

OPTICAL COMMUNICATIONS

Free-Space Propagation:

- Similar to radiowave propagation
- Antenna gain, effective area, path loss expressions unchanged

Devices:

- Detectors (review first recitation)
- Sources—LED's, lasers (next lecture?), amplifiers
- Modulators—amplitude and frequency, mixers, switches
- Passive filters, spectral multiplexers and combiners

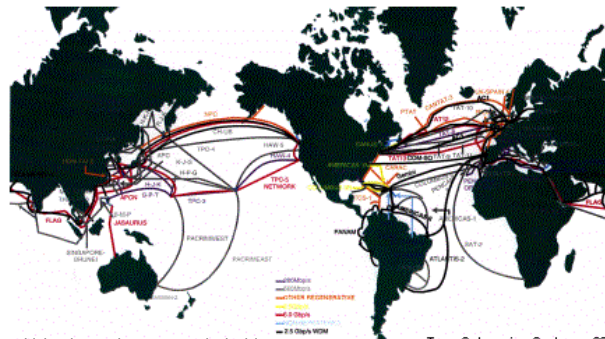
Guided Wave Propagation (including long lines and device interiors):

- Optical fibers trap and guide waves, attenuate little
- Rayleigh scattering is a loss mechanism, $\propto f^4$, favors $\lambda > 1$ -micron
- Rays inside fiber impact wall beyond critical angle
 \Rightarrow total reflection, totally lossless (for smooth walls; unlike mirrors)
- Attenuation $> \sim 1$? DB/km (depends on fiber architecture, materials, f)

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UNDERSEA OPTICAL FIBER CABLES

Fiber Communications Around the Globe



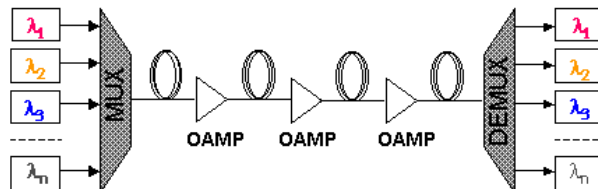
- Virtually all long-distance telecommunication is now by fiber optics
- In-line erbium-doped fiber amplifiers (EDFA's) make transoceanic transmission possible without repeaters – for many wavelengths at the same time in one fiber.
- Without fiber communications there would be no World Wide Web

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WDM MULTIPLEXED LINK

WAVELENGTH DIVISION MULTIPLEXING (WDM):

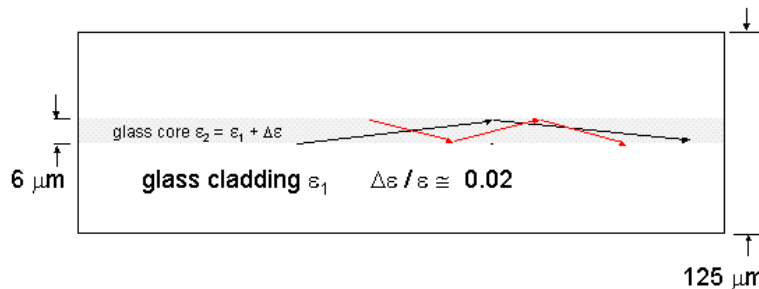
- Multiple wavelengths combined onto one fiber
- All wavelengths amplified simultaneously and independently in each optical amplifier (OAMP)



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WAVES IN FIBERS

Optical Fiber – Simple Picture:



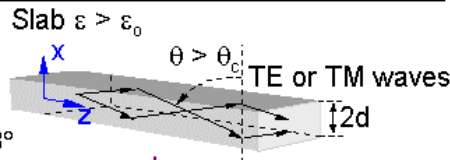
- Light is trapped by total internal reflection in the higher ϵ glass core.
- The small difference in ϵ implies very shallow reflection angles.
- Only certain angles are allowed since the waves must interfere constructively with each reflection \Rightarrow modes.
- Velocity of a mode is determined by the ϵ 's and the core size. (Different modes travel at different velocities.)

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OPTICAL WAVEGUIDES

Dielectric slab waveguide example:

Waves reflect beyond critical angle θ_c
 $\theta_c = \sin^{-1}(n_g^{-1})$ where $n_g \cong 1.5 \Rightarrow \theta_c \cong 41.8^\circ$

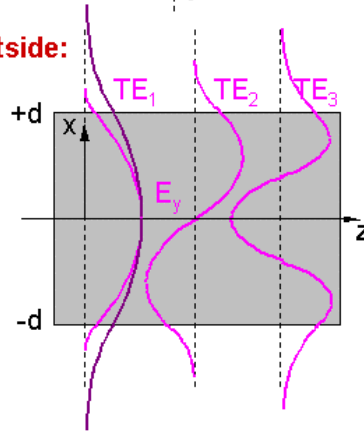


Standing waves inside guide, evanescent outside:

$$\vec{E} = \hat{y}E_0 \begin{Bmatrix} \sin k_x x \\ \cos k_x x \end{Bmatrix} e^{-jk_z z} \quad |x| \leq d$$

and $\vec{E} = \hat{y}\underline{E}_1 e^{-\alpha x - jk_z z}$ for $x > d$,

$$\vec{E} = \pm \hat{y}\underline{E}_1 e^{+\alpha x - jk_z z} \quad \text{for } x < -d$$



Evanescent region:

Decays more rapidly for lower modes and higher frequencies

Boundary conditions:

$$\vec{E}_{\parallel} \text{ and } \partial E_y / \partial x \text{ continuous for TE}_n \quad \nabla \times \vec{E} = -\hat{z} \partial E_y / \partial x - \hat{x} \partial E_y / \partial z = -\partial \vec{H} / \partial t$$

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ELECTROMAGNETIC FIELD DISTRIBUTION

Magnetic Field Distribution: $\vec{H} = -(\nabla \times \vec{E}) / j\omega\mu_0$ (for TE₁, TE₃, etc.)

Inside the slab:

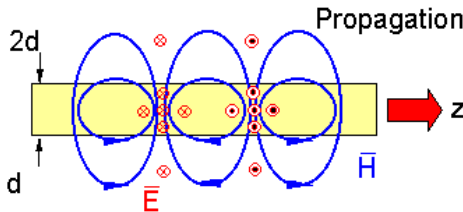
$$\vec{H} = (E_0 / \omega\mu) \left(-\hat{x}k_z \begin{Bmatrix} \sin k_x x \\ \cos k_x x \end{Bmatrix} - \hat{z}jk_x \begin{Bmatrix} -\cos k_x x \\ \sin k_x x \end{Bmatrix} \right) e^{-jk_z z} \quad \text{for } |x| < d$$

Outside the slab:

$$\vec{H} = (E_1 / \omega\mu_0) (-\hat{x}k_z - \hat{z}j\alpha) e^{-\alpha x - jk_z z} \quad \text{for } x > d$$

Matching Boundary Conditions:

Phase: $k_x^2 + k_z^2 = \omega^2\mu\epsilon$ inside the slab, $|x| < d$
 $-\alpha^2 + k_z^2 = \omega^2\mu_0\epsilon_0$ outside, $x > d$



Continuity of \vec{E} at $x = d$: $E_0 \cos k_x d e^{-jk_z z} = E_1 e^{-\alpha d - jk_z z}$ for TE_{1,3,5...}

Continuity of H_z at $x = d$: $(-jk_x E_0 / \omega\mu) \sin k_x d e^{-jk_z z} = -(j\alpha E_1 / \omega\mu_0) e^{-\alpha d - jk_z z}$

Therefore: $k_x d \tan k_x d = \mu\alpha d / \mu_0$ (ratio of continuity equations)
 $k_x^2 + \alpha^2 = \omega^2(\mu\epsilon - \mu_0\epsilon_0)$ (from dispersion equations)

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SOLUTIONS FOR TE_{odd n} DIELECTRIC SLAB WAVEGUIDES

Field Continuity Equations:

$$k_x d \tan k_x d = \mu \alpha d / \mu_0 \quad (\text{ratio of continuity equations})$$

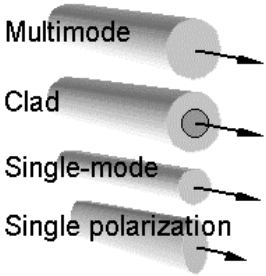
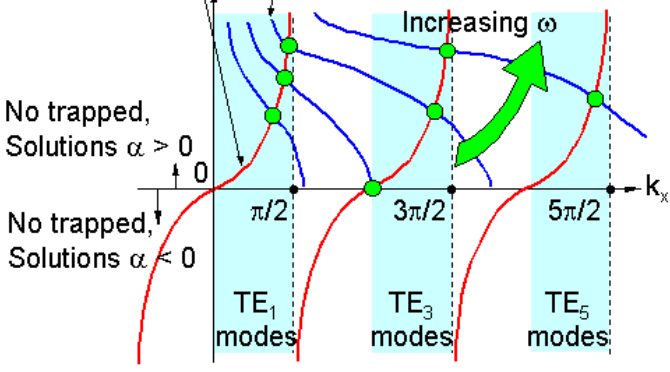
$$k_x^2 + \alpha^2 = \omega^2 (\mu \epsilon - \mu_0 \epsilon_0) \quad (\text{from dispersion equations})$$

Transcendental Equation: $\tan k_x d = (\mu / \mu_0) ([\omega^2 (\mu \epsilon - \mu_0 \epsilon_0) d^2 / k_x^2 d^2] - 1)^{0.5}$

Graphical solution:

Optical Fibers:

Bessel functions
Similar modes

No trapped, Solutions $\alpha > 0$

No trapped, Solutions $\alpha < 0$

TE₁ modes TE₃ modes TE₅ modes

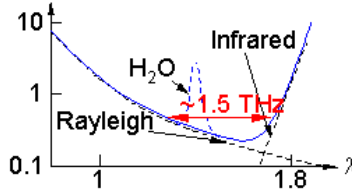
Increasing ω

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FIBER WAVEGUIDE DESIGN

Loss Mechanisms:

Rayleigh scattering from random density fluctuations
Loss $\propto f^4$ (scattering makes sky blue)
Infrared absorption dominates for $\lambda > \sim 1.6$ microns
Minimum total attenuation $\cong 0.2$ dB km⁻¹



Attenuation (dB km⁻¹)

Construction:

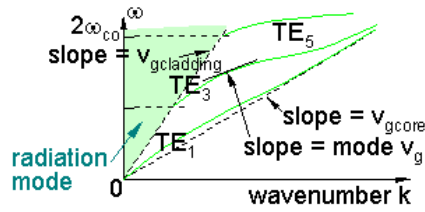
Typical: 10-micron core in 125-micron diameter glass, with 100-micron-thick plastic protective cladding (bundled in cables)

Manufacturing: Solid or hollow preform grown by vapor deposition of SiO₂ and GeO₂ (using e.g. Si(Ge)Cl₄ + O₂ = Si(Ge)O₂ + 2Cl₂)

Pulses Spread Due to Dispersion:

Group Velocity: Want $v_g(f) \cong \text{constant}$, so
Want flat $k(\omega)$ [$n = k/\omega + n_0$]

Dispersion: Determined mostly by $\epsilon(f)$,
modified by $\epsilon(r)$ of fiber



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EFFECTS OF DISPERSION

Pulse Spreading:

Distortion: Square pulse envelope is sum of harmonics--

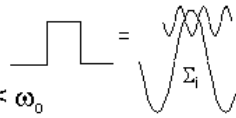
Want all f_i to have same group velocity; $\Delta\omega \ll \omega_0$

Equation: $k = \beta_0\omega_0 + \beta_1(\omega - \omega_0) + \beta_2(\omega - \omega_0)^2 + \dots$ where

$$\beta_0 = k/\omega_0 = v_p^{-1} = n/c$$

$$\beta_1 = dk/d\omega_0 = v_g^{-1} = (1 + [\omega/n]dn/d\omega)n/c$$

$$\beta_2 = d^2k/d\omega^2 = dv_g^{-1}/d\omega = (2dn/d\omega + \omega d^2n/d\omega^2)/c \text{ \{set to 0 at } \omega_0\}$$



Non-linearities:

- Avoid spikes: Large amplitudes generate harmonics at nonlinearities
- Large amplitudes: Desired to lengthen distance between amplifiers
- Nonlinearities: Occur in amplifiers and during propagation
- One remedy: Disperse signals initially (e.g. with grating) so fiber dispersion cancel this initialization over its entire length; Spikes reappear at end when signal is weak