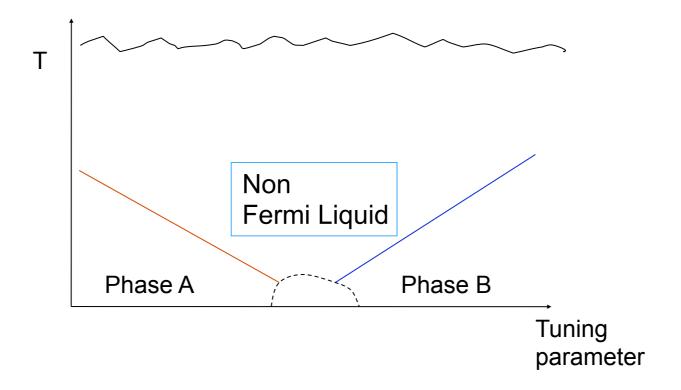


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Some collaborators: Chong Wang, R. Nandkishore, M. Metlitski, A. Vishwanath

## A common phase diagram



NFL due to quantum criticality?

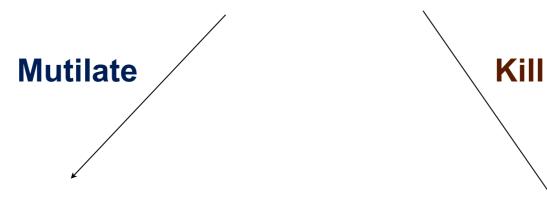
## Approach from Fermi liquid

Crucial: Fate of Fermi surface as a Fermi liquid metal undergoes a quantum phase transition?

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Fermi surface evolves continuously but is distorted in some way. (Ferromagnet, nematic, SDW, CDW,.....)

Original Fermi surface completely disappears. (Mott transition, Kondo breakdown,....)

### Approach from Fermi liquid

Crucial: Fate of Fermi surface as a Fermi liquid metal undergoes a quantum phase transition?



**Mutilate** 

Fermi surface evolves continuously but is distorted in some way. (Ferromagnet, nematic, SDW, CDW,.....)

Many talks at this meeting!

Kill

Original Fermi surface completely disappears. (Mott transition, Kondo breakdown,....)

#### Mutilate the Fermi surface

Model: Electron Fermi surface + X

X = critical order parameter fluctuations

Criticality and non-Fermi liquids at ``low" energy scales?

## Killing a Fermi surface

#### **Expect**

- strong NFL physics near transition
- change of electronic structure happens on electronic energy scales
- theory not described as electron Fermi surface + X

#### How does a Fermi surface die?

Fascinating question fundamental to many issues

1. Quasiparticle residue Z vanishes continuously on approaching transition. (Brinkman, Rice 1970).

Concrete examples: (d = 2,3) TS, Vojta, Sachdev 2004, TS, 2008, Nandkishore, Metlitski, TS (2012); DMFT in d = ∞

2. Fermi surface stays sharp at QCP even though there is no Landau quasiparticle (TS, 2008).

"Critical Fermi Surface"

Concrete examples: (d = 2) TS, 2008, Nandkishore, Metlitski, TS (2012).

Other related problems: Non-fermi liquid *phases*, some (gapless) quantum spin liquids..... Goal: Construct tractable, emergable effective field theories of these phenomena (Emergable: capable of emerging from microscopic lattice models in the `right' physical Hilbert space with the right symmetries.)

## A powerful approach

Slave particles (partons): Fractionalize spin/electron into partons which are then gapless.

Effective theory: Gapless partons + dynamical gauge fields.

# An example Orthogonal Metals: The simplest non-fermi liquids

Nandkishore, Metlitski, TS, 12

'Slave spin' representation:

$$c_{i\alpha} = \tau_i^x f_{i\alpha}$$

 $\tau_i^x$ : Ising spin;  $f_{i\alpha}$ : charge-1, spin-1/2 fermion.

 $Z_2$  gauge redundancy

$$\tau_i^x \to -\tau_i^x$$

$$f_{i\alpha} \to -f_{i\alpha}$$

Phase where  $\langle \tau_i^x \rangle = 0$ , the  $f_{i\alpha}$  form a Fermi surface, and  $Z_2$  gauge field is deconfined:

Electron spectral function gapped (Z = 0) but still a metal: ``Orthogonal metal'' Charge carriers have no overlap with electrons.

(Prior literature: Mis-interpreted as Mott insulator)

#### Comments

- 1. Relation to usual slave bosons/rotors:  $c_{i\alpha} = b_i f_{i\alpha}$ Condense  $\langle b_i^2 \rangle \neq 0$  with  $\langle b_i \rangle = 0$ .
- 2. OM Wavefunction

$$\psi(\vec{r}_i\alpha_i) = \psi_{PSF}(\vec{r}_i)\psi_{Slater}(\vec{r}_i\alpha_i)$$

 $\psi_{PSF}$ : wave function of paired boson superfluid  $(\propto \mathcal{S}[g(\vec{r}_1 - \vec{r}_2)g(\vec{r}_3 - \vec{r}_4)....g(\vec{r}_{N-1} - \vec{r}_N)]).$ 

- 3. OM phase can exist at any density or even in continuum.
- 4. Exactly soluble lattice models easy to write down.

# `Death' of a Fermi surface: FL- Orthogonal Metal transition

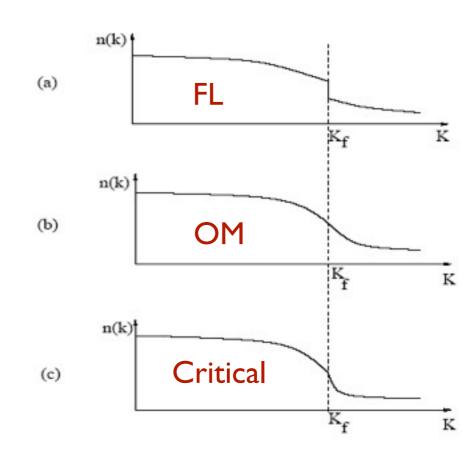
Approach from FL: electron Fermi surface dies.

Slave spin disordering transition

Critical theory: Critical slave spin + f-Fermi surface with energy-energy coupling. (Second order in some situations).

Critical electron Fermi surface: exponents exactly determined by Ising exponents.

$$A(\vec{q},\omega) \sim \omega^{d-2+\eta} g\left(\frac{\omega}{v_f q_\perp}\right)$$

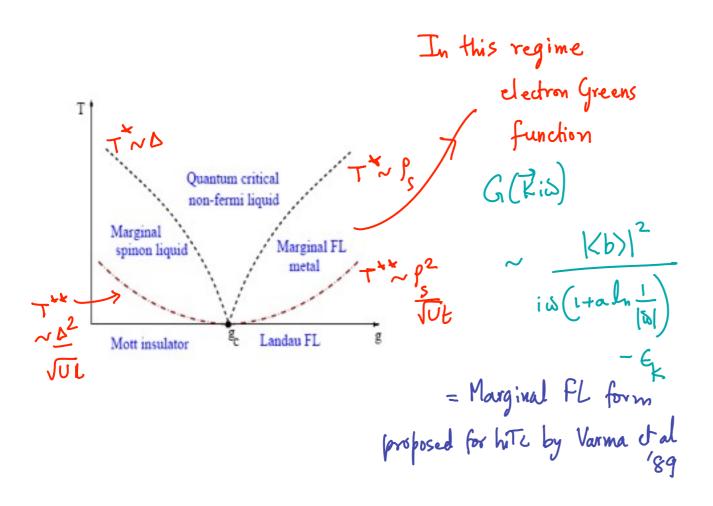


Critical theory  $\neq$  electron Fermi surface + X

#### Similar story: Theory of continuous Mott transitions

Mott metal-insulator transition can be continuous in d = 2 or d = 3 if insulator is spin liquid.

Slave mean field: Florens, Georges 04; Fluctuations, field theory: TS 08; Numerics: Motrunich, Fisher, et al 14.



Many interesting properties:

Eg:Two crossovers

Emergence from criticality # Fermi liquid coherence

#### Comments on slave particle framework

- 1. Conceptually important construction of effective field theories of a class of quantum phases/phase transitions that are beyond standard quasiparticles
- 2. Slave particle effective field theories are emergable, and often tractable.

Theoretical demonstration of many unusual phenomena; some successful contact with experiments (eg, FQHE, 1/2-filled Landau level, some quantum spin liquids)

Essentially only available framework for field theory 'beyond quasiparticles'.

### Beyond (standard) slave particles?

Are there gapless quantum liquids that cannot be understood (easily) within the slave particle framework?

Yes!

Two examples in quantum spin/bose liquids

(i) Quantum vortex liquid phases of quantum XY magnets/frustrated boson systems

Chong Wang, TS to appear.

- (ii) Quantum liquid of fluctuating spiral magnetic orders (inspired by MnSi).
- A. Vishwanath, TS, to appear.

#### Example-I: Gapless quantum vortex liquids

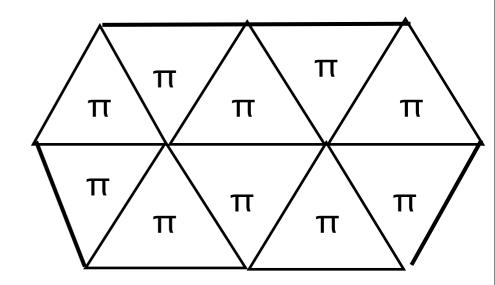
d = 2 XY spin-S triangular lattice quantum antiferromagnet = bosons with frustrated hopping.

Duality: Vortices at 1/2 filling on honeycomb lattice.

Interesting early attempt: (Fisher et al)

Fermionize vortices through flux attachment and let then form a gapless fluid (``Algebraic Vortex Liquid'')

Flaw (Wang, TS 13): With time reversal, this state cannot exist in strict d = 2 (but can exist at boundary of a 3d ``SPT" state)



#### Fractionalizing vortices

Chong Wang, TS 14; also Hermele 10

Dual theory:

$$\mathcal{L} = \mathcal{L}[\Phi_v, a_\mu]$$

 $\Phi_v$ : vortex field on honeycomb lattice

 $a_{\mu}$ : non-compact U(1) gauge field.

Dual partons:

$$\Phi_v = d_1 d_2$$

 $d_{1,2}$ : Fermionic 1/2-vortices.

Dual parton field theory:

Put d in various 'mean field' states; analyze fluctuations.

#### Fractionalized vortex liquids

Dual parton mean field:

Can arrange to break all gauge structure associated with fractionalizing  $\Phi_v$ .

d-fermions gapped: Paired superfluid (spin nematic) of original bosons (spins).

Dirac nodes for d-fermions => new gapless spin liquid phase. Effective theory:

$$\mathcal{L} = \bar{d} \left( \gamma_{\mu} (i\partial_{\mu} - a_{\mu}) d + \frac{1}{g} f_{\mu\nu}^{2} \right)$$

Dual massless non-compact  $QED_3$ .

Power law correlations for many quantities.

#### Physical picture

#### Excitations of paired boson superfluid:

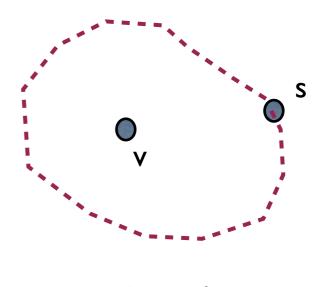
- I. Gapless sound waves
- 2. Gapped neutralized unpaired boson s
- 3. Gapped 1/2-vortices v

s and v are mutual semions.

s-v bound state is a fermionic vortex.

Proliferate to get a gapless fractionalized vortex liquid.

Beyond (standard) slave particles: in terms of original bosons, short distance physics is that of pairing, not fractionalization.



Phase of  $\pi$ 

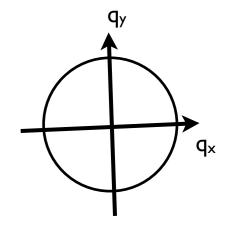
## Example-II: Bose Luttinger Liquid in d > I

Some situations in which bosonic degree of freedom has dispersion minima along surfaces in momentum space:

Eg: Incipient spiral magnetic order (eg MnSi), bosons with Rashba coupling,......

Model action

$$S = \int_{\vec{q},\omega} (\omega^2 + (q^2 - q_0^2)^2 + r) |b(\vec{q},\omega)|^2 + S_{int}$$



Usually interactions will spontaneously pick out some q-vectors for ordering.

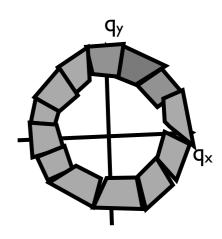
This talk: a fluctuating quantum liquid state without any ordering by 'bosonizing' the bosons.

#### Patch construction

Divide vicinity of Bose surface into discrete patches

Bose field

$$b(\vec{r},\tau) = \frac{1}{\sqrt{N}} \sum_{S} b_{S}(\vec{r},\tau) e^{i\vec{q}_{S}.\vec{r}}$$



Amplitude-phase representation

$$b_S = b_0 e^{i\phi_S}$$

Strategy:

- 1. Assume initially  $b_0$  is uniform on Bose surface but  $\phi_S$  may fluctuate.
- 2. Study "phase-only" theory of  $\phi_S$
- 3. Worry about vortices
- 4. Take patch continuum limit, thermodynamic limit.

#### Phase-only theory on Bose surface

First also ignore inter-patch coupling.

Quadratic action:

$$\mathbf{L}_{\phi}^{(0)} = \frac{1}{N} \sum_{S=1}^{N} \rho_S \left\{ (\partial_{\tau} \phi_S)^2 + (\vec{v}_S \cdot \vec{\nabla} \phi_S)^2 \right\}$$

I+I-d XY field for each patch:

Boson has power law correlations centered at Bose surface with continuously varying exponent.

$$\langle \bar{b}_S(\vec{q},i\omega)b_S(\vec{q},i\omega)\rangle \sim \frac{1}{\left(v^2q_\perp^2+\omega^2\right)^{1-\frac{\eta}{2}}} \qquad \text{``Bose Luttinger Liquid''}$$

#### Stability

Interpatch interactions: As in Fermi Liquid Theory, only `forward' scattering and `BCS' interactions survive.

"BCS": Irrelevant by power counting if  $\eta > 1$ .

"Forward" scattering: Couple together density and currents of XY fields at different patches.

Landau interactions with same qualitative effects as in Fermi Liquid

(similar renormalizations, collective zero sound modes, etc).

Vortices: Only legitimate vortex is in I = 0 phase mode.

This can be suppressed => Bose Luttinger Liquid exists as stable phase of matter.

#### Summary

I. Slave particle framework: effective field theories which concretely demonstrate many interesting phases/phase transitions beyond quasiparticles.

2. There are interesting states beyond slave particles.