

Universality of distribution functions in random matrix theory

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Overview

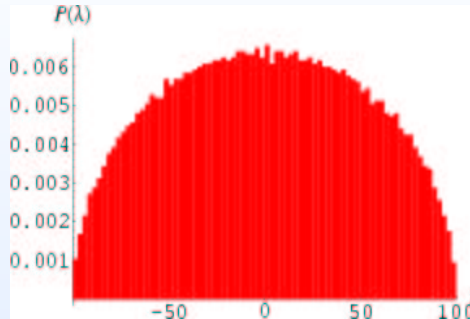
- ▲ **Universality**
 - ▲ **Local eigenvalue statistics**
 - ▲ **Fluctuations of the largest eigenvalue**
- ▲ **Connections outside RMT**
 - ▲ **Zeros of Riemann zeta function**
 - ▲ **Non-intersecting Brownian paths**
 - ▲ **Tiling problem**
- ▲ **Unitary ensembles**
 - ▲ **Determinantal point process**
 - ▲ **Precise formulation of universality**
 - ▲ **Universality in regular cases**
- ▲ **Universality classes in singular cases**
 - ▲ **Singular cases I and II and Painlevé equations**
 - ▲ **Spectral singularity**

Gaussian ensembles

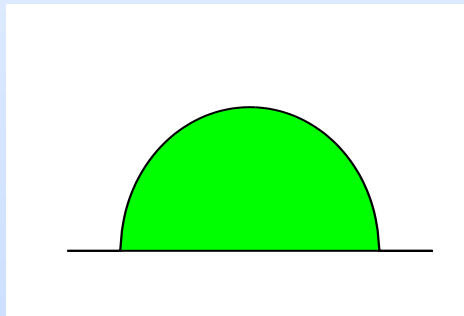
- ▲ Simplest ensembles are **Gaussian ensembles**.
- ▲ Matrix entries have normal distribution with mean zero. The entries are independent up to the constraints that are imposed by the symmetry class.
 - ▲ Gaussian Unitary Ensemble **GUE**: complex Hermitian matrices
 - ▲ Gaussian Orthogonal Ensemble **GOE**: real symmetric matrices
 - ▲ Gaussian Symplectic Ensemble **GSE**: self-dual quaternionic matrices
- ▲ Where are the eigenvalues?

Wigner's semi-circle law

- ▲ Histogram of eigenvalues of large Gaussian matrix, size $10^4 \times 10^4$



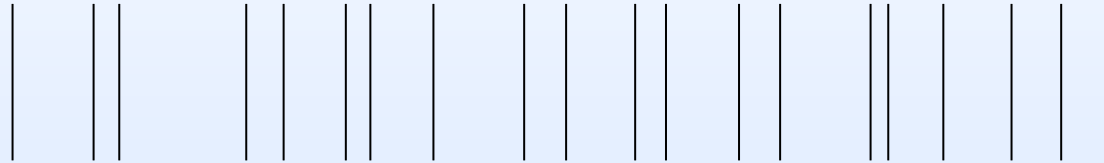
- ▲ After scaling of eigenvalues with a factor \sqrt{n} , there is a limiting mean eigenvalue distribution, known as **Wigner's semi-circle law**



- ▲ This is special for Gaussian ensembles (non-universal). Other limiting distributions for Wishart ensembles, Jacobi ensembles,...

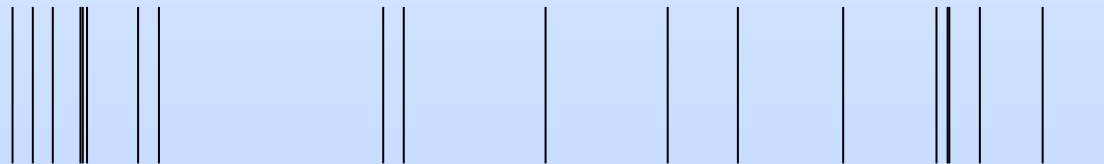
Universality 1: Local eigenvalue statistics

- ▲ Global statistics of eigenvalues depend on the particular random matrix ensemble in contrast to **local statistics**. Distances between consecutive eigenvalues show regular behavior.
- ▲ Rescale eigenvalues around a certain value so that mean distance is one.



plot shows only a few rescaled eigenvalues of a very large GUE matrix

- ▲ This is the same behavior as seen in energy spectra in quantum physics.
- ▲ The **repulsion** between neighboring eigenvalues is very different from Poisson spacings.



Universality 1: Local eigenvalue statistics

- ▲ This local behavior of eigenvalues is not special for GUE.

- ▲ It holds for large class of unitary ensembles
$$\frac{1}{Z_n} e^{-n \text{Tr} V(M)} dM$$

these are ensembles that have the **same symmetry property** as GUE.

Deift, Kriecherbauer, McLaughlin, Venakides, Zhou (1999)

- ▲ Local eigenvalue statistics is different for GOE and GSE which have **different symmetry properties**. Proof of universality for orthogonal and symplectic ensembles is more recent result Deift, Gioev (arxiv 2004)
- ▲ Universality fails at special points, such as **end points** of the spectrum, or points where **eigenvalue density vanishes**.
 - ▲ This gives rise to new universality classes.

Universality 2: Largest eigenvalue

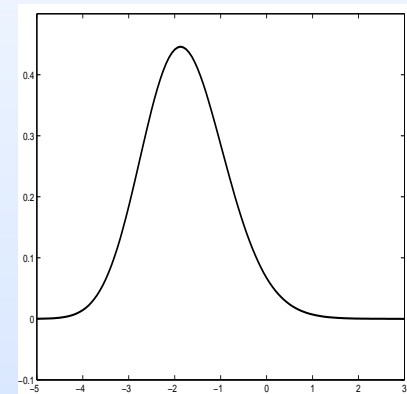
- ▲ Fluctuations of the **largest eigenvalues** of random matrices also show a universal behavior (depending on the symmetry class).
- ▲ For $n \times n$ GUE matrix, the largest eigenvalue grows like $\sqrt{2n}$ and has a standard deviation of the order $n^{-1/6}$.

- ▲ Centered and rescaled largest eigenvalue

$$\sqrt{2n}^{1/6} \left(\lambda_{\max} - \sqrt{2n} \right)$$

converges in distribution as $n \rightarrow \infty$ to a random variable with the **Tracy-Widom distribution**, described by **Tracy, Widom** in 1994.

- ▲ Same limit holds generically for unitary random matrix ensembles.
- ▲ Different TW-distributions for orthogonal and symplectic ensembles.



Density of Tracy-Widom distribution. The density is non-symmetric with top at -1.8 and different decay rates for $x \rightarrow +\infty$ and $x \rightarrow -\infty$

Tracy-Widom distribution

- ▲ There is **no simple formula** for the Tracy-Widom distribution $F(s)$.
- ▲ First formula is as a **Fredholm determinant**:

$$F(s) = \det(I - A_s)$$

where A_s is the integral operator acting on $L^2(s, \infty)$ with kernel

$$\frac{\text{Ai}(x) \text{Ai}'(y) - \text{Ai}'(x) \text{Ai}(y)}{x - y} \quad \text{Airy kernel}$$

and Ai is the Airy function.

- ▲ **Second formula**

$$F(s) = \exp\left(-\int_s^\infty (x - s)q^2(x)dx\right)$$

where $q(s)$ is a special solution of the **Painlevé II equation**

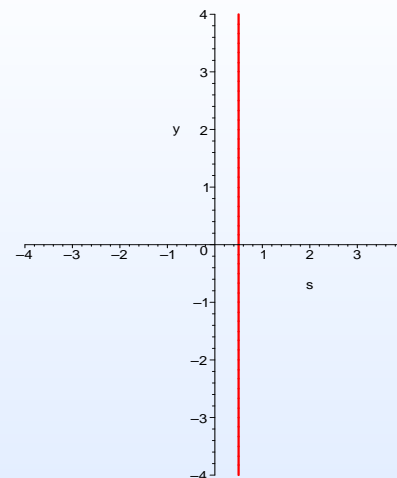
$$q''(s) = sq(s) + 2q^3(s)$$

Limit laws outside RMT

- ▲ Distribution functions of random matrix theory appear in various other domains of mathematics and physics.
 - ▲ **Number theory**
 - ▲ Riemann zeta-function, L -functions, ...
 - ▲ **Representation theory**
 - ▲ Young tableaux, large classical groups, ...
 - ▲ **Random combinatorial structures**
 - ▲ random permutations, random tilings, ...
 - ▲ **Growth models** in statistical physics
 - ▲ last passage percolation, polynuclear growth, ...
- ▲ ..., as well as in applications in statistics, finance, information theory, ...

Riemann zeta function

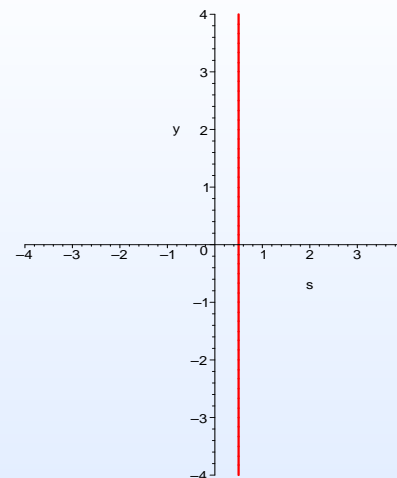
- ▲ The zeta function $\zeta(s) = \sum_{k=1}^{\infty} \frac{1}{k^s}$ has an analytic continuation to the complex plane.
- ▲ The non-trivial zeros of the zeta function are believed to be on the line $\operatorname{Re} s = 1/2$. (Riemann hypothesis)



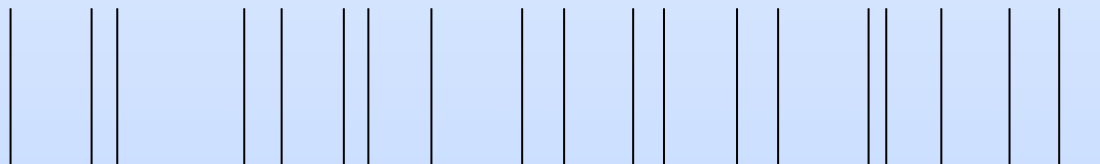
- ▲ **Computational evidence:** no non-real zeros have been found off the critical line.
- ▲ 1,500,000,000 zeros have been found on the critical line.

Riemann zeta function

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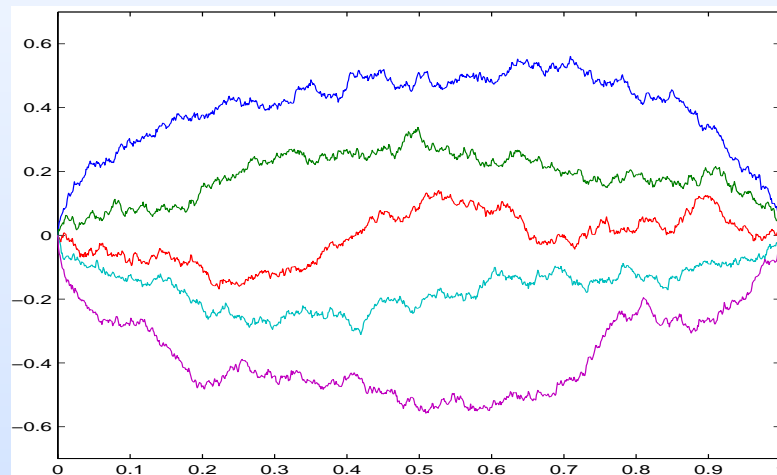
- ▲ **More computational evidence:** Spacings between consecutive zeros on the critical line $\operatorname{Re} s = 1/2$ (after appropriate scaling) show the same behavior as the spacings between eigenvalues of a large GUE matrix



Zeros of $\zeta(s)$ on the critical line have the same local behavior as the eigenvalues of a large random matrix

Non-intersecting Brownian motion paths

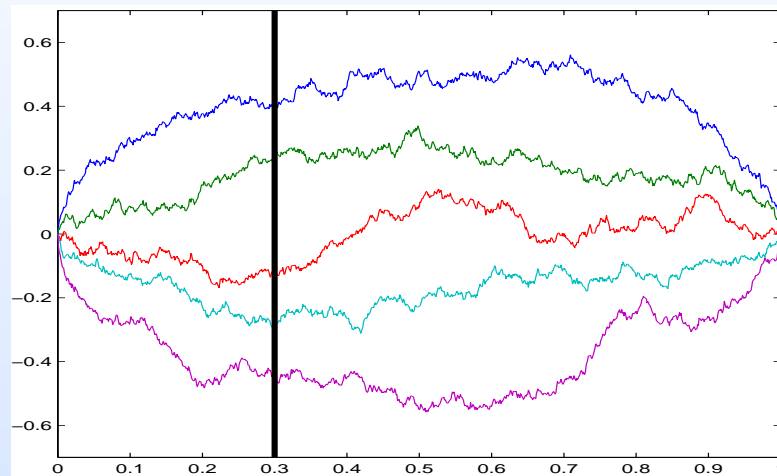
- ▲ Take n independent 1-dimensional Brownian motions with time in $[0, 1]$ conditioned so that:
 - ▲ All paths start and end at the same point.
 - ▲ The paths **do not intersect** at any intermediate time.



Five non-intersecting Brownian bridges

Non-intersecting Brownian motion

- ▲ **Remarkable fact:** At any intermediate time the positions of the paths have **exactly the same distribution** as the eigenvalues of an $n \times n$ **GUE** matrix (up to a scaling factor).

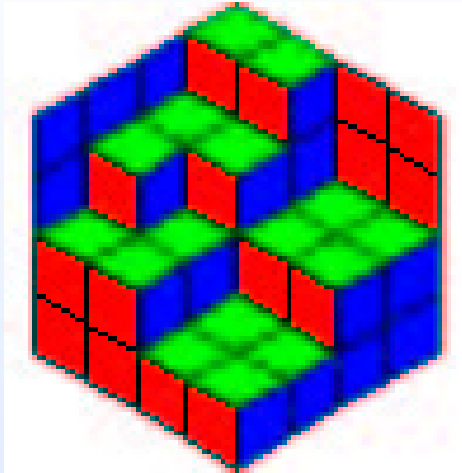


Positions of five non-intersecting Brownian paths behave the same as the eigenvalues of a 5×5 **GUE** matrix

- ▲ This interpretation is basic for the connection of random matrix theory with growth models of statistical physics.

A random tiling problem

- ▲ Hexagonal tiling with rhombi.

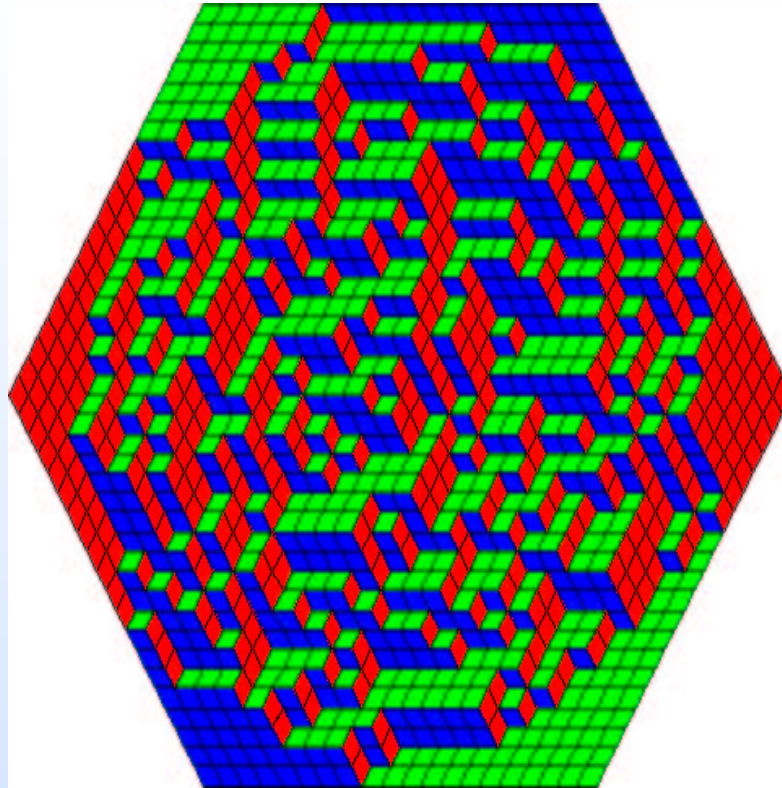


- ▲ May also be viewed as **packing of boxes** in a corner.

- ▲ Take a tiling at random.

- ▲ What does a typical tiling look like, if the number of rhombi increases?

Typical random tiling

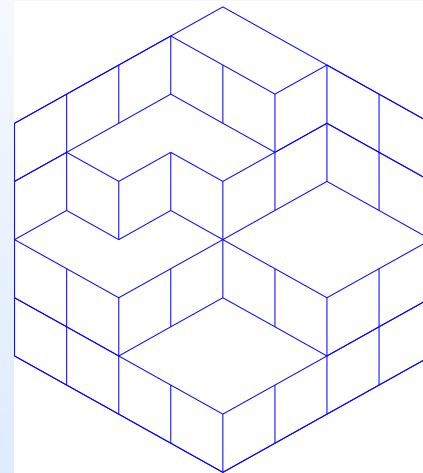
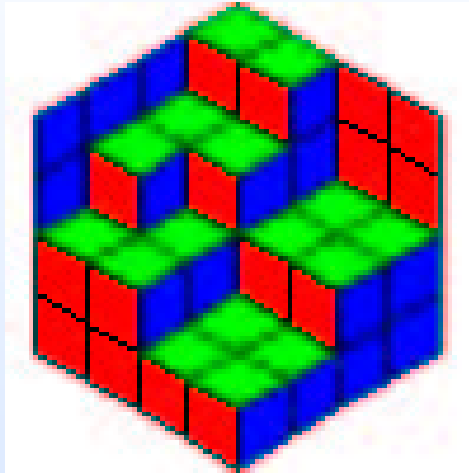


▲ **Observation:**

- ▲ frozen regions near the corners,
- ▲ disorder in the center.

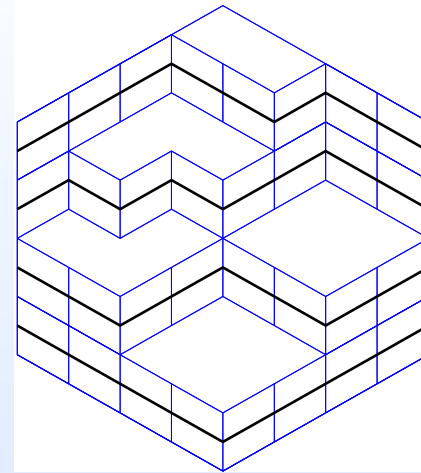
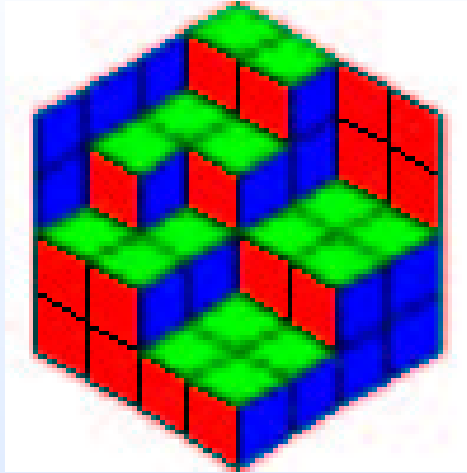
Non-intersecting random walk

- ▲ Consider only blue and red rhombi.



Non-intersecting random walk

- ▲ We can connect the left and the right with non-intersecting paths.



- ▲ A random tiling is the same as a number of non-intersecting random walks.
- ▲ As size increases: **Tracy-Widom distribution** governs the transition between frozen region and disordered region.

Baik, Kriecherbauer, McLaughlin, Miller, (arxiv 2003)

Other examples

- ▲ Longest increasing subsequence of random permutations

Baik, Deift, Johansson (1999)

- ▲ Polynuclear growth model (PNG)

Totally asymmetric exclusion process (TASEP)

Praehofer, Spohn, Ferrari

Imamura, Sasamoto

- ▲ Buses in Cuernavaca, Mexico

Krbalek, Seba

Baik, Borodin, Deift, Suidan

- ▲ Airplane boarding problem

Bachmat

Universality classes in unitary ensembles

▲ Probability measure on $n \times n$ Hermitian matrices

$$\frac{1}{Z_n} e^{-n \operatorname{Tr} V(M)} dM$$

where $dM = \prod_j dM_{jj} \prod_{j < k} d \operatorname{Re} M_{jk} d \operatorname{Im} M_{jk}$

▲ This is a Gaussian ensemble for $V(x) = \frac{1}{2}x^2$

▲ Joint eigenvalue density

$$\frac{1}{\tilde{Z}_n} \prod_{i < j} |x_i - x_j|^2 \prod_{j=1}^n e^{-nV(x_j)}$$

wher

Determinantal point process

- ▲ It is special about unitary ensembles that the eigenvalues follow a **determinantal point process**. This means means that there is a kernel $K_n(x, y)$ so that all eigenvalue correlation functions are expressed as **determinants**

$$\mathcal{R}_m(x_1, x_2, \dots, x_k) = \det [K_n(x_i, x_j)]_{i,j=1,\dots,m}$$

- ▲ $\int_a^b K_n(x, x) dx$ is expected number of eigenvalues in $[a, b]$

- ▲ $\int_a^b \int_c^d \begin{vmatrix} K_n(x, x) & K_n(x, y) \\ K_n(y, x) & K_n(y, y) \end{vmatrix} dx dy$ is expected number of pairs of eigenvalues in $[a, b] \times [c, d]$, etc.

Orthogonal polynomial kernel

- ▲ Let $P_{k,n}(x)$ be the k th degree monic orthogonal polynomial with respect to $e^{-nV(x)}$

$$\int_{-\infty}^{\infty} P_{k,n}(x)P_{j,n}(x)e^{-nV(x)}dx = h_{k,n}\delta_{j,k}.$$

- ▲ Correlation kernel is equal to

$$\begin{aligned} K_n(x, y) &= e^{-\frac{1}{2}n(V(x)+V(y))} \sum_{k=0}^{n-1} \frac{1}{h_{k,n}} P_{k,n}(x)P_{k,n}(y) \\ &= e^{-\frac{1}{2}n(V(x)+V(y))} \frac{h_{n,n}}{h_{n-1,n}} \frac{P_{n,n}(x)P_{n-1,n}(y) - P_{n-1,n}(x)P_{n,n}(y)}{x - y} \end{aligned}$$

Christoffel-Darboux formula

Asymptotical questions

- ▲ All information is contained in the correlation kernel K_n .
- ▲ Asymptotic questions deal with the **global regime**

$$\rho_V(x) = \lim_{n \rightarrow \infty} \frac{1}{n} K_n(x, x)$$

- ▲ and with the **local regime**
 - ▲ Choose x^* and center and scale eigenvalues around x^*

$$\lambda \mapsto (cn)^\gamma (\lambda - x^*)$$

- ▲ This is a determinantal point process with rescaled kernel

$$\frac{1}{(cn)^\gamma} K_n \left(x^* + \frac{x}{(cn)^\gamma}, x^* + \frac{y}{(cn)^\gamma} \right)$$

- ▲ Determine γ and calculate limit of rescaled kernel. Limits turn out to be universal, depending only on the nature of x^* in the global regime.

Global regime

- ▲ In limit $n \rightarrow \infty$, the mean eigenvalue density has a limit $\rho = \rho_V$ which minimizes

$$\iint \log \frac{1}{|x-y|} \rho(x)\rho(y) dx dy + \int V(x)\rho(x) dx$$

among density functions $\rho \geq 0$, $\int \rho(x) dx = 1$.

- ▲ Equilibrium conditions

$$2 \int \log \frac{1}{|x-y|} \rho(y) dy + V(x) = \text{const on support of } \rho$$

$$2 \int \log \frac{1}{|x-y|} \rho(y) dy + V(x) \geq \text{const outside support}$$

Regular and singular cases

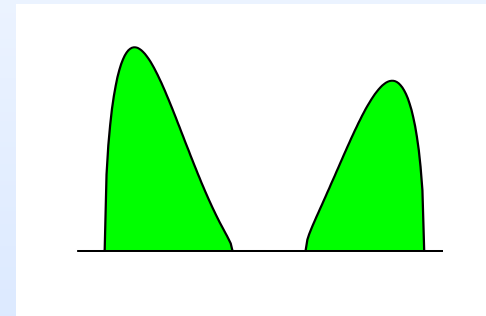
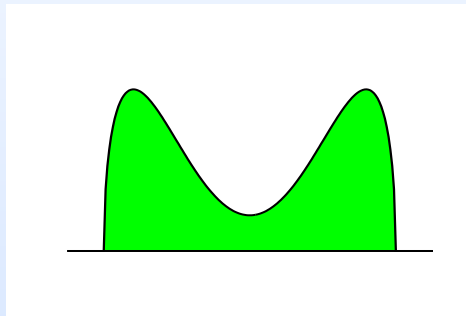
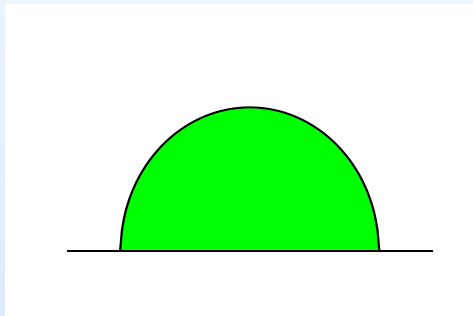
▲ If V is real analytic, then

Deift, Kriecherbauer, McLaughlin (1998)

▲ $\text{supp}(\rho)$ is a finite union of disjoint intervals,

▲ $\rho(x)$ is analytic on the interior of each interval

▲ $\rho(x) \sim |x - a|^{2k+1/2}$ at an endpoint a for some $k = 0, 1, 2, \dots$

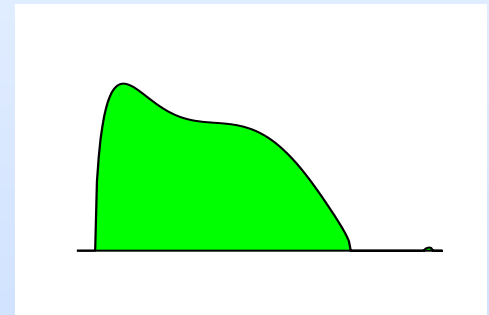
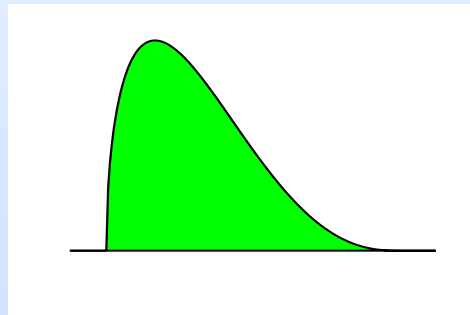
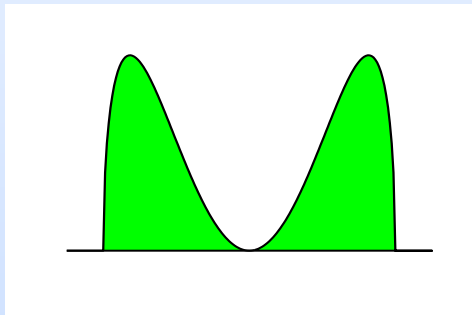


▲ Regular case: positive in interior, square root vanishing at endpoints, and strict inequality in

$$2 \int \log \frac{1}{|x - y|} \rho(y) dy + V(x) > \text{const} \text{ outside the support of } \rho$$

Singular cases

- ▲ Singular cases correspond to possible change in number of intervals if parameters in the external field V change.
 - ▲ Singular case I: ρ vanishes at an interior point
 - ▲ Singular case II: ρ vanishes to higher order than square root at an endpoint.
 - ▲ Singular case III: equality in equilibrium inequality somewhere outside the support



Local regime

▲ Limit of rescaled kernel

$$\frac{1}{(cn)^\gamma} K_n \left(x^* + \frac{x}{(cn)^\gamma}, x^* + \frac{y}{(cn)^\gamma} \right)$$

- ▲ For x^* in the bulk, we take $\gamma = 1$, $c = \rho(x^*)$, and limit is the **sine**

kernel

$$\frac{\sin \pi(x - y)}{\pi(x - y)}$$

Pastur, Shcherbina (1997), Bleher, Its (1999)

Deift, Kriecherbauer, McLaughlin, Venakides, Zhou (1999)

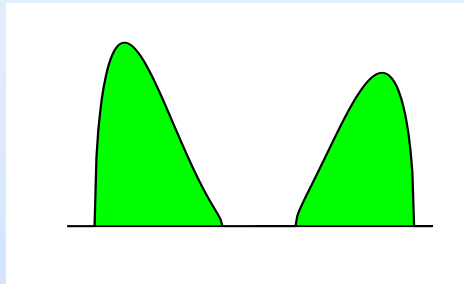
- ▲ Scaling limit of kernel follows from detailed asymptotics of the orthogonal polynomials $P_{n,n}$ and $P_{n-1,n}$ as $n \rightarrow \infty$, which is available in the GUE case since then the orthogonal polynomials are **Hermite polynomials**.
- ▲ For more general cases a powerful new technique was developed: **steepest descent analysis of Riemann-Hilbert problems**

Regular endpoint

- ▲ For a regular endpoint of the support, we take $\gamma = \frac{2}{3}$, and the scaling limit is the **Airy kernel**

$$\frac{\text{Ai}(x) \text{Ai}'(y) - \text{Ai}'(x) \text{Ai}(y)}{x - y}$$

- ▲ This always gives rise to the Tracy Widom distribution for the fluctuations of the extreme eigenvalues.



Singular cases

Unitary ensemble

$$\frac{1}{Z_n} e^{-n \operatorname{Tr} V(M)} dM$$

Reference point x^*

- ▲ Regular interior point:
- ▲ Regular endpoint:
- ▲ Singular case I:
- ▲ Singular case II:
- ▲ Singular case III:

Singular cases

Unitary ensemble

$$\frac{1}{Z_n} e^{-n \operatorname{Tr} V(M)} dM$$

Reference point x^*

- ▲ Regular interior point: sine kernel
- ▲ Regular endpoint: Airy kernel
- ▲ Singular case I:
- ▲ Singular case II:
- ▲ Singular case III:

Singular cases

Unitary ensemble

$$\frac{1}{Z_n} e^{-n \operatorname{Tr} V(M)} dM$$

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- ▲ Regular interior point: sine kernel
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- ▲ Singular case II:
- ▲ Singular case III:

Singular cases

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- ▲ Singular case II: second member of **Painlevé I hierarchy** Claeys, Vanlessen (in progress)
- ▲ Singular case III:

Singular cases

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- ▲ Singular case III: ??

Spectral singularity

Extra factor in random matrix model

$$\frac{1}{Z_n} |\det M|^{2\alpha} e^{-n \operatorname{Tr} V(M)} dM$$

The extra factor does not change the global behavior but it does change the local behavior around the reference point $x^* = 0$

- ▲ Regular interior point:
- ▲ Regular endpoint:

- ▲ Singular case I:
- ▲ Singular case II:
- ▲ Singular case III:

Spectral singularity

Extra factor in random matrix model

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▲ Regular interior point: **Bessel kernel**

AK, Vanlessen (2003)

▲ Regular endpoint:

▲ Singular case I:

▲ Singular case II:

▲ Singular case III:

Spectral singularity

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▲ Regular interior point: **Bessel kernel** AK, Vanlessen (2003)

▲ Regular endpoint: general **Painlevé II** with parameter $2\alpha + \frac{1}{2}$

$$q'' = sq + 2q^3 - 2\alpha - \frac{1}{2}$$

Its, AK, Östenson (in progress)

▲ Singular case I:

▲ Singular case II:

▲ Singular case III:

Spectral singularity

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Its, AK, Östenson (in progress)

▲ Singular case I: **Painlevé II** with parameter α Claeys, AK, Vanlessen (arxiv 2005)

▲ Singular case II:

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▲ Singular case II: ??

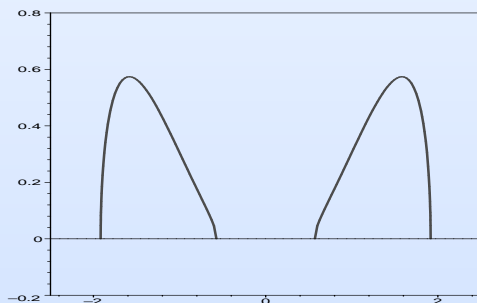
▲ Singular case III: ??

Singular case I

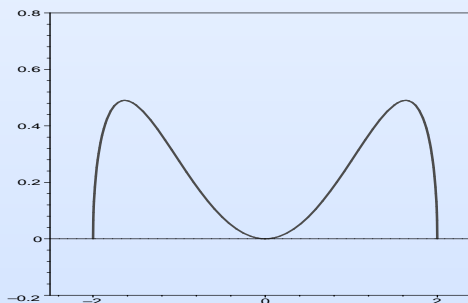
- ▲ Quartic external field $V(x) = \frac{1}{4}x^4 - x^2$ is simplest singular case I.
- ▲ Transition from two-interval to one-interval. If

$$V_t(x) = \frac{1}{t} \left(\frac{1}{4}x^4 - x^2 \right)$$

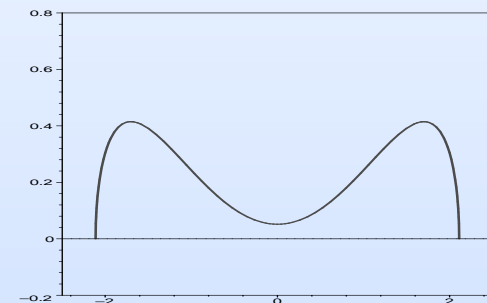
then for $t < 1$: two intervals, and for $t > 1$: one interval



$t = 0.7$



$t = 1$



$t = 1.5$

- ▲ Consider singular case in **double scaling limit** where we rescale eigenvalues $\lambda \mapsto (c_1 n)^{1/3} (\lambda - x^*)$ and we let $t \rightarrow 1$ as $n \rightarrow \infty$ such that $n^{2/3}(t - 1) = c_2 s$

Double scaling limit in singular case I

- ▲ One-parameter family of limiting kernels, depending on s , but independent of V

$$\frac{\Phi_1(x; s)\Phi_2(y; s) - \Phi_2(x; s)\Phi_1(y; s)}{2\pi i(x - y)}$$

- ▲ Φ_1 and Φ_2 satisfy a differential equation

$$\frac{d}{dx} \begin{pmatrix} \Phi_1 \\ \Phi_2 \end{pmatrix} = \begin{pmatrix} -4ix^2 - i(s + 2q^2) & 4xq + 2ir \\ 4xq - 2ir & 4ix^2 + i(s + 2q^2) \end{pmatrix} \begin{pmatrix} \Phi_1 \\ \Phi_2 \end{pmatrix}$$

with parameters s, q and r that are such that $q = q(s)$ satisfies **Painlevé II**:
 $q'' = sq + 2q^3$ and $r = r(s) = q'(s)$. for critical quartic V : [Bleher, Its \(2003\)](#)

for real analytic V : [Claeys, AK \(arxiv 2005\)](#)

for less smooth, even V : [Shcherbina \(arxiv 2006\)](#)

- ▲ Our proof uses the fact that Φ_1 and Φ_2 solve a RH problem that can be used as a local parametrix in the steepest descent analysis.