# Quantum spin liquids and continuous metal-insulator transitions

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TS, PR B 08.

- D. Podolsky, A. Paramekanti, Y.B. Kim, TS, PRL 09.
- D. Mross and TS, PRB 11.
- A. Potter, M. Barkeshli, J. McGreevy, TS, PRL 2012.
- W. Witczak-Krempa, P. Ghaemi, Y.B. Kim, TS, arxiv, 1206.3309

### Plan of talk

Part I. Theory of a continuous Mott metal-insulator transition in d = 2.

Evolution from Fermi liquid to quantum spin liquid insulator: Predictions for transport experiments

Part 2. Metal-insulator transitions in doped semiconductors (Si:P, Si:B).

New questions/insights inspired by quantum spin liquid theory/experiments.

### The electronic Mott transition

Difficult old problem in quantum many body physics

How does a metal evolve into a Mott insulator?

Prototype: One band Hubbard model at half-filling on non-bipartite lattice

AF insulator; Fermi liquid;
No Fermi surface Full fermi surface

### Why hard?

- 1. No order parameter for the metal-insulator transition
- 2. Need to deal with gapless Fermi surface on metallic side
- 3. Complicated interplay between metal-insulator transition and magnetic phase transition

Typically in most materials the Mott transition is first order.

But (at least on frustrated lattices) transition is sometimes only weakly first order - fluctuation effects visible in approach to Mott insulator from metal.

Quantum spin liquid Mott insulators:

Opportunity for progress on the Mott transition - study metal-insulator transition without complications of magnetism.

### Some candidate spin liquid materials

$$\kappa - (ET)_2 Cu_2(CN)_3$$

 $EtMe_3Sb[Pd(dmit)_2]_2$ 

Quasi-2d, approximately isotropic triangular lattice; best studied candidate spin liquids

 $Na_4Ir_3O_8$ 

Three dimensional 'hyperkagome' lattice

 $ZnCu_3(OH)_6Cl_2$ 

Volborthtite, .....

2d Kagome lattice ('strong' Mott insulator)

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Close to pressure driven Mott transition: `weak' Mott insulators

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Volborthtite, .....

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## Some phenomena in experiments

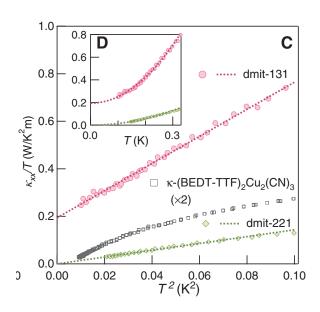
**ALL** candidate quantum spin liquid materials:

**Gapless** excitations down to T << J.

Most extensively studied in organic spin liquids with J ≈ 250 K.

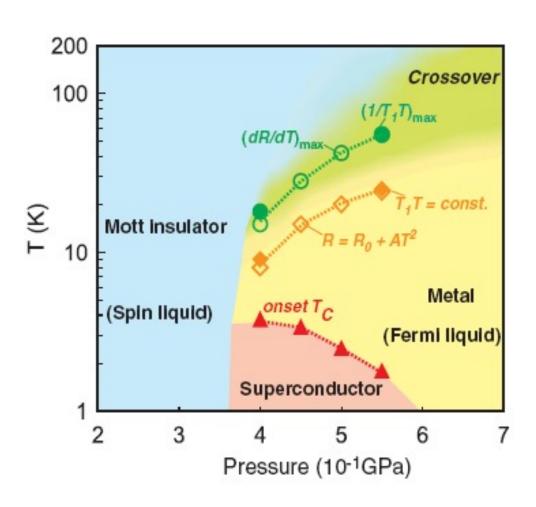
Example: Thermal transport in dmit SL.

Electrical Mott insulator but thermal metal!



M. Yamashita et al, Science 2010.

# Possible experimental realization of a second order(?) Mott transition



### Quantum spin liquids and the Mott transition

#### Some questions:

- 1. Can the Mott transition be continuous?
- 2. Fate of the electronic Fermi surface?

### Slave particle framework

Split electron operator

$$c_{r\sigma}^{\dagger} = b_r^{\dagger} f_{r\alpha}$$

Fermi liquid:  $\langle b \rangle \neq 0$ 

Mott insulator:  $b_r$  gapped

Mott transition:  $b_r$  critical

In all three cases  $f_{r\alpha}$  form a Fermi surface.

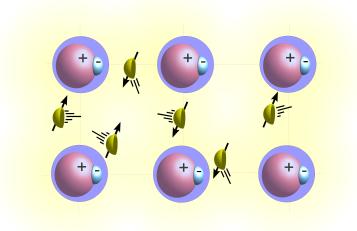
Low energy effective theory: Couple b, f to fluctuating U(1) gauge field.

### Picture of Mott transition

Metal

Electrons swimming in sea of +vely charged ions

Mott spin liquid near metal



Electron charge gets pinned to ionic lattice while spins continue to swim freely.

### Quantum spin liquids and the Mott transition

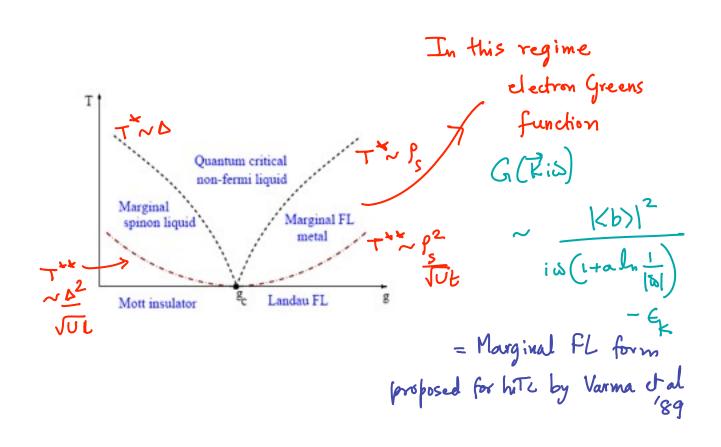
- 1. Can the Mott transition be continuous?
- 2. Fate of the electronic Fermi surface?



Concrete tractable theory of a continuous Mott transition; demonstrate critical Fermi surface at Mott transition; definite predictions for many quantities (TS, 2008).

Example: universal jump of residual resistivity on approaching from metal, weak divergent effective mass,.....

# Finite-T crossovers: emergence of a Marginal Fermi Liquid TS, 2008



### Structure of critical theory

Field theory for critical point

$$S = S[b, a] + S[f_{\alpha}, a]$$

Gauge fluctuations are Landau damped by spinon Fermi surface:

$$S_{eff}[a] = \int_{\mathbf{q},\omega} \left( K_F \frac{|\omega|}{|\mathbf{q}|} + .. \right) |\mathbf{a}(\mathbf{q},\omega)|^2$$

=> at low energies gauge field decouples from critical b fluctuations. Effective critical action

$$S_{eff} = S[b] + S[f, a]$$

S[b]: critical D = 2+1 XY model

S[f]: spinon Fermi surface + Landau damped gauge field with  $z_b = 2$  Both individually understood.

### Non-zero temperature transport/dynamics

$$S_{eff}[a] = \int_{\mathbf{q}} \frac{1}{\beta} \sum_{\omega_n} \left( K_F \frac{|\omega_n|}{|\mathbf{q}|} + .. \right) |\mathbf{a}(\mathbf{q}, \omega_n)|^2$$

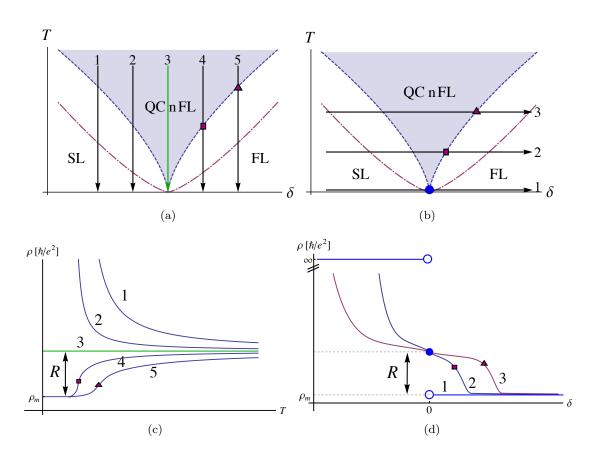
Static gauge fluctuations ( $\omega_n = 0$ ) escape Landau damping, and do not decouple from critical bosons.

Universal transport in a large-N approximation (Witzcak-Krempa, Ghaemi, TS, Y.B. Kim, 2012):

Gauge scattering reduces universal conductivity by factor of  $\approx 8$  from 3D XY result (Damle, Sachdev '97).

Electronic Mott transition: Net resistivity  $\rho = \rho_b + \rho_f$ Universal resistivity jump =  $\rho_b$  enhanced by factor of  $\approx 8$ .

### Non-zero temperature transport



$$\rho - \rho_m = \frac{\hbar}{e^2} G\left(\frac{\delta^{z\nu}}{T}\right) \qquad z = 1, \ \nu \approx 0.672$$

# Part 2: Metal-insulator transitions in doped semiconductors

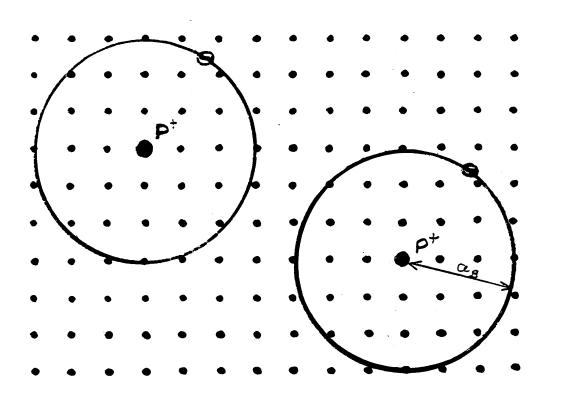
Eg: Si:P, Si:B

Subject of many studies over last 3 decades.

"Anderson-Mott" transition

Is there a quantum spin liquid? (Potter, Barkeshli, McGreevy, TS, PRL 2012)

## Basic picture for Si:P (Si:B, etc)

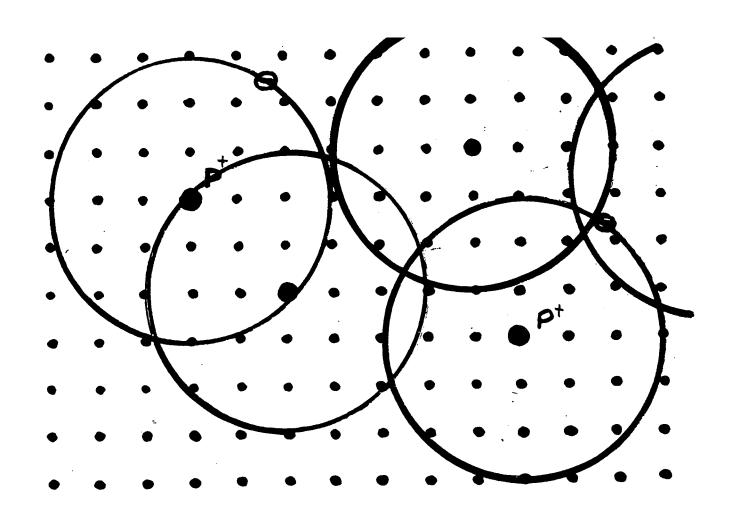


Extra electron of P forms a hydrogen-like state.

 $a_B \approx 20A$ 

Simple model: Randomly placed ``Hydrogen atoms''. Half-filled Hubbard model on random lattice.

# Increase P concentration to get metal



# Local moments in insulator: Random singlets

Bhatt, Lee, 1982

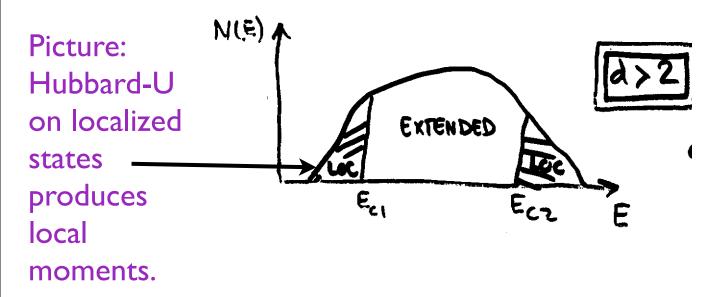
Each local moment forms a singlet bond with a fixed partner.

Broad distribution of singlet bond energies.

Anomalous low-T thermodynamics: diverging spin susceptibility, C/T (dominated by rare weakly coupled spin pairs).

# Metallic phase: persistence of some local moments

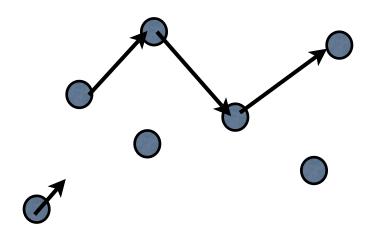
Near transition, some rare fraction of sites retain local moments which then dominate low-T thermodynamics.



# Two fluid phenomenology of metal



Paalanen et al, 1988; Gan, Lee, 86; Milovanovic, Sachdev, Bhatt, 1989; Bhatt, Fisher, 1992



Itinerant electron fluid coexisting with small fraction of local moments.

Near transition, fraction of local moment sites about 15%.

Thermodynamics: independent contribution from both fluids.

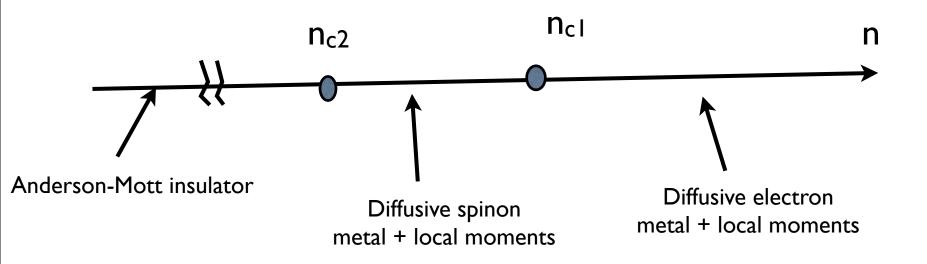
### Evolution across Metal-Insulator Transition (MIT)

What is fate of conducting fluid?

### Three possibilities:

- I. `Conventional wisdom' Itinerant electrons ----> Anderson insulator
- 2. Fraction of sites with local moments increases to approach I at MIT (generically unlikely).
- 3. New possibility: Conducting fluid Mott localizes into a quantum spin liquid with diffusive spinons (Potter, Barkeshli, McGreevy, TS, 2012)

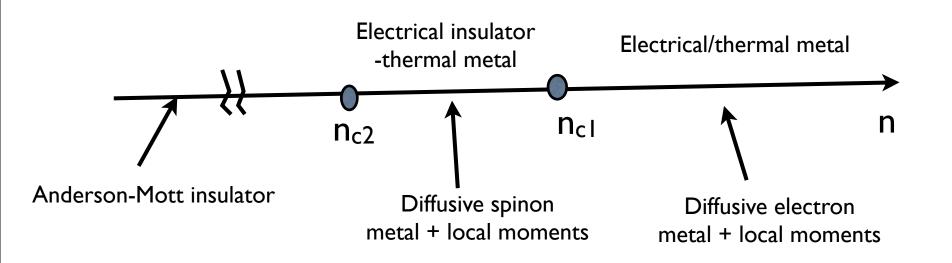
# Possible route to metal-insulator transition



## Some consequences-I

Diffusive spinon metal is electrical insulator but a thermal conductor.

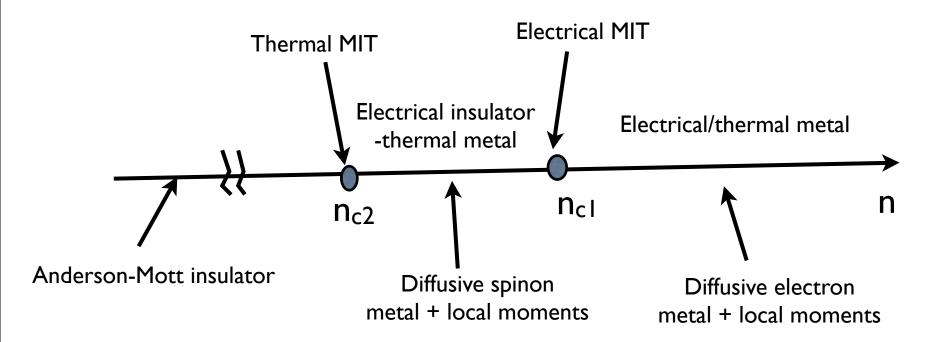
=> electrical metal-insulator transition separated from thermal metal-insulator transition.



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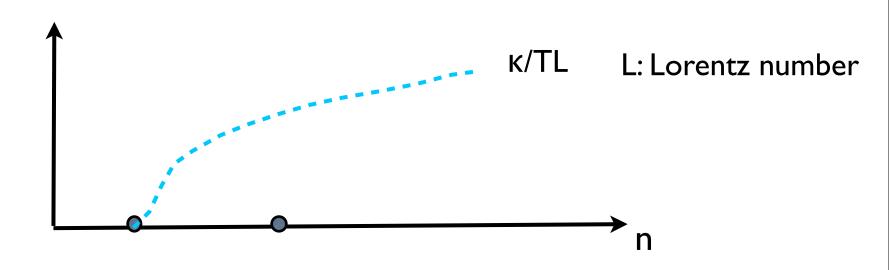


Crucial test: measure thermal transport in insulator.

# Some consequences -II

Existence of diffusive spinon metal will impact critical behavior of electrical conductivity near electrical MIT.

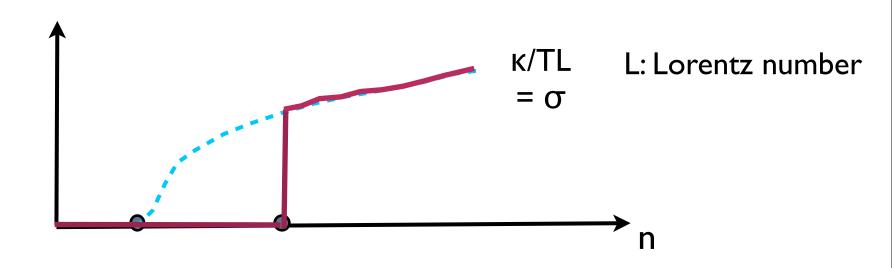
Jump of residual conductivity at MIT; non-monotonic T-dependence.



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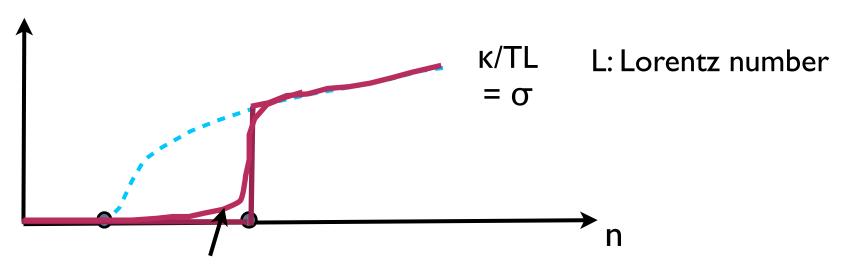
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Thermal rounding of jump; complicates extrapolation to T=0 conductivity.

## Summary

Quantum spin liquids provide an opportunity for progress on classic old problems: Mott and Anderson-Mott metal-insulator transitions.

Clean limit (organics, hyperkagome iridate): Continuous Mott transition possible; several predictions for experiment (eg: universal resistivity jump in d = 2, resistivity peak in d = 3)

Disordered limit (doped semiconductors Si:P, Si:B):

Do electrical and thermal metal-insulator transitions occur simultaneously?