

Massachusetts Institute of Technology  
Physics Department

Physics 8.322  
Quantum Theory II  
*Assignment 11*

Spring 2007  
April 30, 2007

DUE MAY 11, 2007, AT THE END OF THE DAY

**Announcements**

- There will be one more problem set in 8.322, due on May 16. I'll post it as soon as possible

**Reading topics for this period**

- Formal scattering theory; Born Approximation; Separable Potential

**Reading Recommendations 11**

- Sakurai, Formal development and Born Approximation: §7.1 — 7.3, and Method of Partial Waves: 7.5 — 7.6.
- Gottfried, Partial waves, §8.1; Formal development, 8.2(a); Born approximation, 8.3(a),(b).
- Schiff has a lengthy and pretty good section on scattering theory. §5 for basics and §9 for approximation methods.

**Problem Set 11**

**Topics covered in the problems**

- The phase shift and the radial wave function.
- Scattering from atoms in the first Born approximation.
- The Born approximation for a  $\delta$ -shell
- Spin dependent scattering

## Problems

### 1. A formula for the phase shift (postponed from last week.)

In our studying of scattering theory we derived a general formula for elastic scattering amplitude from a potential  $V(\vec{x})$ ,

$$f(\vec{k}', \vec{k}) = \sum_{\ell} (2\ell + 1) f_{\ell}(E) P_{\ell}(\cos \theta_{\vec{k}'\vec{k}}) = -\frac{2m}{4\pi} (2\pi)^3 \langle \vec{k}' | V | \psi_{\vec{k}}^{(+)} \rangle$$

Here,  $E = k^2/2m$ ,  $f_{\ell}(E) = e^{i\delta_{\ell}(E)} \sin \delta_{\ell}(E)/k$ , the state  $|\vec{k}\rangle$  is a plane wave,  $\langle \vec{x} | \vec{k} \rangle = e^{i\vec{k}\cdot\vec{x}}$ , and the state  $|\psi_{\vec{k}}^{(+)}\rangle$  is a solution regular at the origin and obeying outgoing wave boundary conditions. In lecture we showed that for a spherically symmetric potential,

$$\langle \vec{x} | \psi_{\vec{k}}^{(+)} \rangle = \sum_{\ell} (2\ell + 1) i^{\ell} e^{i\delta_{\ell}(k)} P_{\ell}(\cos \theta_{\vec{x}\vec{k}}) R_{\ell}^E(r),$$

where  $R_{\ell}$  is normalized so

$$\lim_{r \rightarrow \infty} R_{\ell}(k, r) = \frac{1}{kr} \sin(kr - \ell\pi/2 + \delta_{\ell}(k)).$$

- (a) From this starting point, derive an exact expression for the phase shift in the  $\ell^{\text{th}}$  partial wave,

$$\sin \delta_{\ell}(k) = -2mk \int_0^{\infty} dr r^2 V(r) j_{\ell}(kr) R_{\ell}(k, r)$$

Note that three different angles appear in this problem:  $\cos \theta_{\vec{x}\vec{k}} = \hat{k} \cdot \hat{x}$ ,  $\cos \theta_{\vec{k}'\vec{x}} = \hat{k}' \cdot \hat{x}$ , and  $\cos \theta_{\vec{k}'\vec{k}} = \hat{k}' \cdot \hat{k}$ . You will have to find identities involving Legendre polynomials and spherical harmonics in order to evaluate the angular integrals that occur here.

- (b) Find an expression for  $\sin \delta_{\ell}(k)$  to *first order* in  $V$ . This is the “first Born approximation”,  $\delta_{\ell}^{\text{1BA}}(k)$ .
- (c) Is  $|\sin \delta_{\ell}^{\text{1BA}}(k)|$  less than unity? If not, why?

### 2. Scattering from a atom or a nucleus in the first Born approximation

When low-energy electrons are scattered off atoms, the motion is non-relativistic and the atomic recoil can be neglected. If the atom is left in its ground state, where it began, the scattering is called *elastic*. Neglect spin. The cross section for scattering from incoming momentum  $\vec{k}$  to outgoing momentum  $\vec{k}'$  can equally well be regarded as function of the angles  $\theta$  and  $\phi$  defined relative to the  $\hat{k}$  axis.

- (a) What is the relation between the momentum transfer  $q \equiv |\vec{q}| = |\vec{k} - \vec{k}'|$  and the angles  $\theta$  and  $\phi$ ?

- (b) Show that *in the first Born approximation*, the differential cross section is given by

$$\frac{d\sigma}{d\phi d\cos\theta} = \left(\frac{2}{a_0 q^2}\right)^2 (Z - F(\vec{q}))^2$$

where  $a_0 = \hbar^2/me^2$ .  $F(\vec{q})$  is known as the *form factor* and contains all the information, measurable in elastic scattering, about the atomic electron charge distribution,  $\rho(\vec{x})$ .

$$F(\vec{q}) = \int d^3x e^{i\vec{q}\cdot\vec{x}} \rho(\vec{x})$$

- (c) Assuming that the charge density is spherically symmetric,  $\rho(\vec{x}) = \rho(r)$ , note  $F(0) = Z$ , and show that

$$\left.\frac{dF}{dq^2}\right|_{q=0} = -\frac{1}{6}\langle r^2 \rangle,$$

where  $\langle r^2 \rangle$  is the mean square charge radius of the atom. What is the *forward* ( $\theta \rightarrow 0$ ) limit of the differential cross section?

- (d) For hydrogen  $\rho(r) = \frac{1}{\pi a_0^3} e^{-2r/a_0}$ . Find  $F(q)$  and  $\langle r^2 \rangle$ .

Now consider scattering from a nucleus in the same approximation: non-relativistic, no spin, no recoil.

- (e) Nuclei have rather sharp surfaces. Model a nucleus with charge  $Z$  and mass number  $A$  as a uniform spherical charge distribution  $\rho(r, Z, A) = \rho_0(Z, A)$  for  $r \leq R(A)$  and  $\rho = 0$  for  $r > R(A)$ . A good estimate for  $R(A)$  is  $R(A) = R_0 A^{1/3}$ , with  $R_0 = 1.2\text{fm}$ . Compute the cross section for scattering in the first Born approximation as a function of  $Z$  and  $A$ .
- (f) Plot the results of the previous section as a function of  $\theta$  at fixed electron energy. Choose a few values of electron energy that show the characteristic structure of the cross section.
- (g) Real nuclei do not have sharp surfaces. Can you describe *qualitatively* what happens to the cross section you computed in the previous part if the surface is smeared out over  $\delta R \ll R(A)$ ?

### 3. Born Approximation for a $\delta$ -shell

Consider scattering from a  $\delta$ -shell potential,

$$2mV(r) = -\lambda\delta(r - a).$$

- (a) Find the first Born approximation to the scattering amplitude,  $f(k, \theta)$ .
- (b) Sketch  $f(k, \theta)$  as a function of  $\theta$  at fixed  $k$ .
- (c) Extract from  $f(k, \theta)$  the first Born approximation to the  $s$ -wave phase shift,  $\delta_0(k)$ .
- (d) Compute the  $s$ -wave phase shift directly by matching interior and exterior solutions. How good is the Born approximation?

#### 4. Spin dependent scattering

A particle with spin-1/2 (mass  $m$ ) scatters from a spin dependent potential. Suppose the spin “up” state is subject to a potential  $V(\vec{x})$ , while the spin “down” state is subject to no potential. In addition, suppose there is a “spin flip” potential  $W(\vec{x})$  which connects spin up and down, so

$$V = \begin{pmatrix} V(\vec{x}) & W(\vec{x}) \\ W(\vec{x}) & 0 \end{pmatrix}$$

in matrix notation. Assume  $V$  to be real, and assume both  $V(\vec{x})$  and  $W(\vec{x})$  to be short range.

- (a) Write down the Schrödinger equation (in terms of a two-component spinor wavefunction) for motion in this potential.
- (b) Introduce a spin index into the notation for scattering states:

$$\begin{aligned} |\phi_{\vec{k}}\rangle &\longrightarrow |\phi_{\vec{k},s}\rangle & s = \uparrow \text{ or } \downarrow \\ |\psi_{\vec{k}}^{(+)}\rangle &\longrightarrow |\psi_{\vec{k},s}^{(+)}\rangle \end{aligned}$$

and write the Lippmann–Schwinger integral equation for  $|\psi_{\vec{k},s}^{(+)}\rangle$

- (c) Calculate the scattering amplitude

$$f_{s,s'}(\vec{k}, \vec{k}') = -\frac{2m}{4\pi} \langle \vec{k}' s' | V | \psi_{\vec{k},s}^{(+)} \rangle$$

to *lowest non-trivial order* in the Born expansion for all four values of  $(s, s')$ . Express your answer in terms of the Fourier transforms of  $V(\vec{x})$  and  $W(\vec{x})$ . Note that the expansion for  $f_{\downarrow\downarrow}$  begins with a term quadratic in  $V$ . Simplify these results in the case that  $V$  and  $W$  depend only on  $r \equiv |\vec{x}|$ .