

Opportunities for Radioactive Molecules in Tests of Fundamental Symmetries and Searches for New Physics

Yevgeny Stadnik

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Conventional Wisdom in the Classification of Atomic/Molecular EDM Experiments

Diamagnetic systems (contain *no* unpaired electrons) are mainly sensitive to **hadronic** sources of CP violation – e.g., **Hg, Xe, n**

Paramagnetic systems (contain *one or more* unpaired electrons) are mainly sensitive to **leptonic** sources of CP violation
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– e.g., **ThO, HfF⁺, YbF, Tl, Cs**

For **semi-leptonic** sources of CP violation, the story is more complicated – the “classification” generally depends on whether the interactions involve mainly **electron spin** or **nuclear spin**

Leptonic CP Violation in Paramagnetic Molecules

Over the past decade, molecular experiments have improved the sensitivity to electron EDM d_e by more than 100-fold:

$$^{232}\text{ThO bound: } |d_e| < 10^{-29} e \text{ cm}$$

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What about sensitivity of paramagnetic systems to hadronic CP violation?

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Hadronic CP Violation in Diamagnetic Atoms

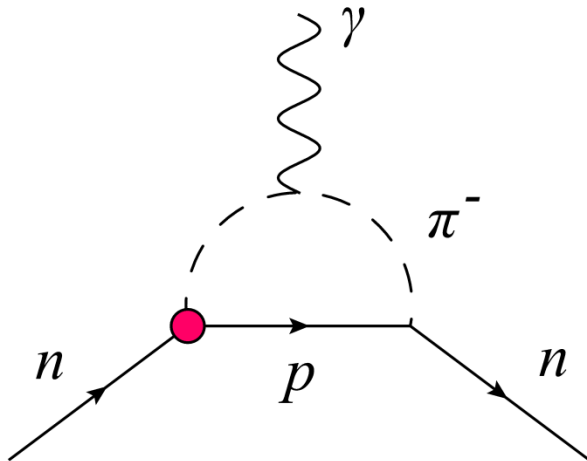
Illustrative example: $\mathcal{L} = \theta \frac{g_s^2}{32\pi^2} G\tilde{G}$

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Nucleon EDMs



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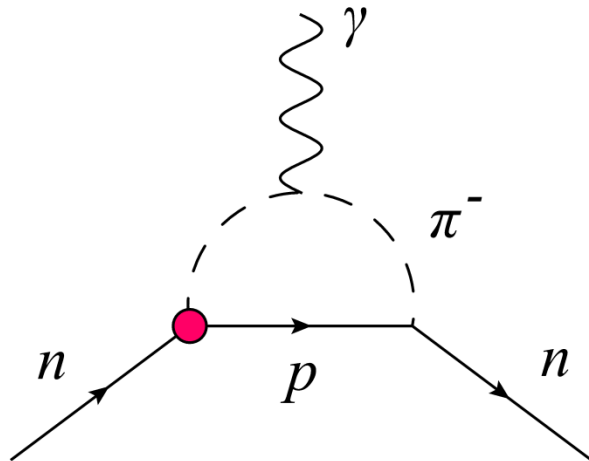
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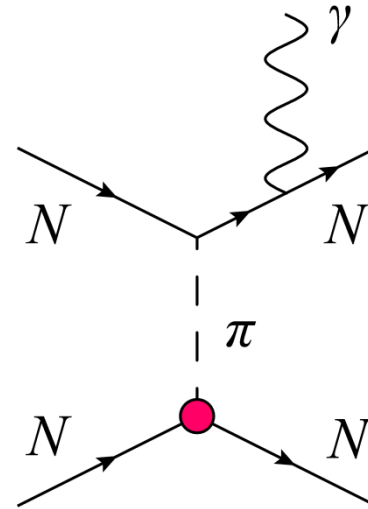
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CP-violating intranuclear forces



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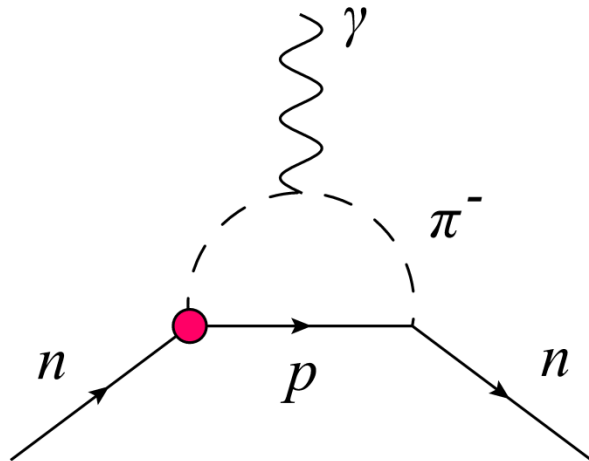
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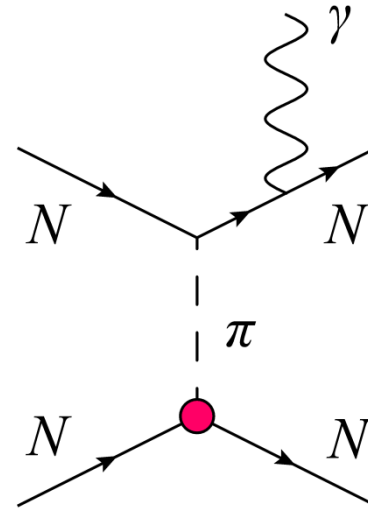
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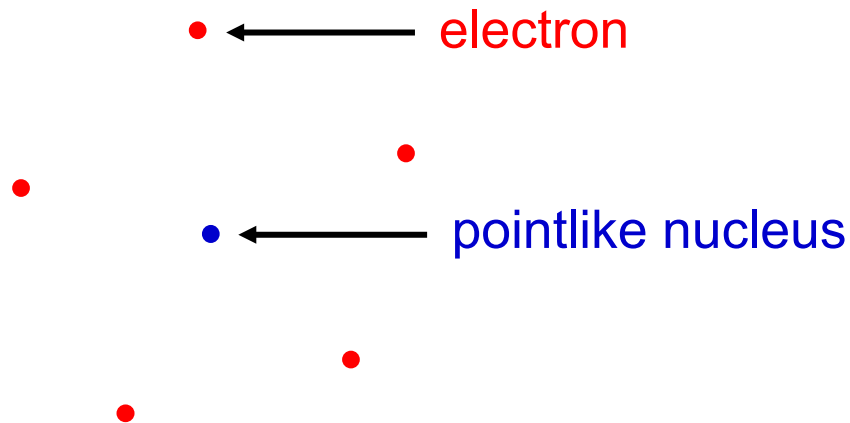


In nuclei, **tree-level** CP-violating intranuclear forces dominate over **loop-induced** nucleon EDMs [loop factor = $1/(8\pi^2)$].

Schiff's Theorem

[Schiff, *Phys. Rev.* **132**, 2194 (1963)]

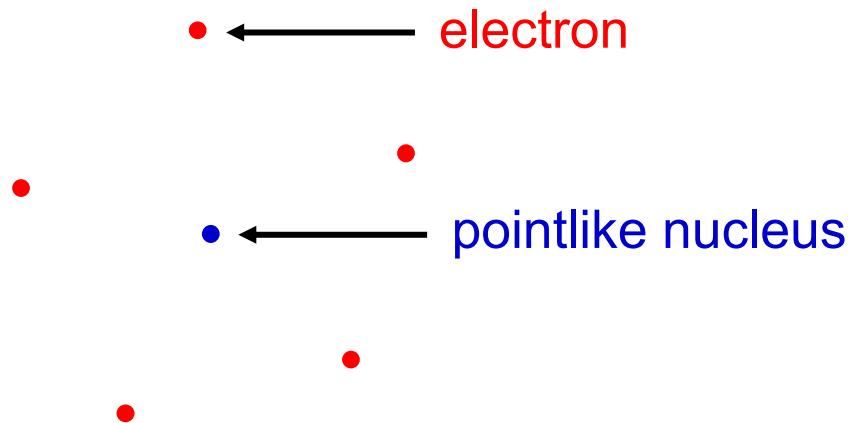
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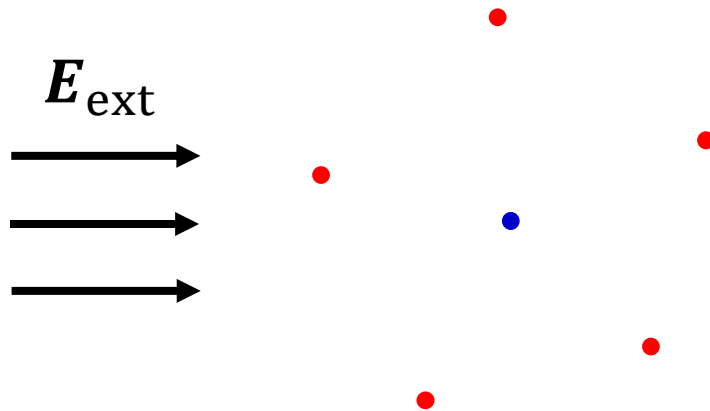


Classical explanation for nuclear EDM: A neutral atom does not accelerate in an external electric field!

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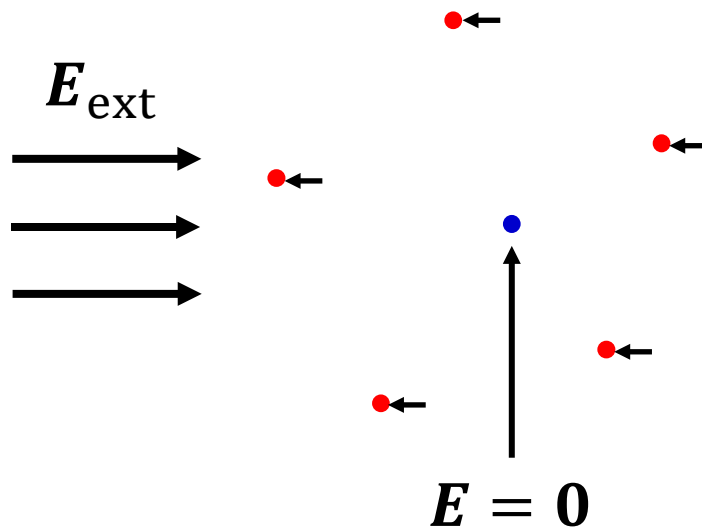


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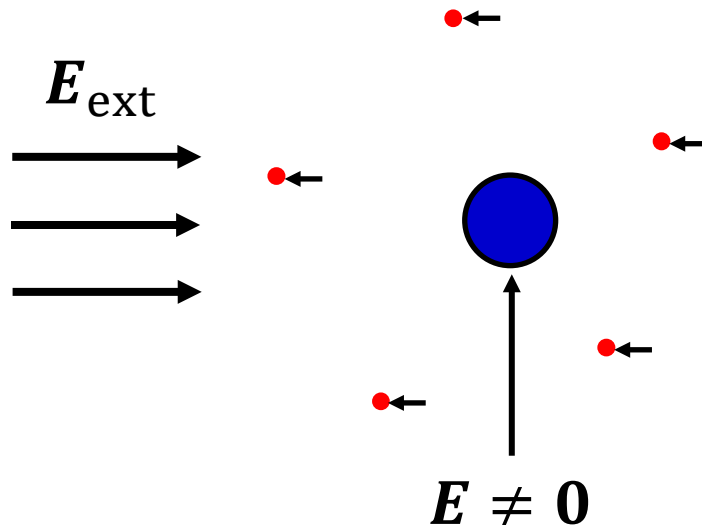
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Lifting of Schiff's Theorem

[Sandars, *PRL* **19**, 1396 (1967)],

[O. Sushkov, Flambaum, Khriplovich, *JETP* **60**, 873 (1984)]

In real (heavy) atoms: Incomplete screening of external electric field due to finite nuclear size, parametrised by **nuclear Schiff moment**.



Hadronic CP Violation in Paramagnetic Molecules

[Flambaum, Pospelov, Ritz, Stadnik, *PRD* **102**, 035001 (2020)]

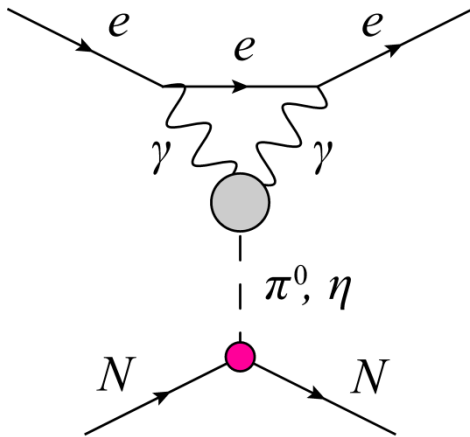
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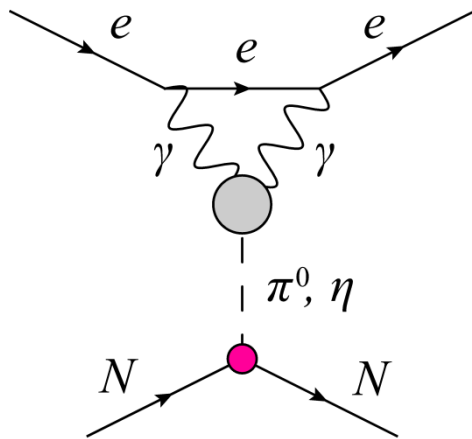
**CP-odd nucleon
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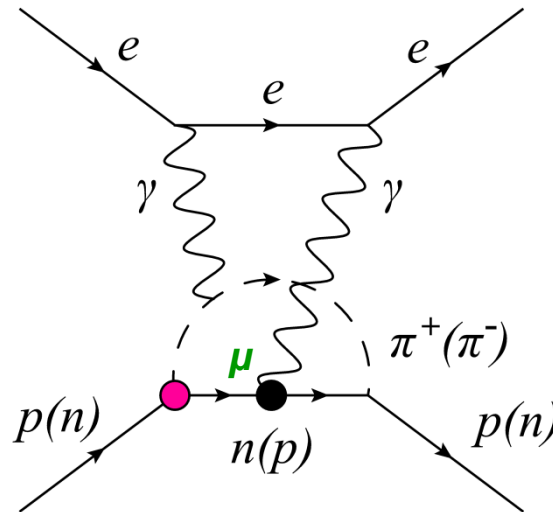
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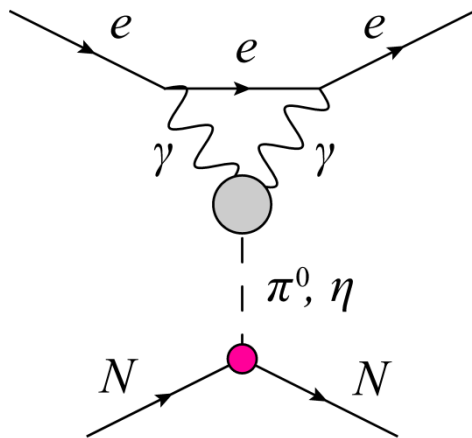
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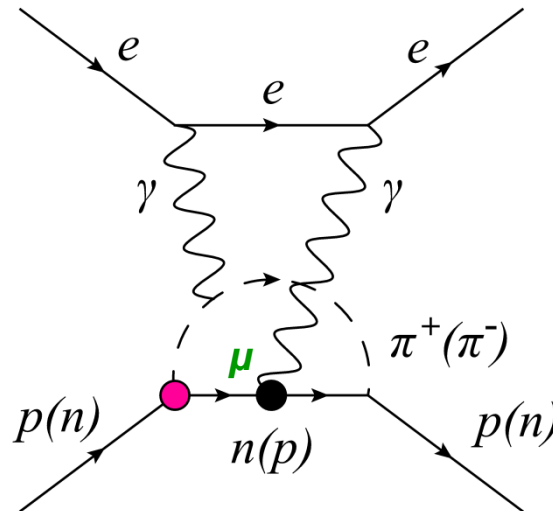
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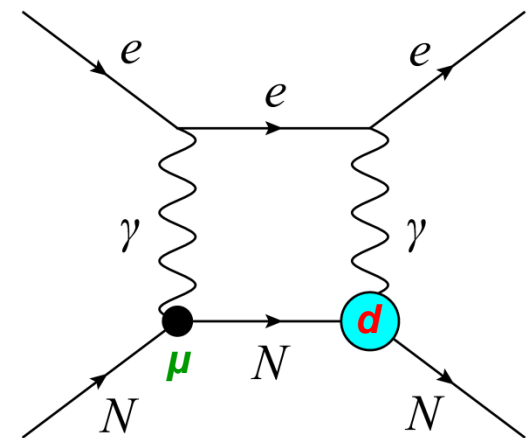
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[Fermi-gas model]



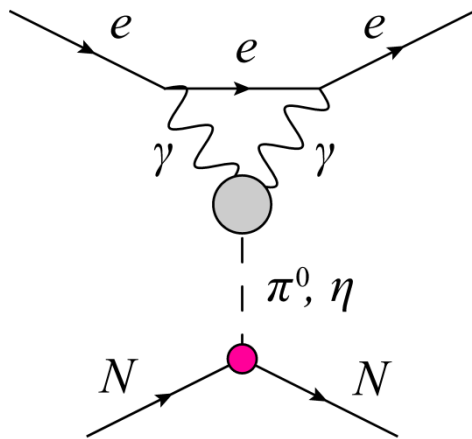
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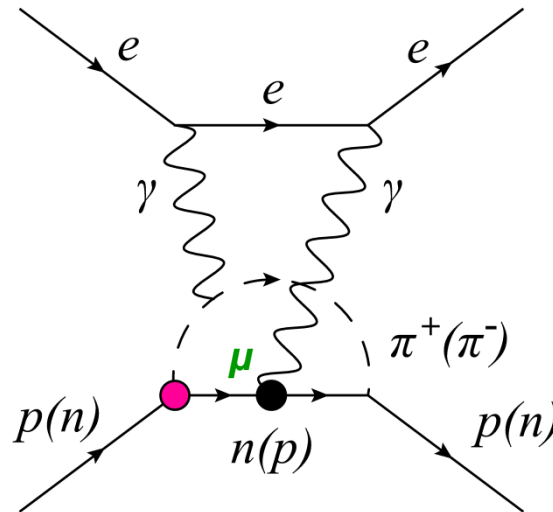
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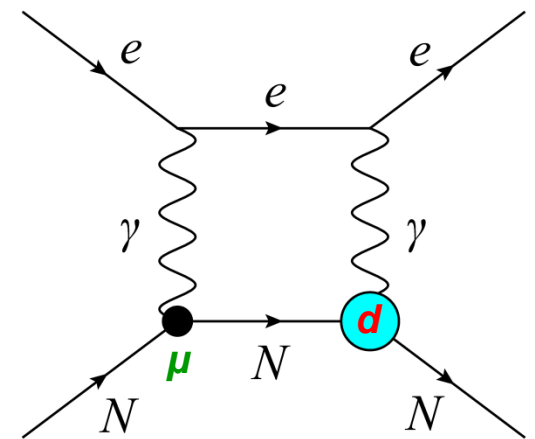
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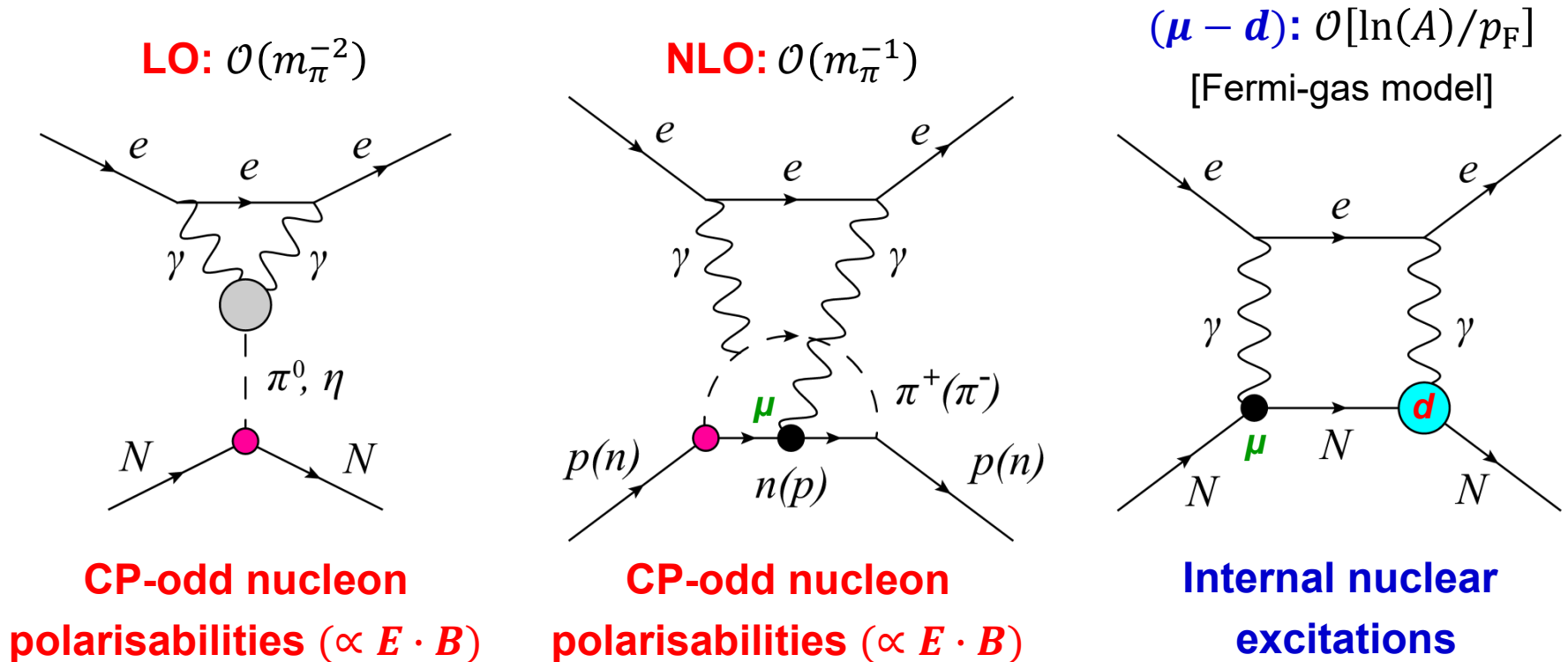


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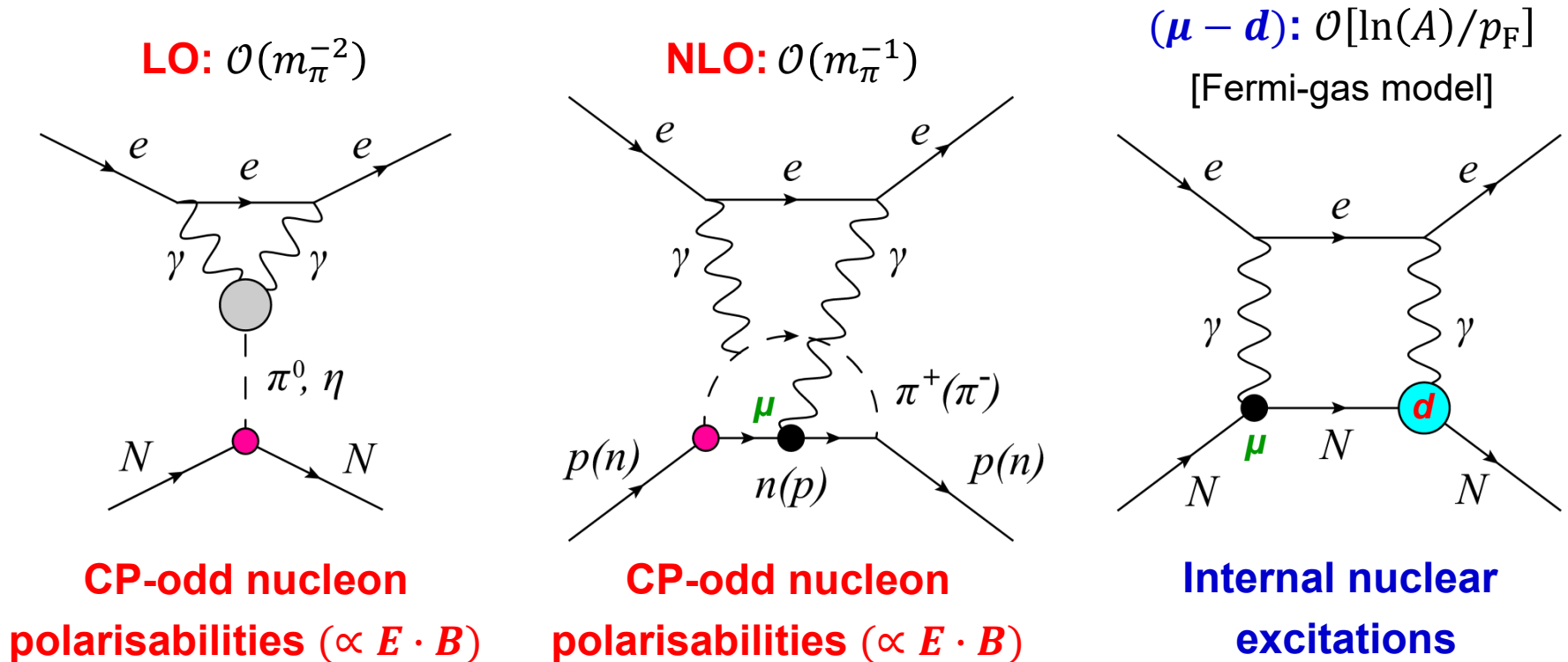
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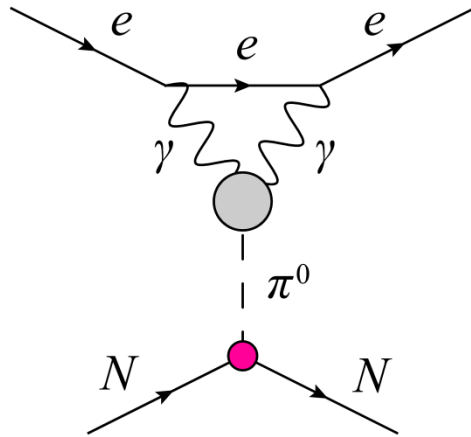
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 - Operative even in *spinless* nuclei (e.g., ^{232}ThO , $^{180}\text{HfF}^+$)



Isoscalar CP-Odd π - N Coupling

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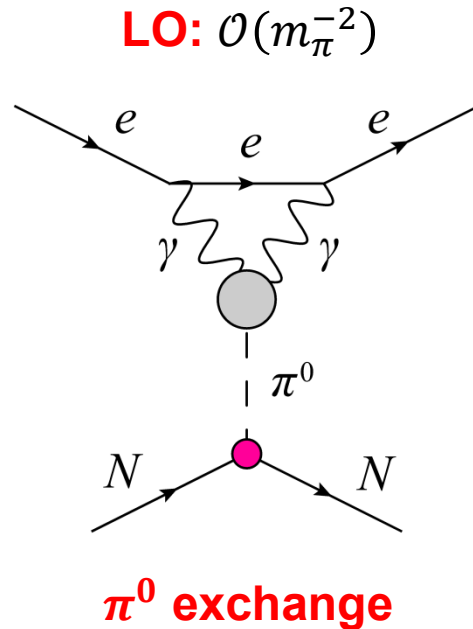


π^0 exchange

$$\mathcal{L} = \bar{g}_{\pi NN}^{(1)} \pi^0 \bar{N} N$$

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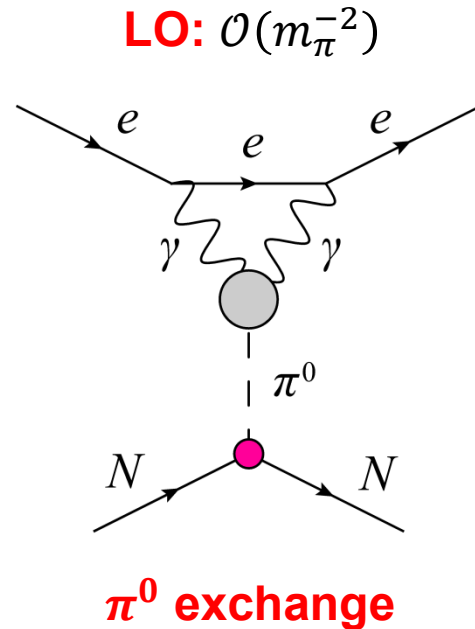


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In molecules with *spinless* nuclei (e.g., ^{232}ThO , $^{180}\text{HfF}^+$), effect dominated by a “**bulk**” property of the nucleus that grows with A in a regular manner, with *no contribution* from the nuclear Schiff moment mechanism (needs $I \neq 0$)

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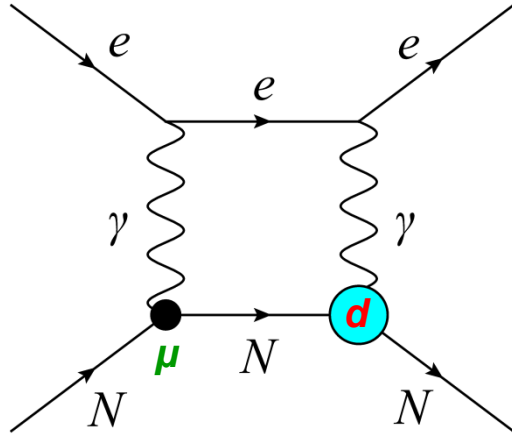
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=> Clean bounds, since less sensitivity to details of nuclear structure

(cf. strong sensitivity of ^{199}Hg Schiff moment to assumptions about underlying nuclear structure – different models give different signs for sensitivity coefficient)

Nucleon EDMs $\mathcal{L} = -\frac{i}{2} d_N \bar{N} F_{\mu\nu} \sigma^{\mu\nu} \gamma_5 N$

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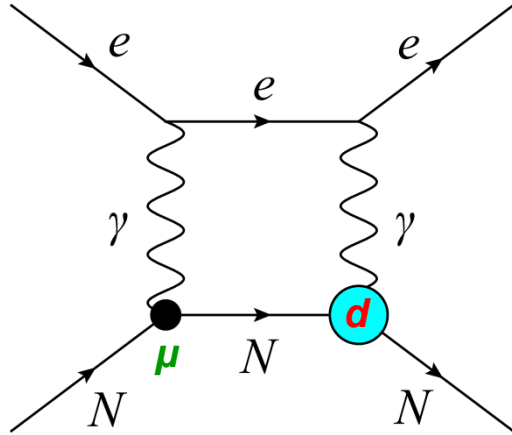


Internal nuclear excitations

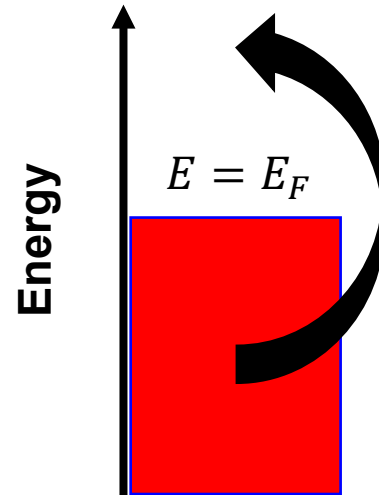
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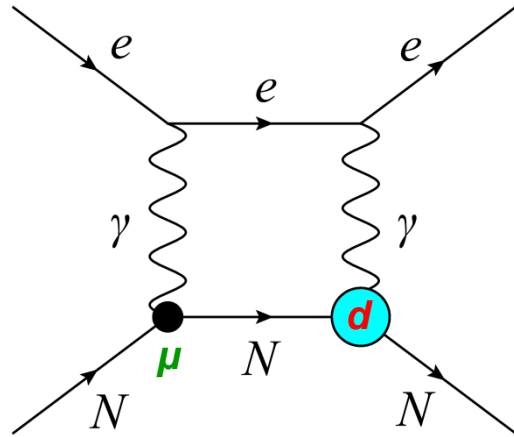


Continuum

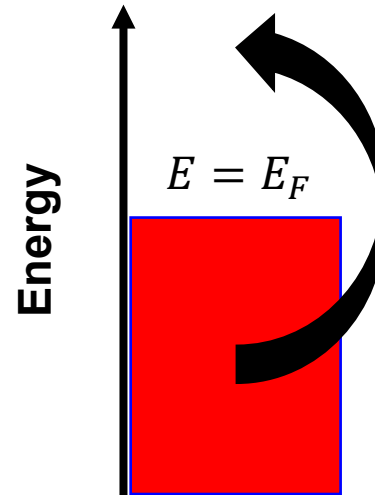
Excitations to continuum above Fermi surface: $\sim \ln(A)/p_F$ [Fermi-gas model]

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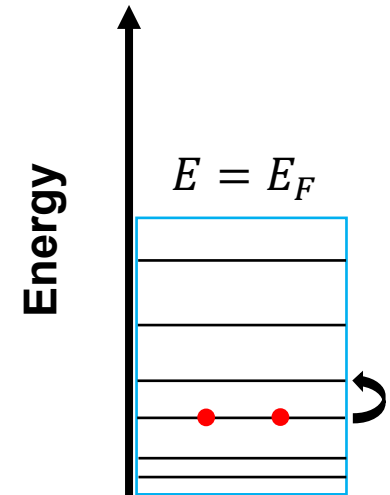
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Internal nuclear excitations



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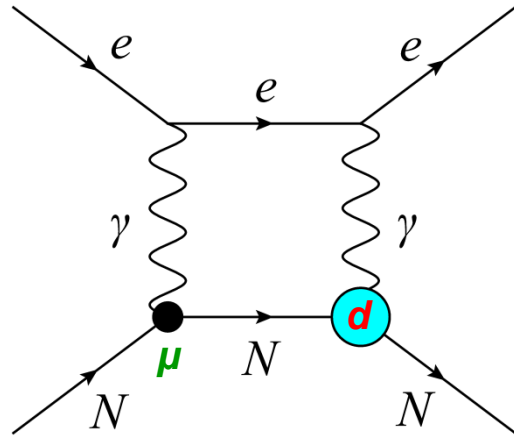
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Discrete transitions between L-S doublets: $\sim [\mathcal{O}(10)/A] \times (1/\Delta E_{\text{nucl}})$

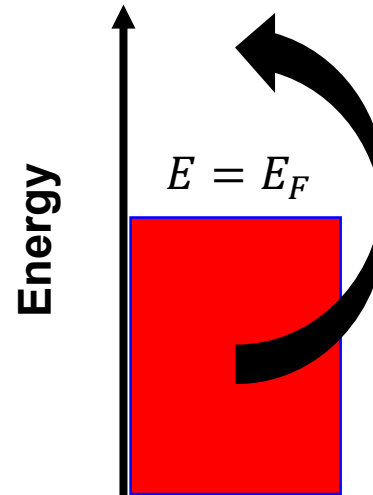
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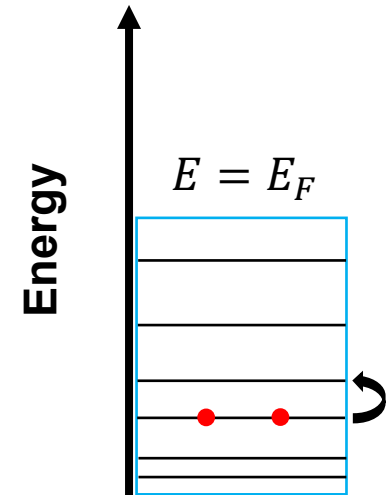
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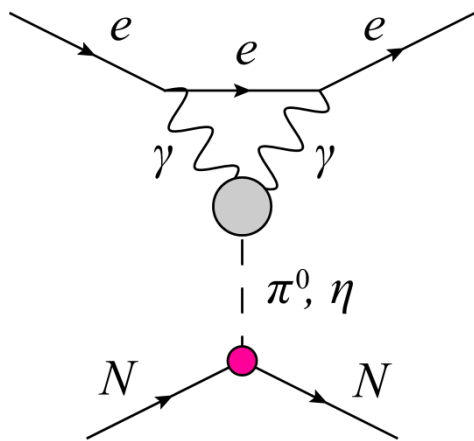
For $A \sim 200$ and $\Delta E_{\text{nucl}} \sim$ several MeV, the two contributions are comparable in size (and of the same sign)

QCD Vacuum Angle

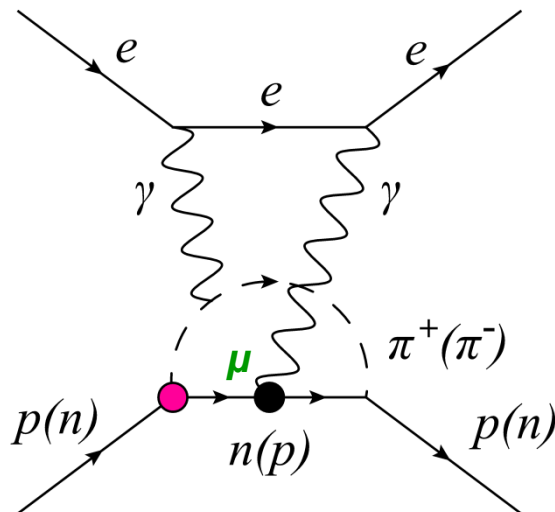
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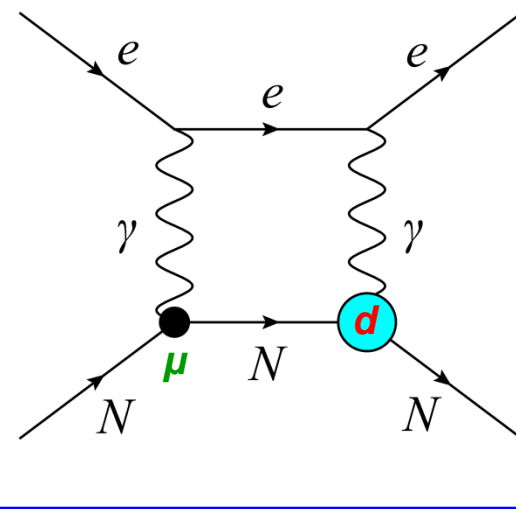


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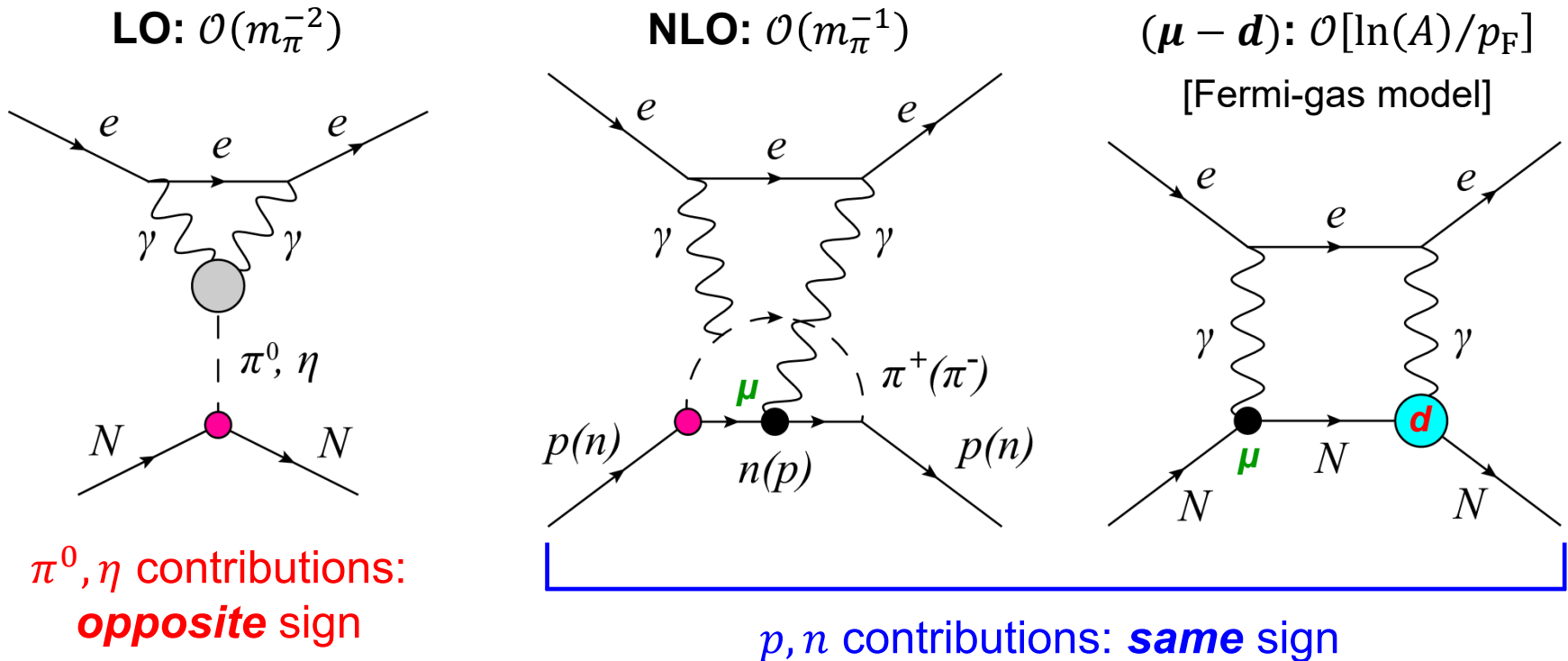
π^0, η contributions:
opposite sign
 (near-cancellation
 in heavy nuclei)

p, n contributions: **same** sign

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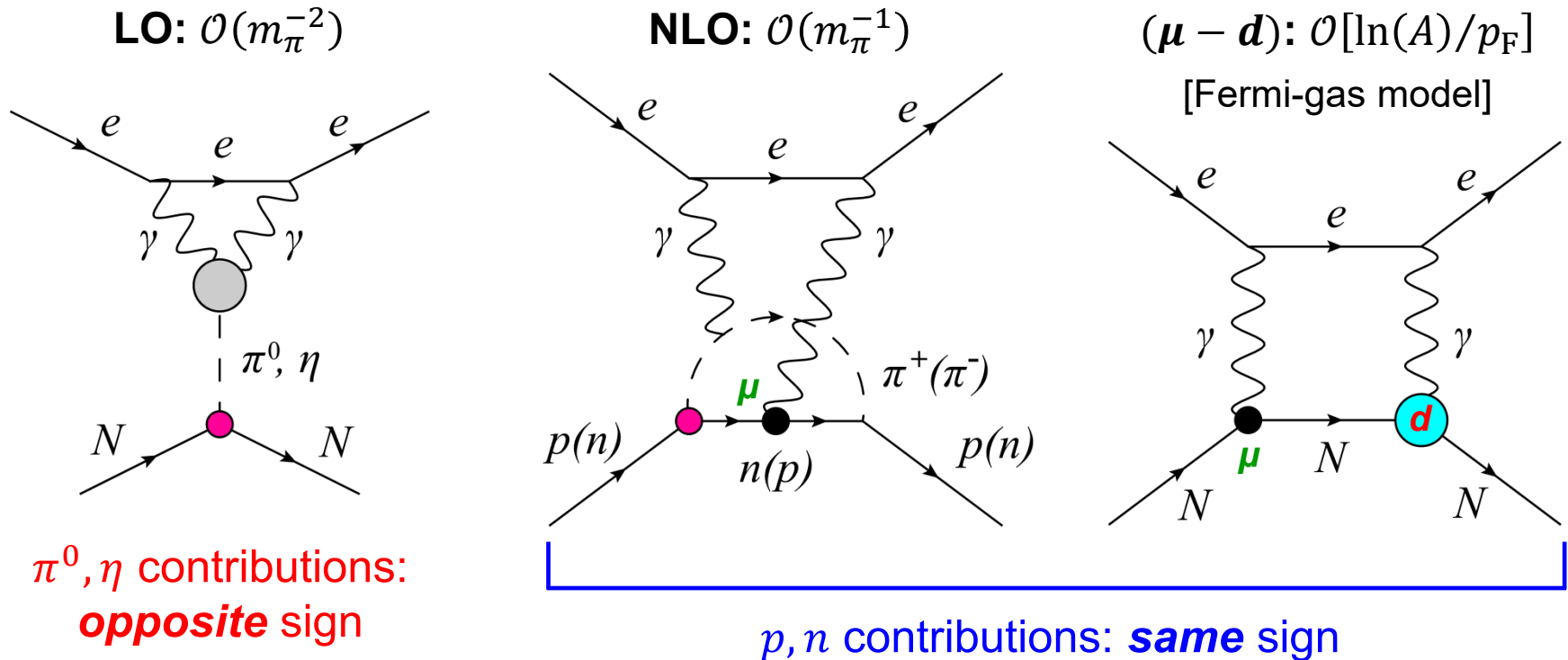
For $Z \sim 80$ & $A \sim 200$: $C_{\text{SP}}(\theta) \approx [0.1_{\text{LO}} + 1.0_{\text{NLO}} + 1.7_{(\mu-d)}] \times 10^{-2} \theta \approx 0.03\theta$

$$\mathcal{L}_{\text{contact}} = - \frac{G_F C_{\text{SP}} \bar{N} N \bar{e} i \gamma_5 e}{\sqrt{2}}$$

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Future work: η' contribution and other N²LO contributions, nuclear in-medium effects (NLO process), nuclear structure effects [$(\mu - d)$ process]

Bounds on Hadronic CP Violation Parameters

ThO bounds: [Flambaum, Pospelov, Ritz, Stadnik, *PRD* **102**, 035001 (2020)]

System	$ \bar{g}_{\pi NN}^{(1)} $	$ \tilde{d}_u - \tilde{d}_d $ (cm)	$ d_p $ (e cm)	$ \theta $
ThO	4×10^{-10}	2×10^{-24}	2×10^{-23}	3×10^{-8}
<i>n</i>	1.1×10^{-10}	5×10^{-25}	—	2.0×10^{-10}
Hg	1×10^{-12}	5×10^{-27}	2.0×10^{-25}	1.5×10^{-10}
Xe	6.7×10^{-8}	3×10^{-22}	3.2×10^{-22}	3.2×10^{-6}

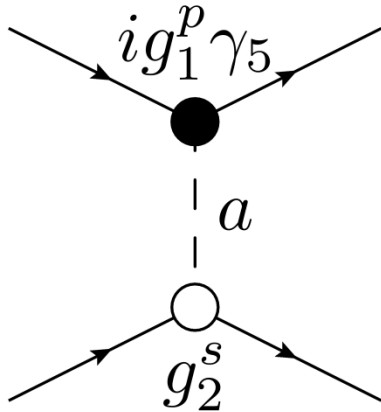
* These limits can formally be null within nuclear uncertainties

Current bounds from molecules are $\sim 10 - 100$ times weaker than from Hg & *n*, but are $\sim 10 - 100$ times stronger than bounds from Xe

P,T-Violating Forces Mediated by Dark Bosons

[Stadnik, Dzuba, Flambaum, *PRL* **120**, 013202 (2018)],

[Dzuba, Flambaum, Samsonov, Stadnik, *PRD* **98**, 035048 (2018)]



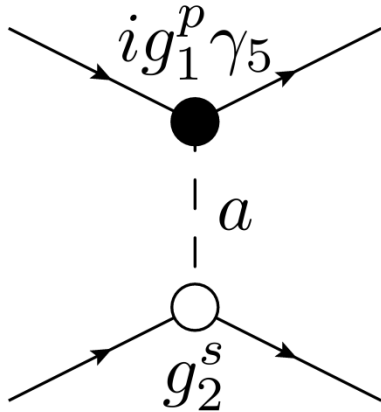
$$\mathcal{L}_{\text{int}} = a \bar{f} \left(g_f^s + i g_f^p \gamma_5 \right) f$$

$$V(r) \approx \frac{g_1^p g_2^s}{8\pi m_1} \boldsymbol{\sigma}_1 \cdot \hat{\mathbf{r}} \left(\frac{m_a}{r} + \frac{1}{r^2} \right) e^{-m_a r}$$

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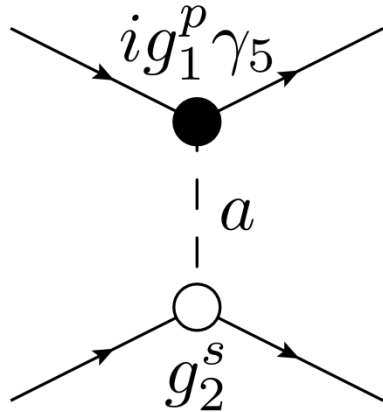
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P,T-violating forces => Atomic and Molecular EDMs

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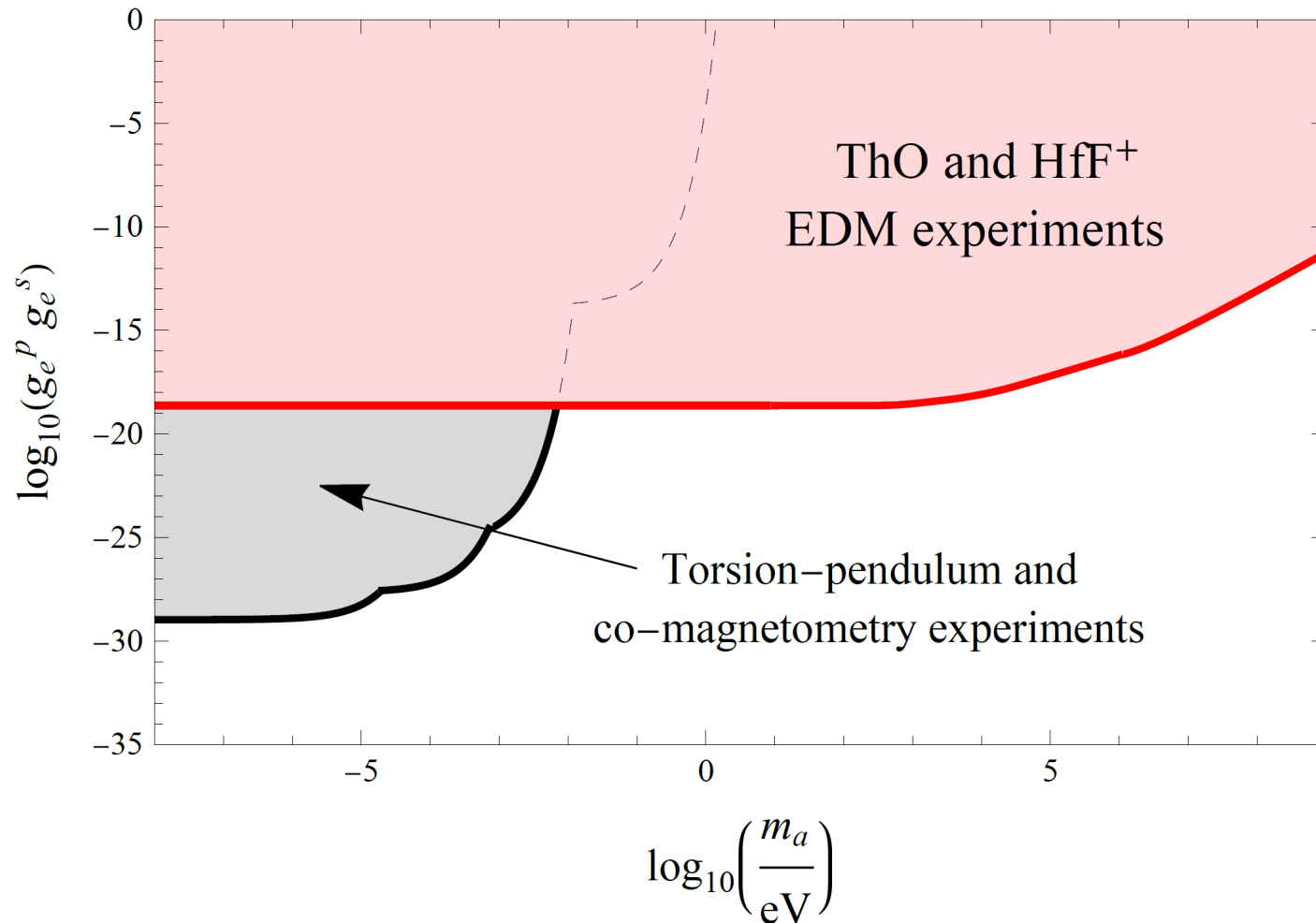
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P,T-violating forces => Atomic and Molecular EDMs

If exchanged boson is sufficiently low-mass, then P,T-violating forces are long-range on the scale of atom/molecule, and the non-vanishing contribution arises from the Thomas-Fermi length scale $r \sim a_B/Z^{1/3}$

Constraints on Scalar-Pseudoscalar Electron-Electron Interaction

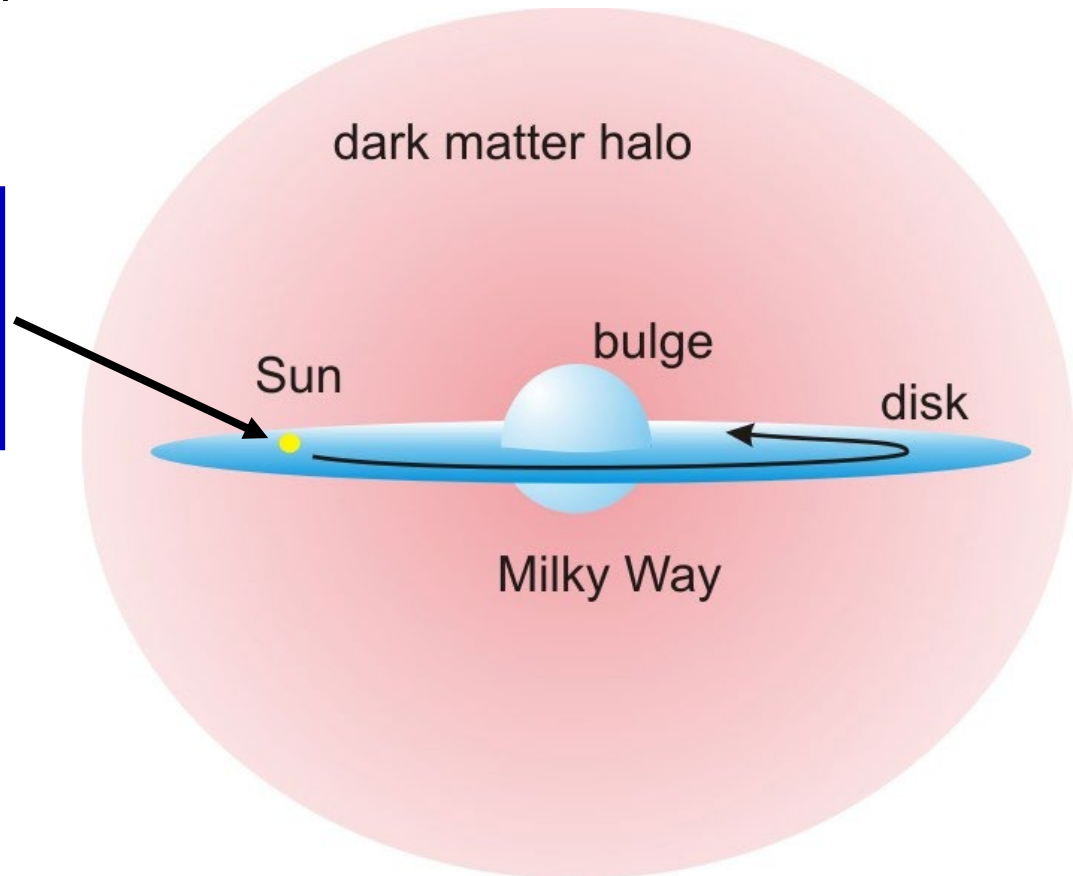
EDM constraints: [Stadnik, Dzuba, Flambaum, *PRL* **120**, 013202 (2018); arXiv:1708.00486]



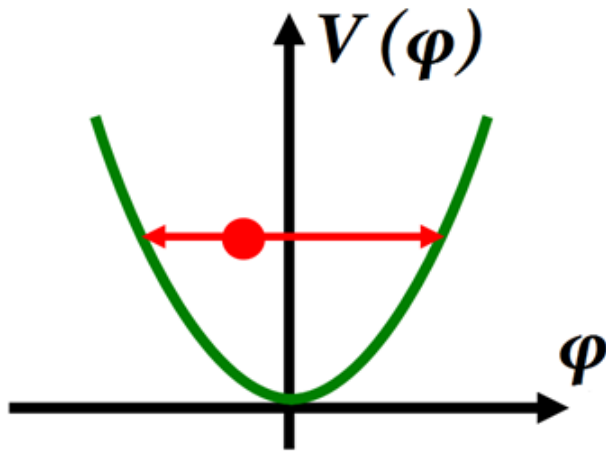
Dark Matter

Strong astrophysical evidence for existence of **dark matter** (~5 times more dark matter than ordinary matter)

$$\rho_{\text{DM}} \approx 0.4 \text{ GeV/cm}^3$$
$$v_{\text{DM}} \sim 300 \text{ km/s}$$



Oscillating Electric Dipole Moments



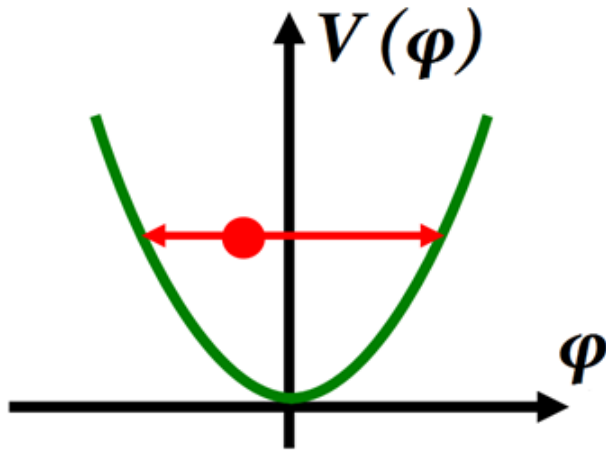
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$$\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$$

Oscillating Electric Dipole Moments

Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)]

Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)];
[Flambaum, Pospelov, Ritz, Stadnik, *PRD* **102**, 035001 (2020)]



$$\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar),$$

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$$\mathcal{L} = \frac{C_G g_s^2}{32\pi^2 f_a} \varphi_0 \cos(m_\varphi t) G \tilde{G} \Rightarrow$$

$$\mathbf{d}(t) \propto J \cos(m_\varphi t),$$

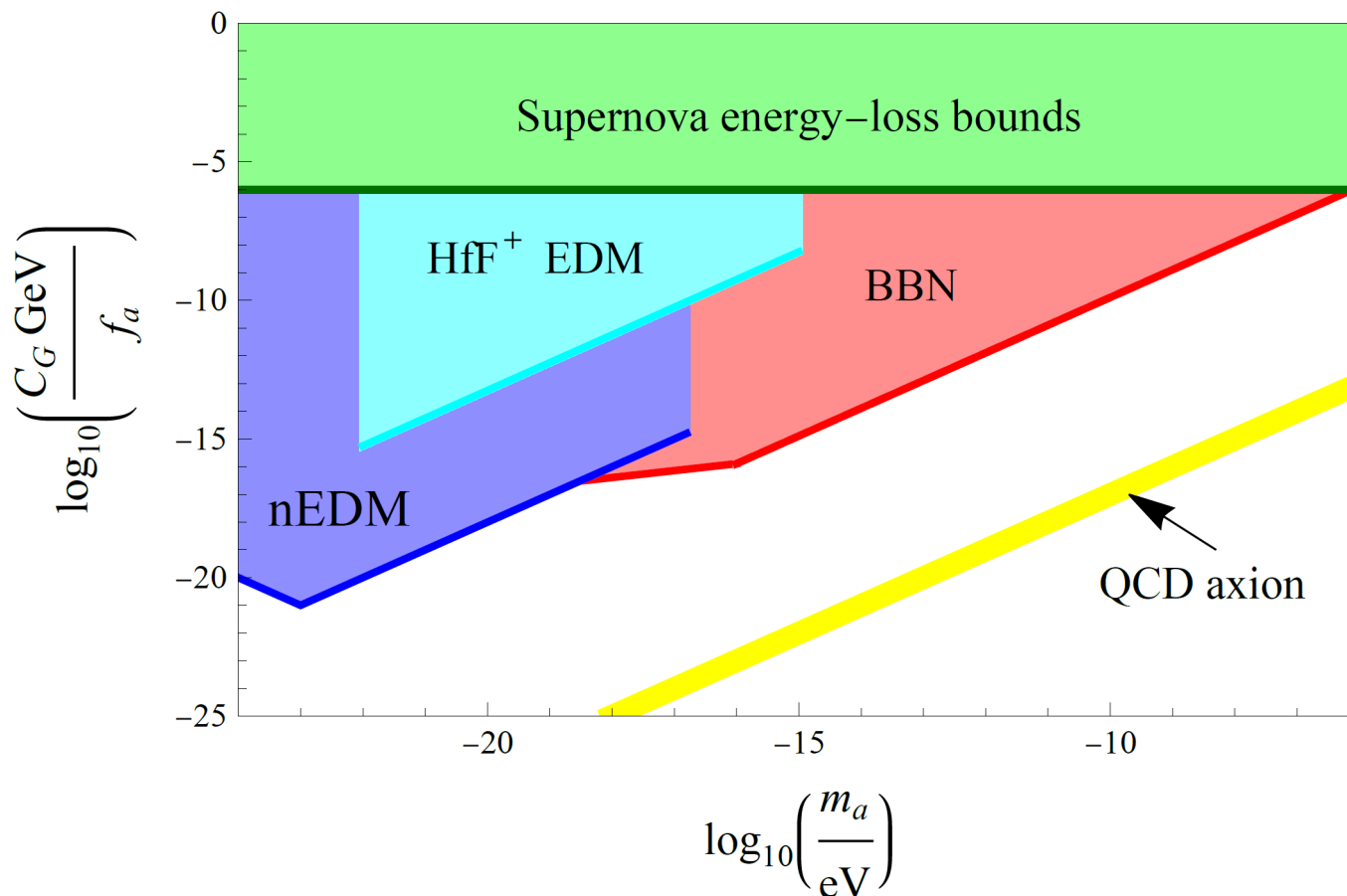
$$H_{\text{EDM}}(t) = \mathbf{d}(t) \cdot \mathbf{E}$$

$$\text{cf. } \mathcal{L} = \theta \frac{g_s^2}{32\pi^2} G \tilde{G} \Rightarrow \theta \leftrightarrow C_G \varphi_0 \cos(m_\varphi t) / f_a$$

Constraints on Interaction of Axion Dark Matter with Gluons

nEDM constraints: [nEDM collaboration, *PRX* **7**, 041034 (2017)]

HfF⁺ EDM constraints: [Roussy *et al.*, *PRL* **126**, 171301 (2021)]



Summary

- Plethora of opportunities for radioactive molecules in tests of fundamental symmetries and searches for new physics:
 - Paramagnetic molecules sensitive to hadronic sources of CP violation via two-photon-exchange processes (*regardless* of nucleus spin; i.e., operative for *spinless* nuclei, such as in ^{232}ThO , $^{180}\text{HfF}^+$)
 - Exchange of low-mass dark bosons within molecules can induce “long-range” P,T-violating forces, generating permanent EDMs
 - Low-mass bosonic dark matter can induce oscillating-in-time EDMs

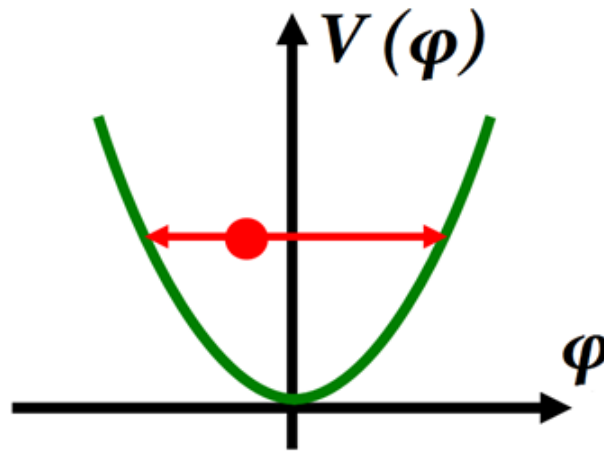
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 - Low-mass bosonic dark matter can induce oscillating-in-time EDMs
- Further work:
 - *For two-photon-exchange processes*: nuclear structure effects, nuclear in-medium effects, η' and other N²LO contributions
 - *For nucleon-spin-dependent phenomena*: improved knowledge of proton and neutron spin contributions in heavy nuclei
 - Improved calculations of Schiff moments of heavy nuclei

Back-Up Slides

Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)



$$V(\varphi) = \frac{m_\varphi^2 \varphi^2}{2}$$

$$\ddot{\varphi} + m_\varphi^2 \varphi \approx 0$$

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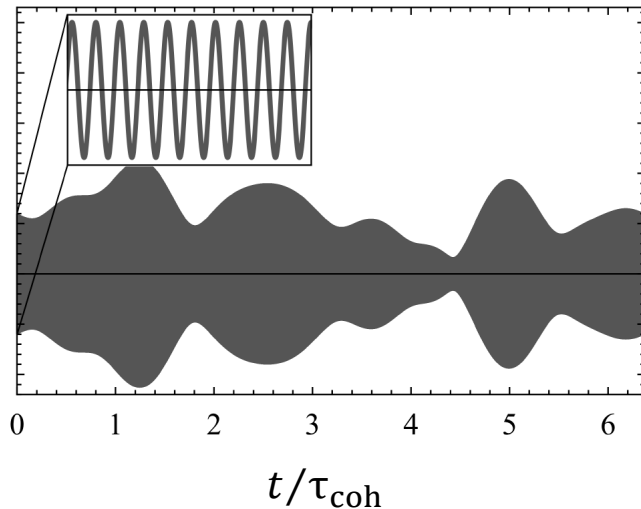
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\uparrow
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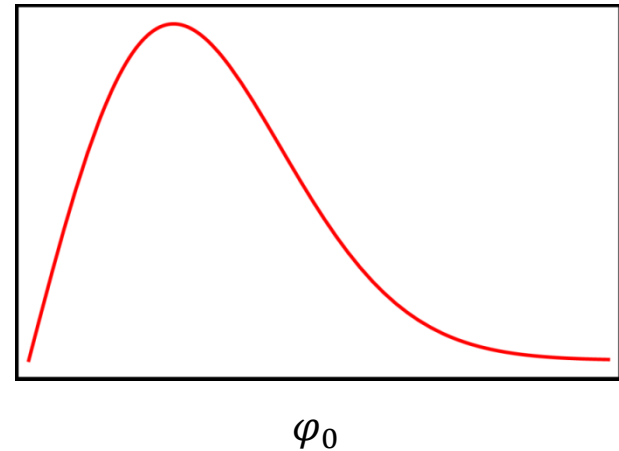
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Evolution of φ_0 with time



Probability distribution function of φ_0
(e.g., Rayleigh distribution)



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$T_{\text{osc}} \sim 1 \text{ month}$

IR frequencies

Lyman- α forest measurements [suppression of structures for $L \lesssim \mathcal{O}(\lambda_{\text{dB},\varphi})$]

[Related figure-of-merit: $\lambda_{\text{dB},\varphi} / 2\pi \leq L_{\text{dwarf galaxy}} \sim 100 \text{ pc} \Rightarrow m_\varphi \gtrsim 10^{-21} \text{ eV}$]

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- **Wave-like signatures** [cf. *particle-like* signatures of WIMP DM]