This notes accompany the discussion of the "generalized gauge transformation" given in Lecture 17, page 4. Begin with a spacetime whose metric we represent as some background $\hat{g}_{\alpha\beta}$ plus a perturbation $h_{\alpha\beta}$; these are all expressed in terms of the coordinate x^{γ} . We wish to change to a new coordinate y^{γ} which is related to x^{γ} by an infinitesimal shift ξ^{γ} :

$$x^{\alpha} = y^{\alpha} - \xi^{\alpha}(y^{\gamma}) \; .$$

This coordinate transformation changes the metric according to

$$g_{lphaeta}(y^{\gamma}) = g_{\gamma\delta}(x^{\gamma}) rac{\partial x^{\gamma}}{\partial y^{lpha}} rac{\partial x^{\delta}}{\partial y^{eta}}$$

Let's now expand this out:

$$g_{\alpha\beta}(y^{\gamma}) = (\hat{g}_{\gamma\delta} + h_{\gamma\delta}) \left(\delta^{\gamma}{}_{\alpha} - \partial_{\alpha}\xi^{\gamma}\right) \left(\delta^{\delta}{}_{\beta} - \partial_{\beta}\xi^{\delta}\right)$$

$$= \hat{g}_{\alpha\beta} + h_{\alpha\beta} - \hat{g}_{\gamma\beta}\partial_{\alpha}\xi^{\delta} - \hat{g}_{\alpha\delta}\partial_{\beta}\xi^{\gamma} + \mathcal{O}[h\,\partial\xi,\,(\partial\xi)^{2}]$$

$$= \hat{g}_{\alpha\beta} + h_{\alpha\beta} - \hat{g}_{\gamma\beta}(\hat{\nabla}_{\alpha}\xi^{\gamma} - \xi^{\delta}\hat{\Gamma}^{\gamma}_{\alpha\delta}) - \hat{g}_{\alpha\delta}(\hat{\nabla}_{\beta}\xi^{\delta} - \xi^{\gamma}\hat{\Gamma}^{\delta}_{\beta\gamma})$$

$$= \hat{g}_{\alpha\beta} + h_{\alpha\beta} - \hat{\nabla}_{\alpha}\xi_{\beta} - \hat{\nabla}_{\beta}\xi_{\alpha} + \xi^{\gamma}(\hat{\Gamma}_{\beta\alpha\gamma} + \hat{\Gamma}_{\alpha\beta\gamma})$$

$$= \hat{g}_{\alpha\beta} + h_{\alpha\beta} - \hat{\nabla}_{\alpha}\xi_{\beta} - \hat{\nabla}_{\beta}\xi_{\alpha} + \xi^{\gamma}\partial_{\gamma}\hat{g}_{\alpha\beta}.$$

On going from the fourth to the fifth line, we've used the definition of the Christoffel symbol. On that last line, the background in this expression is currently expressed as a function of the old coordinates, x^{γ} . We want it to appear as a function of the new coordinates, y^{γ} :

$$\hat{g}_{\alpha\beta}(x^{\gamma}) = \hat{g}_{\alpha\beta}(y^{\gamma} - \xi^{\gamma}) = \hat{g}_{\alpha\beta}(y^{\gamma}) - \xi^{\gamma} \partial_{\gamma} \hat{g}_{\alpha\beta}$$

Note, we could do this expansion for the perturbation as well: $h_{\alpha\beta}(x^{\gamma}) = h_{\alpha\beta}(y^{\gamma} - \xi^{\gamma}) = h_{\alpha\beta}(y^{\gamma}) - \xi^{\gamma}\partial_{\gamma}h_{\alpha\beta}$. This means that $h_{\alpha\beta}(x^{\gamma}) = h_{\alpha\beta}(y^{\gamma}) + \text{terms that are of order "small squared," which we neglect.$

Substituting the result for $\hat{g}_{\alpha\beta}(x^{\gamma})$ into $g_{\alpha\beta}(y^{\gamma})$ yields

$$g_{lphaeta}(y^{\gamma}) \;\;=\;\; \hat{g}_{lphaeta} + h_{lphaeta} - \hat{
abla}_{lpha} \xi_{eta} - \hat{
abla}_{eta} \xi_{lpha} \;,$$

from which we read off

$$h_{\alpha\beta} \to h_{\alpha\beta} - \hat{\nabla}_{\alpha}\xi_{\beta} - \hat{\nabla}_{\beta}\xi_{\alpha}$$