

Recap:

- Found that velocity addition formula broke momentum conservation. Fix: Tweak the definition: $\mathbf{p} = \gamma(u) m \mathbf{u}$ is conserved during interactions among bodies.

- This notion of momentum makes kinetic energy more complicated. Integrating the work done on a body, we find that a body with speed u has kinetic energy

$$K = (\gamma(u) - 1) mc^2$$

- Interpreting the total energy as $E = \gamma(u) mc^2$, this implies a rest energy $E_0 = mc^2$.

Suppose we have a body with energy E and momentum \underline{p} as measured in our lab.

What are the energy E' and momentum \underline{p}' as measured by an observer whose velocity relative to the lab is $\underline{v} = v \underline{e}_x$?

Recipe to follow:

1. Compute the velocity of the body in the lab from \underline{p} and E .
2. Using the velocity addition formulas, compute the velocity in the new ("moving") frame.
3. From the velocity, assemble E' & \underline{p}' .

Result:

$$E' = \gamma (E - v p^x)$$

$$p^{x'} = \gamma (p^x - v E/c^2)$$

$$p^{y'} = p^y, \quad p^{z'} = p^z$$

or

$$\begin{bmatrix} E'/c \\ p^{x'} \\ p^{y'} \\ p^{z'} \end{bmatrix} = \begin{bmatrix} \gamma & -\gamma\beta & 0 & 0 \\ -\gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} E/c \\ p^x \\ p^y \\ p^z \end{bmatrix}$$

→ Energy & momentum form a set of quantities that transform under a Lorentz transformation!

Recall $\Delta s^2 = -c^2 \Delta t^2 + \Delta x^2 + \Delta y^2 + \Delta z^2$ is a "Lorentz invariant": All Lorentz frames agree on its value. Can we do something similar with energy and momentum?

Inspired by the spacetime invariant, let's think of energy as the "time-like" piece of momentum. Using it, we'll construct something similar:

$$-\frac{E^2}{c^2} + |\mathbf{p}|^2 = ?$$

Plug in $E^2 = \gamma^2 m^2 c^4 = \frac{m^2 c^4}{1 - u^2/c^2}$

$$|\mathbf{p}|^2 = \gamma^2 m^2 u^2 = \frac{m^2 u^2}{1 - u^2/c^2}$$

So,
$$-\frac{E^2}{c^2} + |\mathbf{p}|^2 = \frac{-m^2 c^2 + m^2 u^2}{1 - u^2/c^2} = -m^2 c^2 \left[\frac{1 - u^2/c^2}{1 - u^2/c^2} \right]$$

$$- E^2 + |\mathbf{p}|^2 c^2 = m^2 c^4$$

or
$$E^2 = |\mathbf{p}|^2 c^2 + m^2 c^4$$

Notice: If $m=0$, we have $|\mathbf{p}| = E/c$

→ Momentum carried by massless bodies. This corresponds perfectly to energy and momentum carried by electromagnetic radiation (c.f. the Poynting vector.)

Not difficult to show that this relationship is an invariant: $E^2 = |\mathbf{p}|^2 c^2 + m^2 c^4$ in all Lorentz frames.

We can make this invariance a little clearer by recognizing that $E/c, p^x, p^y, p^z$ are, by their Lorentz transformation rule, components of a 4-vector.

We've briefly met this notion already in the context of the displacement 4-vector:

$$\Delta \vec{x} = \sum_{\mu=0}^3 \Delta x^\mu \vec{e}_\mu$$

$$\Delta x^0 = \Delta x^t = c \Delta t$$

$$\Delta x^1 = \Delta x^x = \Delta x$$

$$\Delta x^2 = \Delta x^y = \Delta y$$

$$\Delta x^3 = \Delta x^z = \Delta z$$

Already saw that under our rules for transforming components and unit vectors, this geometric object is invariant as we change between Lorentz frames:

$$\begin{aligned} \Delta \vec{x} &= \sum_{\mu=0}^3 \Delta x^\mu \vec{e}_\mu \\ &= \sum_{\mu'=0}^3 \Delta x^{\mu'} \vec{e}_{\mu'} \end{aligned}$$

We similarly define the "4-momentum" to be the geometric object in spacetime whose components transform between Lorentz frames according to the rules we have found:

$$\vec{p} = \sum_{\mu=0}^3 p^{\mu} \vec{e}_{\mu}$$

$$p^0 = E/c$$

$$p^x = p^x, \quad p^y = p^y, \quad p^z = p^z$$

Very useful definition and theorem:

Given two 4-vectors, $\vec{a} = \sum_{\mu=0}^3 a^{\mu} \vec{e}_{\mu}$ and $\vec{b} = \sum_{\mu=0}^3 b^{\mu} \vec{e}_{\mu}$, we define

$$\vec{a} \cdot \vec{b} = -a^t b^t + a^x b^x + a^y b^y + a^z b^z$$

Notice that Δs^2 , the invariant interval, is equal to $\Delta \vec{x} \cdot \Delta \vec{x}$. Likewise, we can see that $\vec{p} \cdot \vec{p} = -m^2 c^2$.

Theorem: For any two 4-vectors, $\vec{a} \cdot \vec{b}$, the scalar product $\vec{a} \cdot \vec{b}$ is a Lorentz invariant.

Forgot to mention on the previous page: The 4-momentum gives us a good tool to combine conservation of momentum and energy together. Our combined rule reads as follows:

Given N_i bodies that interact and produce N_f bodies, 4-momentum is conserved:

$$\sum_{i=1}^{N_i} \vec{p}_i^{\text{init}} = \sum_{i=1}^{N_f} \vec{p}_i^{\text{final}}$$

\vec{p}_i^{init} = initial 4-momentum of particle i

\vec{p}_i^{final} = final 4-momentum of particle i

Corollary: Given a 4-vector \vec{a} , the scalar product of \vec{a} with itself is a value that all observers agree upon. If

$$\vec{a} \cdot \vec{a} < 0$$

We say \vec{a} is "timelike": There exists a Lorentz frame in which \vec{a} points purely in the timelike direction. If

$$\vec{a} \cdot \vec{a} > 0$$

We say \vec{a} is "spacelike": There exists a Lorentz frame in which \vec{a} points purely in a spatial direction. If

$$\vec{a} \cdot \vec{a} = 0$$

\vec{a} is "lightlike" or "null": In all Lorentz frames, \vec{a} points along light cones.

Notice: $\vec{p} \cdot \vec{p} = -m^2 c^2$: \vec{p} is timelike if $m > 0$.

Another definition: $\vec{u} = \frac{1}{m} \vec{p}$

We call \vec{u} the "4-velocity." What does it mean? Look at its components:

$$u^t = p^t / m = E / mc = \gamma mc^2 / mc = \gamma c$$

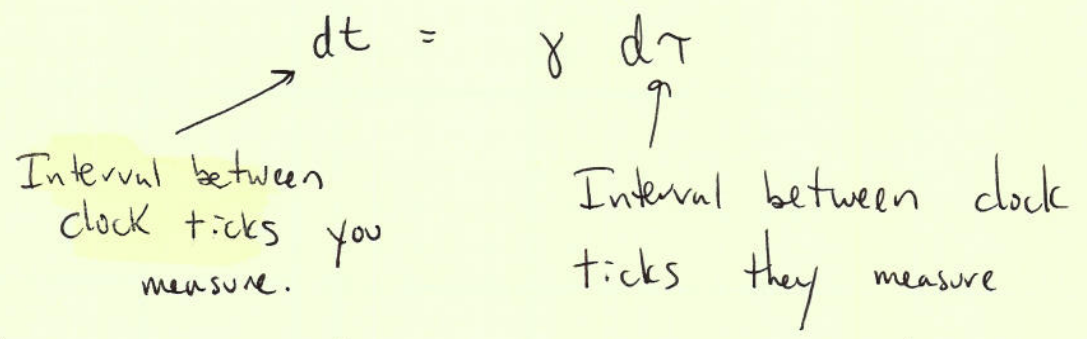
$$u^x = p^x / m = \gamma (\underline{u})^x$$

$$u^y = p^y / m = \gamma (\underline{u})^y$$

$$u^z = p^z / m = \gamma (\underline{u})^z$$

[1] Notation:
 $(\underline{u})^i$ = the i -component of 3-vector \underline{u} .

The spatial components look just like "normal" velocity, but with an annoying factor of γ . To interpret this, consider someone passing by you with 3-velocity \underline{u} . That person's clocks run slow according to you:



(In other words, the moving clock runs slow: In their time interval $d\tau$, we see $\gamma d\tau$ go past.)

Definition: $d\tau$ is an interval of "proper" time, time as measured by the person we say is moving.

"Proper" comes from French "propre," and refers to the fact that it is the moving observer's own time. (German sometimes calls it "eigenzeit.")

The key reason that this is useful is that $d\tau$, properly interpreted, is a Lorentz invariant:

All observers agree that the person in motion measures a time interval $d\tau$. Indeed, $d\tau$ is the time interval as measured in the frame in which the observer is at rest!

With this in mind, re-examine components of the

4-velocity: $u^x = \gamma \frac{dx}{dt} = \frac{dx}{d\tau}$

$$u^y = \gamma \frac{dy}{dt} = \frac{dy}{d\tau}$$

$$u^z = \gamma \frac{dz}{dt} = \frac{dz}{d\tau}$$

Last component: $u^t = c \frac{dt}{d\tau}$. Modulo the factor of c , this tells us how the two observers interpret intervals of time.

Compare with the displacement 4-vector:

$$\vec{u} = d\vec{x} / d\tau$$

$\rightarrow \vec{u}$ is spacetime displacement per unit proper time.

4-velocity versus 3-velocity

We now have two important ways to characterize a moving body's motion:

3-velocity: $\underline{u} = \frac{d\underline{x}}{dt}$ → Motion through SPACE - as seen by some observer - per unit TIME - as seen by some observer.

→ Clearly strongly depends on the frame in which it is measured!

4-velocity: $\underline{u} = \frac{d\underline{x}}{d\tau}$ → Motion through SPACETIME - as seen by all observers - per unit proper time - as agreed to by all observers.

→ Frame-independent, geometric object.

Components that a given observer assigns to 4-velocity are assembled from the 3-velocity:

$$u^0 = \gamma c$$

$$u^1 = \gamma (\underline{u})^x$$

$$u^2 = \gamma (\underline{u})^y$$

$$u^3 = \gamma (\underline{u})^z$$

$$\gamma = \left(1 - |\underline{u}|^2 / c^2\right)^{-1/2}$$

Key difference is how we describe them in different Lorentz frames:

4-vel: Lorentz transformation.

$$u^{\alpha'} = \sum_{\mu=0}^3 \Lambda^{\alpha'}_{\mu} u^{\mu}$$

$\Lambda^{\alpha'}_{\mu}$ \equiv elements of transformation matrix relating two Lorentz frames.

$$\vec{u} = \sum_{\mu=0}^3 u^{\mu} \vec{e}_{\mu} = \sum_{\alpha'=0}^3 u^{\alpha'} \vec{e}_{\alpha'}$$

\rightarrow same object, but different components.

3-vel: Velocity addition formula.

$$(\underline{u})^{x'} = \frac{(\underline{u})^x + v}{1 + (\underline{u})^x v / c^2}$$

$$(\underline{u})^{y'} = \frac{\underline{u}^y}{\gamma (1 + (\underline{u})^x v / c^2)}$$

$$(\underline{u})^{z'} = \frac{\underline{u}^z}{\gamma (1 + (\underline{u})^x v / c^2)}$$

\rightarrow Actually a different 3-vector!