

# Lecture 1: Random Walks and 1-D Diffusion

- A localized high concentration of a substance will spread out to make the concentration spatially homogeneous (2<sup>nd</sup> Law of thermo)
- Molecules "move" from high to low density
- Individual particles are randomly battered by solvent molecules  $\Rightarrow$  causes a random walk
- Many random walks made by many molecules is diffusion
- Diffusion is dominant at less  $\sim 1 \mu\text{m}$  length scale  $\rightarrow$

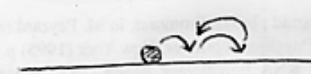
Show Applet Movie of Random Walk  
Beach ball at concert example


seems remarkable that  $\text{H}_2\text{O}$  collisions perturb the sphere  
cannot see little jerks, but can see the large displacements  
Einstein 1905: while distracted from his thesis


1<sup>st</sup> observation: pollen particles Robert Brown 1828  
the concept of molecules was controversial  
'Brownian motion' was mistaken for a life process

Single molecule measurements: Random walk

Bulk measurements: Diffusion

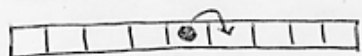
1D  transcription factor sliding on DNA  
growth of actin filaments

2D  protein traveling in lipid membrane

3D  protein traveling in cytoplasm  
cell motility

# One-Dimensional Random Walk

(2)



- flip a coin to determine whether you go left or right
- how far are you likely to go in a random walk
- cannot solve for one walk, but we can determine the average displacement

Enumeration of the possibilities for  $N=3$  steps:

$x=0$ $L \leftarrow$ step length	Distance ( $x_3$ )	$\frac{x_3^2}{}$
	+3L	+9L <sup>2</sup>
	+1	+1
	+1	+1
	-1	+1
	+1	+1
	-1	+1
	-1	+1
	-3	+9
	$\langle x_3 \rangle = 0$	$\langle x_3^2 \rangle = 3L^2$ (mean-square displacement)

for  $N$  steps:

$$\langle x_N \rangle = 0 \quad \langle x_N^2 \rangle = NL^2$$

- the mean displacement of a random walk is zero
- that does not mean that you don't go anywhere!

• Add time to the process

$$N = \frac{t}{\Delta t} \leftarrow \text{total time}$$

$$N = \frac{t}{\Delta t} \leftarrow \text{amount of time per step}$$

- Define diffusion constant  $D \equiv \left[ \frac{\text{length}^2}{\text{time}} \right]$  eg,  $\frac{\mu\text{m}^2}{\text{s}}$

$$D = \frac{\langle x_N^2 \rangle}{2t} = \frac{NL^2}{2t} = \frac{L^2}{2\Delta t}$$

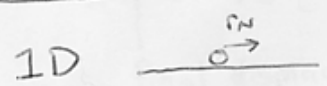
$$\langle x_N^2 \rangle = 2Dt \quad (\text{mean-square displacement increases linearly with time})$$

$\uparrow$  completely arbitrary/convenient later

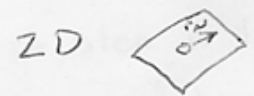
Any individual random walk will not conform to the diffusion law, even approximately!

- we will observe excursions of any size  $X$  as long as we are willing to wait a time  $\sim O\left(\frac{X^2}{2D}\right)$
- can easily experimentally measure  $D$ : wait time  $t$ , measure displacement, and repeat

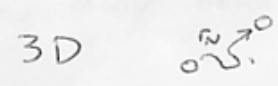
Higher-dimensions



$$\langle r_N^2 \rangle = \langle x_N^2 \rangle = 2Dt$$



$$\langle r_N^2 \rangle = \langle x_N^2 \rangle + \langle y_N^2 \rangle = 4Dt$$



$$\langle r_N^2 \rangle = \langle x_N^2 \rangle + \langle y_N^2 \rangle + \langle z_N^2 \rangle = 6Dt$$

Example

About how long does it take for a sudden supply of sugar molecules to spread uniformly?

$\langle r_N^2 \rangle = 6Dt$        $\overset{\text{sugar}}{D} \sim 500 \frac{\mu\text{m}^2}{\text{s}}$

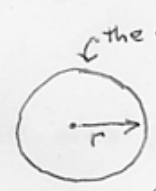
$$t \sim \frac{\langle r_N^2 \rangle}{6D}$$

$$t_{\text{coli}} \sim \frac{(1 \mu\text{m})^2}{6(500 \frac{\mu\text{m}^2}{\text{s}})} = 0.0003 \text{ s}$$

$$t_{\text{neut}} \sim \frac{(10)^2}{6(500)} = 0.03 \text{ s}$$

$$t_{\text{nerve}} \sim \frac{(1000)^2}{6(500)} = 300 \text{ s}$$

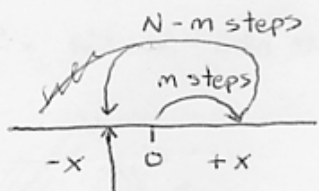
Compare with  $D \sim 10 \frac{\mu\text{m}^2}{\text{s}}$  for an average-sized protein in the cytoplasm.



- $r = 1 \mu\text{m}$  for bacterium
- $r = 10 \mu\text{m}$  for neutrophil
- $r = 1000 \mu\text{m}$  for nerve cell

# Explicit calculation of the probability distribution

- 1D random walk



final displacement probability  $P(m, N)$

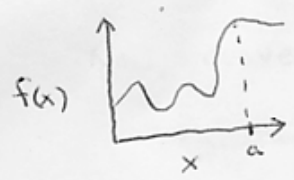
- each step is like flipping a coin  $P_+ = P_- = 1/2$
- How many heads in 100 flips?

- Binomial probability distribution:  $P(m, N) = \binom{N}{m} \left(\frac{1}{2}\right)^N \frac{N!}{m!(N-m)!}$

- factorials are mathematically intractable
- want to approximate function P

- Calculate a Taylor series around the most probable end point  $m^*$

reminder: purpose of a Taylor Series is to approximate a function with a simpler mathematical expression



$$f(x) \approx f(a) + \left(\frac{df}{dx}\right)_{x=a} (x-a) + \frac{1}{2} \left(\frac{d^2f}{dx^2}\right)_{x=a} (x-a)^2 + \dots$$

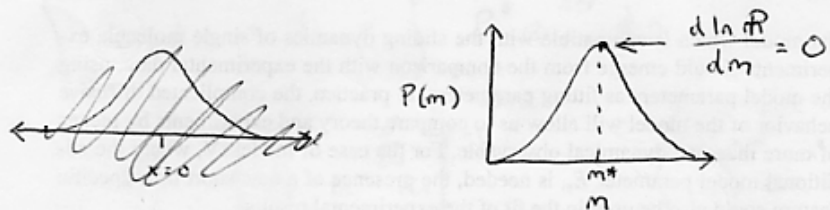
- Approximate  $P(m, N)$  - work w/ log is easier

$$\ln P(m) = \ln P(m^*) + \left(\frac{d \ln P}{dm}\right)_{m^*} (m-m^*) + \frac{1}{2} \left(\frac{d^2 \ln P}{dm^2}\right)_{m^*} (m-m^*)^2 + \dots$$

Need to evaluate the derivatives  $\left(\frac{d \ln P}{dm}\right)$  and  $\left(\frac{d^2 \ln P}{dm^2}\right)$  (4)

Take log of (1) and use Stirling's Approximation ( $\ln N! \approx N \ln N$ )

$$\ln P = N \ln N - m \ln m - (N-m) \ln (N-m) + N \ln \left(\frac{1}{2}\right)$$



$$(A) \quad \left(\frac{d \ln P}{dm}\right)_{m^*} = 0 = -1 - \ln m^* + \ln (N-m^*) + 1$$

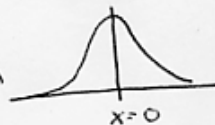
solve for  $m^*$ ,  $m^* = \frac{N}{2}$  (makes sense, right?)

$$(B) \quad \left(\frac{d^2 \ln P}{dm^2}\right)_{m^* = N/2} = \left(-\frac{1}{m} - \frac{1}{N-m}\right)_{m^* = N/2} = -\frac{4}{N}$$

Put (A) & (B) into Taylor Series

$$\ln P = \ln P^* + 0 + \frac{1}{2} \left(-\frac{4}{N}\right) (m-m^*)^2$$

$$P = P^* e^{-\frac{2(m-m^*)^2}{N}} \quad \left. \vphantom{P} \right\} \text{- form of Gaussian Distribution}$$



Now, convert from  $m$  to  $x$

~~forward~~ displacement progress is forward-reverse steps

$$\lambda = m - (N-m) = 2m - N$$

$$m = \frac{x+N}{2}$$

$$P(x) = P^* e^{-\frac{x^2}{2N}}$$

↑ at  $x=0$



• Now, calculate  $P^*$

⑤

- this is a probability distribution that needs to be normalized

$$\int_{-\infty}^{\infty} P(x) dx = 1 = \int_{-\infty}^{\infty} P^* e^{-\frac{x^2}{2N}} dx$$

$$P^* = \left[ \int_{-\infty}^{\infty} e^{-\frac{x^2}{2N}} dx \right]^{-1}$$

(can look up  $\int_{-\infty}^{\infty} e^{-ax^2} dx = \sqrt{\frac{\pi}{a}}$ )

$$= \frac{1}{\sqrt{2\pi N}} = (2\pi N)^{-\frac{1}{2}}$$

$$P(x) = \frac{1}{\sqrt{2\pi N}} e^{-\frac{x^2}{2N}}$$

a Gaussian of this form  $\langle x^2 \rangle = N$

from the Diffusion law

$$\langle x^2 \rangle = 2Dt = N$$

$$P(x) = \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}}$$

- can calculate  $D$  if you have a distribution  $P(x)$  at ~~time~~ of displacements at time  $t$