

## 8.08 Statistical Physics II

### Random Walks and Diffusion

(Dated: May 2, 2011)

A simplified model of diffusion consists of a one-dimensional lattice, with lattice spacing  $a$ , in which an "impurity" makes a random walk from one lattice site to an adjacent one, making jumps at time intervals  $\tau$ .

1. After  $N$  jumps have been made, find the probability that the atom has moved a distance  $d$  from its starting point, in the limit of large  $N$ .
2. The diffusion coefficient is defined by the differential equation  $D \frac{\partial^2 f}{\partial x^2} = \frac{\partial f}{\partial t}$ , where  $f$  is the concentration of the impurity. Find an expression for  $D$  in the model described above.

#### Solution:

1. Suppose that the particle made  $N_+$  steps forward and  $N_-$  steps backward. The distance traveled is  $d = a(N_+ - N_-) = as$ . We have the relations  $N_+ + N_- = n$  and  $N_+ - N_- = s$ . By solving these we get  $N_+ = (n+s)/2$  and  $N_- = (n-s)/2$ .

The probability of traveling a distance  $d$  is the probability of making  $N_+$  jumps to the right, so

$$P(d) = \frac{N!}{N_+!N_-!} = \frac{N!}{(N+s)/2!(N-s)/2!}$$

Note that this is un-normalized. We will normalize it when we go to the continuum limit.

Using Stirling's approximation we get

$$\ln P(d) = N \ln 2 - \frac{1}{2}(N+s) \ln(1+s/N) - \frac{1}{2}(N-s) \ln(1-s/N)$$

$P(d)$  will be sharply peaked around  $s = 0$ , so we Taylor expand around  $s/N = 0$  and get

$$P(d) = 2^N e^{-s^2/2N}$$

By setting  $x = as$  and normalizing we get

$$P(x) = \frac{1}{a\sqrt{2\pi N}} e^{-x^2/2Na^2}$$

2. The number of jumps is the time over the time it takes to make one jump,  $N = t/\tau$ . So

$$P(x, t) = \sqrt{\frac{\tau}{2\pi a^2 t}} e^{-x^2 \tau / 2a^2 t}$$

The density of concentration of impurities in the lattice is  $f(x, t)$ . If we start with  $M$  impurities at time  $t = 0$  in positions  $x_0^1, x_0^2, \dots, x_0^M$  then  $f(x, t) = \sum_{i=1}^M P(x - x_0^i, t)$ .

From linearity substitute  $P(x)$  into

$$D \frac{\partial^2 f}{\partial x^2} = \frac{\partial f}{\partial t}$$

and get

$$D = a^2/2\tau$$

3. Compute the ratio of thermal conductivity of helium gas at  $P = 0.1atm$  and  $300K$  to that at  $P = 0.5atm$  and  $300K$ . Compute also the ratio of viscosities at the two pressures

**Solution:** The rate of energy transfer through a unit area in the plane  $Z = const$  is proportional to

$$(nv)(-\lambda dE/dT) = (nv\lambda dE/dT)(-dT/dZ)$$

The first factor is the flux of molecules in either direction across the plane. The second term is the difference in average molecular energy in a distance of one mean free path ( $\lambda$ ). The coefficient of  $(-dT/dZ)$  is the heat conductivity,  $K \propto nv\lambda C$ , where  $C$  is the molecular specific heat.

The mean free path is  $\lambda = 1/n\sigma$ , where  $\sigma$  is the collision cross section. In addition,  $C = 3k/2$  for a monoatomic gas and  $v = \sqrt{3kT/M}$ . Finally

$$K \propto \sqrt{T}$$

independent of pressure, so the desired ratio is 1.

The viscosity may be calculated by considering the net transverse momentum transfer across a unit area of a plane  $Z = const$ , in unit time. In the same notation, this is proportional to  $-(nv)(m\lambda du/dZ)$ , where  $m$  is the molecular mass and  $u$  is the transverse velocity. By definition, the coefficient of  $du/dZ$  is the viscosity  $\eta$ . Thus

$$\eta \propto \sqrt{T}$$

independent of pressure, so the ratio is 1. Furthermore

$$\eta/K = const$$

4. In a hot plasma, all the atoms may be regarded as completely ionized. Although the ions have long-range forces due to Coulomb interactions, macroscopically the plasma is electrically neutral. This suggests that the Coulomb interactions are screened, and so become short-range. Estimate this range, making suitable approximations.

**Solution:**

The electrical potential in the vicinity of an ion is  $\phi(r)$ . The density near the ion is

$$n(r) = ne^{-e\phi(r)/kT}$$

Poisson's equation is

$$\nabla^2\phi(r) = -4\pi \sum_{\alpha} e_{\alpha}n_{\alpha}$$

For a hot plasma

$$n_{\alpha}(r) = n_{\alpha}(1 - e_{\alpha}\phi/kT)$$

Substituting into the Poisson equation

$$\nabla^2\phi(r) = 4\pi \left( \sum_{\alpha} \frac{n_{\alpha}e_{\alpha}^2}{kT} \right) \phi = \kappa^2\phi$$

which has a solution

$$\phi(r) = e_{ion} \frac{e^{-\kappa r}}{r}$$

$\kappa^{-1}$  is a characteristic length scale of the plasma and is called the Debye radius.