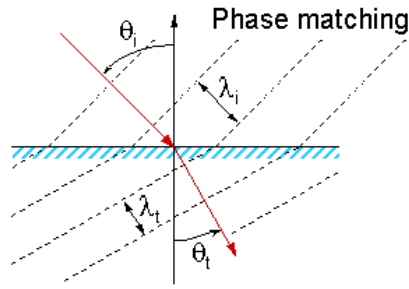
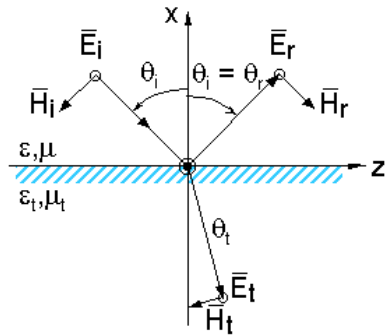


PLANE WAVES AT PLANAR BOUNDARIES

Amplitudes and phases of reflection and transmission:

First consider TE waves



$$\bar{E}_i = \hat{y} E_0 e^{+jk_x x - jk_z z}$$

$$\bar{E}_r = \hat{y} \Gamma E_0 e^{-jk_x x - jk_z z}$$

$$\bar{E}_t = \hat{y} \Gamma E_0 e^{+jk_x x - jk_z z}$$

where

$$k_{x_i} = k_i \cos \theta_i$$

$$k_i = \omega \sqrt{\mu \epsilon}, \quad k_t = \omega \sqrt{\mu_t \epsilon_t}$$

Apply boundary conditions at $x \triangleq 0$

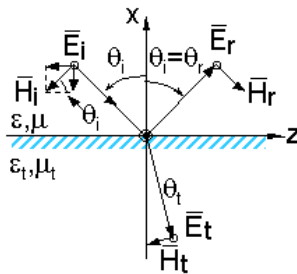
$\bar{E}_{//}$ continuous: $E_0 + \Gamma E_0 = \Gamma E_0$

$$1 + \Gamma = \Gamma$$

LS-1

PLANE WAVES AT PLANAR BOUNDARIES (2)

(1) $\bar{E}_{//}$ continuous: $1 + \Gamma = \Gamma$



(2) $\bar{H}_{//}$ continuous: $-\cos \theta_i (E_0 / \eta_0) + \cos \theta_i \Gamma (E_0 / \eta_0) = -\cos \theta_t \Gamma (E_0 / \eta_t)$

$$1 - \Gamma = \frac{\cos \theta_t \eta_0}{\cos \theta_i \eta_t} \Gamma \quad (\text{Note: } \eta \neq n)$$

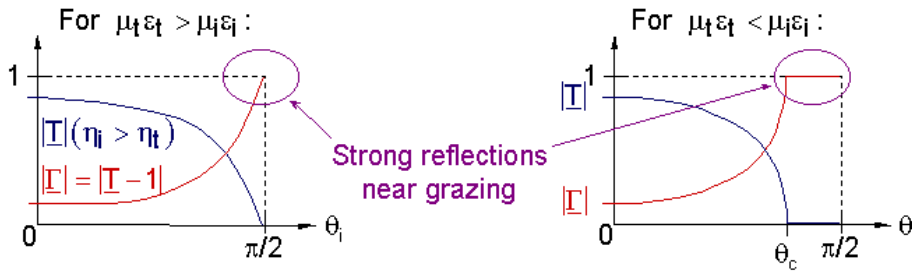
(1) & (2) $\Rightarrow \Gamma = 2 / [1 + (\eta_0 \cos \theta_t / \eta_t \cos \theta_i)] \rightarrow 0$ as $\theta_i \rightarrow \frac{\pi}{2}$

(1) $\Rightarrow \Gamma = \Gamma - 1 \rightarrow -1$ as $\theta_i \rightarrow \frac{\pi}{2}$, and $\rightarrow \begin{cases} 1 & \text{for } \eta_0 \rightarrow \ll \eta_t \\ -1 & \text{for } \eta_0 \gg \eta_t \end{cases}$ for $\theta_i = 0$

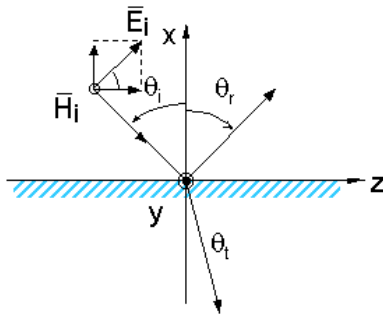
where $\cos \theta_t = \sqrt{1 - \sin^2 \theta_t}$; $\sin \theta_t = (c_t / c_i) \sin \theta_i$ (Snell's Law)

LS-2

PLANE WAVES AT PLANAR BOUNDARIES (3)



Consider TM Waves



$$\bar{H}_i = \hat{y} H_0 e^{+jk_{x1}x - jk_z z}$$

$$\bar{E}_i = (\hat{x} \sin \theta_i + \hat{z} \cos \theta_i) H_0 \eta_i e^{+jk_{x1}x - jk_z z}$$

Can use same method as for TE waves (see above) or "Duality"

L8-3

METHOD OF DUALITY

Note: for $\rho = \bar{J} \equiv 0$ we have:

$\nabla \times \bar{E} = -j\omega\mu\bar{H}$	$\nabla \times \bar{H} = j\omega\epsilon\bar{E}$
$\nabla \times \bar{H} = j\omega\epsilon\bar{E}$	$\nabla \times \bar{E} = -j\omega\mu\bar{H}$
$\nabla \cdot \mu\bar{H} = 0$	$\nabla \cdot \epsilon\bar{E} = 0$
$\nabla \cdot \epsilon\bar{E} = 0$	$\nabla \cdot \mu\bar{H} = 0$

Claim: if given valid solution in zone where $\rho = \bar{J} \equiv 0$

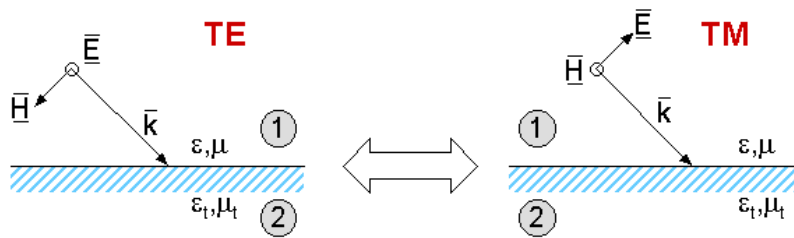
and we write $\bar{E} \rightarrow \pm\bar{H}, \bar{H} \rightarrow \mp\bar{E}$

$$\mu \rightarrow \epsilon, \epsilon \rightarrow \mu$$

This "dual solution" is valid (if the new fields satisfy the new boundary conditions)

L8-4

EXAMPLE OF DUALITY



- A) If TE is solution in ①, then TM is too if $\epsilon \rightarrow \mu, \vec{E} \rightarrow \vec{H}$
 $\mu \rightarrow \epsilon, \vec{H} \rightarrow -\vec{E}$
- B) In region ②, use same transformation to find TM solution, always

C) TM solutions in ①, ② match boundary conditions (B.C.) to **only if B.C. are also dual.**

Here: B.C. are $E_{||}, H_{||}$ continuous; they satisfy duality ($E_{||} \leftrightarrow \pm H_{||}$)
 Therefore duality yields TM solution immediately when B.C. satisfy duality

L8-5

EXAMPLE OF DUALITY (2)

TE Solution:

$$\Gamma_{TE} = 2 \left/ \left(1 + \frac{\cos \theta_t \sqrt{\mu_o / \epsilon_o}}{\cos \theta_i \sqrt{\mu_t / \epsilon_t}} \right) \right.$$

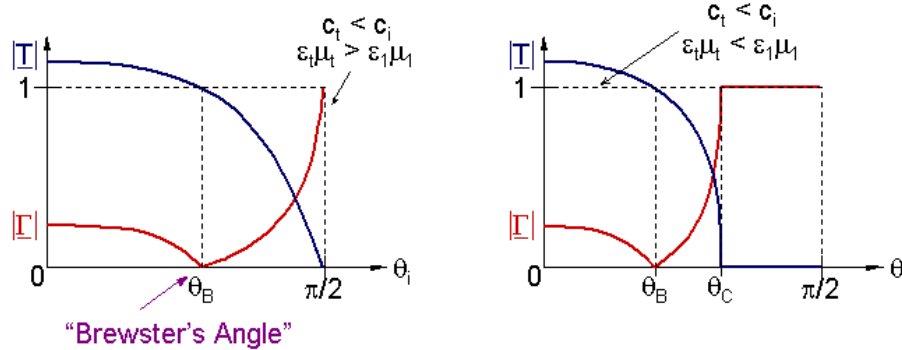
$$\Gamma_{TE} = \Gamma_{TE} - 1$$

Dual TM Solution:

$$\Gamma_{TM} = 2 \left/ \left(1 + \frac{\cos \theta_t \sqrt{\epsilon_o / \mu_o}}{\cos \theta_i \sqrt{\epsilon_t / \mu_t}} \right) \right.$$

$$\Gamma_{TM} = \Gamma_{TM} - 1$$

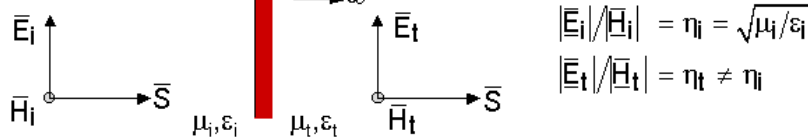
Example: TM reflection and transmission magnitudes



L8-6

PHYSICAL INTERPRETATIONS OF BREWSTER'S ANGLE

Normal Incidence:



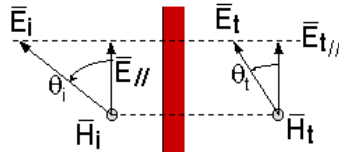
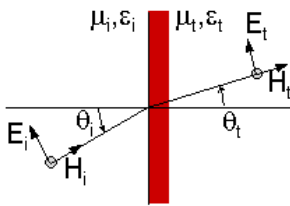
$$|\vec{E}_i|/|\vec{H}_i| = \eta_i = \sqrt{\mu_i/\epsilon_i}$$

$$|\vec{E}_t|/|\vec{H}_t| = \eta_t \neq \eta_i$$

For NO reflection $\vec{E}_{//}, \vec{H}_{//}$ must be continuous.

For normal incidence this is possible only IF $\eta_i = \eta_t$

Angular Incidence:



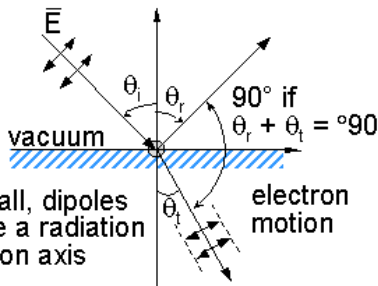
If $|\vec{H}_t| = |\vec{H}_i|$, and $|\vec{E}_t| < |\vec{E}_i|$, then there is an incidence angle where impedances match.

e.g. If: $\mu_i = \mu_t$ and $\epsilon_t > \epsilon_i$

Then: $\eta_t < \eta_i = \sqrt{\mu_i/\epsilon_i}$, $|\vec{E}_t| < |\vec{E}_i|$

L8-7

ALTERNATIVE INTERPRETATION OF θ_B



If $\mu_t = \mu_i$; $\epsilon_i \neq \epsilon_t \Rightarrow$ null for TM wave when $\theta_r + \theta_t = 90^\circ$ (then $\theta_i = \theta_B$).

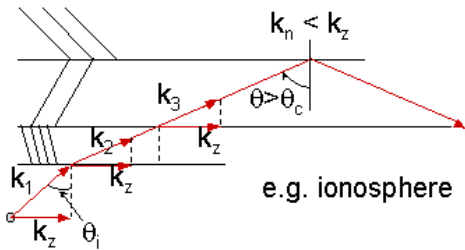
What happens when $\epsilon_i = \epsilon_t$, $\mu_t \neq \mu_i$?

Magnetic dipoles are included, and null for TE wave when $\theta_i + \theta_t = 90^\circ$ (then $\theta_i = \theta_B$)

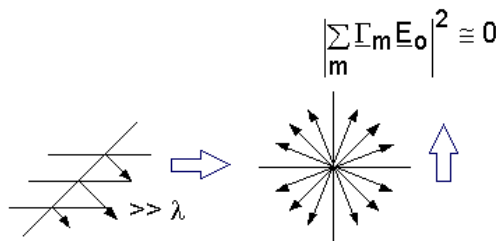
Recall, dipoles have a radiation null on axis

electron motion

Multilayer refraction:



e.g. ionosphere



reflection phases cancel \Rightarrow no losses, 100% transmission

$$\left| \sum_m \Gamma_m E_o \right|^2 \cong 0$$

L8-8