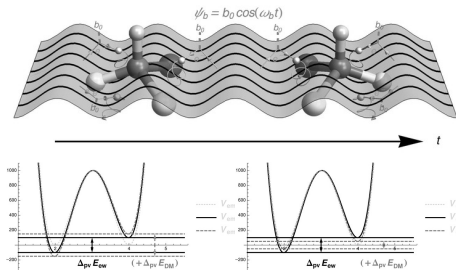
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Sources of $\mathcal{P}, \mathcal{T}$ -violation in molecules

From elementary particles to atoms and molecules

Fundamental Theory

θ , CKM, SUSY, Multi Higgs, LR-symmetry, etc.

Wilson coefficients

d_e

C_{qe}, C_{qq}

$\theta, d_q, \tilde{d}_q$

Elementary Particles

High-energy
physics

Sources of $\mathcal{P}, \mathcal{T}$ -violation in molecules

From elementary particles to atoms and molecules

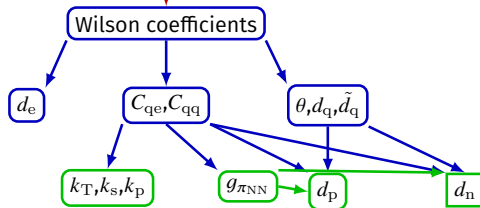
Fundamental Theory

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Elementary Particles

Hadrons

High-energy
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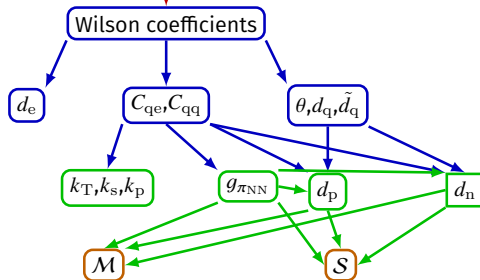
Elementary Particles

Hadrons

Nuclei

High-energy
physics

Precision
Spec-
troscopy

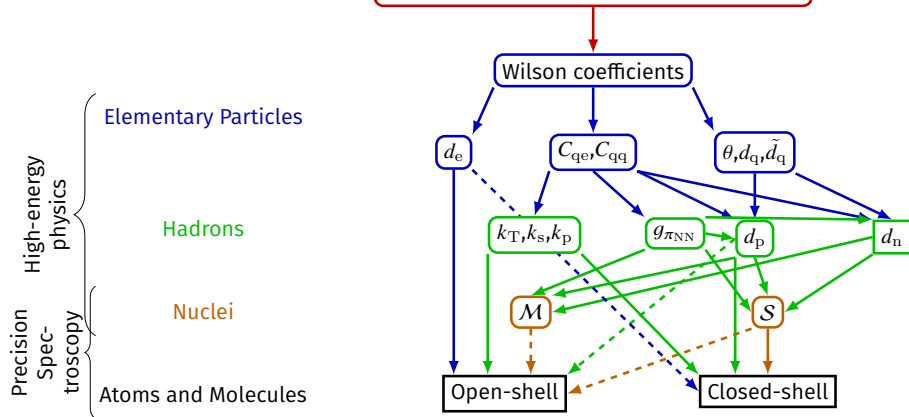


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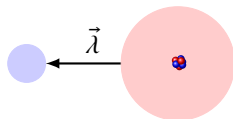
Sources of $\mathcal{P}, \mathcal{T}$ -violation in molecules

Effective molecular spin-rotational Hamiltonian

Polar linear heavy-elemental molecules

⇒ Strong internal fields

⇒ Easy to polarize



Effective spin-rotational Hamiltonian for heavy nucleus

$$H_{\text{sr}} =$$

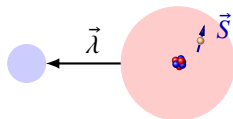
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Effective spin-rotational Hamiltonian for heavy nucleus

$$H_{\text{sr}} = \underbrace{\vec{\lambda} \cdot \vec{S}'}_{\Omega} (W_{\text{d}} d_{\text{e}} + W_{\text{s}} k_{\text{s}})$$

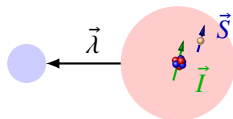
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Effective spin-rotational Hamiltonian for heavy nucleus

$$H_{\text{sr}} = \underbrace{\vec{\lambda} \cdot \vec{S}'}_{\Omega} (W_{\text{d}} d_{\text{e}} + W_{\text{s}} k_{\text{s}}) + \underbrace{\vec{\lambda} \cdot \vec{I}}_I (W_{\text{T}} k_{\text{T}} + W_{\text{p}} k_{\text{p}} + W_{\text{s}}^{\text{m}} k_{\text{s}} + W_{\text{S}} S + (W_{\text{m}} + W_{\text{S}} R) d_{\text{p}} + W_{\text{d}}^{\text{m}} d_{\text{e}})$$

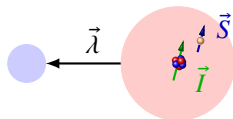
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$$H_{\text{sr}} = \underbrace{\vec{\lambda} \cdot \vec{S}'}_{\Omega} (W_d d_e + W_s k_s) + \underbrace{\vec{\lambda}^T \cdot \mathbf{T} \cdot \vec{S}'}_{\Theta} W_M \tilde{M} + \text{higher moments...}$$

$$+ \underbrace{\vec{\lambda} \cdot \vec{I}}_I (W_T k_T + W_P k_P + W_s^m k_s + W_S S + (W_m + W_S R) d_P + W_d^m d_e)$$

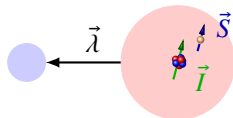
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⇒ **Large effective electronic structure factors $W \Leftrightarrow$ high sensitivity(?)**

Numerical calculation and analytical models

Molecular *ab initio* calculations

- Quasi-relativistic Hartree–Fock and DFT calculations
- Sufficiently accurate computation of W (within 10 % for most)
- Can compute arbitrary properties
- Excited states via SCF-MOM

Analytical atomic models

- One-electron atom model
- Laws for *relativistic enhancement*
→ Scaling with nuclear charge Z

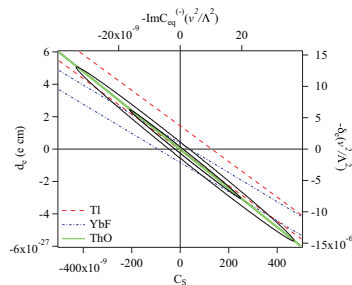
Numerical results for molecules with $\Omega = 1/2, \mathcal{I} = 0$

$$H_{\text{srA}} = \Omega (W_{\text{dA}} d_{\text{e}} + W_{\text{sA}} k_{\text{s}})$$

- electron EDM (eEDM) d_{e}
- scalar-pseudoscalar nucleon-electron current k_{s}

$$h \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} = \Omega \begin{pmatrix} W_{\text{d},1} & W_{\text{s},1} \\ W_{\text{d},2} & W_{\text{s},2} \end{pmatrix} \begin{pmatrix} d_{\text{e}} \\ k_{\text{s}} \end{pmatrix}$$

⇒ Need **at least two** experiments



Numerical results for molecules with $\Omega = 1/2, \mathcal{I} = 0$

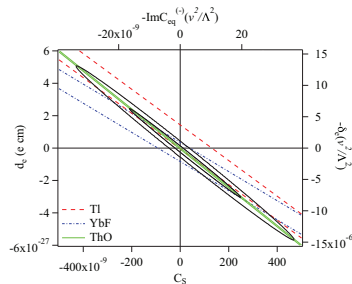
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⇒ Need **at least two** experiments

- $\mathcal{P}, \mathcal{T}$ -violation in diatomic molecules with $^2\Sigma_{1/2}$ -ground state
 - *Ab initio* calculation of many diatomic molecules vs. atomic models
⇒ Chemical enhancement
 - Dependence of the coverage area A on Z and chemical influence



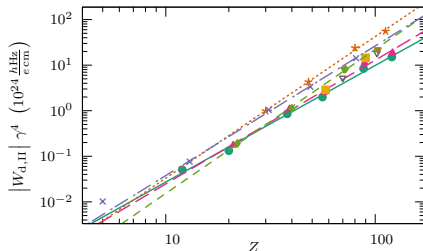
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eEDM interactions in $^2\Sigma_{1/2}$ -ground state molecules

Chemical influence on scaling behavior of eEDM interactions

Empirical atomic model $W_d \sim \frac{\alpha^2 Z^3}{\gamma^4}$; ZORA-cGKS-B3LYP calculations

$$\log_{10} \left\{ |W_d| \gamma^4 \times 10^{-24} \frac{\text{e cm}}{\text{h Hz}} \right\} = b_{d,\text{FS}} + a_{d,\text{FS}} \log_{10} \{Z\}$$



(Mg-E120)F

$10^{-4.12} Z^{2.56}$

(Sc-E121)O

$10^{-4.32} Z^{2.71}$

(Ce-Th)N

(Yb-No)F

(Lu-Lr)O



(Ti-Hf)N

$10^{-4.97} Z^{3.16}$

(Cd-Cn)H

$10^{-4.59} Z^{3.11}$

(B-Tl)O

$10^{-4.27} Z^{2.85}$



- ⇒ Chemically suppressed Z-scaling in group 2 and 3
- ⇒ Chemical enhancement in group 4 and 12

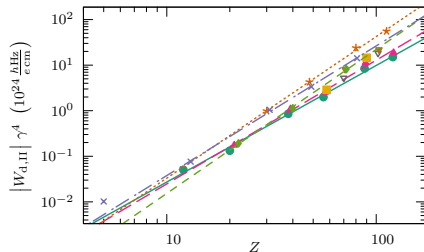
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(Mg-E120)F	●	(Ti-Hf)N	◆
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(Sc-E121)O	▲	(Cd-Cn)H	★
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(Ce-Th)N	■	(B-Tl)O	×
(Yb-No)F	▼	$10^{-4.27} Z^{2.85}$	---
(Lu-Lr)O	▼		

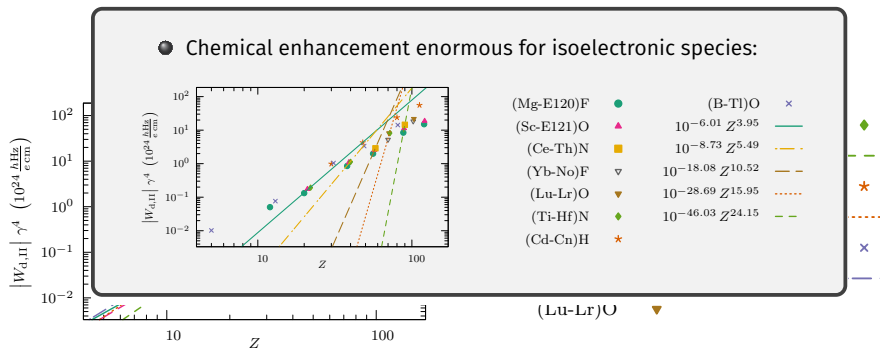
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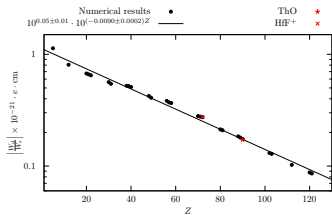


⇒ Chemically suppressed Z-scaling in group 2 and 3

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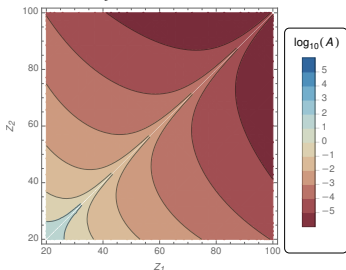
$\mathcal{P}, \mathcal{T}$ -odd ratio and coverage region in $d_e - k_s$ parameter space



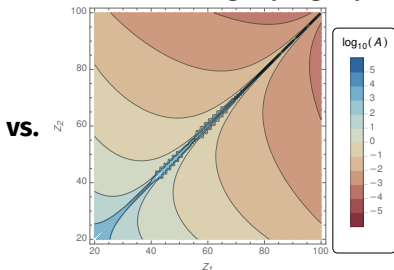
$$A \sim W_{d,1}^{-1} W_{d,2}^{-1} \left| \frac{W_{s,1}}{W_{d,1}} - \frac{W_{s,2}}{W_{d,2}} \right|^{-1}$$

$$A \approx \frac{k_p^2 \pi |u(v_1)u(v_2)|}{10^{b_{d,1}+b_{d,2}} \frac{Z_1^{a_{d,1}} Z_2^{a_{d,2}}}{\gamma_1^4 \gamma_2^4} 0.9 \cdot |1.02Z_1 - 1.02Z_2| \times 10^{27} \frac{\text{Hz}^2}{\text{e} \cdot \text{cm}}}$$

Analytical atomic model



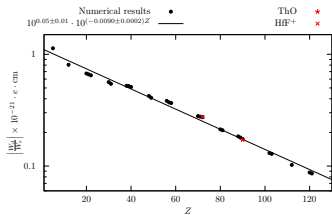
Ab initio model: group 2, group 2



vs.

Numerical results for molecules with $\Omega = 1/2, \mathcal{I} = 0$

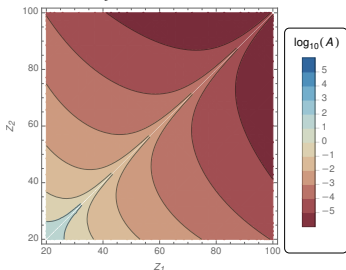
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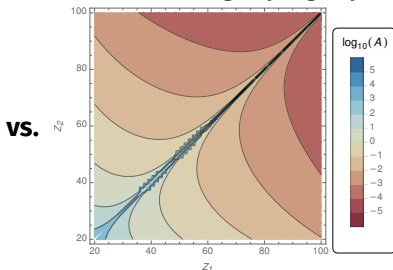
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Analytical atomic model



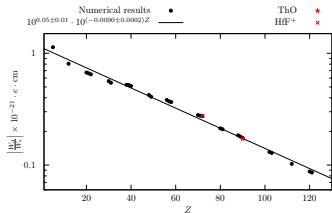
Ab initio model: group 2, group 12



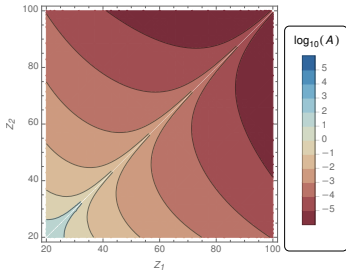
vs.

Numerical results for molecules with $\Omega = 1/2, \mathcal{I} = 0$

$\mathcal{P}, \mathcal{T}$ -odd ratio and coverage region in $d_e - k_s$ parameter space

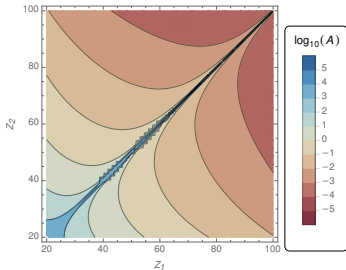


Analytical atomic model



vs.

Ab initio model: group 12, group 12



$$A \sim W_{d,1}^{-1} W_{d,2}^{-1} \left| \frac{W_{s,1}}{W_{d,1}} - \frac{W_{s,2}}{W_{d,2}} \right|^{-1}$$

$$A \approx \frac{k_p^2 \pi |u(v_1)u(v_2)|}{10^{b_{d,1}+b_{d,2}} \frac{Z_1^{a_{d,1}} Z_2^{a_{d,2}}}{\gamma_1^4 \gamma_2^4} 0.9 \cdot |1.02 Z_1 - 1.02 Z_2| \times 10^{27} \frac{\text{Hz}^2}{\text{e} \cdot \text{cm}}}$$

Numerical results for molecules with $\Omega = 1/2, \mathcal{I} = 0$

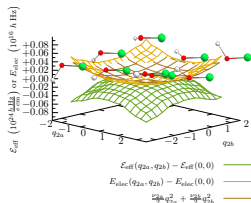
Linear polyatomic molecules: MOH

- W_d, W_s and $\frac{W_s}{W_d}$ similar to MF

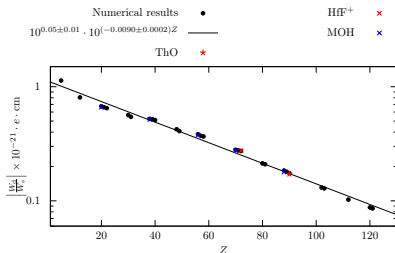
⇒ No advantages for disentanglement

- **Complementary future searches:**

- ▶ **More complex polyatomic molecules.**
- ▶ **Diamagnetic systems**



H-bending modes in YbOH



Laser-coolable 1st excited vibrational state:

- Experiment: Internal co-magnetometer states
- Calculations: Vibrational corrections to $\mathcal{E}_{\text{eff}}$ negligible ($< 1\%$)

⇒ Promising candidates for experiment

Numerical results for molecules with $\Omega = 1/2, \mathcal{I} = 0$

Linear polyatomic molecules: MOH

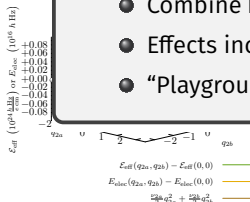
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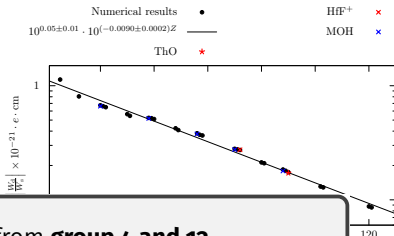
- **Complementary future searches:**

▶ **More complex polyatomic molecules**

- Largest effects in molecules from **group 4 and 12**
- Combine **heavy** molecules that **differ** in Z
- Effects increase by changing chemical environment
- “Playground” of diatomic/linear molecules restricted ?



H-bending modes in YbOH



negligible ($< 1\%$)

⇒ Promising candidates for experiment

Global $\mathcal{P}, \mathcal{T}$ -odd measurement model

- Six relevant *fundamental* parameters in a molecule

$$\hbar\vec{\omega} = \mathbf{W}\vec{x}_{\mathcal{P},\mathcal{T}},$$

$$\mathbf{W} = \begin{pmatrix} & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \end{pmatrix}$$

Global $\mathcal{P}, \mathcal{T}$ -odd measurement model

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$$\hbar \vec{\omega} = \mathbf{W} \vec{x}_{\mathcal{P}, \mathcal{T}},$$
$$\mathbf{W} = \begin{pmatrix} \Omega_1 W_{\text{d},1} & \Omega_1 W_{\text{s},1} \\ \Omega_2 W_{\text{d},2} & \Omega_2 W_{\text{s},2} \end{pmatrix}$$

- | open-shell molecules ($\Omega > 0$) with a closed-shell nucleus with $\mathcal{I} = 0$,

Global $\mathcal{P}, \mathcal{T}$ -odd measurement model

- Six relevant *fundamental* parameters in a molecule

$$\hbar\vec{\omega} = \mathbf{W}\vec{x}_{\mathcal{P},\mathcal{T}},$$

$$\mathbf{W} = \begin{pmatrix} \mathcal{I}_1 W_{\text{d},1}^{\text{m}} & \mathcal{I}_1 W_{\text{s},1}^{\text{m}} & \mathcal{I}_1 W_{\text{T},1} & \mathcal{I}_1 W_{\text{S},1} & \mathcal{I}_1 (R_1 W_{\text{S},1} + W_{\text{m},1}) \\ \mathcal{I}_2 W_{\text{d},2}^{\text{m}} & \mathcal{I}_2 W_{\text{s},2}^{\text{m}} & \mathcal{I}_2 W_{\text{T},2} & \mathcal{I}_2 W_{\text{S},2} & \mathcal{I}_2 (R_2 W_{\text{S},2} + W_{\text{m},2}) \\ \mathcal{I}_3 W_{\text{d},3}^{\text{m}} & \mathcal{I}_3 W_{\text{s},3}^{\text{m}} & \mathcal{I}_3 W_{\text{T},3} & \mathcal{I}_3 W_{\text{S},3} & \mathcal{I}_3 (R_3 W_{\text{S},3} + W_{\text{m},3}) \\ \mathcal{I}_4 W_{\text{d},4}^{\text{m}} & \mathcal{I}_4 W_{\text{s},4}^{\text{m}} & \mathcal{I}_4 W_{\text{T},4} & \mathcal{I}_4 W_{\text{S},4} & \mathcal{I}_4 (R_4 W_{\text{S},4} + W_{\text{m},4}) \\ \mathcal{I}_5 W_{\text{d},5}^{\text{m}} & \mathcal{I}_5 W_{\text{s},5}^{\text{m}} & \mathcal{I}_5 W_{\text{T},5} & \mathcal{I}_5 W_{\text{S},5} & \mathcal{I}_5 (R_5 W_{\text{S},5} + W_{\text{m},5}) \end{pmatrix}$$

- I open-shell molecules ($\Omega > 0$) with a closed-shell nucleus with $\mathcal{I} = 0$,
- II closed-shell molecules ($\Omega = 0$) with an open-shell nucleus with $\mathcal{I} = 1/2$,

Global $\mathcal{P}, \mathcal{T}$ -odd measurement model

- Six relevant *fundamental* parameters in a molecule

$$\hbar\vec{\omega} = \mathbf{W}\vec{x}_{\mathcal{P},\mathcal{T}},$$

$$\mathbf{W} = \begin{pmatrix} \Omega_1 W_{d,1} + \mathcal{I}_1 W_{d,1}^m & \Omega_1 W_{s,1} + \mathcal{I}_1 W_{s,1}^m & \mathcal{I}_1 W_{T,1} & \mathcal{I}_1 W_{S,1} & \mathcal{I}_1 (R_1 W_{S,1} + W_{m,1}) \\ \Omega_2 W_{d,2} + \mathcal{I}_2 W_{d,2}^m & \Omega_2 W_{s,2} + \mathcal{I}_2 W_{s,2}^m & \mathcal{I}_2 W_{T,2} & \mathcal{I}_2 W_{S,2} & \mathcal{I}_2 (R_2 W_{S,2} + W_{m,2}) \\ \Omega_3 W_{d,3} + \mathcal{I}_3 W_{d,3}^m & \Omega_3 W_{s,3} + \mathcal{I}_3 W_{s,3}^m & \mathcal{I}_3 W_{T,3} & \mathcal{I}_3 W_{S,3} & \mathcal{I}_3 (R_3 W_{S,3} + W_{m,3}) \\ \Omega_4 W_{d,4} + \mathcal{I}_4 W_{d,4}^m & \Omega_4 W_{s,4} + \mathcal{I}_4 W_{s,4}^m & \mathcal{I}_4 W_{T,4} & \mathcal{I}_4 W_{S,4} & \mathcal{I}_4 (R_4 W_{S,4} + W_{m,4}) \\ \Omega_5 W_{d,5} + \mathcal{I}_5 W_{d,5}^m & \Omega_5 W_{s,5} + \mathcal{I}_5 W_{s,5}^m & \mathcal{I}_5 W_{T,5} & \mathcal{I}_5 W_{S,5} & \mathcal{I}_5 (R_5 W_{S,5} + W_{m,5}) \end{pmatrix}$$

- I open-shell molecules ($\Omega > 0$) with a closed-shell nucleus with $\mathcal{I} = 0$,
- II closed-shell molecules ($\Omega = 0$) with an open-shell nucleus with $\mathcal{I} = 1/2$,
- III open-shell molecules with an open-shell nucleus with $\mathcal{I} = 1/2$,

Global $\mathcal{P}, \mathcal{T}$ -odd measurement model

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- IV open-shell molecules with an open-shell nucleus with $\mathcal{I} \geq 1$.

$$\Rightarrow \text{We need to minimize coverage volume } V = P^3 \frac{\pi^3}{6} \frac{\prod_{i=1}^6 |u(\omega_i)|}{|\det(\mathbf{W})|}$$

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- IV open-shell molecules with an open-shell nucleus with $\mathcal{I} \geq 1$.

⇒ We need to minimize coverage volume $V = P^3 \frac{\pi^3}{6} \frac{\prod_{i=1}^6 |u(\omega_i)|}{|\det(\mathbf{W})|}$

- **Radioactive molecules** with long-isotope chains
- ⇒ **2-4 classes with "one" molecule**

Radioactive molecules for $\mathcal{P}$, $\mathcal{T}$ -violation

Radium monofluoride

- RaF has a simple electronic structure
- RaF is laser-coolable
- RaF: Class I (^{226}RaF), [Class II ($^{223,225}\text{RaF}^+$)], Class III (^{225}RaF) and Class IV (^{223}RaF).
- But no direct sensitivity to d_p (no unpaired valence proton)?
- Other candidates: Iso-electronic Actinides

*	89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No
*														

Radioactive molecules for $\mathcal{P}$, $\mathcal{T}$ -violation

Actinide molecules for search for $\mathcal{P}$, $\mathcal{T}$ -violation

Nuclear structure effects:

- $^{225,227}\text{Ac}$, ^{229}Th , ^{229}Pa Supposed to have octupole deformation
 - Ac and Pa have an open proton shell.
 - Pa possibly has a parity doublet with splitting $\Delta E \sim 60 \text{ eV}$ (^{229}Th has $\Delta E < 130 \text{ keV}$)
- ⇒ Collective Schiff moment two orders larger than in ^{229}Th
- Embedding ^{229}Pa in a molecule may help to understand nuclear structure!

I. Ahmad et al., *Phys. Rev. Lett.* **1982**, 49, 1758–1761, V. V. Flambaum, *Phys. Rev. C* **2019**, 99, 035501.

I. Ahmad et al., *Phys. Rev. C* **2015**, 92, 024313.

J. T. Singh, *Hyperfine Interact.* **2019**, 240, 29.

Radioactive molecules for $\mathcal{P}$, $\mathcal{T}$ -violation

Actinide molecules for search for $\mathcal{P}$, $\mathcal{T}$ -violation

Electronic structure:

- ThO most successful eEDM experiment so far, ThF⁺ experiment planned
- Large enhancement of eEDM and k_s in AcO and ThN, both may have $^2\Sigma_{1/2}$ ground state (ZORA-cGHF)
- Schiff moment enhancement in AcO ($\sim -5800/a_0^4$) larger than in RaF ($\sim -4000/a_0^4$) (in ThN ($\sim -2200/a_0^4$) smaller)
- Cationic species: AcF⁺ (and similar ones), ThO⁺
- What about Pa?

V. Andreev et al., *Nature* **2018**, 562, 355–360.

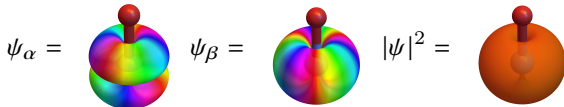
D. N. Gresh et al., *J. Mol. Spec.* **2016**, 319, 1–9.

K. Gaul et al., *Phys. Rev. A* **2019**, 99, 032509.

Radioactive molecules for $\mathcal{P}$, $\mathcal{T}$ -violation

Protactinium containing molecules?

- No neutral molecules iso-electronic to RaF
 - PaN or PaO⁺ would be iso-electronic to ThO
- ⇒ f-electron dominated → strongly mixed ground state?
- PaN⁺ and PaO²⁺ were experimentally observed
 - Theoretical studied of PaO²⁺ by Kovacs *et.al.* (CASPT2)
 - $\Omega = 5/2$ ground state with $\sim 80\%$ Φ and $\sim 20\%$ Δ characterized by a non-bonding f-electron (ZORA-cGKS-B3LYP)



- Schiff moment strongly enhanced ($W_S \lesssim -11\,000/a_0^4$)
- Electron spin dependent effects suppressed (f-electron).
- Sympathetic cooling (mass to charge ratio in range of Ba⁺)?

M. Santos *et al.*, *J. Phys. Chem. A* **2006**, 110, PMID: 16640369, 5751–5759.

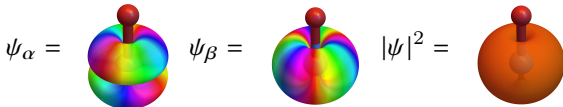
A. Kovács *et al.*, *Struct. Chem.* **2013**, 24, 917–925.

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 - $\Omega =$ non

Continuation follows soon...



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Conclusion

- Many good molecular $\mathcal{P}, \mathcal{T}$ -violation experiments needed
- Robust bounds on $\mathcal{P}, \mathcal{T}$ -Violation require molecules with fundamentally different spin-rotational Hamiltonian
 - Many experiments involving non-zero spin nuclei are required
 - ⇒ Radioactive molecules most interesting here
 - Have to find balance between large enhancement and different ratio of enhancement factors
 - Polyatomic molecules can reduce experimental uncertainty
 - Non-linear polyatomic molecules have a different spin-rotational Hamiltonian!
- PaO^{2+} has potential to be used for precision spectroscopy
- More interesting Pa candidates and a detailed study will follow soon...

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Thank You for Your Attention!