Wrap up of black hole orbits: Motion of light; orbits in Kerr

Going beyond symmetry.

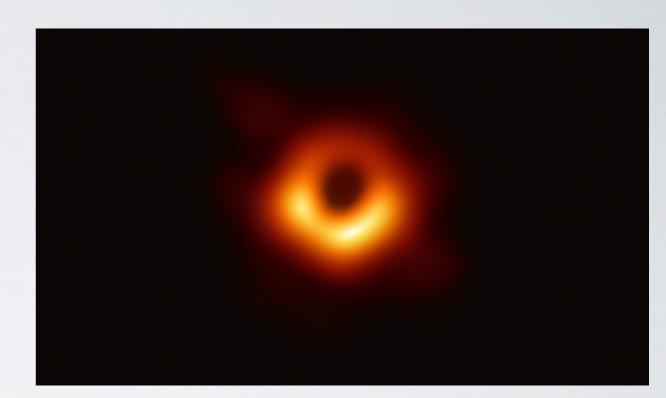


Image of radiation emitted by material flowing onto the black hole at the core of galaxy M87. Image produced using large-scale radio interferometry project known as the Event Horizon Telescope

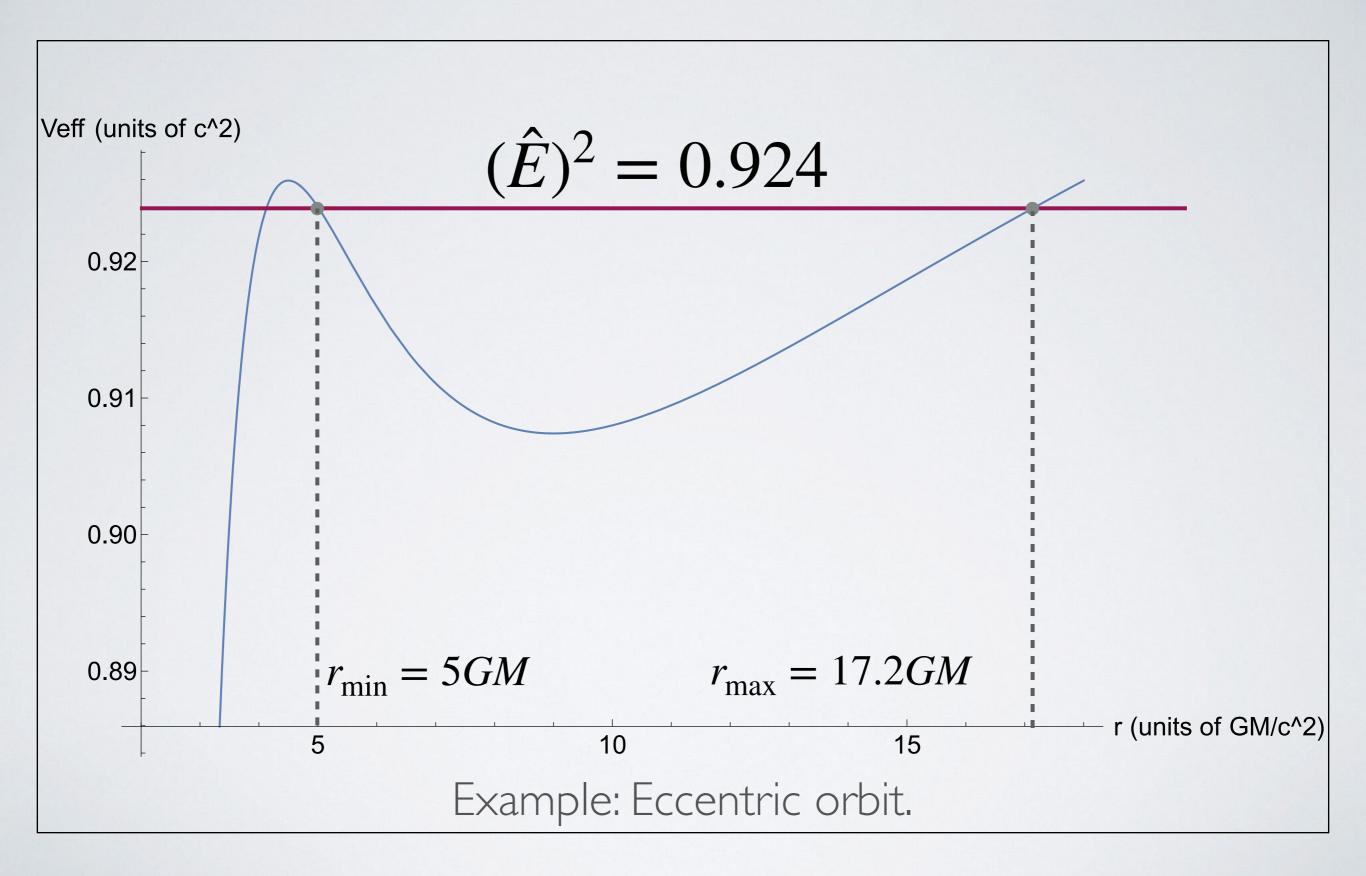
**Recap:** Examined motion of a massive body in Schwarzschild spacetime ... It is governed by the equations

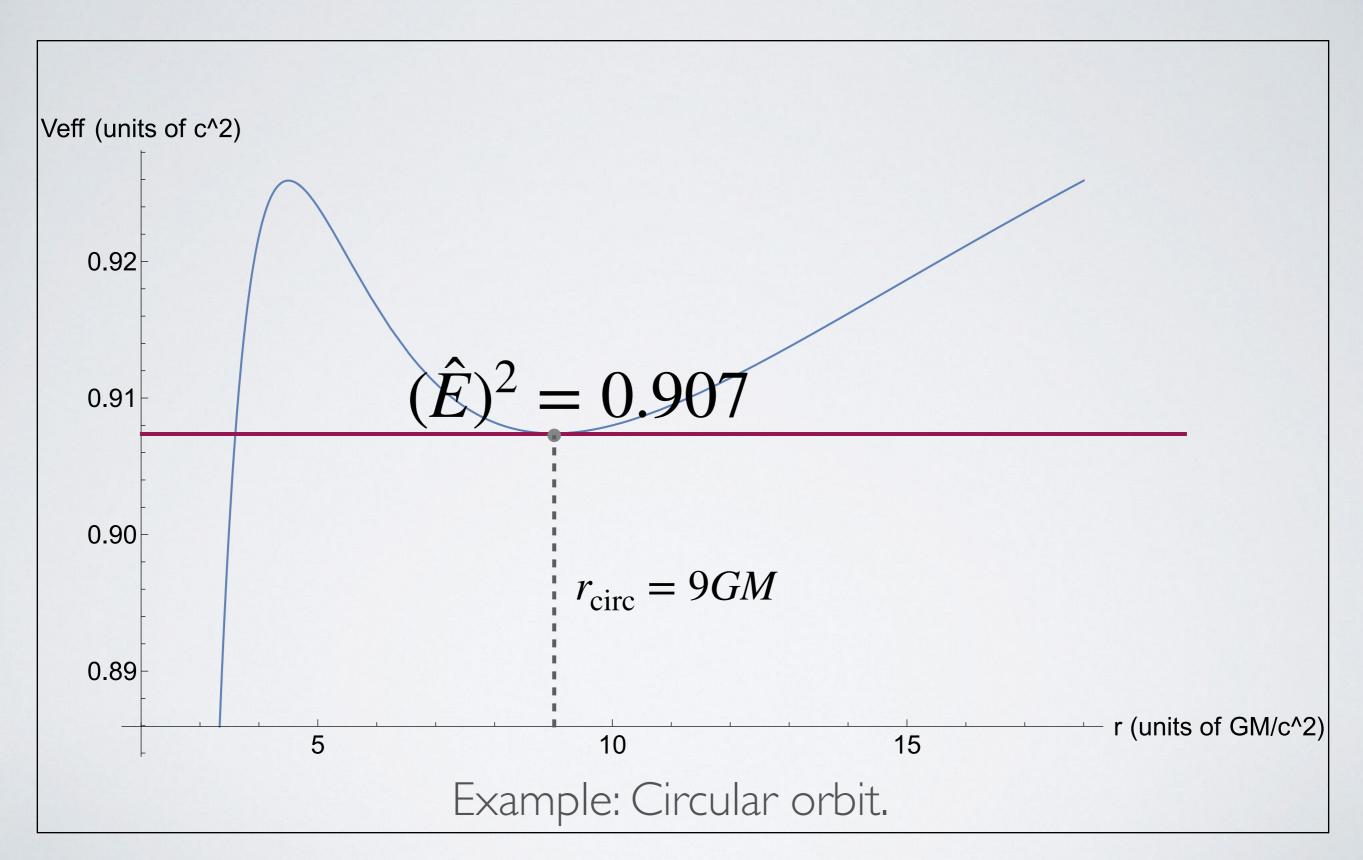
$$\frac{dt}{d\tau} = \frac{E}{1 - 2GM/r}$$

$$\frac{d\phi}{d\tau} = \hat{L}/r^2$$

$$\left(\frac{dr}{d\tau}\right)^2 = \hat{E}^2 - \left(1 - \frac{2GM}{r}\right)\left(1 + \frac{\hat{L}^2}{r^2}\right) \equiv \hat{E}^2 - V_{\text{eff}}(r)$$

Pick energy and angular momentum; characteristics of effective potential and energy determine character of the orbit.





Circular condition is that the orbit have  $dr/d\tau = 0$  and sit at the minimum of  $V_{\text{eff}}(r)$ ; imposing these conditions leads to the solutions

$$\hat{E} = \frac{1 - 2GM/r}{\sqrt{1 - 3GM/r}}$$

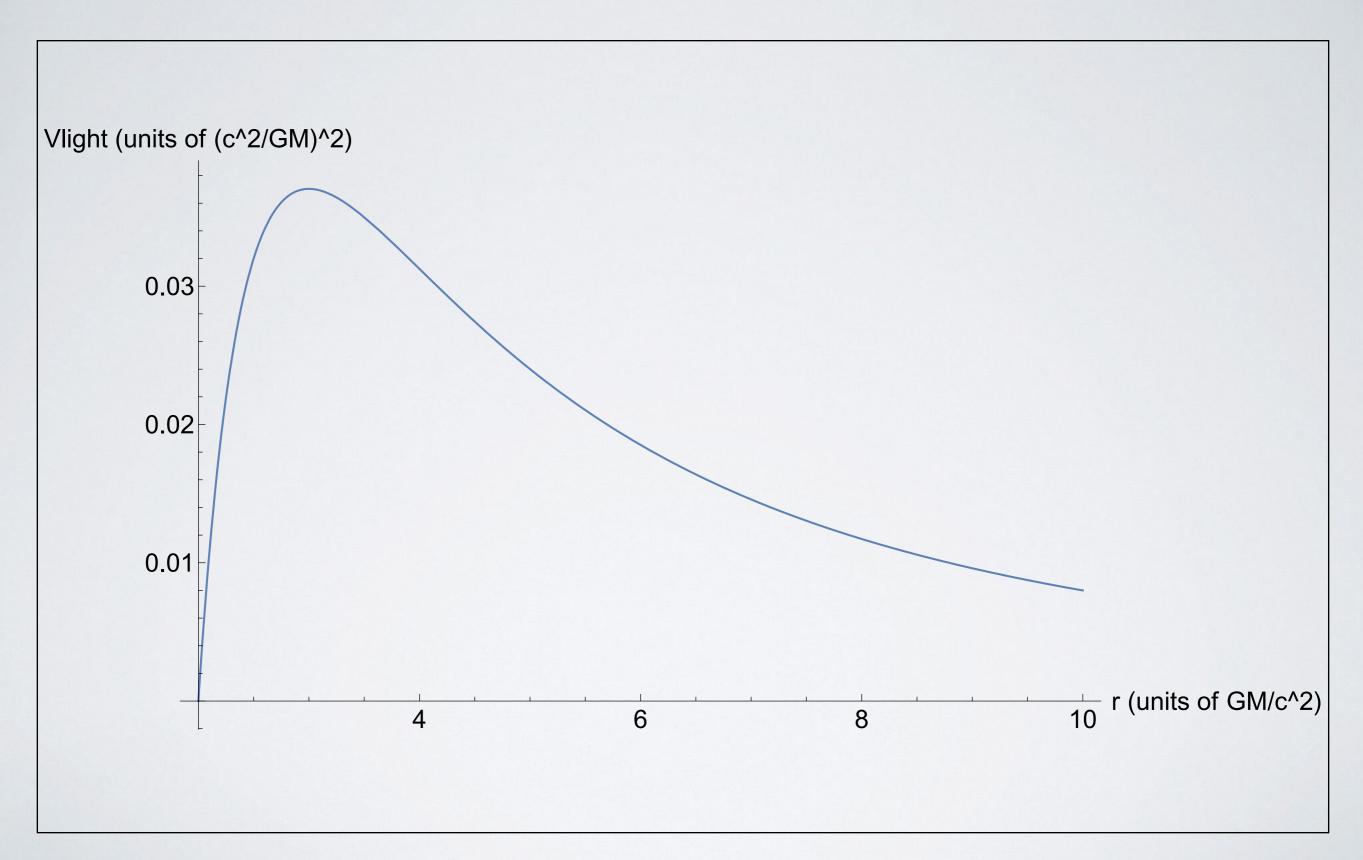
$$\hat{G}Mr$$

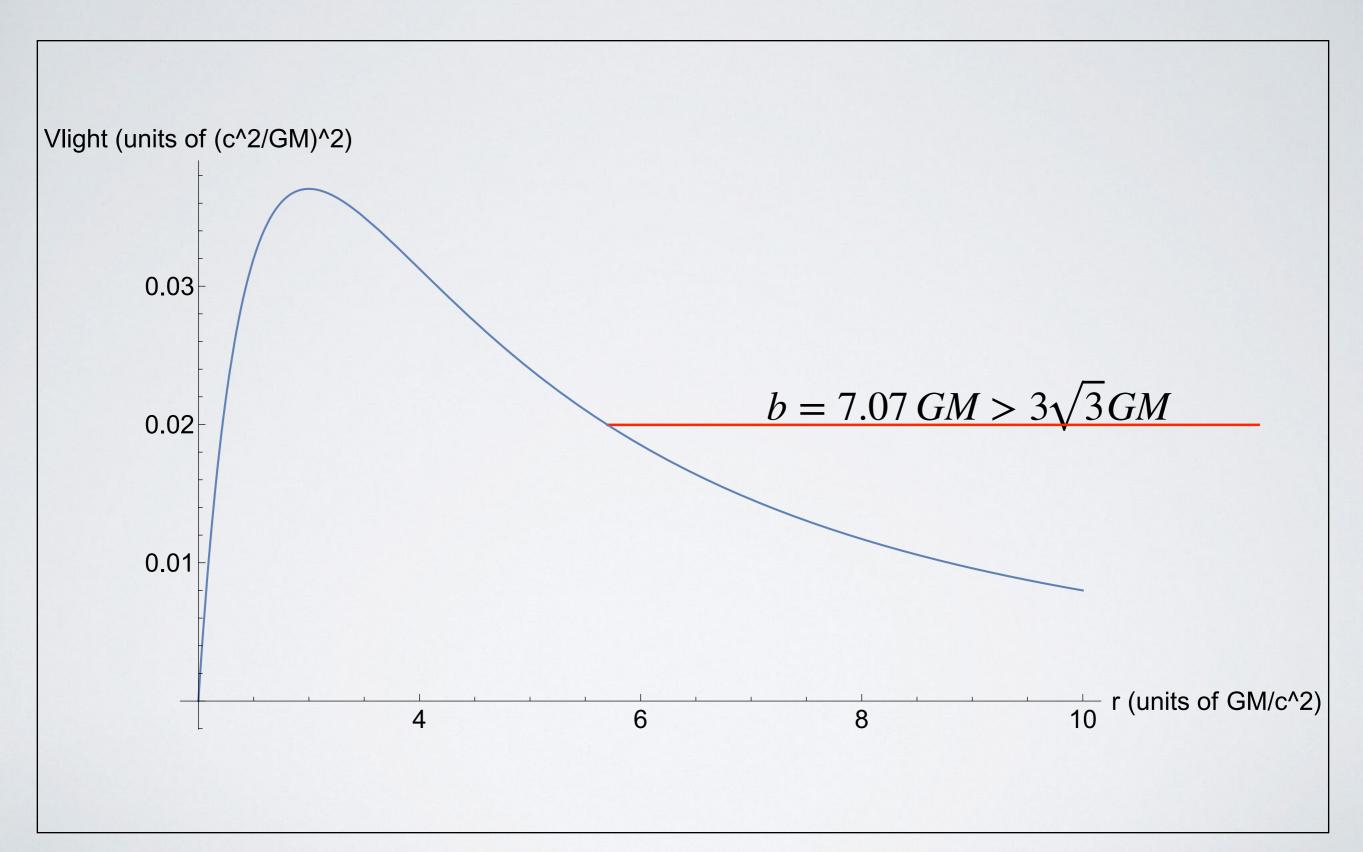
$$\hat{L} = \pm \sqrt{\frac{GMr}{1 - 3GM/r}}$$

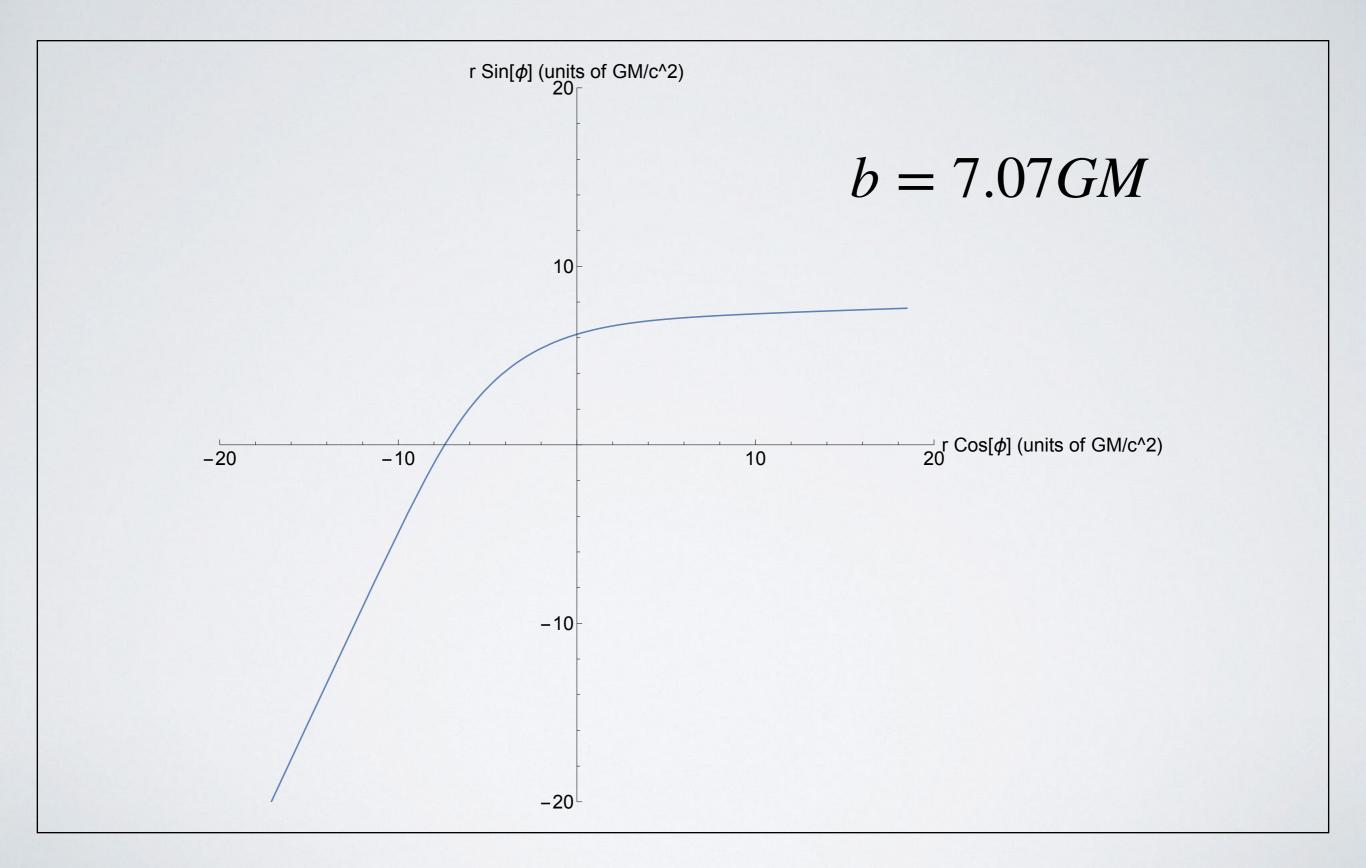
Require in addition that the minimum be stable:

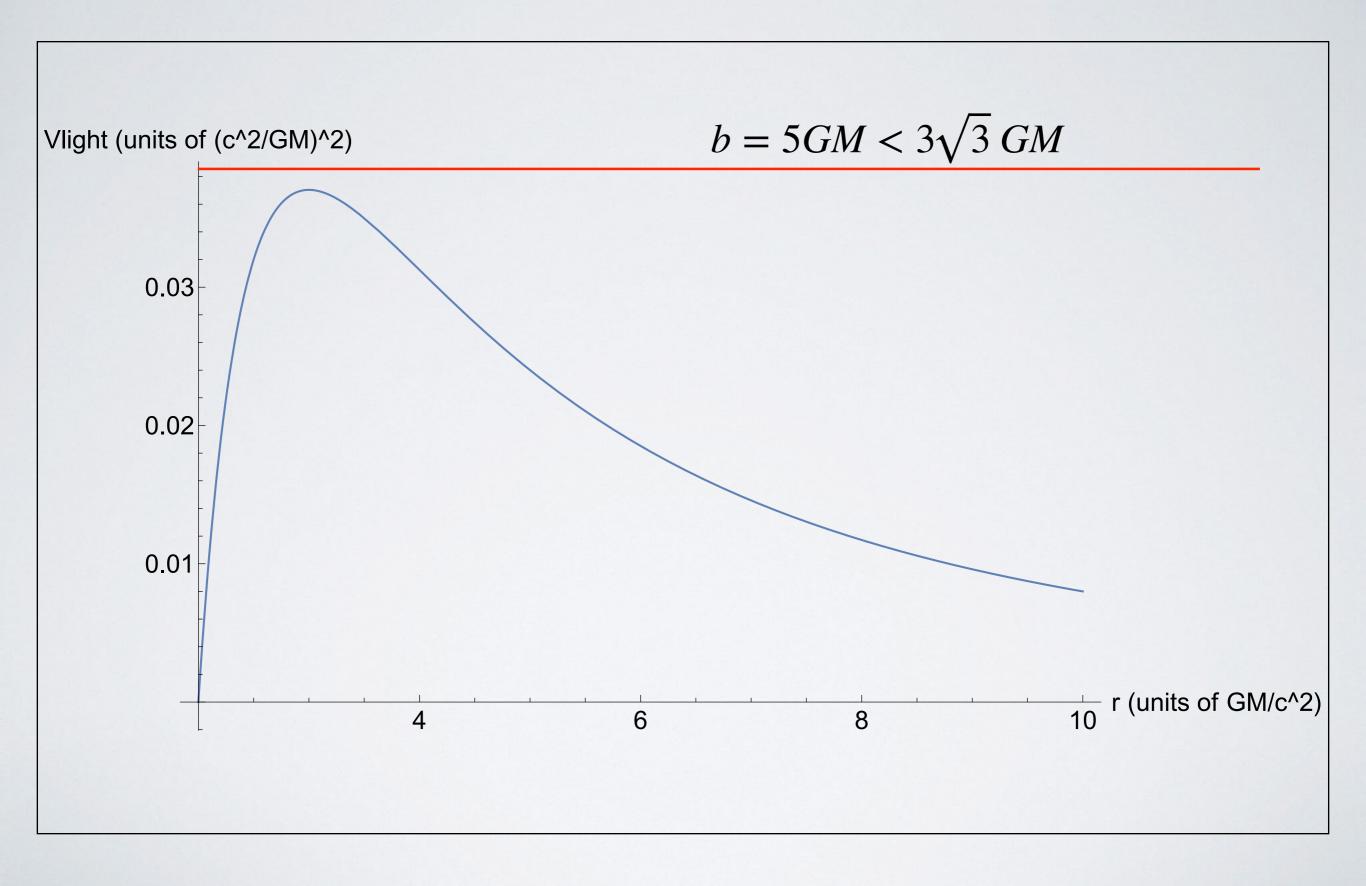
$$\frac{\partial^2 V_{\text{eff}}}{\partial r^2} > 0 \longrightarrow r > 6GM.$$

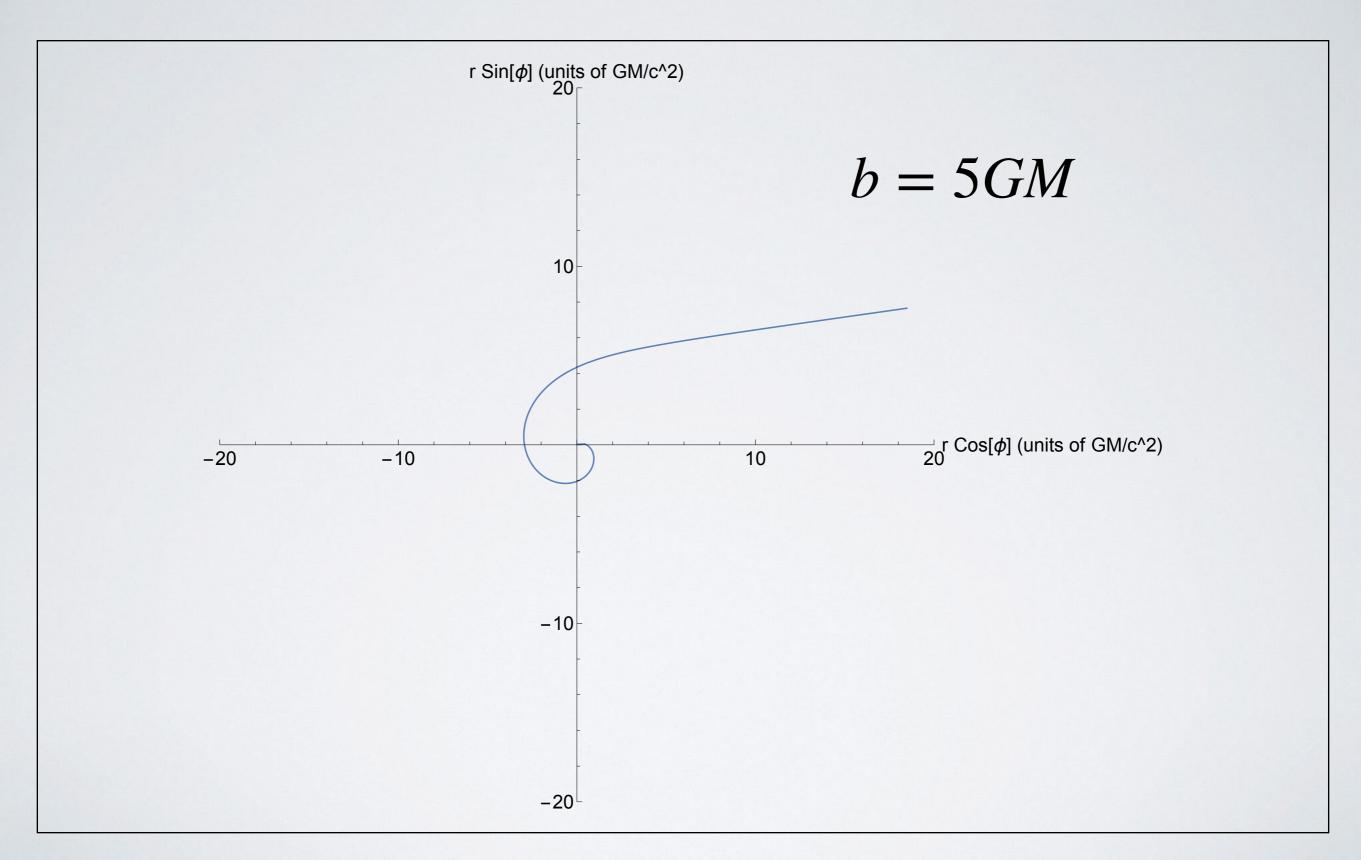
Leads to a very non-Newtonian characteristic: an "innermost stable orbit." No stable circular orbits exist inside the radius r = 6GM.

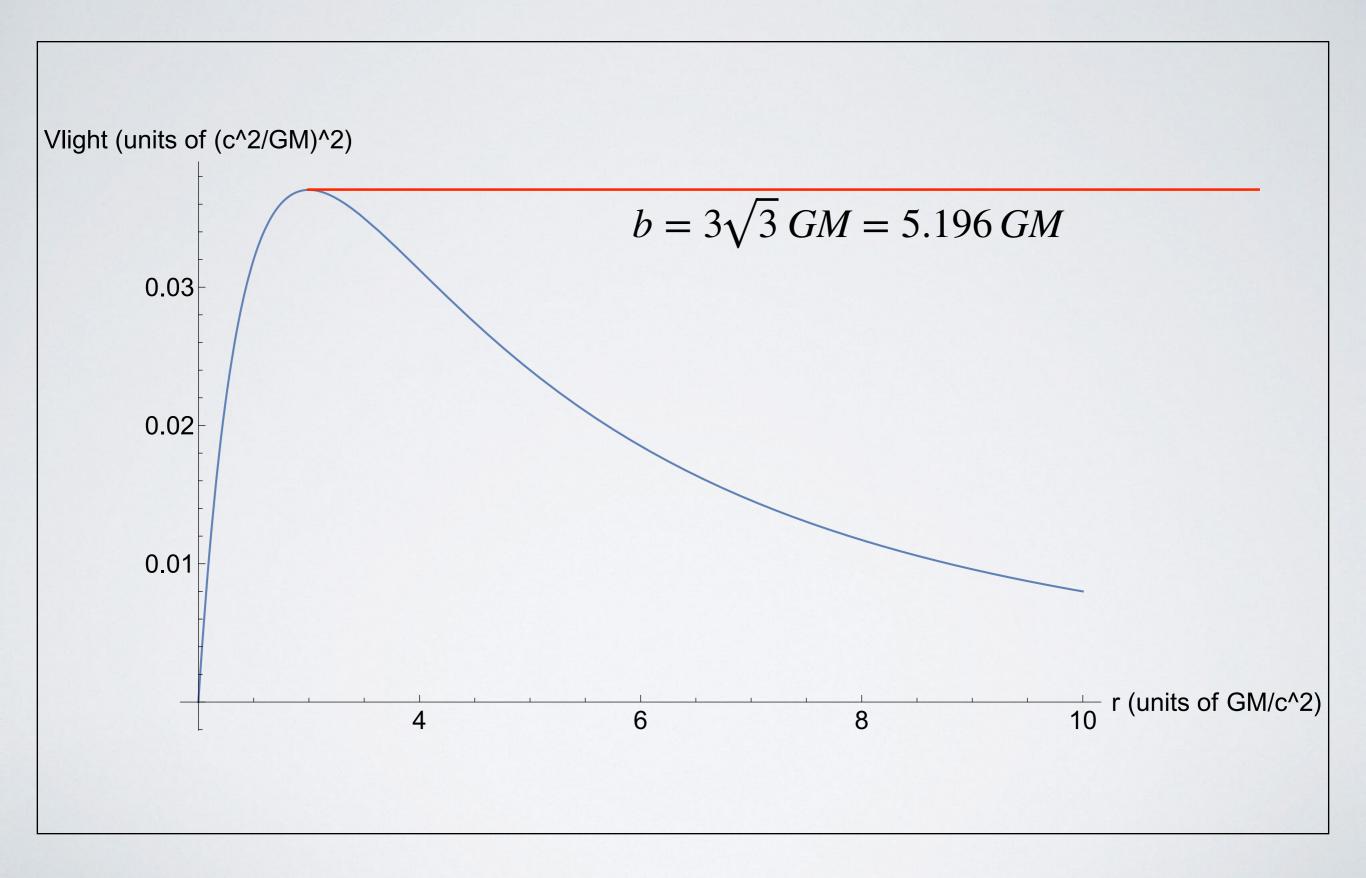


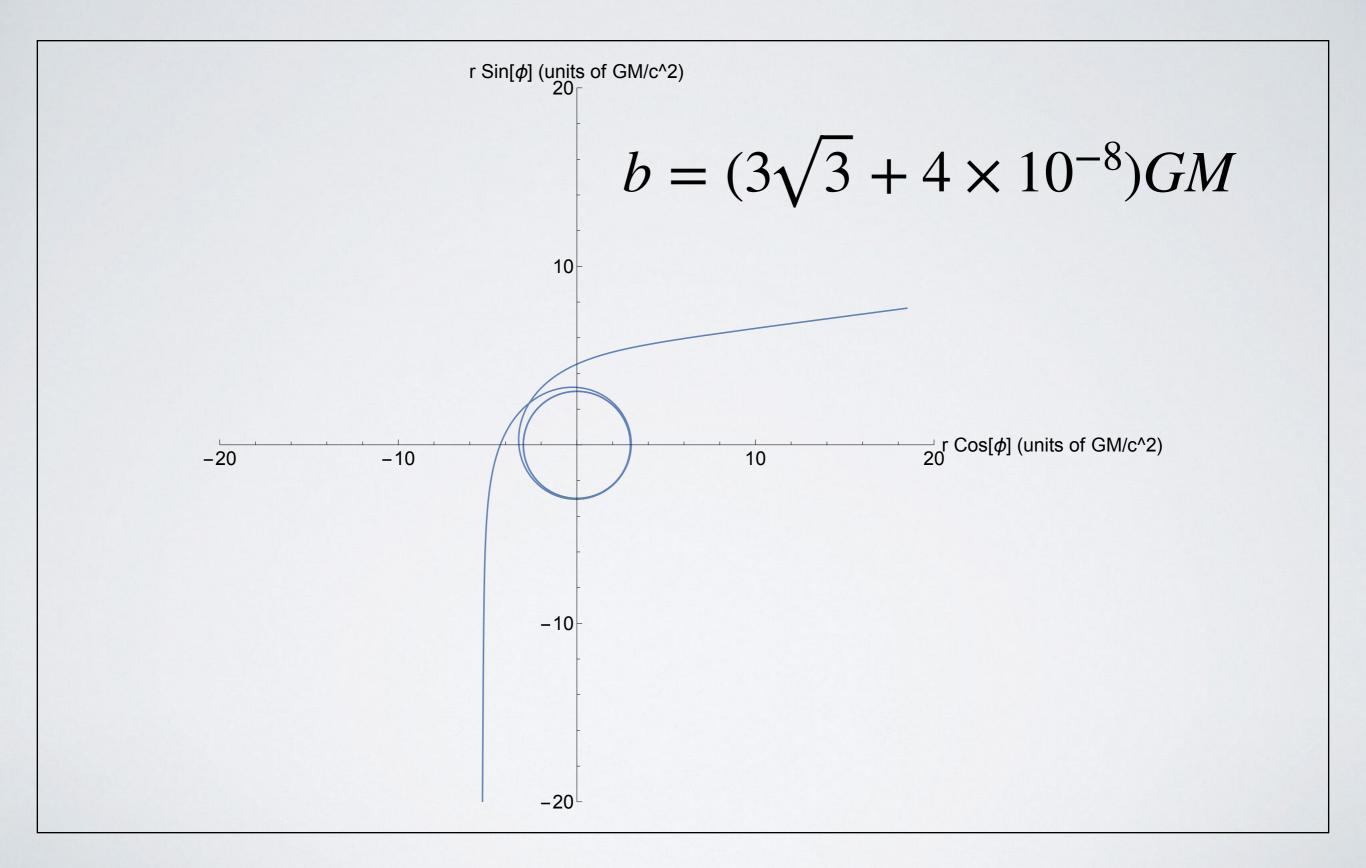














Radiation emitted by material accreting onto two famous BHs

Kerr motion: Conserved constants make it possible to separate the equations of motion. Using  $\Sigma = r^2 + a^2 \cos^2 \theta$ ,  $\Delta = r^2 - 2GMr + a^2$ , plus the definition  $a = |\mathbf{S}|/GM$ , we have

$$\Sigma^2 \left(\frac{dr}{d\tau}\right)^2 = \left(\hat{E}(r^2 + a^2) - a\hat{L}_z\right)^2 - \Delta\left(r^2 + (\hat{L}_z - a\hat{E})^2 + \hat{Q}\right) \equiv R(r)$$

$$\Sigma^{2} \left(\frac{d\theta}{d\tau}\right)^{2} = \hat{Q} - \hat{L}_{z}^{2} \cot^{2}\theta - a^{2} \cos^{2}\theta (1 - \hat{E}^{2}) \equiv \Theta(\theta)$$

$$\Sigma \left(\frac{d\phi}{d\tau}\right) = \csc^2\theta \hat{L}_z + \frac{a}{\Delta}(2Mr\hat{E} - a\hat{L}_z)$$

$$\Sigma \left(\frac{dt}{d\tau}\right) = \hat{E} \left[ \frac{(r^2 + a^2)^2}{\Delta} - a^2 \sin^2 \theta \right] - \frac{2Mra}{\Delta} \hat{L}_z$$

Pick energy, axial angular momentum, "Carter constant" Q: determines the properties of bound orbits.

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Useful subset: Q = 0 requires  $\theta = \pi/2$ , so these are equatorial orbits. If we further set R = 0, dR/dr = 0, we have circular equatorial orbits.

Stable circular equatorial orbits further have  $d^2R/dr^2 > 0$ ; solve for the system (R = 0, dR/dr = 0,  $d^2R/dr^2 = 0$ ), generalize the innermost stable circular orbit to Kerr. Result:

$$r_{\rm ISCO}/GM = 3 + Z_2 \mp \left[ (3 - Z_1)(3 + Z_1 + 2Z_2) \right]^{1/2}$$

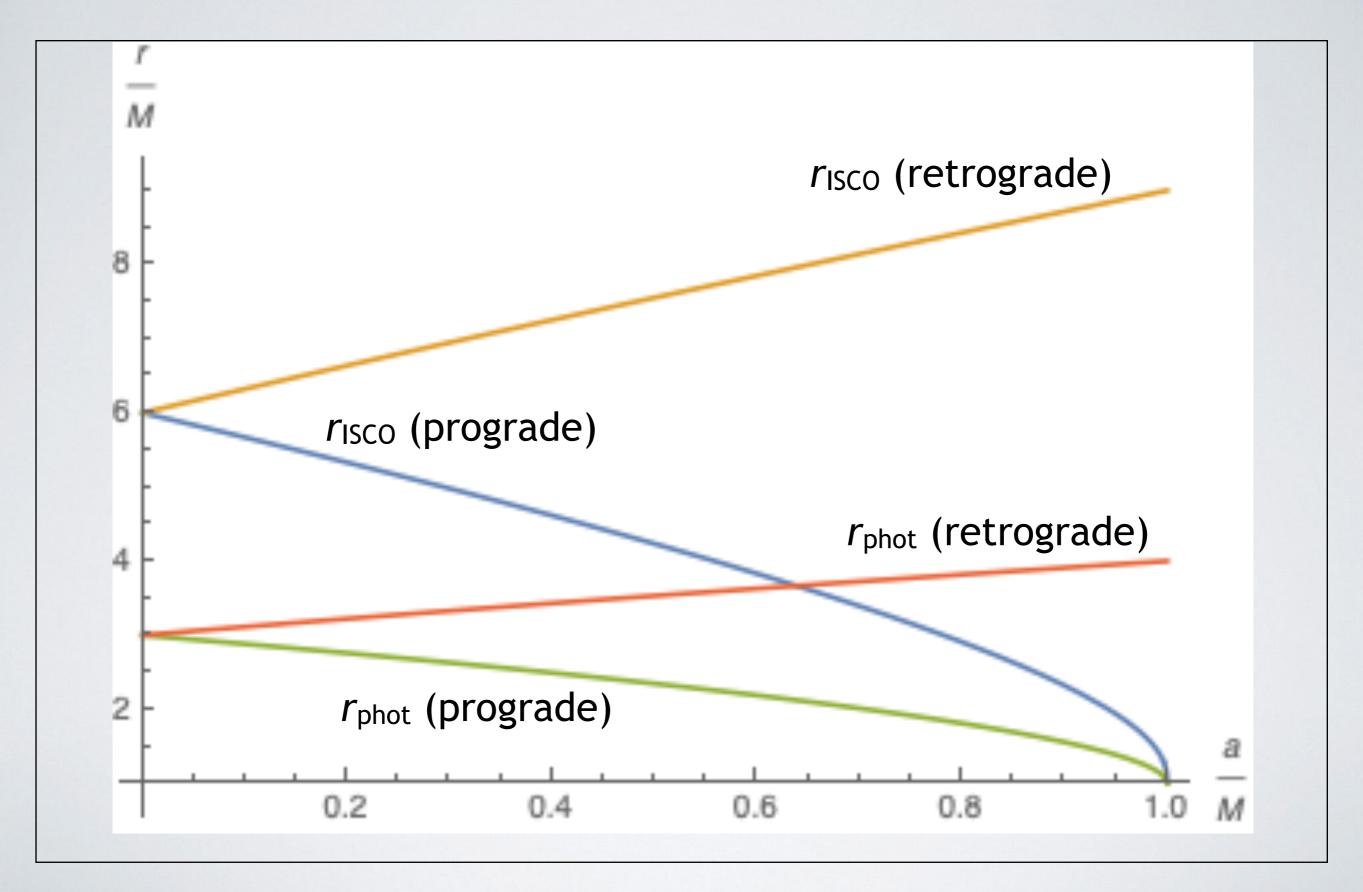
$$Z_1 = 1 + \left( 1 - q^2 \right)^{1/3} \left[ (1 + q)^{1/3} + (1 - q)^{1/3} \right]$$

$$Z_2 = \left( 3q^2 + Z_1^2 \right)^{1/2} \qquad (q \equiv a/GM)$$

Do this exercise for a light-like trajectory:

$$r_{\text{phot}}/GM = 2 + 2\cos\left[\frac{2}{3}\cos^{-1}(\mp q)\right]$$

Reference: Bardeen, Press, Teukolsky, Astrophys. J. 178, 347 (1972).



Quantity which appears in the "iterative" form of the field equations that we use to define post-Newtonian theory:

$$\begin{split} \Lambda^{\alpha\beta} &= -h^{\mu\nu}\partial^2_{\mu\nu}h^{\alpha\beta} + \partial_\mu h^{\alpha\nu}\partial_\nu h^{\beta\mu} + \frac{1}{2}g^{\alpha\beta}g_{\mu\nu}\partial_\lambda h^{\mu\tau}\partial_\tau h^{\nu\lambda} \\ &- g^{\alpha\mu}g_{\nu\tau}\partial_\lambda h^{\beta\tau}\partial_\mu h^{\nu\lambda} - g^{\beta\mu}g_{\nu\tau}\partial_\lambda h^{\alpha\tau}\partial_\mu h^{\nu\lambda} + g_{\mu\nu}g^{\lambda\tau}\partial_\lambda h^{\alpha\mu}\partial_\tau h^{\beta\nu} \\ &+ \frac{1}{8}\big(2g^{\alpha\mu}g^{\beta\nu} - g^{\alpha\beta}g^{\mu\nu}\big)\big(2g_{\lambda\tau}g_{\epsilon\pi} - g_{\tau\epsilon}g_{\lambda\pi}\big)\partial_\mu h^{\lambda\pi}\partial_\nu h^{\tau\epsilon} \,. \end{split}$$