The Influence of Transient Extensional Viscosity on Flow Through an Axisymmetric Contraction-Expansion

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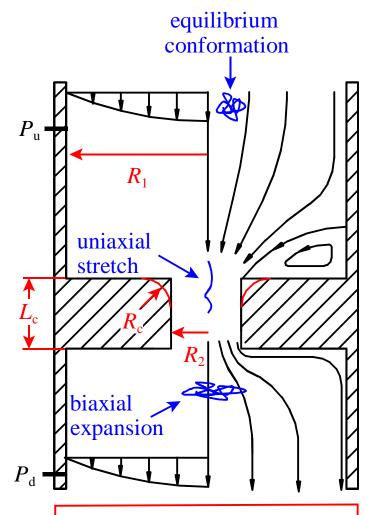




The Axisymmetric Contraction-Expansion



(Cartalos & Piau, 1992; Szabo et al., 1997; Rothstein & McKinley,1999)



$$R_1 / R_2 = \beta = 2, 4, 8$$

 $L_c / R_2 = 0.5, 1, 2$
 $R_c = 0, 0.1R_2, 0.2R_2, 0.5R_2$

 Analog of 'orifice plate' in Newtonian fluid mechanics!

Sampson (1891)
Dagan et. al (1982)
$$\Delta P_N = \frac{3\eta Q}{\pi R^3} \left(1 + \frac{8}{3} \frac{L_c}{R_2} \right)$$

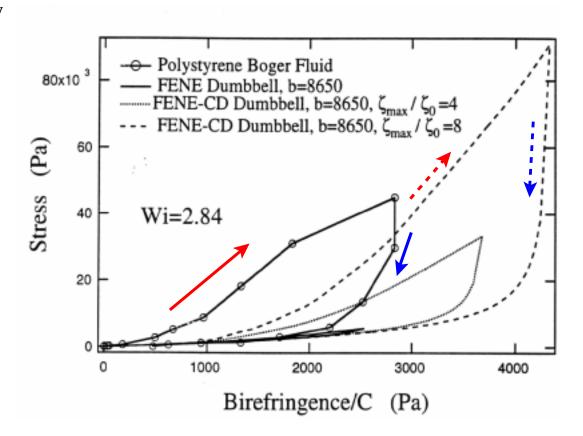
- Upstream contraction flow is numerical and experimental benchmark problem
- Complex flow containing mixture of shear near walls and extensional effects in vicinity of contraction
- Characteristic Deborah number

$$De = \lambda \gamma = \lambda Q / \pi R_2^3$$

- Determine extra pressure drop $\Delta P_{\rm ext}$ across orifice
 - Insight into existence of purely dissipative 'internal' stress

Motivation

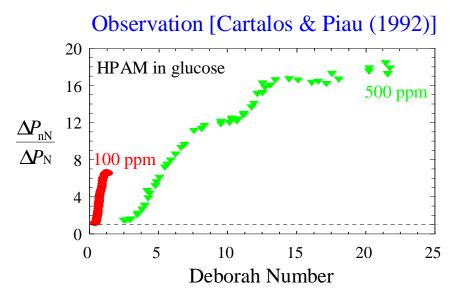
- Existence of a purely dissipative 'internal' stress in homogeneous uniaxial elongation resulting in a *stress-conformation hysteresis* proposed by Ryskin(1987) and Doyle *et al.*(1998)
- Stress increases along one pathway as fluid filament is extended and relaxes down a very different pathway once flow is removed
- Result of *stress-conformation hysteresis* is large energy loss
- FENE dumbbell models demonstrate little hysteresis
- FENE dumbbell models with additional conformation dependent drags can increase hysteresis but cannot quantitatively match experimental data



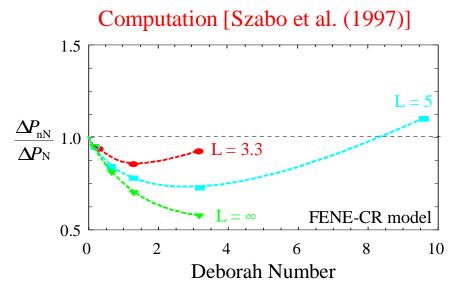
Manifestation of Hysteresis in Complex Flows



• We wish to observe impact on prototypical non-homogeneous flow





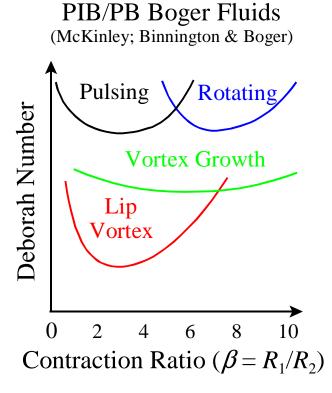


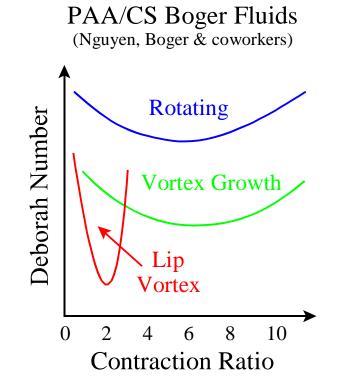
Entropic elasticity and energy storage
 ⇒ Reduced pressure drop

Question: Is this discrepancy due to 'dissipative' or 'internal' stresses arising from non- equilibrium conformation?

- Similar discrepancies between observations and computation present in other complex flows
 - ° C_D of spheres [Solomon & Muller (1996)]
 - o Pressure drop across banks of cylinders [Khomami & Moreno (1997)]

Overview of Flow Stability in Abrupt Contractions

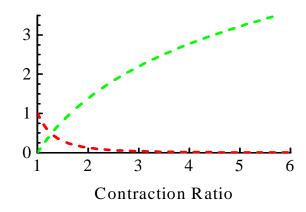




PS/PS Boger Fluids (Rothstein & McKinley)

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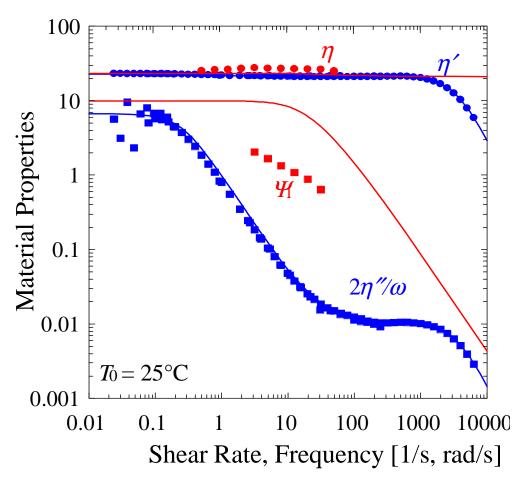
- Why do vortex growth dynamics change with contraction ratio and Boger fluid?
- Competing roles of upstream shear rate $y_1/y_2 = \beta^{-3}$ and accumulated strain $\varepsilon = \ln(\beta^2)$
- Measurements by Shelley Anna (SF3) of transient extensional viscosity may yield insights



Fluid Rheology



- Monodisperse polystyrene $Mw = 2.03 \times 10^6$ g/mol dissolved in oligomeric polystyrene
- Dilute solution with concentration c = 0.025wt% $\Rightarrow c/c^* = 0.23$
- Model viscoelastic fluid to probe hysteresis in the absence of polydispersity effects



• Small amplitude oscillatory shear rheology well fit by Rouse-Zimm bead-spring model

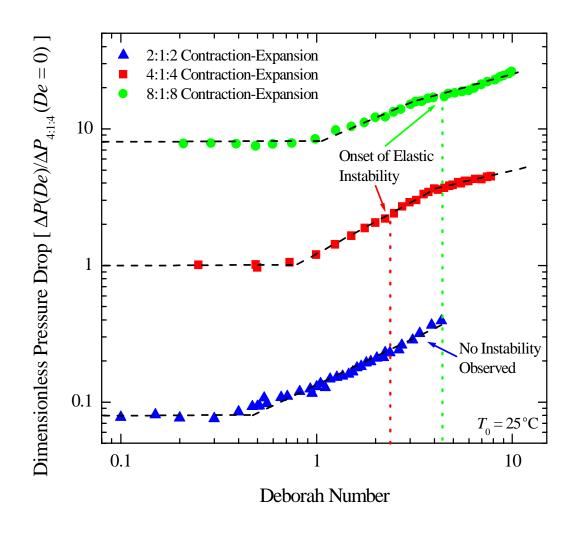
$$h^* \simeq 0.1 \Rightarrow$$
 dominant hydrodynamic interactions

$$\lambda_1 = 3.08s$$
 $\beta = \eta_s/\eta_0 = 0.92$
 $\lambda_i = \lambda_1/i^{1.77}$

- Weakly elastic polymeric solvent $\lambda_{ps} = 2.5 \times 10^{-4} \text{ s}$
- Steady shear data poorly fit by FENE-P model with a finite extensibilty of L = 88
- Temperature variations described by shift factor of WLF theory.

Pressure Drop Across Contraction/Expansions of β = 2,4 and 8

- All contraction ratios show similar pressure growth trends.
- Pressure drop increase with *De* before onset of elastic flow instability.



• Newtonian-like response

$$\Delta P_{\rm nN} \sim Q$$

$$0 < De \lesssim 1$$

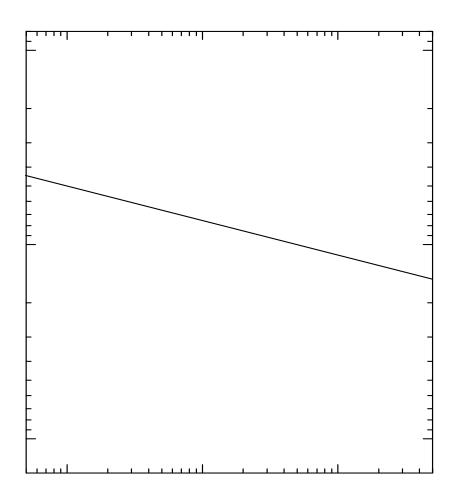
Pressure drop increases quadratically with flow rate

$$\Delta P_{\rm nN} \sim Q^2$$
 $1 \leq De \leq 5$

 Pressure drop increases more slowly approaching a linear dependence on flow rate

$$\Delta P_{\rm nN} \sim Q$$

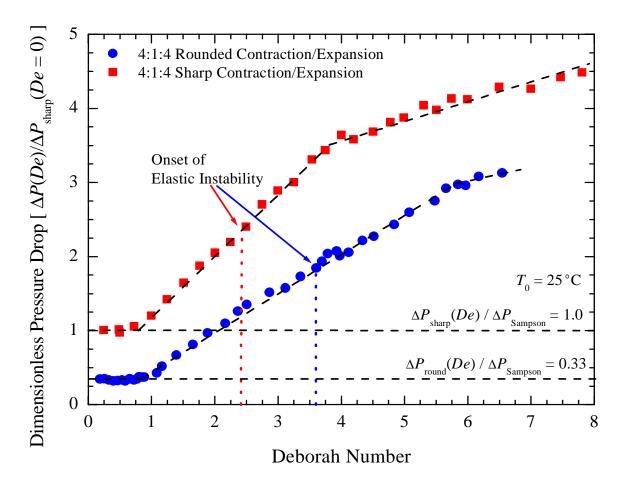
$$De \gtrsim 5$$



Effect of Re-entrant Lip Curvature on Pressure Drop



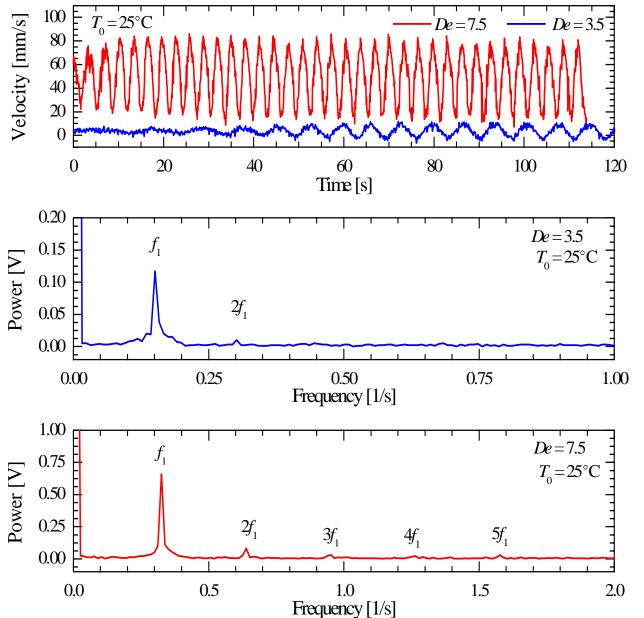
- Sharp and rounded entrance lip pressure drop data follow same trends
- Small changes in lip curvature have huge effect on pressure drop



• Rounding lip entrance delays onset of each pressure drop transition and elastic flow instability

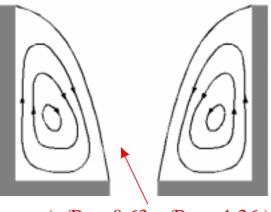
Onset of Elastic Instability





- Small amplitude
 oscillations seen in Laser
 Doppler Velocimetry
 (LDV) measurements at
 large De
- Higher harmonics of fundamental frequency appear as *De* increases

Location of LDV Probe



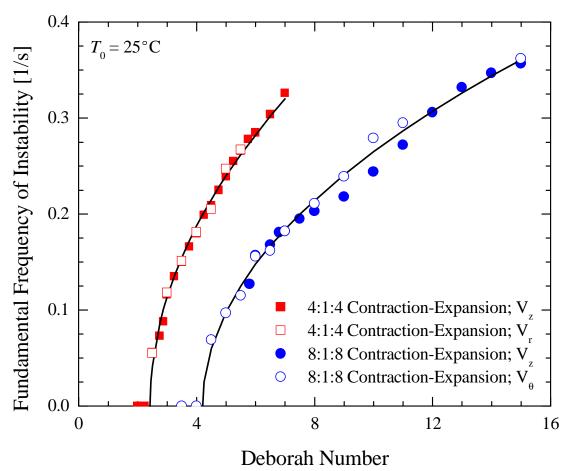
 $(r/R_2 = 0.63, z/R_2 = -1.26)$

Characterization of Elastic Instability



	Critical Deborah Number for Onset of Instability		
4:1:4 Sharp Contraction/Expansion	2.4 ± 0.1		
4:1:4 Round Contraction/Expansion	3.6 ± 0.2		
8:1:8 Sharp Contraction/Expansion	4.2 ± 0.1		

• Lip curvature delays the onset of the elastic instability



 Solid lines denote theoretical result for supercritical Hopf bifurcation

$$f_1 \propto (De - De_{crit})^{1/2}$$

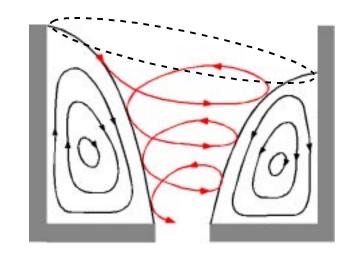
• Dimensionless frequency of elastic flow instability

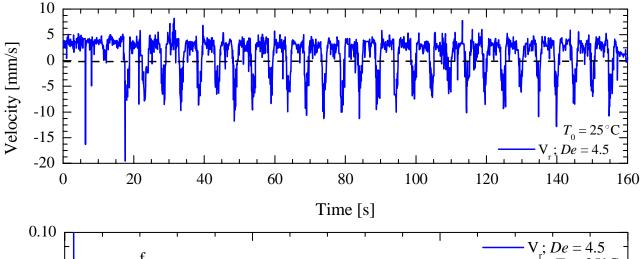
$$\lambda_{\rm z} f_1 = \frac{\lambda_{\rm z} \omega}{2\pi} \sim 1$$

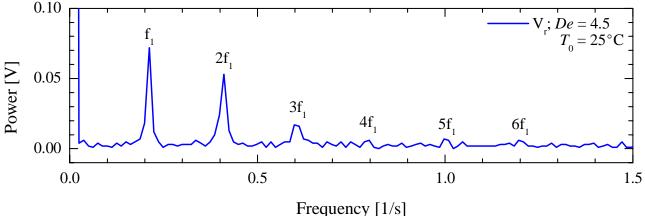
$$\lambda_z = 3.08s$$

Nonlinear Dynamics: Jetting

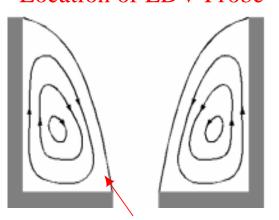
- Localized elastic jetting instability seen in 4:1:4 contraction-expansion after onset of vortex asymmetry
- We postulate that instability is helical jet of high speed fluid traveling down the interior of vortex structure







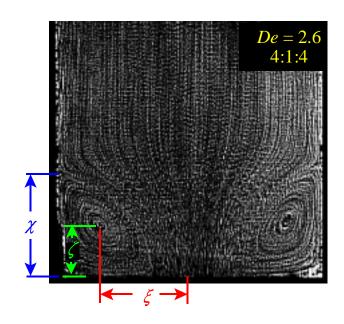
Location of LDV Probe

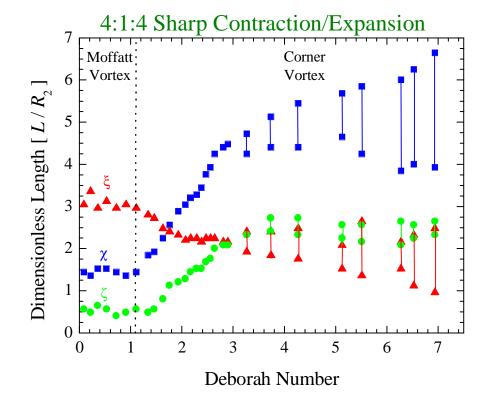


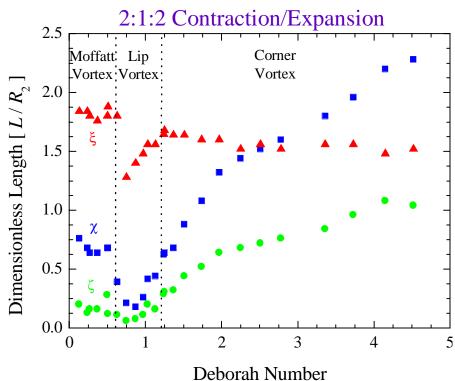
 $(r/R_2 = 1.57, z/R_2 = -0.16)$

Vortex Growth Dynamics

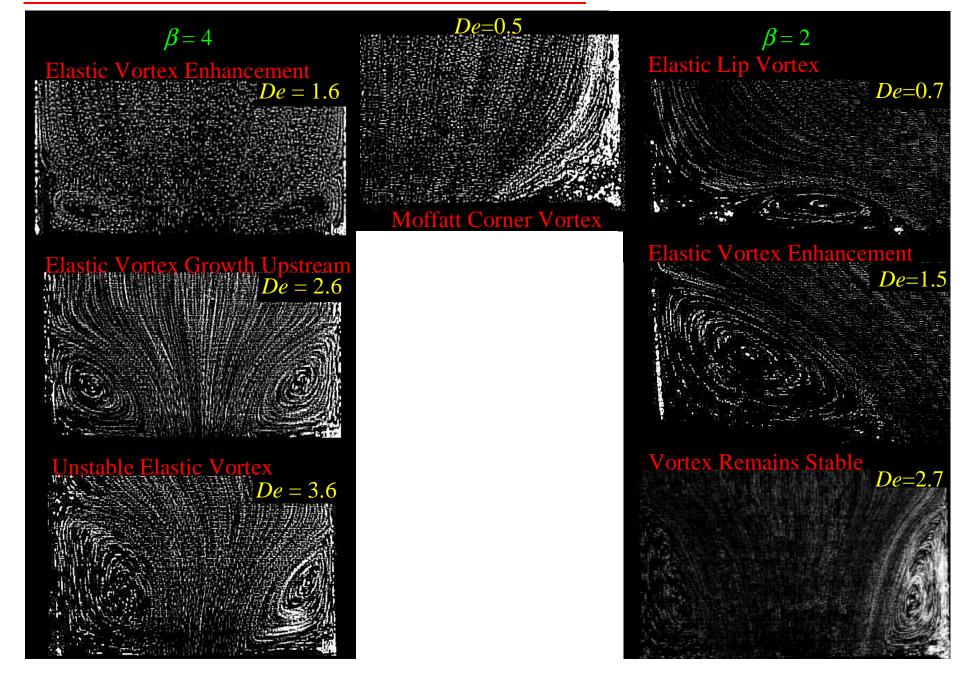
- At low *De*, Newtonian-like Moffatt vortex exist
- At moderate De, two distinct patterns of vortex growth exist
 - For $\beta \ge 4$ vortex grows out from salient corner, grows upstream and eventually becomes unstable
 - For $\beta = 2$ new lip vortex emerges near re-entrant corner, grows toward the salient corner and then proceeds to grow upstream



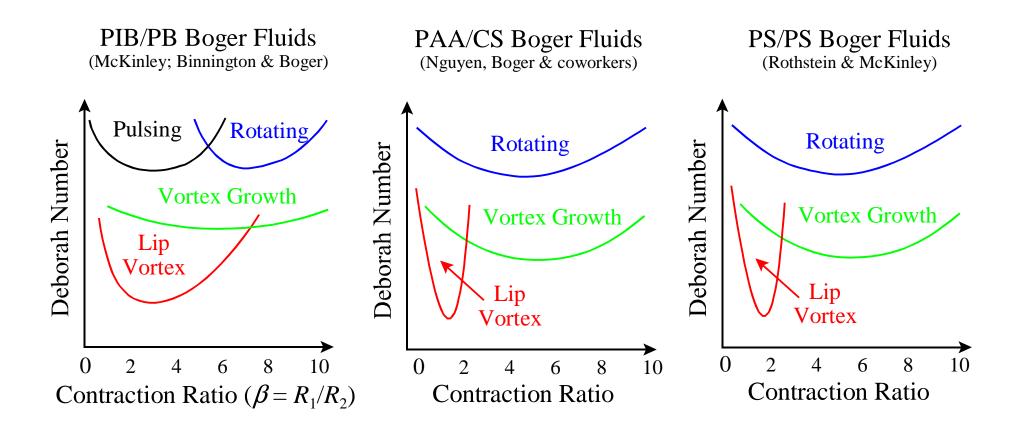




A Proper Taxonomy of Vortex Structure



Flow Stability Diagram for PS/PS Boger Fluid



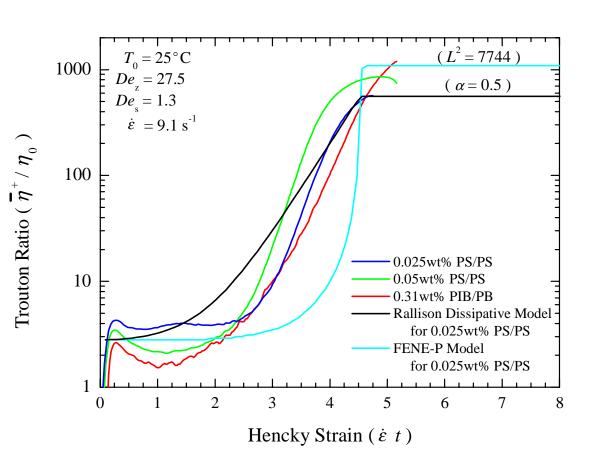
• Flow stability diagram of PS/PS Boger fluid is very similar to PAA/CS Boger fluid and dissimilar to PIB/PB Boger fluid

Transient Extensional Rheology (In collaboration with Shelley Anna)



• Why do two fluids with similar shear properties behave so differently in complex flows?

- PIB and PS Boger fluids exhibit markedly different stress growth
- Results $\bar{\eta}^+(\dot{\varepsilon_0}) / \eta_0$ are approximately independent of $De_z = \lambda_z \dot{\varepsilon_0}$ (for $De_z > 1$)
- How much stress growth do we get in an axisymmetric contraction?



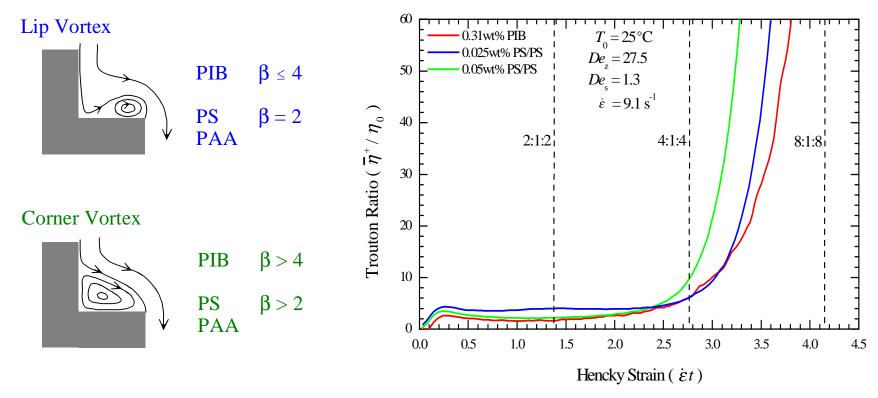
Accumulated Strain Along Centerline

$$\varepsilon = \int_{0}^{t_1} \dot{\varepsilon} dt = \int_{v_z(-\infty)}^{v_z(z = 1/2L_c)} \frac{dv_z}{v_z} = \ln\left(\frac{R_1}{R_2}\right)^2 = 2\ln\beta$$

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Comparison of PS and PIB Boger Fluids





- Experiments suggest that the lip vortex is present for smaller contractions (with larger shear component) while the corner vortex is present in flows dominated by extension.
- However, extensional viscosity alone still cannot differentiate between the PS and PIB Boger fluids if we consider only the pure extension along the centerline.
- Does the shearing near the wall increase the effective Hencky strain? Can we investigate this by preshearing the fluid before stretching it?

Concluding Remarks



- Observed large monotonic increase in pressure drop not correlated with the onset of elastic flow instability and not qualitatively affected by changes in lip curvature
 - ⇒ Is this the result of the *stress-conformation hysteresis* observed in transient uniaxial extension?
- To answer this question more concretely, molecular simulations in transient inhomogeneous flows need to be performed

- Evolution in flow structure of PS/PS Boger fluid similar to PAA/CS Boger fluids and dissimilar to PIB/PB Boger fluids
 - → Are variations in vortex growth dynamics the result of differences in transient extensional rheology of fluids?
- Results of simple transient uniaxial extension are inconclusive
- Preliminary results of preshear to transient uniaxial extension show considerable impact on Trouton ratio.