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Environ. Sci. Technol. Lett., **Just Accepted Manuscript** • DOI: 10.1021/acs.estlett.5b00303 • Publication Date (Web): 03 Nov 2015

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Scalable and Continuous Water Deionization by Shock Electrodialysis

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Abstract

1
2 Rising global demand for potable water is driving innovation in water treatment
3 methods. Shock electrodialysis is a recently proposed technique that exploits deioniza-
4 tion shock waves in porous media to purify water. In this letter, we present the first
5 continuous and scalable shock electrodialysis system and demonstrate the separation
6 of sodium, chloride, and other ions from a feed stream. Our prototype continuously
7 removes over 99% (and up to 99.99%) of salt from diverse electrolytes over a range
8 of concentrations (1 mM, 10 mM, and 100 mM). The desalination data collapses with
9 dimensionless current, scaled to charge advection in the feed stream. Enhanced water
10 recovery with increasing current (up to 79%) is a fortuitous discovery, which we at-

11 tribute to electro-osmotic pumping. The results suggest the feasibility of using shock
12 electrodialysis for practical water purification applications.

13 **Introduction**

14 Access to potable water is a critical global challenge. It is estimated that over one billion
15 people currently do not have reliable access to clean and safe water.¹ Unfortunately, the
16 purification of non-potable water is infrastructure-intensive and expensive, as many countries
17 currently supplement their water supply with seawater desalination by large-scale reverse
18 osmosis (RO), or in some cases electrodialysis (ED).^{2,3} Moreover, widespread interest in
19 purifying wastewater has intensified with the ascent of hydraulic fracturing for shale gas
20 extraction, which is difficult to treat due to the presence of large amounts of dissolved salts.⁴⁻⁶
21 For these and other applications, extensive research is underway to develop improved water
22 treatment methods, and the unique capabilities of electrochemical systems have attracted
23 renewed attention.⁷⁻¹⁰

24 One of the new electrochemical approaches to water purification is shock electrodial-
25 ysis (SED),¹¹⁻¹³ which is based on the emerging science of deionization shocks in porous
26 media.¹⁴⁻²³ The SED process involves flowing feed water through a weakly charged porous
27 slab with micron-sized pores that is placed between two ion-selective elements, such as
28 ion-exchange membranes or electrodes. When current is passed through the device, zones
29 of ion depletion and enrichment are formed in order to maintain electroneutrality near the
30 ion-selective surfaces. In the classical picture, a diffusion-limited current is reached whenever
31 the salt concentration approaches zero, but it is well established that overlimiting currents
32 are possible in *bulk* electrolytes due to electrokinetic or electrochemical phenomena,¹⁰ such
33 as the Rubinstein-Zaltzman electro-osmotic instabilities²⁴⁻²⁷ or current induced membrane
34 discharge.²⁸ Recently, it has been shown that two new mechanisms¹⁸ – surface conduction
35 (electromigration)^{19,20} and surface convection (electro-osmosis)^{21,29,30} – are responsible for

36 over-limiting current when the electrolyte is *confined* in a microchannel^{18,23,31} or porous
37 medium,^{12,32} whose surface charge is opposite to the active ionic species. The transient re-
38 sponse to an applied overlimiting current (discovered first in 2009^{15,16}) is the propagation of a
39 deionization shock wave through the microchannel or porous medium with a sharp boundary
40 between concentrated and depleted zones.^{14,17,19,20,33} Stationary deionization shocks were first
41 observed at nanochannel junctions³⁴ and can cause localized seawater desalination within a
42 microchannel.²² For SED in porous media,^{11,12} the flow of water can also be separated into
43 brine and depleted streams by a physical splitter that is placed within the location of the
44 shock in a pressure-driven cross flow. The basic physics of SED was recently demonstrated
45 in a non-continuous copper electrodeposition cell that was able to reduce the concentration
46 of copper sulfate by five orders of magnitude in two passes¹² and perform other separations
47 and disinfection¹³ through a silica glass frit.

48 In this letter, we report the first continuous and scalable SED system for arbitrary feed
49 streams. We characterized its voltage response using IV-curves and tested its capability to
50 desalinate NaCl solutions at concentrations of 1 mM, 10 mM, and 100 mM, as well as KCl,
51 KNO₃, and Na₂SO₄ solutions at 10 mM. We also made the fortuitous discovery that electro-
52 osmotic (EO) pumping leads to greatly enhanced water recoveries, in contrast to existing
53 theoretical SED models.^{12,14,15,17-19}

54 **Materials and Methods**

55 Following previously introduced concepts,^{11,12,19} the prototype was designed to be continuous
56 and scalable with a novel cross-flow architecture (Fig. 1(a)) that employs two identical ion-
57 exchange membranes and a porous frit, which acts as a “leaky membrane” of the same
58 polarity. When an overlimiting current is applied to the system, a deionization shock (sharp
59 concentration gradient) propagates away from the cathode-side membrane and bends in the
60 imposed cross flow. The deionized solution in its wake is extracted by splitting the flow

61 into two streams at sufficiently high currents (or low flow rates) for the shock to span the
62 fresh water outlet. This structure could be easily repeated by making a stack of alternating
63 glass frits and membranes; thereby easily scaling the device. Some pictures of the device are
64 shown in Fig. 1(b).

65 The main body of the prototype was made from extruded acrylic sheets (W. W. Grainger)
66 and GORE® gasket sheets. These materials were cut using a laser cutter to obtain 2" × 2" × 1.5"
67 overall device dimensions. The electrodes were made of platinum mesh connected to plat-
68 inum wires (Alfa Aesar). The porous material was a silica glass frit (Adams & Chittenden
69 Scientific Glass, ultra-fine, pore size 0.9-1.4 μm , BET internal area, $a_m = 1.75 \text{ m}^2/\text{g}$, mass
70 density, $\rho_m = 1.02 \text{ g}/\text{cm}^3$, porosity of 0.31, and dimensions 20 × 10 × 2.7 mm), and the mem-
71 branes were Nafion® (Ion Power, N115, thickness ~ 127 microns).

72 Before assembly, the membranes were treated chemically to remove any impurities and
73 to activate them. The frit was glued into its acrylic frame using Devcon 2 Ton® Epoxy
74 from McMaster-Carr before assembly. The gasketing material was used to seal the device
75 and house the electrode channels (each about 0.8 mm thick during operation) that had been
76 cut into the gaskets. The electrode channels were open and pressurized during operation
77 using downstream pressure tubing in order to hold the membranes flat against the frit. The
78 splitter was also made from teflon gasketing material that was compressed against the end
79 of the frit using the outlet port plate. (See Supporting Information for more details).

80 During device operation, flow rates through three syringes were controlled by Harvard
81 Apparatus syringe pumps. The electrode channels were supplied with electrolyte solution of
82 the same concentration and composition as the feed to the frit. The flow rates were verified
83 manually at the device outlets before testing. Current was then applied using a Keithley
84 Instruments Model 2450 SourceMeter. During desalination tests, the system was allowed to
85 reach steady state before samples were taken from the outlets for impedance measurements
86 using a Gamry Reference 3000 Potentiostat to determine their conductivities.

87 Results and Discussion

88 The current-voltage relation of the device was obtained by slow linear sweeps in current,
89 since traditional voltage sweeps (cyclic voltammetry) were found to be less stable in the
90 over-limiting regime. Typical polarization curves at a sweep rate of 25 nA/s ($125 \mu\text{A}/\text{m}^2\text{s}$)
91 for 10 mM NaCl are shown in Fig. 2. Although the current is sustained by water splitting, the
92 onset voltage for significant current is consistently close to 1 V, which is less than the stan-
93 dard potential for water electrolysis (1.229 V), likely because of the very low partial pressures
94 of oxygen and hydrogen gas, since no gases are fed to the electrode channels. The role of sur-
95 face transport in over-limiting current is demonstrated by comparing polarization curves for a
96 bare, negatively charged frit (used for desalination experiments) and for a positively charged
97 frit, whose surfaces were treated by layer-by-layer deposition of a charged polymer,^{12,32} poly-
98 diallyldimethylammonium chloride (pDADMAC). Importantly, the two curves diverge above
99 the same current where strong desalination is observed in the bare frit (see below), which is
100 well above the theoretical diffusion-limited current ($j_{lim} = 2zeD_{eff}c_0/L \approx 0.18 \text{ A}/\text{m}^2$ for 10
101 mM NaCl, where $D_{eff} = 4.01 \times 10^{-10} \text{ m}^2/\text{s}$ and $L = 4.3 \text{ mm}$) due to convection.¹² After this
102 divergence, for the bare frit, there is a nearly linear increase in current, as predicted theo-
103 retically^{18,19} and observed experimentally in previous experiments on overlimiting current in
104 negatively charged porous media^{12,32} and microfluidic devices.²³ The overlimiting conduc-
105 tivity (slope * thickness of frit) of $9.79 \times 10^{-3} \text{ S}/\text{m}$ is quantitatively consistent with previous
106 experiments with the same porous silica glass frits and Nafion membranes that attributed
107 this phenomenon to electro-osmotic surface convection.¹² At still higher voltages, a nonlinear
108 increase in current is observed, likely associated with concentration polarization and water
109 splitting in the electrode channels, as in traditional electro dialysis systems.^{7,8,10} Bubble for-
110 mation is observed in the effluent of the electrode streams above $11 \text{ A}/\text{m}^2$ ($> 50j_{lim}$), and
111 the voltage becomes unstable and unreliable.

112 As noted above, it seems better to operate shock electrochemical systems by controlling
113 the current, rather than the voltage. Recent experiments¹² and simulations³⁵ have shown

114 that voltage sweeps tend to overshoot and oscillate around the limiting current plateau. In
115 steady state, galvanostatic operation also helps to ensure the formation of a stable deion-
116 ization shock by maintaining over-limiting current.^{15–17,19,20,23,33} In contrast, potentiostatic
117 operation leads to variable shock speeds in the over-limiting regime,^{19,36} and unstable cur-
118 rents at high voltage (from bubble formation and electro-osmotic instabilities in the electrode
119 channels) could disrupt the shock structure required for stable desalination.

120 Desalination tests were carried out over a wide range of currents, flow rates, and concen-
121 trations with several different electrolytes. Figure 3(a) shows a remarkable collapse of all the
122 data on a single dimensionless curve by plotting the percentage of salt removed $1 - \tilde{c}$, where
123 $\tilde{c} = c/c_0$ is the ratio of fresh to feed salt concentration, versus the applied current scaled to
124 the rate of positive charge advection into the device, $\tilde{I} = I/z_+c_+eQ$, where Q is the inlet
125 flow rate. The master curve is approximately exponential, $\tilde{c} \approx e^{-\gamma\tilde{I}}$, $\gamma = 0.619$, as shown by
126 the semilog plot in Fig 3(b). The data collapse with dimensionless current further supports
127 the benefits of galvanostatic operation noted above.

128 In dimensionless variables, the desalination performance of SED is thus mainly controlled
129 by the properties of the porous medium (macroscopic dimensions, surface charge, and mi-
130 crostructure) and does not explicitly depend on the ion type (subject to minor differences
131 in ionic mobilities), salt concentration, current, or flow rate. Simple models of SED pre-
132 dict similar data collapse and improved desalination by increasing the dimensionless ratio
133 $\tilde{\rho}_s = q_s a_m \rho_m / 2\varepsilon_p z_+ e c_0$ of the surface charge to the feed cation charge per macroscopic vol-
134 ume.¹⁹ This trend is consistent with the poor performance of a frit with larger pores and
135 $\sim 5x$ smaller $\tilde{\rho}_s$ (Fig. S6 of Supporting Information), although more experiments are needed
136 to test the predicted scaling.

137 Counter-intuitively, the placement of the splitter is not the sole factor that determines the
138 water recovery (defined as the ratio of flow of desalinated water to total flow of water into the
139 frit), in contrast to existing theoretical SED models.^{12,14,15,17–19} Since our splitter was placed
140 midway along the frit's downstream edge, the recovery was expected to be approximately

141 50%, with only small variations due to random pore structure or uncertainty in the splitter
142 placement (since it is made of non-rigid gasketing material). However, as the data in Fig.
143 3(c) shows, the water recovery actually increased from $\sim 45\%$ to 79% with increasing current
144 and decreasing flow rate. We hypothesize that this observation can be explained by an
145 increasingly significant contribution of electroosmotic pumping to the total flow (i.e., the
146 larger the ratio of electroosmotic flow to the applied flow is, the higher the water recovery
147 would be). Porous glass frits have been used as electroosmotic pumps,³⁷ and it is known that
148 water pumping is towards the cathode at near neutral pH. In our device with impermeable
149 membranes, however, existing models predict that any transverse electro-osmotic pumping
150 would be opposed by pressure-driven back flow, yielding no change in water recovery.

151 The resolution of this paradox may come from electro-osmotic surface convection behind
152 the shock,^{18,21} which could dissipate the pressure in vortices and redirect the net electro-
153 osmotic flow into the fresh outlet. Such convection has already been implicated as the
154 dominant mechanism for over-limiting conductance in our glass frits¹² and in microchannels
155 with similar vortex sizes.²³ Indeed, the classical model of electro-osmotic pumping without
156 any back pressure³⁸ predicts velocities at high currents comparable to the applied velocity.
157 To test this hypothesis, we replot the water recovery data in Fig. 3(d) versus the ratio of the
158 electro-osmotic pumping rate to the applied flow rate, $\hat{I} = Q_{EO}/Q$, where $Q_{EO} = \epsilon\zeta I/(\mu\sigma)$
159 is an estimate of the electro-osmotic flow, ϵ the permittivity, ζ the zeta potential (set to
160 100 mV^{37}), μ the viscosity, and σ the conductivity.³⁹ This leads to a reasonable collapse of
161 the data on a straight line, $R_w = \alpha + \beta\hat{I}$, $\alpha = 0.454$, $\beta = 0.092$, consistent with a simple
162 derivation in the supporting information. Statistical analysis gives a reduced χ^2 value of
163 1.22, indicating a good fit of the data. The small dimensionless slope (0.092) suggests that
164 electro-osmotic flow is significantly retarded by back pressure or surface charge regulation¹²
165 in the depleted region. Overall, the results strongly suggest that electro-osmotic pumping
166 leads to increased water recovery, which could be exploited in future system designs.

167 All desalination methods exhibit a tradeoff between salt removal, water recovery, and

168 energy efficiency. Plots of total (electrical and hydraulic pumping) energy consumption are
169 shown in Fig. S4 of the supporting information for all the conditions in Fig. 3. In this
170 first SED prototype, the impressive salt removal and water recovery come with significant
171 energy costs, in the range of 10^{-1} to 10^3 kWh/m³ (hydraulic pumping makes up about 0.5
172 kWh/m³). The energy efficiency (ratio of the thermodynamic limit set by osmotic pressure
173 to the actual total energy consumption) varied between 0.3% and 0.003%. The current
174 efficiency peaks at 57.7% at the onset of strong deionization and decays with increasing
175 current toward an apparent lower limit around 10% (see Fig. S5). This novel behavior,
176 compared to current efficiencies $> 80\%$ in electrodialysis,⁴⁰ suggests that surface current
177 in the deionized region has a larger transference number for protons than for cations. In
178 future SED systems, the energy efficiency may be dramatically improved by using a stack
179 of frits and membranes (to reduce the fraction of the total voltage lost to electrolysis) and
180 by varying properties of the porous medium, such as surface charge, proton affinity, matrix
181 microporosity (to enhance salt removal and shock formation) and anisotropic macroporosity
182 (to lower hydraulic resistance in the flow direction). As in traditional desalination plants,
183 the overall energy efficiency could also be improved by harvesting lost heat or re-using the
184 brine and electrode streams in related processes.

185 In summary, we have shown for the first time that a concentration shock can be stable
186 and propagate despite continuous cross-flow through a random porous medium. Thereby,
187 we have demonstrated that SED can be harnessed to continuously and scalably remove
188 99+% of ions from electrolyte solutions with water recovery up to 79%, biased by a new
189 mechanism of electro-osmotic pumping. Furthermore, this system can be scaled analogously
190 to electrodialysis systems by repeating our unit cell between two electrodes. In the future,
191 the design of the prototype will be altered to give more precise control of the water recovery
192 by electro-osmotic flow and to improve energy efficiency, guided by the scaling laws revealed
193 in Fig. 3. Furthermore, in the future, we will investigate the recycling of the anolyte and
194 catholyte solutions, keeping in mind that the pH of these solutions, when unbuffered, can

195 deviate significantly during operation. Possible applications include wastewater recycling,
196 ultrapure water production, treatment of produced water from hydraulic fracturing, and
197 water disinfection.¹³ The enormous gradients in salt concentration and electric field that
198 arise in SED also provide new opportunities for chemical or biological separations.

199 Acknowledgement

200 This research was supported by Weatherford International, the MIT Energy Initiative, the
201 SUTD-MIT Fellowship Program, and the USA-Israel Binational Science Foundation (grant
202 2010199).

203 Associated Content

204 The supporting material includes a derivation of the scaling of water recovery with electro-
205 osmotic pumping, more details on materials and methods, and further plots on desalination
206 performance, energy consumption, and current efficiency.

207 This material is available free of charge via the Internet at <http://pubs.acs.org/>

208 Notes

209 The authors declare no competing financial interest.

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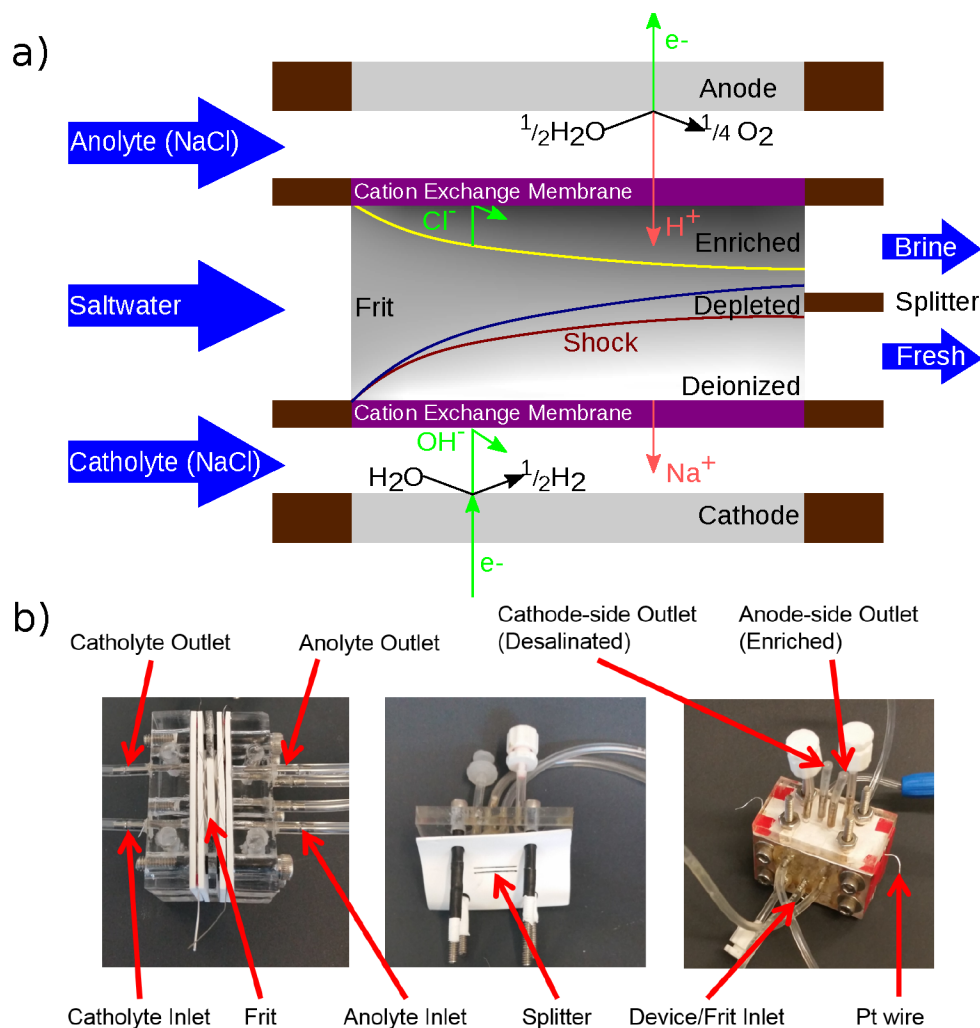


Figure 1: (a) Diagram of the SED prototype, showing its operating principles. In the presence of chloride ions, oxidation of chloride ions to chlorine gas may take place in addition to oxygen evolution at the anode. (b) Photographs of the SED prototype, both assembled (right) and unassembled (left and center) to show the frit placement and splitter design.

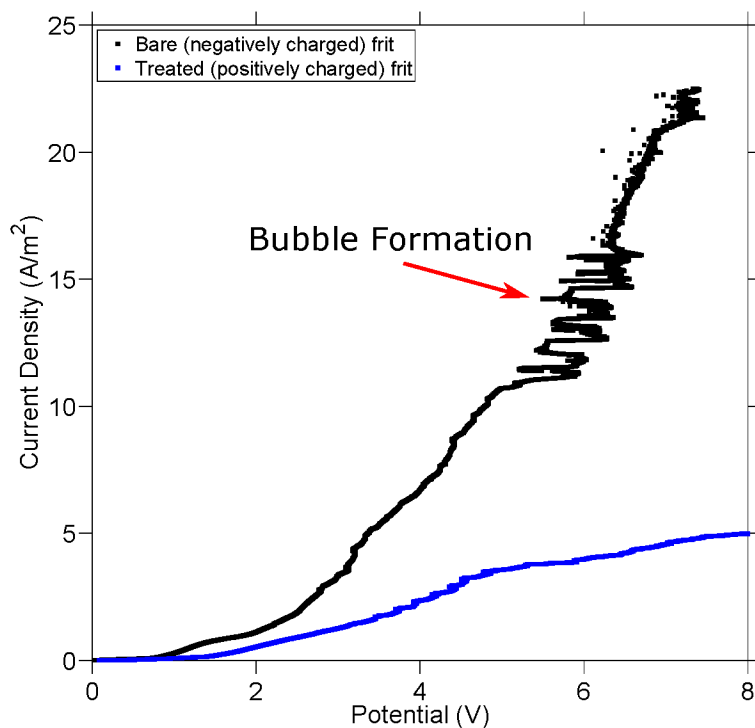


Figure 2: Typical current-voltage relations for the SED prototype measured at a sweep rate of 25 nA/s ($125 \mu\text{A}/\text{m}^2\text{s}$) for 10 mM NaCl ($j_{lim} \approx 0.18 \text{ A}/\text{m}^2$ without convection) with the bare, negatively charged silica glass frit used for desalination (below), as well as a frit whose surfaces are coated with positively charged polymers to show the role of negative surface charging in sustaining over-limiting current.^{12,32} The over-limiting conductance¹⁸ is consistent with electro-osmotic surface convection,¹² and unstable voltage associated with water splitting is observed above $11 \text{ A}/\text{m}^2$ ($> 50j_{lim}$).

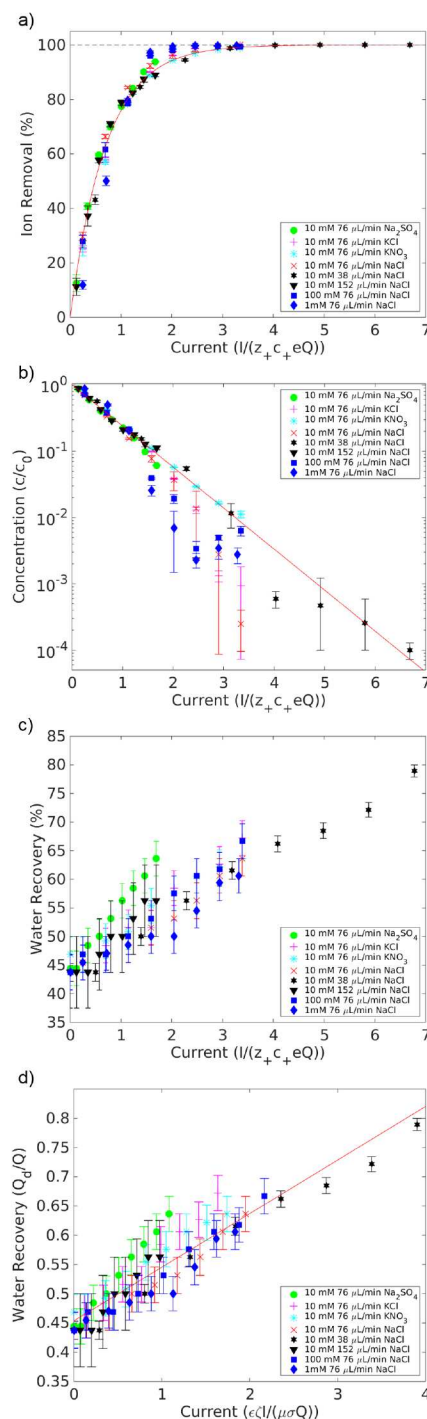


Figure 3: Desalination performance of the SED prototype. (a) Data collapse of the percentage of ion removal, $1 - \tilde{c}$, where $\tilde{c} = c/c_0$ is the ratio of fresh to feed salt concentrations, versus dimensionless current, $\tilde{I} = I/(z_+c_+eQ)$, scaled to the rate of positive charge advection into the device. (b) This semilog plot of \tilde{c} versus \tilde{I} shows excellent data collapse around the red line, $\log \tilde{c} = -0.619 * \tilde{I}$, especially below 95% salt removal. (c) Data for water recovery R_w versus dimensionless current \tilde{I} . As shown in (d), data collapse is quite good by scaling the applied flow rate to an estimate of the transverse electroosmotic flow, $\hat{I} = Q_{EO}/Q = \epsilon\zeta I/(\mu\sigma Q)$. The red line, $R_w = 0.454 + 0.092\hat{I}$, gives a reasonable fit of all the data ($\chi^2_{reduced} = 1.22$).

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