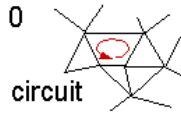


## KIRCHOFF'S VOLTAGE LAW

**Kirchoff's Voltage Law: (KVL)**

Around any loop:  $\sum_i V_i = 0$

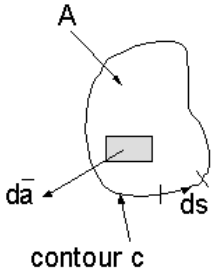


**Faraday's Law:**

$\nabla \times \vec{E} = -\partial \vec{B} / \partial t$  [differential form]

Integral form is:

$$\int_c \vec{E} \cdot d\vec{s} = -\int_A (\partial \vec{B} / \partial t) \cdot d\vec{a}$$



[Recall Stoke's Theorem:  $\int_A (\nabla \times \vec{G}) \cdot \hat{n} da = \int_c \vec{G} \cdot d\vec{s}$ ]

If  $\partial \vec{B} / \partial t = 0$  in A, then Kirchoff's voltage law must be satisfied

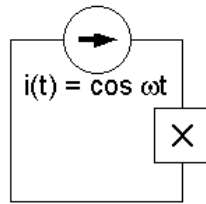
More generally:

KVL assumes all magnetic energy is stored inside circuit elements

L10-1

## KIRCHOFF'S VOLTAGE LAW (2)

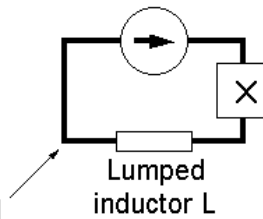
**Undefined Circuit:**



Current in loop varies, so  $(\partial \vec{B} / \partial t) \neq 0 \Rightarrow$  voltage!

**Defined Circuit:**

We assume voltage drops occur only across elements, and ignore H fields generated by currents through them.



We use small loops, thick wires, and high-impedance lumped elements

L10-2

## KIRCHOFF'S CURRENT LAW (KCL)

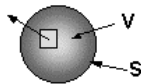
**Kirchoff's Current Law: (KCL)**  $\sum_i i_i = 0 =$  Total current into any node

**Ampere's Law:**  $\nabla \times \vec{H} = \vec{J} + \partial \vec{D} / \partial t$  [differential form]  
 $\nabla \cdot (\nabla \times \vec{H}) = \nabla \cdot (\vec{J} + \partial \vec{D} / \partial t) = 0$   
 $= \nabla \cdot \vec{J} + \partial (\nabla \cdot \vec{D}) / \partial t$   
 $= \nabla \cdot \vec{J} + \partial \rho / \partial t$

**Conservation of Charge:**  $\nabla \cdot \vec{J} = -\partial \rho / \partial t = 0$  at each node

[Recall Gauss's Divergence Theorem:  $\int_V (\nabla \cdot \vec{G}) dv = \int_S \vec{G} \cdot d\vec{a}$ ]

**Current Law follows from:**  $\int_S \vec{J} \cdot d\vec{a} = \sum_i i_i = 0$



More generally: KCL assumes nodes store no charge, or that all electric energy is stored inside circuit elements

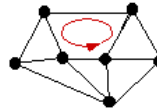
L10-3

## SOLVING CIRCUITS

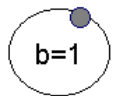
### Generic Circuit Topology:

Assume b branches, n nodes, and p unique loops

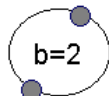
$b=12, p=6, n=7$



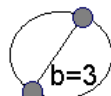
We can show  $b = n + p - 1$ .



$n = p = 1$



$n = 2$



$n = p = 2$

Every time we add a branch, the number of nodes or loops (meshes) increases by one.

Voltage sources are branches characterized by  $v = V_o$   
 Current sources are branches characterized by  $i = I_o$

L10-4

## SOLVING CIRCUITS (2)

**Number of Unknowns:** Number of unknowns is  $2b$  ( $v, i$  for each branch)

**Number of Equations:** We have one device equation for each branch relating or specifying  $v$  and  $i$

We also have  $n - 1$  independent node equations (KCL) and  $p$  loop equations (KVL), for a total of:

$$b + (n - 1) + p = 2b \text{ equations} = \text{number of unknowns}$$

Given initial conditions and all sources we can solve the equations analytically for simple circuits, or by simulation for any circuit.

(Non-unique exceptions: indeterminate flip-flop or chaotic circuits)

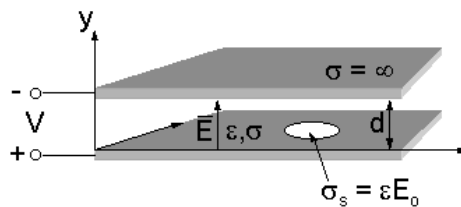
L10-5

## SIMPLE CIRCUIT ELEMENTS, R AND C

### Electrostatics:

In general:	Everywhere	$\nabla \times \vec{E} = 0, \quad \nabla \cdot \vec{E} = \rho/\epsilon$
	At conductor	$\hat{n} \cdot \vec{E} = \sigma_s/\epsilon, \quad \vec{E}_{\parallel} = 0$
	Between conductors,	$\nabla \cdot \vec{E} = \rho/\epsilon = 0$

### Parallel-Plate Devices:



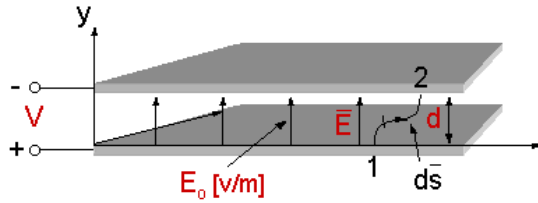
### Field Solution:

$$\vec{E} = \hat{y}E_0 \text{ for rectangular geometry here}$$

$$\sigma_s = \epsilon E_0$$

L10-6

## SIMPLE CAPACITOR



### Relating Fields to Potentials:

Since:  $\nabla \times \vec{E} = 0$

We define  $\vec{E} = -\nabla\Phi$      $\Phi = \int_1^2 \vec{E} \cdot d\vec{s}$

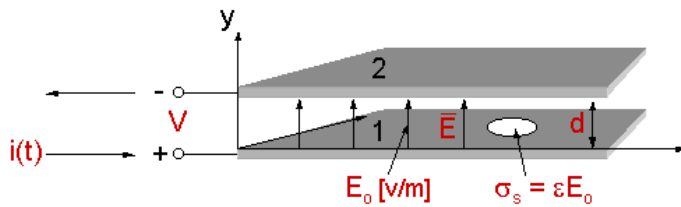
Where  $\Phi$  is the electrostatic potential relative to [2]

Therefore here:  $E_0 d = V$  ,  $E_0 = V/d$

In general, absolute potential  $\Phi = 0$  at infinity, by definition

L10-7

## SIMPLE CAPACITOR (2)



### Capacitor Charge Q:

We define:  $Q = \int_A \sigma_s da = A\epsilon E_0 = A\epsilon V/d$

We also define:  $Q = CV$

### Capacitance C:

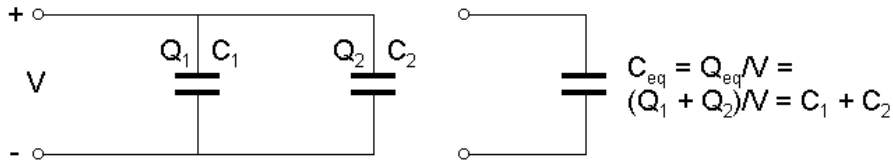
Therefore:  $C = A\epsilon/d$

We know:  $q(t) = \int_{-\infty}^t i(t) dt = Cv(t)$

Therefore:  $v(t) = (1/C) \int_{-\infty}^t i(t) dt$     Also,  $i = C dv/dt$

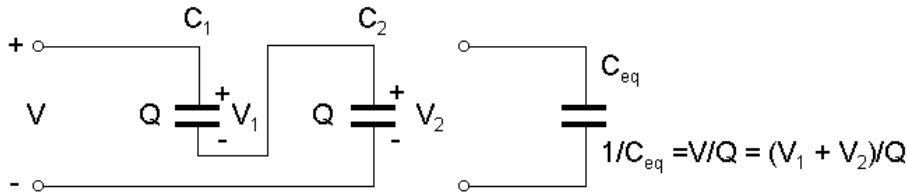
L10-8

## CAPACITORS IN SERIES AND PARALLEL



**Capacitors in Parallel:**

$$C_{eq} = C_1 + C_2$$



**Capacitor in Series:**

$$1/C_{eq} = 1/C_1 + 1/C_2$$

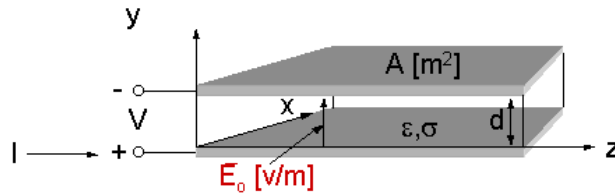
L10-9

## SIMPLE RESISTORS

**Conductance  $\sigma$ :**

$$\vec{J} = \sigma \vec{E} \quad [\text{am}^{-2}]$$

$$I = AJ = A\sigma E_0 = A\sigma V/d$$



**Resistance R:**

$$R = V/I = d/A\sigma$$

**Resistors in Series:**

$$R_{eq} = R_1 + R_2$$

$$R_{eq} = (V_1 + V_2)/I$$

**Resistors in Parallel:**

$$1/R_{eq} = 1/R_1 + 1/R_2$$

$$1/R_{eq} = (I_1 + I_2)/V$$

L10-10

## CHARGE RELAXATION

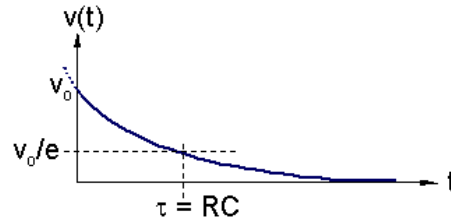
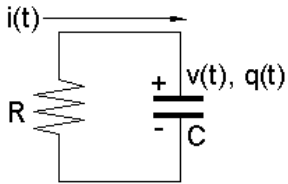
### RC Circuits:

$$q(t) = Cv(t) = \int i(t)dt + q_0$$

$$v(t) = -i(t)R = -RC dv(t)/dt$$

Try:  $v(t) = v_0 e^{-t/RC}$  [ $v_0$  is initial condition]

It works:  $v_0 e^{-t/RC} = (-RC)(-1/RC)v_0 e^{-t/RC}$



### Dielectric Relaxation:

$$R = d/A\sigma, C = \epsilon A/d$$

$$\tau = RC = \epsilon/\sigma \text{ seconds "Relaxation time constant"}$$

independent  
of geometry